

Effects of Water Deuteration on Thermodynamic and Structural Properties of Proteins and Biomembranes

Carmelo Tempa, Victor Cruces Chamorro, and Pavel Jungwirth*



Cite This: *J. Phys. Chem. B* 2023, 127, 1138–1143



Read Online

ACCESS |



Metrics & More

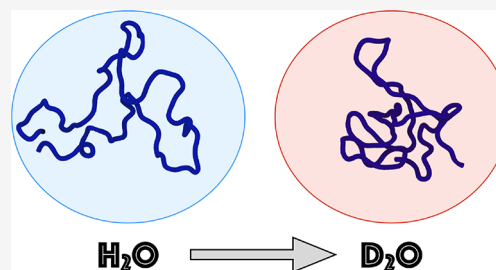


Article Recommendations



Supporting Information

ABSTRACT: Light and heavy water are often used interchangeably in spectroscopic experiments with the tacit assumption that the structure of the investigated biomolecule does not depend too much on employing one or the other solvent. While this may often be a good approximation, we demonstrate here using molecular dynamics simulations incorporating nuclear quantum effects via modification of the interaction potential that there are small but significant differences. Namely, as quantified and discussed in the present study, both proteins and biomembranes tend to be slightly more compact and rigid in D_2O than in H_2O , which reflects the stronger hydrogen bonding in the former solvent.



INTRODUCTION

Properties of light (H_2O) and heavy (D_2O) water are very similar to each other, save for a trivial $\sim 10\%$ difference in density due to the higher mass of the D over the H isotope.¹ This similarity is the rationale behind using the two water isotopes interchangeably as biomolecular solvents in spectroscopic experiments.^{2,3} Nevertheless, concerning nontrivial differences the two solvents vary by several degrees in melting points, by 0.4 pH (or pD) units in the autoionization equilibrium constant, and by about 20% in viscosity. Also, the number density (i.e., the number of molecules per unit volume) of D_2O is not equal to, but is actually lower than, that of H_2O . These variations in turn translate to a slightly different behavior of biomolecules dissolved in H_2O vs in D_2O . In particular, soluble proteins tend to be somewhat more compact and rigid in heavy water and so do phospholipid bilayers.^{4,5} As we have discussed recently, an intriguing consequence of nuclear quantum effects in water is the observed sweet taste of heavy water, as contrasted to a taste-neutral light water.⁶ More precisely, it is the reduction of nuclear quantum effects upon moving from H_2O to D_2O that triggers the activation of the human sweet taste receptor.

Small differences between light and heavy water can be related to slightly stronger hydrogen bonds in the latter liquid. These in turn can be traced back to nuclear quantum effects. Namely, zero point motions along as well as perpendicular to the direction of the water–water hydrogen bond are more pronounced in H_2O over in D_2O with a net effect of a slight hydrogen bond destabilization.⁷ Rigorously, computationally demanding quantum simulations such as path integral molecular dynamics (PIMD) should be employed to recover these effects. While feasible for neat water, such simulations become prohibitively expensive when large biomolecules are added to the solution. However, as already demonstrated by

Feynmann and Hibbs, zero point energy effects can effectively be incorporated into classical simulations by modifying the interaction potential.^{8,9} We have recently employed this approach to develop, based on an earlier model,¹⁰ a classical force field for heavy water. Here, we use this approach to quantify the differences in thermodynamic and structural properties of amino acids, proteins, and phospholipid membranes in light vs heavy water, comparing the simulation result to experiment whenever possible and providing a molecular interpretation of the observed phenomena. Focusing on differences between bulk properties of light vs heavy water, secondary effects of deuteration of exchangeable hydrogens of the biomolecules in heavy water have been neglected in the present study.

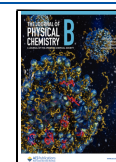
METHODS

Force Fields. The force fields used here to simulate H_2O and D_2O are the commonly used SPC/E model¹¹ for the former and our recently developed SPCE-HW parametrization¹⁰ for the latter. Amino acids, proteins, lipids, and ions were modeled using the CHARMM36 topology generated by the CHARMM-GUI web interface.^{12,13} Classical equations of motion were solved numerically with a 2 fs integration time step using the Verlet-list algorithm.¹⁴ Long-range electrostatic interactions were accounted for using the Particle Mesh Ewald scheme^{15,16} employing a short-range cutoff of 1.2 nm. For the

Received: November 25, 2022

Revised: January 23, 2023

Published: February 1, 2023



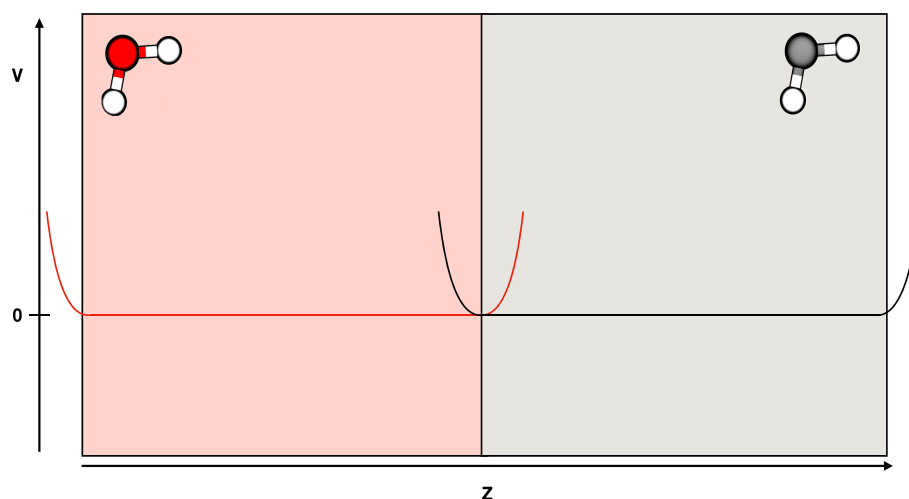


Figure 1. Simplified description of the setup used for calculating the free energy of transfer from H_2O to D_2O . Red (gray) color indicates D_2O (H_2O). The flat-bottomed potentials are also indicated in the figure.

van der Waals interaction, a force-switching algorithm from 1.0 to 1.2 nm was employed. Simulations were run in the isothermal–isobaric (NpT) ensemble with the velocity-rescale thermostat¹⁷ and the Parrinello–Rahaman barostat¹⁸ imposing a temperature of 298 K and a pressure of 1 atm, with coupling constants of 5 and 1 ps, respectively.

Free Energy of Amino Acid Transfer. A box $6 \times 6 \times 12$ nm^3 unit cell was filled with adjacent equally sized slabs of light (SPCE) and heavy (SPCE-HW) water, each containing 7203 water molecules. A flat-bottomed potential along the long (z) axis (see Figure 1) with a force constant of 100 kJ/mol and a distance from the center of each slabs R_i of 3 nm was applied to each of the cubes to keep the light and heavy water molecules separated from each other.

The system was energy minimized and then equilibrated in the NpT ensemble for 10 ns. All the essential amino acids were one by one placed in the center of the SPCE water slab and energy minimized. The umbrella sampling technique was then applied to compute the Potential of Mean Force (PMF) along the z axis (moving from the H_2O slab to D_2O), and the free energy of transfer is extracted as the difference of the PMF at the bulk of the two solvents. Thirty windows were generated along the z axis, each separated by 0.2 nm, and a force constant of $1000 \text{ kJ mol}^{-1} \text{ nm}^{-2}$ was applied in each window. Free energies in the individual windows were connected using the weighted histogram analysis method (WHAM), with the associated statistical error evaluated using the bootstrap method.^{19,20}

Protein and Phospholipid Membrane Simulations.

Three representative globular proteins have been chosen for the present study: azurine,²¹ lactoglobuline,²² and ribonuclease T1.²³ The initial PDB structures were processed and solvated in a water box extending at least 2 nm from the protein to the edges of the unit cell using the CHARMM-GUI web server.^{12,13} The CHARMM-GUI default water model (i.e., TIP3P) was changed to SPCE or SPCE-HW. Sodium or chloride counterions¹³ were added to neutralize the systems. The obtained systems were then energy minimized and equilibrated in the NpT ensemble for 10 ns, after which a production run of 1 μs followed for each of the three proteins. In addition, for ribonuclease a set of extra simulations in a range of different temperatures was performed in order to

simulate the melting of the protein. For each temperature, a 1.7 μs trajectory was generated with the first μs taken as equilibration and discarded from the analysis (for further details see the Supporting Information, Tables S2 and S3).

To explore the effect of water deuteration on biological membranes, a bilayer containing 200 phospholipids (POPC) was constructed using CHARMM-GUI.^{12,13} The total amount of water molecules (SPCE or SPCE-HW) added was 15180. After energy minimization and equilibration in the NpT ensemble of 10 ns, the systems were run for 200 ns. Furthermore, a patch of a dipalmitoylphosphatidylcholine (DPPC) membrane was built using CHARMM-GUI^{12,13} and simulated to evaluate the effect of the employed water models on the temperature of phase transition from the gel phase to the liquid phase. A bilayer composed of 64 lipids was solvated with a total of 2600 water molecules and simulated with temperature annealing from 325 to 305 K in 2 μs .

RESULTS AND DISCUSSION

The free energies of transfer ΔG from H_2O to D_2O are summarized for all the amino acids in Figure 2. All the calculated free energies are positive, which means the amino acids are less stable in the heavy water than in light water. It is worth noting that the ΔG values, which vary between 0.7 and 2.2 kcal/mol do not follow the hydrophobicity scale of amino acids. The results presented in Figure 2 rather point to the molecular size as the main factor governing the ΔG values: the larger the amino acid, the more unfavorable the transfer from H_2O to D_2O . The always positive free energy of transfer between the two solvents indicates that, compared to H_2O , D_2O has a higher propensity to form water–water hydrogen bonds than water–amino acid hydrogen bonds. This explains why the free energy of hydration depends on the excluded volume, which for small molecules like amino acids correlates well with the molecular weight. We indeed see a very good linear correlation between the molar mass and the free energy of transfer (Figure 3).

The above results concerning amino acids indicate that D_2O may be a somewhat worse solvent than H_2O for proteins, hence inducing also more compact structures with a reduced radius of gyration. To test this, we modeled and analyzed the behavior of three globular proteins (azurine, lactoglobuline,

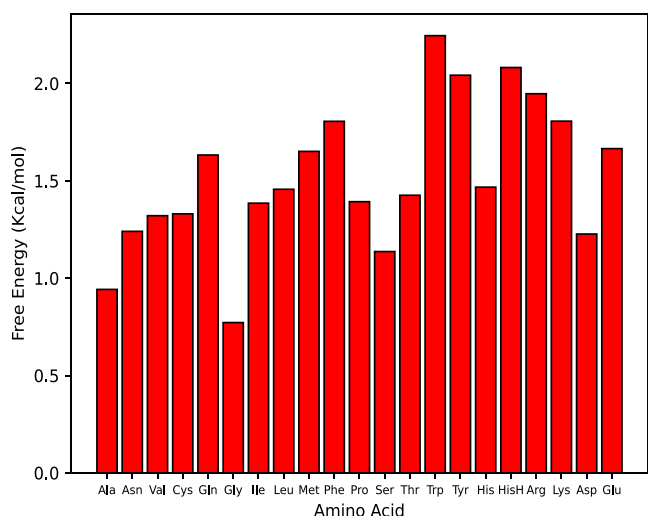


Figure 2. Free energy of transfer from SPC/E (H_2O) to SPCE-HW (D_2O). Errors are not reported because too small. A table of the values with associated errors can be found in the [Supporting Information](#) (Table S1).

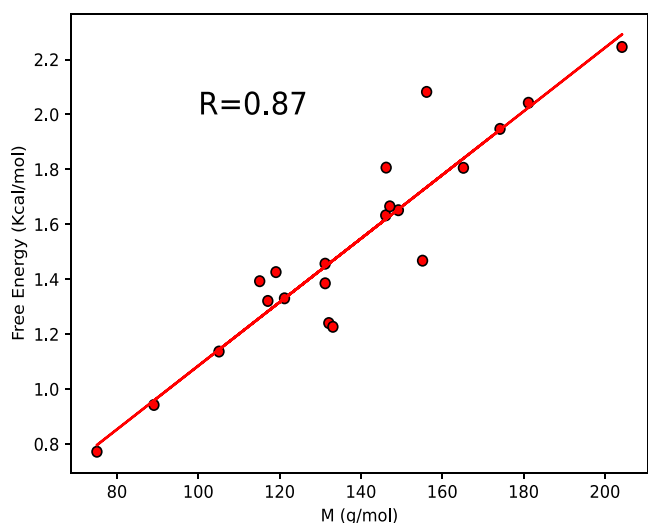
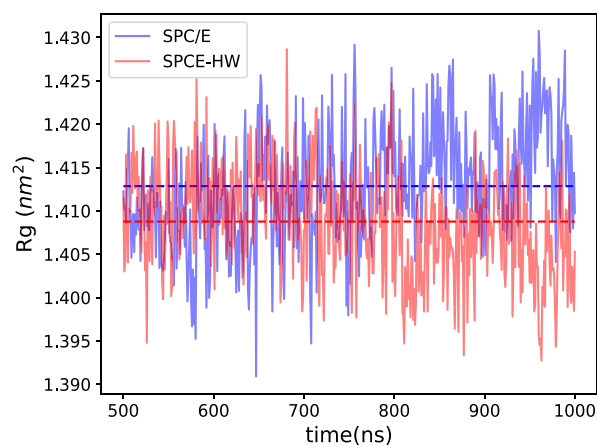


Figure 3. Free energy of transfer from SPC/E (H_2O) to SPCE-HW (D_2O) as a function of the amino acid molar mass showing a very good linear correlation.

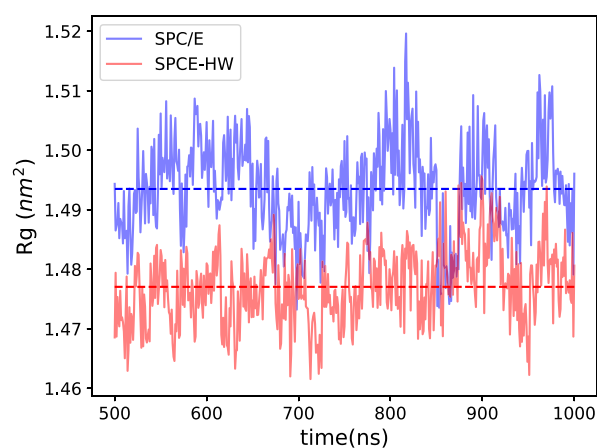
and ribonuclease) in D_2O vs H_2O . All three proteins show a small but consistent decrease in the radius of gyration when moving from H_2O to D_2O ; see [Figure 4](#). The same trend is also observed for the solvent accessible surface area (SASA); see [Figure S1](#) in the Supporting Information. Simulations thus show that water deuteration is making the proteins tighter, which is consistent with the positive free energy of transfer presented above, as well as with the generally tightening effect of D_2O found in experiment.²⁴

For ribonuclease, we also modeled the effect of deuteration on the protein melting temperature. From the results presented in [Figure 5](#), we see that the tightening of the protein structure upon water deuteration also leads to stabilization and increase of the protein melting temperature.

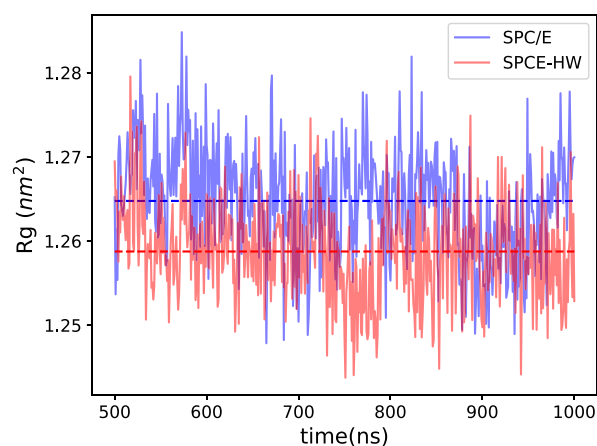
This is in line with experimental observations,^{27,28} and it is consistent with the sign of the free energies of transfer from H_2O to D_2O of individual amino acids in [Figure 2](#).



(A)



(B)



(C)

Figure 4. Radius of gyration of azurine (A), lactoglobuline (B), and ribonuclease (C) in SPC/E (H_2O , blue line) and SPCE-HW (D_2O , red line). The dashed lines represent average values over the production runs of 500 ns.

The tightening effect of the heavy water is not limited to proteins as demonstrated on the areas per lipid (APL) calculated for a POPC bilayer, yielding a value of 0.63 nm²

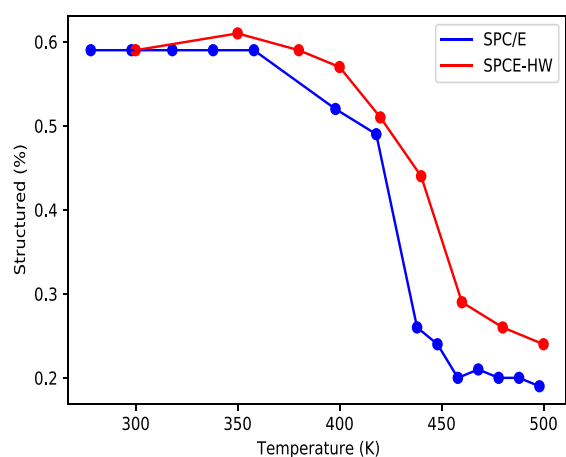


Figure 5. Annealing of ribonuclease. The percentage of structured protein is plotted versus the temperature. The percentage of structured protein was calculator using the gromacs tool “gmx do dssp”, which uses the DSSP algorithm.^{25,26} The blue line is SPC/E (H₂O), and the red line is SPCE-HW (D₂O). Error bars are not reported as they are too small to be visible in the figure (i.e., below 1%).

for H₂O and 0.59 nm for D₂O. The effect of the solvent on the POPC membrane APL is qualitatively in line (albeit more pronounced) with previous simulations.²⁹ At the same time, we observe a small (about 3%) bilayer thickening, as deduced from the density profile of phosphate (see the [Supporting Information](#)). These results are consistent with experimental findings.³⁰

In addition, the behavior of a DPPC bilayer around the melting point was investigated by monitoring the APL as a function of temperature. The results are shown in [Figure 6](#).

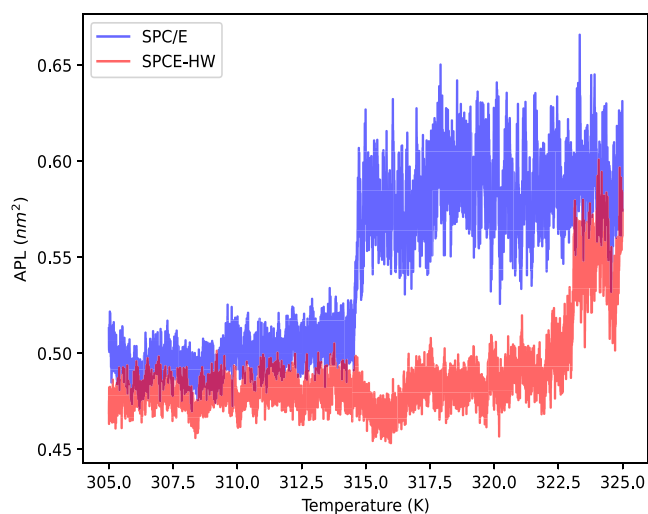


Figure 6. Area per lipid of a DPPC bilayer as a function of temperature in SPC/E H₂O (blue) and SPCE-HW D₂O (red).

The first thing to notice is that the APL in D₂O is systematically smaller than that in H₂O, which is consistent with the above results for POPC. DPPC in H₂O exhibits a melting temperature of 314.5 K, which is in very good agreement with both the experimentally determined value of 314.15 K³¹ and the value from MD simulations using the default TIP3P charmm model.³² The system simulated in D₂O using the present model shows a clear upward shift of the

melting temperature of almost 10 K. Qualitatively, this is in accord with experimental findings,³¹ although the measured shift in melting temperature is smaller (less than 1 K).

CONCLUSION

In this work, we quantified the effect of water deuteration on amino acids, proteins, and phospholipid membranes. This was done using classical molecular dynamics simulations employing models that account for differences between H₂O and D₂O in an effective way, incorporating nuclear quantum effects into the intermolecular potential. In particular, we focused on differences in structural properties such as the compactness of the biomolecules and thermodynamic effects like melting temperatures and free energies of transfer of solutes from light to heavy water. To the former, our results reveal small but systematic structural effects on proteins. Namely, we observe a decrease in radii of gyration of less than 1% upon moving from H₂O to D₂O. Interestingly, structural effects on phospholipid membranes are larger than on proteins; in particular, upon deuteration we observed a decrease of the area per lipid by more than 10% and thickening of the bilayer by about 3%. To the latter, our results show that all amino acids are slightly less soluble in heavy vs light water. Also, moving from H₂O to D₂O, we observe an upward shift by several degrees of the melting point of a model protein (ribonuclease). The same affect of increasing the melting temperature is found for DPPC. Altogether the simulations show that the structural effect on a globular protein might be small, but the thermodynamic effect (melting) on protein and membranes can be important, especially if an experiment is conducted close to the phase transition temperature, where even a small shift can change the physical-chemical properties. Comparison to available experimental data shows that our simple models capture well the principal effect, namely, that D₂O is a somewhat worse solvent for biomolecules than H₂O. This also implies that association between proteins or between a protein and a biomembrane may be positively affected by water deuteration. Finally, protein domains that are intrinsically disordered and thus very sensitive to the balance between protein–solvent and solvent–solvent interaction, may show a high sensitivity to the H₂O to D₂O substitution.

ASSOCIATED CONTENT

Supporting Information

The Supporting Information is available free of charge at <https://pubs.acs.org/doi/10.1021/acs.jpcb.2c08270>.

Free energy of transfer (with associated error) between water and heavy water for each single amino acid, solvent accessible surface area (SASA) of azurine, lactoglobulin, and ribonuclease in water and heavy water. Detailed data about annealing of ibonuclease T1. Phosphate density profile of POPC in water and heavy water ([PDF](#))

AUTHOR INFORMATION

Corresponding Author

Pavel Jungwirth – Institute of Organic Chemistry and Biochemistry of the Czech Academy of Sciences, 160 00 Prague 6, Czech Republic; orcid.org/0000-0002-6892-3288; Email: pavel.jungwirth@uochb.cas.cz

Authors

Carmelo Tempra – Institute of Organic Chemistry and Biochemistry of the Czech Academy of Sciences, 160 00 Prague 6, Czech Republic; orcid.org/0000-0002-2890-6993

Victor Cruces Chamorro – Institute of Organic Chemistry and Biochemistry of the Czech Academy of Sciences, 160 00 Prague 6, Czech Republic

Complete contact information is available at:

<https://pubs.acs.org/10.1021/acs.jpcb.2c08270>

Notes

The authors declare no competing financial interest.

ACKNOWLEDGMENTS

P.J. thanks the Czech Science Foundation (EXPRO grant no. 19-26854X). C.T. and V.C.C. are grateful for support from the Faculty of Science of the Charles University, Prague, where they are enrolled as Ph.D. students. C.T. and V.C.C. thank the International Max Planck Research School for Many-Particle Systems in Structured Environments hosted by the Max Planck Institute for the Physics of Complex Systems, Dresden, Germany.

REFERENCES

- (1) Ceriotti, M.; Fang, W.; Kusalik, P. G.; McKenzie, R. H.; Michaelides, A.; Morales, M. A.; Markland, T. E. Nuclear Quantum Effects in Water and Aqueous Systems: Experiment, Theory, and Current Challenges. *Chem. Rev.* **2016**, *116*, 7529–7550.
- (2) De Meutter, J.; Goormaghtigh, E. Evaluation of protein secondary structure from FTIR spectra improved after partial deuteration. *Eur. Biophys. J.* **2021**, *50*, 613–628.
- (3) Hume, S.; Hithell, G.; Greetham, G. M.; Donaldson, P. M.; Towrie, M.; Parker, A. W.; Baker, M. J.; Hunt, N. T. Measuring proteins in H₂O with 2D-IR spectroscopy. *Chemical science* **2019**, *10*, 6448–6456.
- (4) Cioni, P.; Strambini, G. B. Effect of heavy water on protein flexibility. *Biophysical journal* **2002**, *82*, 3246–3253.
- (5) Beranová, L.; Humpolíčková, J.; Šykora, J.; Benda, A.; Cwiklik, L.; Jurkiewicz, P.; Gröbner, G.; Hof, M. Effect of heavy water on phospholipid membranes: experimental confirmation of molecular dynamics simulations. *Phys. Chem. Chem. Phys.* **2012**, *14*, 14516–14522.
- (6) Ben Abu, N.; Mason, P. E.; Klein, H.; Dubovski, N.; Ben Shoshan-Galeczki, Y.; Malach, E.; Pražienková, V.; Maletínská, L.; Tempra, C.; Chamorro, V. C.; et al. Sweet taste of heavy water. *Commun. Biol.* **2021**, *4*, 440.
- (7) Ceriotti, M.; Fang, W.; Kusalik, P. G.; McKenzie, R. H.; Michaelides, A.; Morales, M. A.; Markland, T. E. Nuclear quantum effects in water and aqueous systems: Experiment, theory, and current challenges. *Chem. Rev.* **2016**, *116*, 7529–7550.
- (8) Guillot, B.; Guissani, Y. Quantum effects in simulated water by the Feynman–Hibbs approach. *J. Chem. Phys.* **1998**, *108*, 10162–10174.
- (9) Sesé, L. M. Feynman–Hibbs quantum effective potentials for Monte Carlo simulations of liquid neon. *Mol. Phys.* **1993**, *78*, 1167–1177.
- (10) Chamorro, V. C.; Tempra, C.; Jungwirth, P. Heavy Water Models for Classical Molecular Dynamics: Effective Inclusion of Nuclear Quantum Effects. *J. Phys. Chem. B* **2021**, *125*, 4514–4519.
- (11) Berendsen, H. J. C.; Grigera, J. R.; Straatsma, T. P. The missing term in effective pair potentials. *J. Phys. Chem.* **1987**, *91*, 6269–6271.
- (12) Jo, S.; Kim, T.; Iyer, V. G.; Im, W. CHARMM-GUI: a web-based graphical user interface for CHARMM. *Journal of computational chemistry* **2008**, *29*, 1859–1865.
- (13) Lee, J.; Cheng, X.; Swails, J. M.; Yeom, M. S.; Eastman, P. K.; Lemkul, J. A.; Wei, S.; Buckner, J.; Jeong, J. C.; Qi, Y.; et al. CHARMM-GUI input generator for NAMD, GROMACS, AMBER, OpenMM, and CHARMM/OpenMM simulations using the CHARMM36 additive force field. *J. Chem. Theory Comput.* **2016**, *12*, 405–413.
- (14) Swope, W. C.; Andersen, H. C.; Berens, P. H.; Wilson, K. R. A computer simulation method for the calculation of equilibrium constants for the formation of physical clusters of molecules: Application to small water clusters. *J. Chem. Phys.* **1982**, *76*, 637–649.
- (15) Darden, T.; York, D.; Pedersen, L. Particle mesh Ewald: An N log(N) method for Ewald sums in large systems. *J. Chem. Phys.* **1993**, *98*, 10089–10092.
- (16) Essmann, U.; Perera, L.; Berkowitz, M. L.; Darden, T.; Lee, H.; Pedersen, L. G. A smooth particle mesh Ewald method. *J. Chem. Phys.* **1995**, *103*, 8577–8593.
- (17) Bussi, G.; Donadio, D.; Parrinello, M. Canonical sampling through velocity rescaling. *J. Chem. Phys.* **2007**, *126*, 014101.
- (18) Parrinello, M.; Rahman, A. Polymorphic transitions in single crystals: A new molecular dynamics method. *J. Appl. Phys.* **1981**, *52*, 7182–7190.
- (19) Kumar, S.; Rosenberg, J. M.; Bouzida, D.; Swendsen, R. H.; Kollman, P. A. The weighted histogram analysis method for free-energy calculations on biomolecules. I. The method. *Journal of computational chemistry* **1992**, *13*, 1011–1021.
- (20) Hub, J. S.; De Groot, B. L.; Van Der Spoel, D. g_wham A Free Weighted Histogram Analysis Implementation Including Robust Error and Autocorrelation Estimates. *J. Chem. Theory Comput.* **2010**, *6*, 3713–3720.
- (21) Nar, H.; Messerschmidt, A.; Huber, R.; Van de Kamp, M.; Canters, G. W. Crystal structure of *Pseudomonas aeruginosa* azurin at 1.85 Å resolution. *Febs Letters* **1992**, *306*, 119–124.
- (22) Brownlow, S.; Cabral, J. H. M.; Cooper, R.; Flower, D. R.; Yewdall, S. J.; Polikarpov, I.; North, A. C.; Sawyer, L. Bovine β -lactoglobulin at 1.8 Å resolution—still an enigmatic lipocalin. *Structure* **1997**, *5*, 481–495.
- (23) Kostrewa, D.; Choe, H. W.; Heinemann, U.; Saenger, W. Crystal structure of guanosine-free ribonuclease T1, complexed with vanadate (V), suggests conformational change upon substrate binding. *Biochemistry* **1989**, *28*, 7592–7600.
- (24) Cioni, P.; Strambini, G. B. Effect of heavy water on protein flexibility. *Biophysical journal* **2002**, *82*, 3246–3253.
- (25) Joosten, R. P.; Te Beek, T. A.; Krieger, E.; Hekkelman, M. L.; Hooft, R. W.; Schneider, R.; Sander, C.; Vriend, G. A series of PDB related databases for everyday needs. *Nucleic acids research* **2011**, *39*, D411–D419.
- (26) Kabsch, W.; Sander, C. Dictionary of protein secondary structure: pattern recognition of hydrogen-bonded and geometrical features. *Biopolymers: Original Research on Biomolecules* **1983**, *22*, 2577–2637.
- (27) Pica, A.; Graziano, G. Effect of heavy water on the conformational stability of globular proteins. *Biopolymers* **2018**, *109*, No. e23076.
- (28) Makhatadze, G. I.; Clore, G. M.; Gronenborn, A. M. Solvent isotope effect and protein stability. *Nature structural biology* **1995**, *2*, 852–855.
- (29) Rog, T.; Murzyn, K.; Milhaud, J.; Karttunen, M.; Pasenkiewicz-Gierula, M. Water isotope effect on the phosphatidylcholine bilayer properties: a molecular dynamics simulation study. *J. Phys. Chem. B* **2009**, *113*, 2378–2387.
- (30) Matsuki, H.; Okuno, H.; Sakano, F.; Kusube, M.; Kaneshina, S. Effect of deuterium oxide on the thermodynamic quantities associated with phase transitions of phosphatidylcholine bilayer membranes. *Biochimica et Biophysica Acta (BBA)-Biomembranes* **2005**, *1712*, 92–100.
- (31) Matsuki, H.; Okuno, H.; Sakano, F.; Kusube, M.; Kaneshina, S. Effect of deuterium oxide on the thermodynamic quantities associated with phase transitions of phosphatidylcholine bilayer membranes.

Biochimica et Biophysica Acta (BBA)-Biomembranes **2005**, *1712*, 92–100.

(32) Khakbaz, P.; Klauda, J. B. Investigation of phase transitions of saturated phosphocholine lipid bilayers via molecular dynamics simulations. *Biochimica et Biophysica Acta (BBA)-Biomembranes* **2018**, *1860*, 1489–1501.