

Article

Free-Radical Propagation Rate Coefficients of Diethyl Itaconate and Di-n-Propyl Itaconate Obtained via PLP-SEC

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Abstract: The propagation step is one of the key reactions in radical polymerization and knowledge about its kinetics is often vital for understanding and designing polymerization processes leading to new materials or optimizing technical processes. Arrhenius expressions for the propagation step in free-radical polymerization of diethyl itaconate (DEI) as well as di-n-propyl itaconate (DnPI) in bulk, for which propagation kinetics was yet unexplored, were thus determined via pulsed-laser polymerization in conjunction with size-exclusion chromatography (PLP-SEC) experiments in the temperature range of 20 to 70 °C. For DEI, the experimental data was complemented by quantum chemical calculation. The obtained Arrhenius parameters are $A = 1.1 \text{ L}\cdot\text{mol}^{-1}\cdot\text{s}^{-1}$ and $E_a = 17.5 \text{ kJ}\cdot\text{mol}^{-1}$ for DEI and $A = 1.0 \text{ L}\cdot\text{mol}^{-1}\cdot\text{s}^{-1}$ and $E_a = 17.5 \text{ kJ}\cdot\text{mol}^{-1}$ for DnPI.

Keywords: propagation rate coefficient; k_p ; PLP-SEC; itaconate; kinetics; radical polymerization



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1. Introduction

Esters of itaconic acid (2-methylenesuccinic acid) show the potential to replace styrene as a comonomer in the synthesis of polymer resins [1] or as composite material with cotton [2]. This is especially interesting given the fact that itaconic acid can be produced by many microbiological organisms such as *Aspergillus terreus* and thus can be obtained from biomass while styrene is petrol based [3–5].

In order to fully utilize itaconic acid esters in free-radical polymerizations, the exact knowledge of the reaction rate coefficients is highly advantageous. Here, we are focusing on the propagation rate coefficient, k_p . While in a number of publications [6–13], data for k_p for a few members of the series of itaconic acid homo-diester are already available, this series is lacking some data, most notably for the diethyl and the di-n-propyl ester. This work intends to fill these gaps by applying pulsed-laser polymerization in conjunction with the size-exclusion chromatography (PLP-SEC) method as well as quantum chemical prediction in order to obtain Arrhenius expressions for k_p of these monomers.

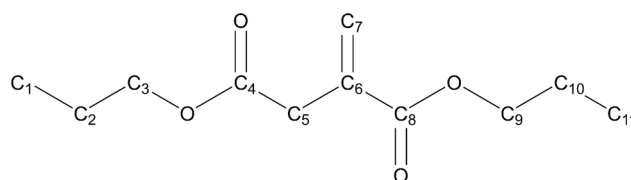
PLP is a highly specialized technique for studying the kinetics of radical polymerization reactions. It involves using a pulsed laser to initiate polymerization in short bursts, allowing for precise control over reaction conditions and reaction kinetics. One of the key advantages of PLP is that it allows for the measurement of propagation rate coefficients of radical polymerization with high precision and accuracy, and is thus IUPAC-recommended. The propagation rate coefficient, k_p , is one of the kinetic key parameters that describe how fast the chain-growth reaction proceeds and how sensitive it is to changes in reaction conditions such as temperature and concentration of reactants. Measuring kinetic coefficients in radical polymerization is important for understanding and optimizing radical polymerization processes, as it allows to predict and control the behavior of the reaction under different conditions and in addition enables to predict the properties of the resulting polymer. The kinetic information can thus be used to design new polymers with specific properties, optimize reaction conditions for industrial-scale production, and troubleshoot problems that arise during polymerization.

2. Materials and Methods

Chemicals: Diethyl itaconate (DEI, $M = 186.21$ g/mol) (98%) was purchased from Tokyo Chemical Industry (TCI). The inhibitor 4-tert-butylcatechol (TBC) was removed via column chromatography over alkaline aluminum oxide.

Di-n-propyl itaconate (DnPI, $M = 214.26$ g/mol) was obtained by esterification of itaconic acid (TCI, 99%) with 1-propanol (Sigma-Aldrich, 99.5%) according to a literature procedure [14]. 1 eq. of itaconic acid was dissolved in 10 eq. of 1-propanol and acidified with 0.4 eq. conc. H_2SO_4 . The mixture was boiled under reflux for 5 h. Then, the alcohol was removed using a rotary evaporator. The residue was poured in a 5-times excess of ice water, the phases were separated, and the organic phase was neutralized with a 10 wt.%-solution of $NaHCO_3$ in water. Afterwards, the product was dried over $CaCl_2$. The purity was determined by NMR and was found to be >95 %.

1H NMR (300 MHz, $CDCl_3$, δ , see Scheme 1): 6.31 (s, 1H; C7), 5.68 (s, 1H; C7), 4.11 (t, $J = 6.8$ Hz, 2H; C3), 4.04 (t, $J = 6.8$ Hz, 2H; C9), 3.33 (s, 2H; C5), 1.65 (m, 4H; C2+C10), 0.92 (q $J = 7.3$ Hz, 6H; C1+C11); ^{13}C NMR (75 MHz, $CDCl_3$, δ): 170.9 (C4), 166.4 (C8), 134.3 (C6), 128.2 (C7), 66.72 (C3/C9), 66.6 (C3/C9), 37.9 (C5), 22.06 (C2+C10), 10.5 (C1+C11).



Scheme 1. Carbon atom numbering for DnPI used in the NMR assignment.

The photoinitiator 2,2-dimethoxy-2-phenylacetophenone (DMPA, Aldrich, 99%) was used as received.

PLP experiments: All samples were degassed by a gentle flow of argon for at least 10 min. Pulsed-laser-polymerizations were performed in silica glass cells which were tempered by a heating bath for at least 10 min prior to the start of irradiation. The laser used for irradiation is an ATLEX 1000 I XeF exciplex laser operating at 351 nm with a pulse energy of 1–7 mJ. Repetition rates varied from 0.25 to 2 Hz. The samples were prepared by dissolving 20 mM to 75 mM of DMPA in the respective itaconate ester. No further solvent was added.

After irradiation, the mixture was used directly for SEC analysis without isolation of the polymer.

Density measurements: A gas pycnometry system (AccuPyc II from Micromeritics) was used to measure room temperature densities of the monomers. Since the instrument did not have a temperature control system, room temperature densities were used for the whole temperature range. The following densities were obtained: 1.044 g·mL $^{-1}$ (DEI) and 1.025 g·mL $^{-1}$ (DnPI).

Size-exclusion chromatography: Size-exclusion chromatography was performed using an Agilent 1260 G1310B iso pump with a PSS degasser. Injection was performed with an Agilent 1260 ALS G1329B autosampler onto three PSS SDV 5 μ m bead columns ($10^6, 10^4, 10^3$ Angstrom) which was embedded in a PSS TCC6000 column oven at 35 °C. THF at a flow rate of 1 mL·min $^{-1}$ was used as the solvent. The detector was an Agilent 1260 RID G1362A. Access to absolute molecular weight distributions was obtained by a PS-PMMA calibration in conjunction with the Mark-Houwink-constants of DEI and DnPI, which have been determined previously [15]. These values refer to the solvent toluene, but may also be used for THF, as suggested by Szablan et al. [7].

Quantum chemical calculations: A propagating DEI radical was mimicked by a DEI molecule with a methyl group bound to the former olefinic double bond. Minimum structures of the reactants (consisting of the radical and a monomer) and the product were found using the crest program by Grimme [16]. Repeated runs on different conformers of reactant and product were performed and the global minima were identified. These

global minima were optimized with ORCA 4.2.1 and subsequently a NEB-TS-calculation was performed to find the TS structure [17,18].

3. Results

The Arrhenius parameters for the propagation step of the sterically hindered monomer DEI were calculated using quantum chemistry and transition state theory (TST). In order to obtain experimental values for k_p , PLP-SEC experiments were performed for DEI in bulk in the temperature range between 20 and 70 °C and for DnPI in the temperature range of 30 to 70 °C. For the basic principles of the IUPAC-recommended PLP-SEC method and for experimental peculiarities, please also refer to [19,20].

3.1. Prediction of Arrhenius Parameters for DEI

k_p as function of temperature, $k_p(T)$, can be calculated from the partition functions Q_n of the involved molecules n and the electronic barrier height E_0 (see Equation (1)). κ denotes a correction factor in order to account for tunneling (assumed to be equal to 1 here and will be left out in the following), c is the inverse of the volume used in the translational partition function, k_B is the Boltzmann constant, R the universal gas constant, and m is the molecularity of the reaction [21].

$$k_p(T) = \kappa \cdot c^{(1-m)} \cdot \frac{k_B T}{h} \cdot \frac{Q^\ddagger}{\prod_n Q_n} \cdot \exp\left(-\frac{E_0}{RT}\right) \quad (1)$$

This expression can also be reformulated using the thermodynamical properties entropy of activation, ΔS^\ddagger , and enthalpy of activation, ΔH^\ddagger :

$$k_p(T) = c^{(1-m)} \cdot \frac{k_B T}{h} \exp\left(\frac{\Delta S^\ddagger}{R}\right) \cdot \exp\left(-\frac{\Delta H^\ddagger}{RT}\right) \quad (2)$$

From Equation (2), the Arrhenius parameters for k_p can be calculated directly by Equations (3) and (4). The formulae for calculating the individual entropy components can be found in the Supporting Information.

$$A = c^{(1-m)} \cdot \exp(m) \cdot \frac{k_B T}{h} \exp\left(\frac{\Delta S^\ddagger}{R}\right) \quad (3)$$

$$E_a = E_0 + ZPVE + \Delta\Delta H^\ddagger + mRT \quad (4)$$

ZPVE is the vibrational zero-point energy and $\Delta\Delta H^\ddagger$ a temperature correction, which is calculated for a single species by Equation (5) where ν_j denotes the frequency of the j -th vibration and h is the Planck constant.

$$\Delta\Delta H = R \cdot \sum_j \frac{h\nu_j}{\exp(h\nu_j) - 1} + \frac{5}{2}RT + \frac{3}{2}RT \quad (5)$$

Itaconates are 1,1-disubstituted ethylene derivatives and thus four different propagation variants are possible. The attacking radical as well as the resulting macroradical can each be centered on the tertiary or the primary carbon. Radical stability favors the radical to be centered on the tertiary carbon, while steric hindrance favors an attack on the primary carbon, resulting in the formed radical also to be tertiary. Consequently, only this reaction has been considered in this work. One might improve the results presented here by considering all variants and Boltzmann-weighting them.

Geometries and frequencies are rather insensitive to the level of theory, so using a cheap method is recommended. The electronic barrier, on the other side, is typically very sensitive to the level of theory used, so a high-quality method should be used. However, for reasons of computational power, in this work only UHF/6-31G(d) and B3LYP/def2-TZVP were used.

The calculations were carried out in the gas phase. More accurate results can be obtained by including a solvent model to include interactions between the solute and solvent. These can influence for example electronic energy levels and vibrational frequencies [22]. Since the monomer in question is rather exotic, no implicit solvent model parameters exist for it, to our best knowledge. Consequently, an explicit solvent model would have to be used. However, this comes at a largely increased computational cost and is far beyond the available resources. Consequently, no solvent model could be used.

The vibrational frequencies were used unscaled as well as scaled. For UHF/6-31G(d), the average of the scaling factors reported by Scott et al. [23] was used, which amounts to 0.90064. For B3LYP/def2-TZVP, the average of the scaling factors reported by Kesharwani et al. [24] was used, which amounts to 0.9870.

The obtained minimum structures (on UHF/6-31G(d) level of theory) are depicted in the Appendix B and their .xyz-coordinates are also given. It is clearly visible that both molecules are approximately orthogonally aligned to each other in the reactant structure, in order to minimize steric repulsion. In the transition state (TS), however, both molecules are aligned in parallel, in order to facilitate overlap of the involved orbitals. The product structure is very similar to the TS structure, indicating a late TS. This is unexpected since exothermic reactions as these typically exhibit an early TS.

Obtained electronic barrier heights and Arrhenius parameters are shown in Table 1. The Arrhenius activation energy is very similar for both levels of theory. Likely, a high level of theory such as CCSD(T) would increase the accuracy significantly. The Arrhenius prefactors differ by approximately one order of magnitude for both levels of theory. Both, however, are significantly lower compared to other typical classes of monomers such as acrylates, where they are in the order of $10^6 \text{ L} \cdot \text{mol}^{-1} \cdot \text{s}^{-1}$ [25]. This can be explained by the bulky substituents at the olefinic bond, which hinder the approach of a new monomer for reaction.

Table 1. Calculated electronic barriers, E_0 , and Arrhenius parameters (prefactor, A , and activation energy, E_A) for the propagation step of DEI at 20 °C. The frequencies used for the calculation were used unscaled as well as scaled by their respective scaling factors.

Method	E_0 [kJ·mol ⁻¹]	A [L·mol ⁻¹ ·s ⁻¹]	E_A [kJ·mol ⁻¹]
UHF/6-31G(d)	12.7	$2.0 \cdot 10^3$	22.6
UHF/6-31G(d) scaled	12.7	$3.3 \cdot 10^3$	22.3
B3LYP/def2-TZVP	13.8	$1.4 \cdot 10^4$	22.7
B3LYP/def2-TZVP scaled	13.8	$1.5 \cdot 10^4$	22.7

3.2. PLP Experiments

The PLP-SEC-method introduced by Olaj and coworkers [26] offers a very easy way to determine propagation rate coefficients in radical polymerization. A sample consisting of monomer, photoinitiator, and solvent (if applicable) is subjected to repeated, short laser pulses. The laser pulse initiates polymerization by instantaneous formation of radicals. In between the pulses, the radicals propagate (and also terminate to some extent) but upon the following laser pulse, the radical concentration rises sharply again, and the termination probability is thus greatly enhanced. The length L of the chains which were initiated by a laser pulse and were terminated in the moment of the following laser pulse corresponds to $L = k_p \cdot c_M \cdot t_0$ with t_0 being the time between two consecutive laser pulses and c_M the monomer concentration. However, not all radicals terminate at the first following laser pulse, but may be stopped at the second or third (i -th) subsequent pulse. The corresponding chain length is then $L_i = i \cdot k_p \cdot c_M \cdot t_0$, provided the monomer conversion is kept relatively low. In any case, a polymer with that specific chain length of $L, 2L, 3L, \dots$ is more frequently occurring than that having other chain lengths, and thus, structured chain-length distributions with multiple maxima are typically observed. The structure can be evaluated to obtain estimates for L_i and the inflection point at the low molecular weight side of the peak typically yields the most precise results [26]. The

inflection points were determined by evaluating the maxima of the derivative of the molar mass distribution, as shown in Figure 1.

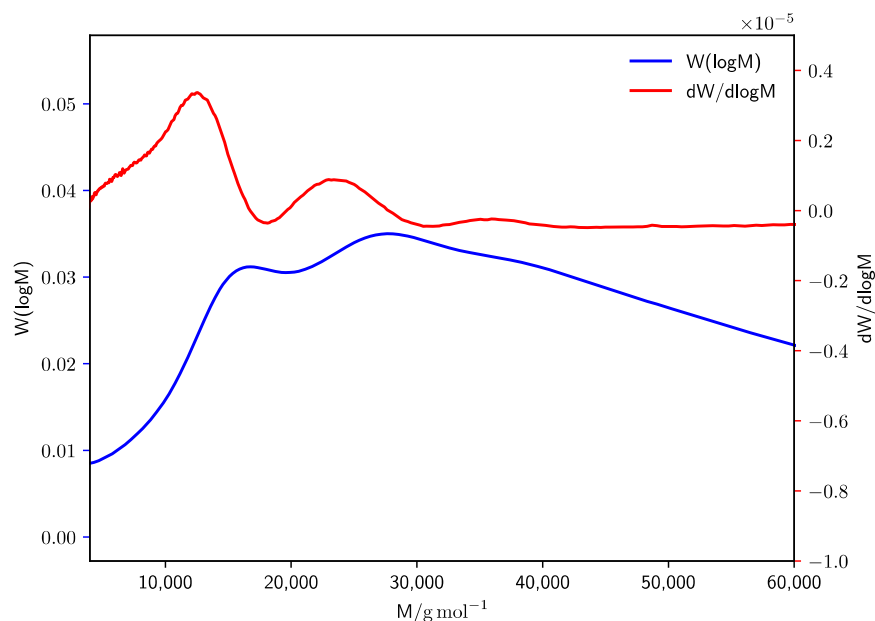


Figure 1. Detail of a representative molar mass distribution (**blue**) from PLP of DEI and its derivative (**red**). The molar masses L_i were obtained from the maxima of the derivative curve.

Itaconates propagate slowly and also terminate slowly [8], which presents a challenge for the PLP method. In this work, we addressed this challenge by choosing an extremely slow laser repetition rate, which allowed us to obtain successful results.

3.2.1. Diethyl Itaconate

The PLP-SEC data obtained for DEI are presented in Table 2. For each temperature, two separate solutions of the photoinitiator 2,2-dimethoxy-2-phenylacetophenone (DMPA) in DEI were prepared and from each solution two experiments were performed. The molar masses at the inflection points, M_{Inf} , were extracted from complete molar mass distributions obtained via size-exclusion chromatography.

To verify the integrity of the data, the obtained values of k_p shall be independent of the initiator concentration and the repetition rate. Furthermore, the position of the second additional PLP peak shall be at double the molar mass of the first peak and analogously the third peak at triple the molar mass of the first peak (or 1.5 times of the second peak). This requirement rests on the assumption that k_p is chain-length independent and is called the “consistency criterion” of PLP-SEC. It was introduced for distinguishing the additional PLP peaks from any other occurring peaks that are not kinetically relevant.

As can be seen in Figure 2, the ratio of the second peak position divided by the first is systematically lower than 2, indicating a significant chain length dependence of propagation for chain lengths up to ca. 150 to 200. For the ratio between the third and second peak, however, the agreement with the expected ratio of 1.5 is good. This observation has been made for many other systems studied via PLP-SEC [27–30] and there still is discussion, whether this effect—which is clearly going beyond the generally accepted chain length dependency of k_p for short macroradicals up to chain lengths of ca. 5 [31,32]—is reflecting a real chain-length dependency up to longer chain-lengths or is an SEC artifact [33,34]. PLP-SEC alone is apparently not capable of deciding this open question and it is now generally accepted that the k_p data obtained are estimates subject to uncertainty, due to such as yet unexplained effects.

Table 2. PLP-SEC data for DEI: initiator concentration, c_{Ini} , temperature T , total number of pulses, n_{pulses} , laser pulse frequency, ν , molecular weight at the various inflection points, M_{inf} , and the average k_p of all obtained inflection points.

c_{Ini} [mol·L ⁻¹]	T/K	n_{pulses}	ν [Hz]	M_{inf1} [g·mol ⁻¹]	M_{inf2} [g·mol ⁻¹]	M_{inf3} [g·mol ⁻¹]	k_p [L·mol ⁻¹ ·s ⁻¹]
0.022	343.15	2000	1	25,400	42,800	-	22.4
0.022	343.15	2000	2	15,400	27,700	-	28.0
0.040	343.15	2000	1	26,600	42,800	-	23.0
0.040	343.15	2000	2	15,200	28,000	-	28.0
0.026	333.15	1000	0.5	43,000	73,800	108,800	18.5
0.026	333.15	4000	1	25,100	44,900	65,600	22.2
0.047	333.15	2400	0.5	44,700	79,100	116,500	19.7
0.047	333.15	4000	1	23,600	44,500	-	22.0
0.022	323.15	2000	1	20,000	36,400	54,300	18.0
0.022	323.15	2000	0.5	35,700	63,100	96,100	15.9
0.061	323.15	2000	1	19,780	36,400	-	18.2
0.061	323.15	2000	0.5	35,100	61,120	-	15.7
0.033	313.15	2000	1	16,320	30,050	-	15.0
0.051	313.15	2000	1	15,700	29,800	-	14.7
0.051	313.15	2000	0.5	29,100	52,300	-	13.2
0.027	303.15	2000	1	12,510	23,310	35,790	11.5
0.027	303.15	2000	0.5	23,500	42,900	65,000	10.6
0.046	303.15	2000	1	12,490	23,420	-	11.6
0.046	303.15	2000	0.5	23,100	42,900	64,200	10.5

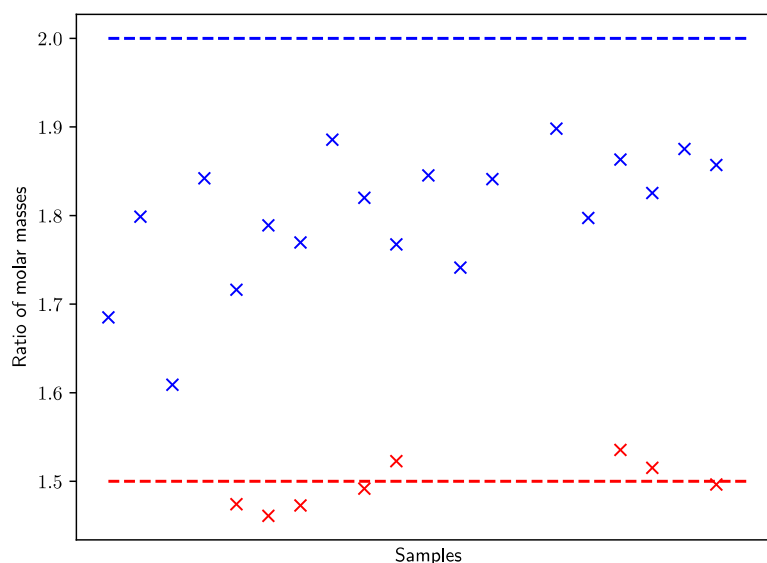


Figure 2. Ratio of molar masses of inflection points for all DEI samples. Blue points show the ratio of the second inflection point to the first one, while the red points show the ratio of the third inflection point to the second one.

Figure 3 shows the obtained k_p values for different temperatures and different laser pulse frequencies as a function of initiator concentration. It is clearly visible that k_p is independent of initiator concentration—a strict criterion for a successful PLP-SEC experiment—, but, however, is not independent of the repetition rate, which alters the chain length of the produced polymer, as discussed above. With increasing repetition rate, k_p increases.

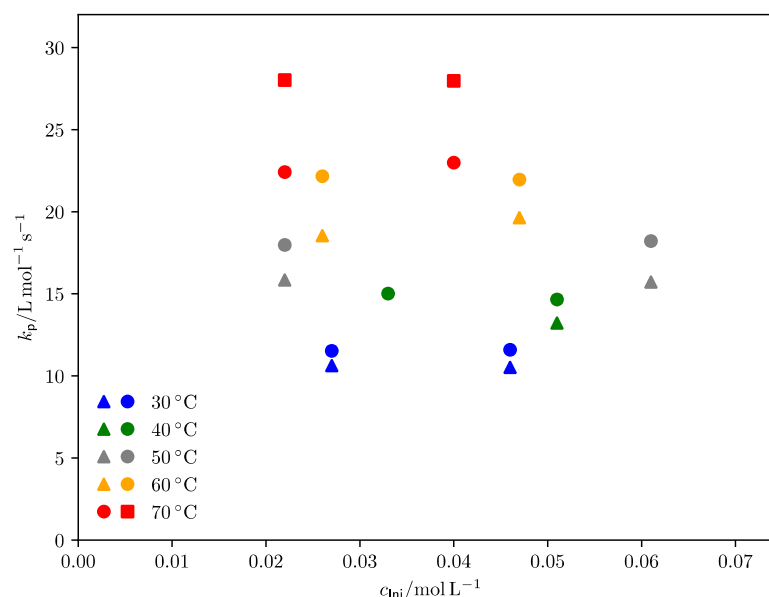


Figure 3. Obtained propagation rate coefficients, k_p , for DEI as function of initiator concentration, c_{ini} . An average of the k_p values extracted from the different inflection points is used here. Triangles refer to a laser repetition rate of 0.5 Hz, while circles refer to 1 Hz, and squares to 2 Hz.

In order to extract the Arrhenius parameters from the obtained k_p data, a regular Arrhenius fit is employed (see Figure 4). The observed chain-length dependency of k_p is not taken into account, but is accepted as uncertainty; consequently, the Arrhenius parameters, given in Table 3, will be average values over the whole chain length regime. It can already be noted here that these Arrhenius parameters fit very well to the ones obtained for other itaconates, but this will be discussed in greater detail below.

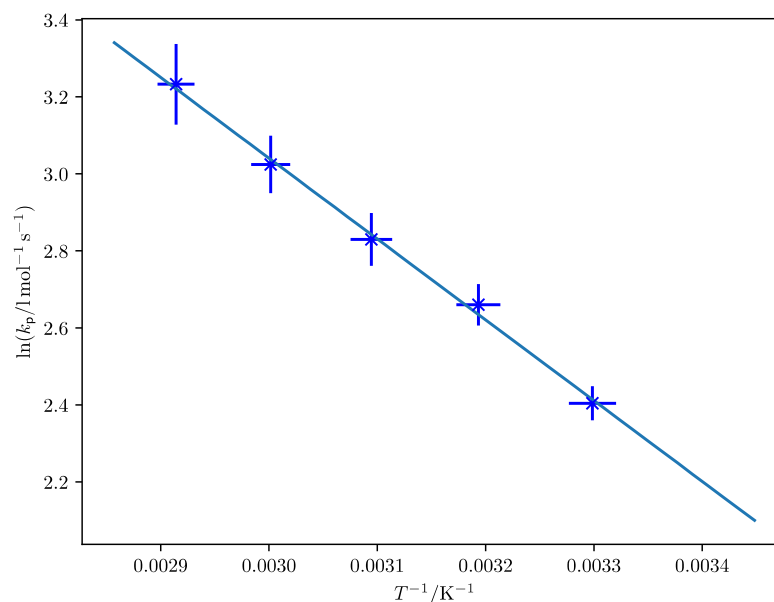


Figure 4. Arrhenius fit for k_p of DEI. The indicated uncertainties in $\ln(k_p)$ result from the standard deviation of the measurements for the respective temperature, while for the uncertainty in T^{-1} an uncertainty of $\Delta T = 2$ K was assumed. The obtained Arrhenius parameters are shown in Table 3.

Table 3. Experimental Arrhenius parameters for chain-length-averaged k_p of DEI, obtained via PLP-SEC.

A [$\text{L}\cdot\text{mol}^{-1}\cdot\text{s}^{-1}$]	E_A [$\text{kJ}\cdot\text{mol}^{-1}$]
$(1.1 \pm 0.3) \cdot 10^4$	17.5 ± 0.6

3.2.2. Di-n-Propyl Itaconate

Other than for DEI, only two stock solutions of DnPI with different concentrations of DMPA were prepared. The pulse energy was not recorded explicitly during the irradiation but was in a similar range as for the DEI experiments. For all samples, 1000 pulses were irradiated onto the sample, guaranteeing that the overall monomer conversion was low, as required for a successful PLP-SEC-experiment. The resulting PLP-data are collated in Table 4.

Table 4. PLP data for DnPI: initiator concentration, c_{ini} , temperature T , total number of pulses, n_{pulses} , laser pulse frequency, ν , molecular weight at the various inflection points, M_{inf} , and the average k_p of all obtained inflection points.

c_{ini} [$\text{mol}\cdot\text{L}^{-1}$]	T/K	ν [Hz]	M_{inf1} [$\text{g}\cdot\text{mol}^{-1}$]	M_{inf2} [$\text{g}\cdot\text{mol}^{-1}$]	M_{inf3} [$\text{g}\cdot\text{mol}^{-1}$]	k_p [$\text{L}\cdot\text{mol}^{-1}\cdot\text{s}^{-1}$]
0.074	333.15	0.25	70,800	129,100	-	16.5
0.049	333.15	0.25	70,200	127,900	193,900	16.2
0.074	333.15	0.5	41,300	72,800	111,900	18.7
0.049	333.15	0.5	41,100	72,500	110,000	18.5
0.074	323.15	0.25	61,800	115,200	-	14.6
0.049	323.15	0.25	61,200	113,500	-	14.4
0.074	323.15	0.5	33,930	61,600	-	15.8
0.049	323.15	0.5	35,000	61,000	-	16.0
0.074	313.15	0.25	47,200	92,800	-	11.4
0.049	313.15	0.25	50,500	93,400	136,800	11.6
0.074	313.15	0.5	26,200	49,800	-	12.5
0.049	313.15	0.5	28,400	48,400	-	12.8
0.049	303.15	0.5	19,400	39,100	-	9.5
0.074	293.15	0.25	12,100	22,400	32,700	2.8
0.049	293.15	0.25	12,400	20,700	32,100	2.7
0.074	293.15	0.5	16,100	29,600	-	7.5
0.049	293.15	0.5	14,300	30,100	-	7.1

The same consistency criteria as for DEI were checked (see Figure 5). When inspecting the ratio of molar masses of the inflection points, clearly a higher scattering is observed compared to DEI. The data for the ratio of peaks 2 and 1 also shows that most samples are well below the anticipated value of 2, again pointing towards a chain-length dependency of k_p .

As can be seen in Figure 6, just as for DEI, the obtained k_p values are independent of initiator concentration and increase with increasing laser repetition rate. The same reasoning as for DEI can be applied here.

Since for 30 °C only one data point exists, the uncertainty for the Arrhenius fit, which is shown in Figure 7, was assumed to be similar to that of the other data points.

Table 5. Experimental Arrhenius parameters for chain-length-averaged k_p of DnPI, obtained via PLP-SEC.

A [$\text{L}\cdot\text{mol}^{-1}\cdot\text{s}^{-1}$]	E_A [$\text{kJ}\cdot\text{mol}^{-1}$]
$(1.0 \pm 0.5) \cdot 10^4$	17.5 ± 1.2

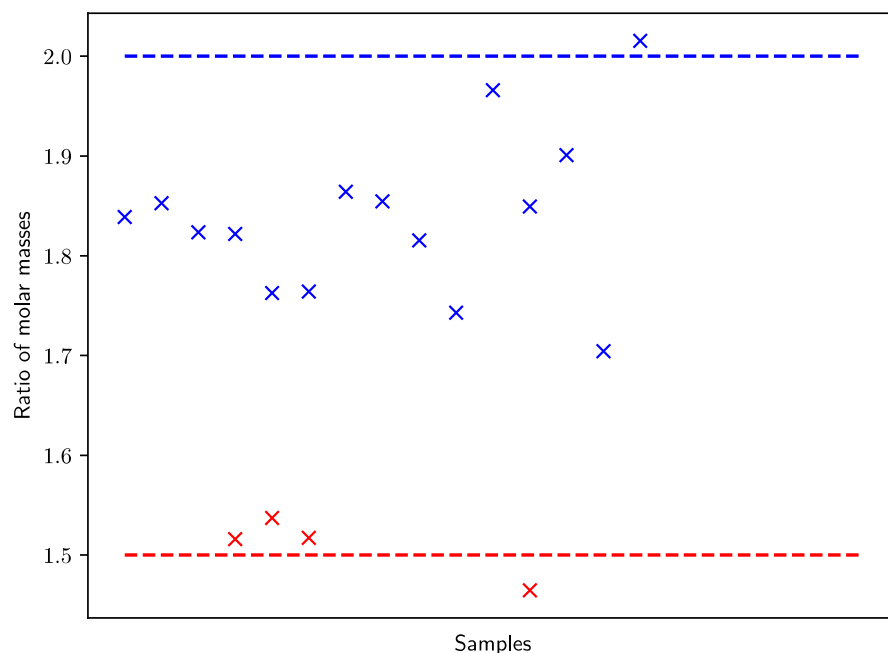


Figure 5. Ratio of molar masses of inflection points for all DnPI samples. Blue points show the ratio of the second inflection point to the first one, while the red points show the ratio of the third inflection point to the second one. Data points that significantly deviate from the expected ratios have been omitted from the analysis according to the consistency criterion of PLP-SEC.

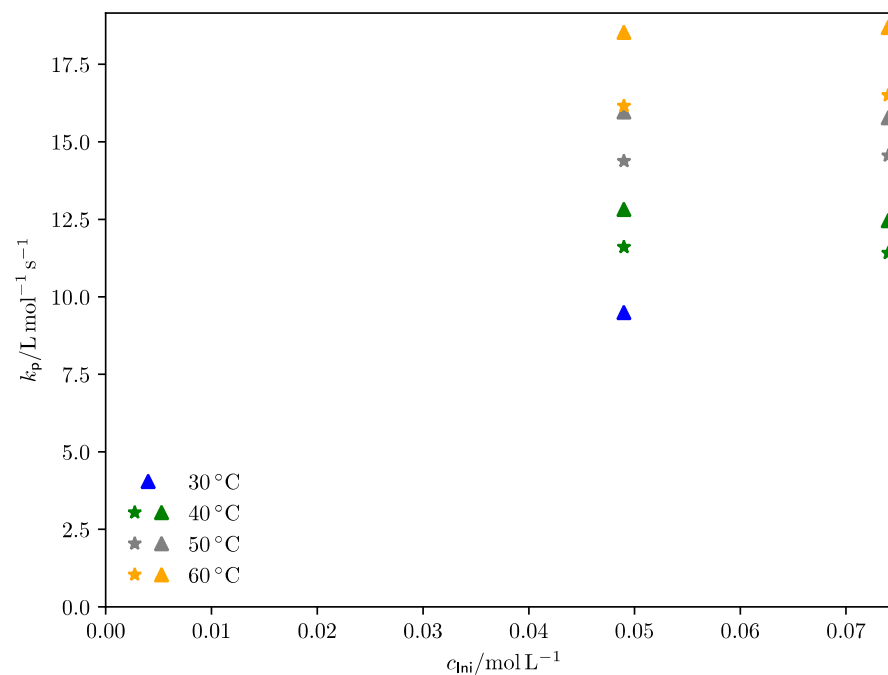


Figure 6. Obtained propagation rate coefficients for DnPI as a function of initiator concentration, c_{Ini} . An average of the k_p values extracted from the different inflection points is used here. Stars refer to a laser repetition rate of 0.25 Hz, while triangles refer to 0.5 Hz.

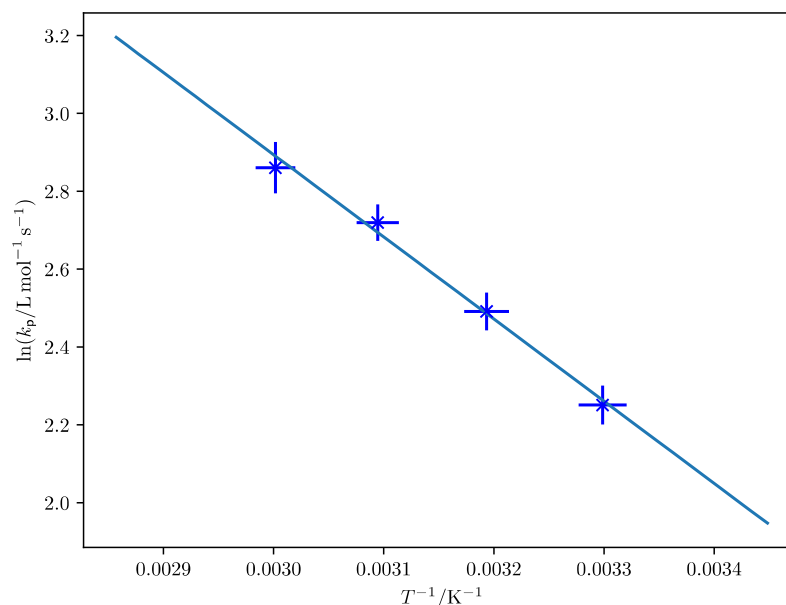


Figure 7. Arrhenius fit for k_p of DEI. The indicated uncertainties in $\ln(k_p)$ result from the standard deviation of the measurements for the respective temperature, while for the uncertainty in T^{-1} an uncertainty of $\Delta T = 2$ K was assumed. For 30 °C, only one data point exists, consequently the uncertainty in $\ln(k_p)$ was assumed to be similar to the other values. The obtained Arrhenius parameters are shown in Table 5.

4. Discussion

Inspecting the experimental results of the Arrhenius fits, one notices that the results for DEI and DnPI are extremely close to each other. A correlation analysis for the Arrhenius parameters has not been performed, however. Within the series of itaconate esters (see Table 6), DEI and DnPI are on the lower end of prefactor and activation energy values.

Table 6. Comparison of Arrhenius prefactors and activation energies for k_p of different itaconic acid esters. DMI = Dimethyl itaconate, DnBI = Di-n-butylitaconate and DCHI = Dicyclohexylitaconate.

Substance	A [$10^5 \text{ L} \cdot \text{mol}^{-1} \cdot \text{s}^{-1}$]	E_A [$\text{kJ} \cdot \text{mol}^{-1}$]
DMI [6]	2.2	24.9
DMI [7]	7.3	27.8
DEI (this work)	0.11	17.5
DnPI (this work)	0.10	17.6
DnBI [7]	0.33	21.3
DCHI [7]	0.99	26.5
DCHI [8]	0.17	22.0

The quantum chemical calculations show good agreement with the experimental results. Compared to the experimental values, the prefactors match nearly perfectly on B3LYP/def2-TZVP level. The predicted activation energy is about 5 kJ/mol too high which is still a very good agreement, considering the rather low level of theory. This good agreement is likely due to error compensation.

Haehnel et al. [35] investigated the family behavior of linear alkyl methacrylate monomers. They stated that high frequency factors are accompanied by high values of the activation energy, making it difficult to identify any trends with Arrhenius parameters. Therefore, the actual values of k_p at a given temperature were used for revealing family behavior of monomers. In the case of linear alkyl methacrylates, k_p (at 50 °C) increases linearly with the length of the side chain [34]. For itaconates, we find an opposite trend: as indicated in Figure 8, k_p clearly decreases with increasing size of the ester groups. This

can possibly be reasoned by a pre-structuring effect of the monomer, as has also been put forward by Haehnel et al. [35]. For alkyl methacrylates, longer side chains lead to a more structured monomer bulk. Due to stronger dispersive interactions, the longer side chains align more pronouncedly with each other. This effect also aligns the olefinic bonds such that propagation is easily possible and promoted with ester groups size. Itaconates, however, have two side chains. This allows for two different patterns of alignment. The first is identical to the one discussed for methacrylates. The second pattern is not stacked, but shifted, in a zig-zag motif. This alignment increases the distance of the olefinic bonds, leading to decreased propagation rates. Apparently for itaconates, the second pattern is preferred. This interesting effect might be confirmed by MD simulations in the future.

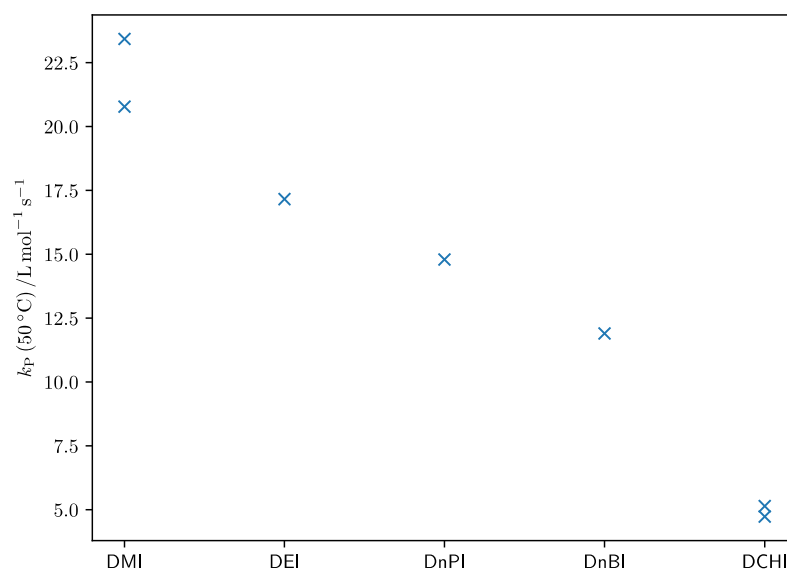


Figure 8. Values of k_p at 50 °C for the family of itaconates using Arrhenius parameters from Table 6.

5. Conclusions

Arrhenius expressions for the propagation step in free-radical polymerization of diethyl itaconate (DEI) as well as di-n-propyl itaconate (DnPI) in bulk, for which propagation kinetics was yet unexplored, were determined via pulsed-laser polymerization in conjunction with size-exclusion chromatography (PLP-SEC) experiments in the temperature range of 20 to 70 °C. The values were found to be very similar for both monomers. Together with literature data for other itaconate monomers, a family behavior was found, where k_p clearly decreases with increasing size of the ester groups. This behavior is different to other monomers and can possibly be explained by a pre-structuring effect of these rather sterically hindered monomers.

Author Contributions: Conceptualization, E.M.; methodology, E.M. and T.W.; software, E.M.; validation, E.M.; formal analysis, E.M. and T.W.; investigation, E.M. and T.W.; writing—original draft preparation, E.M.; writing—review and editing, P.V.; visualization, E.M.; supervision, P.V. All authors have read and agreed to the published version of the manuscript.

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Data Availability Statement: The data presented in this study are available in this work.

Conflicts of Interest: The authors declare no conflict of interest.

Appendix A

Formulae for Calculating the Individual Entropy Components

$$S_{trans} = R \left(\ln \left(\left(\frac{2\pi M k_B T}{h^2} \right)^{\frac{3}{2}} \frac{k_B T}{p} \right) + 1 + \frac{3}{2} \right)$$

$$S_{rot} = R \left(\ln \left(\frac{\pi^{\frac{1}{2}}}{\sigma_r} \left(\frac{T^{\frac{3}{2}}}{\Theta_{r,x} \Theta_{r,y} \Theta_{r,z}} \right) \right) + \frac{3}{2} \right)$$

$$S_{vib} = R \sum_i \left(\frac{\frac{h\nu_i}{k_B T}}{\exp\left(\frac{h\nu_i}{k_B T}\right) - 1} - \ln \left(1 - \exp\left(\frac{h\nu_i}{k_B T}\right) \right) \right)$$

Appendix B

Structures of the Reactants, Transition State and Product on UHF/6-31g(d) Level of Theory

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Reactants UHF/6-31G(d)

C -4.40972401531386 -3.71503806114475 -4.19211604628649
 C -3.68133265382649 -2.50113964837613 -3.82846057294931
 C -2.18555069312768 -2.58287955867867 -3.74895926164922
 H -1.81956361353836 -1.90927486262867 -2.97994166704005
 C -1.55123355560259 -2.21500407854842 -5.08428872443647
 H -1.85981769088834 -3.58155274742932 -3.48373455760964
 O -3.87993817666339 -4.76861766065154 -4.41207592473984
 O -2.19867266439477 -2.02984556122397 -6.06256526758899
 O -5.72557494763414 -3.54020745632567 -4.24515591923209
 C -6.54017447797737 -4.64564347728663 -4.61958058930150
 O -0.23107825111582 -2.11988533201135 -5.14587436878297
 C 0.62860159593704 -2.24965352535422 -4.01615485238005
 C -6.84381429855741 -5.54313533752586 -3.43397987468031
 H -6.05382188123488 -5.19375753995276 -5.41217021980953
 H -7.44408010230698 -4.19722006485559 -5.00595428455019
 C 2.05126359981739 -2.21283246002602 -4.53134556977455
 H 0.43228935067523 -3.17915492013272 -3.50505209750336
 H 0.44776061679191 -1.42471687501624 -3.33574782757450
 H -7.52779037380381 -6.33091526848121 -3.73515231333415
 H -7.30921982574788 -4.97521989160534 -2.63537015695070
 H -5.93929548340384 -6.00294157450675 -3.05616604113733
 H 2.22599638910386 -3.03416206606115 -5.21539910098320
 H 2.74329238575857 -2.31283602695059 -3.70356746094692
 H 2.24579827549218 -1.28111638750133 -5.05035417824159
 C 1.41023686313188 -7.39479693429179 -3.21670814323402
 C 1.76218588985207 -7.00859229226371 -1.80537284102727
 C 0.40030108509587 -6.50567479832831 -3.86586886936411
 C 2.51335328251719 -5.68767506802261 -1.74174766720343
 H 2.39075946646450 -7.77227607069698 -1.35759309022347
 H 0.86155241563142 -6.94414051075366 -1.20845866996825
 C 1.90936515216638 -8.44863568539955 -3.83552909442938
 O 2.96604170334251 -5.15779889207647 -2.69980821647521
 O -0.13492953713609 -5.62162351698555 -3.26298085202178
 O 0.15592786208118 -6.79625093352468 -5.12018673812008
 C -0.78815096521645 -5.98171011233685 -5.82956890531792
 O 2.69679684031063 -5.13549053655022 -0.54522744993166
 C 2.15076670664725 -5.65150973860652 0.65601647843522

```

C -0.99407495047473 -6.61281636474743 -7.18859925965537
H -0.37738582773625 -4.98579544880722 -5.91623978867437
H -1.70513446640562 -5.92923065237836 -5.26468399659097
C 2.67742937800976 -4.79518178786613 1.78889817293130
H 2.44840934259808 -6.68468532830934 0.79052106358879
H 1.06902617152998 -5.60801958172911 0.61572335119833
H -0.06041114415965 -6.66737370844859 -7.73713970390276
H -1.69456867264684 -6.01422677361150 -7.76061057091692
H -1.39699203958323 -7.61500923783392 -7.09171340373149
H 3.75936486793610 -4.83910902452768 1.82657350671995
H 2.38120460531527 -3.76220599922446 1.65136599612789
H 2.28125101852981 -5.14748058512089 2.73574103607832
C -4.37617783051305 -1.18068970392636 -3.63246950175260
H 2.62660274641475 -9.09041746684623 -3.35615376276326
H 1.62192414224581 -8.69793813042399 -4.83783698934764
C -4.86717639532277 -0.97403239860796 -2.19209531994862
H -5.21608532866781 -1.10596935176055 -4.31090971943407
H -3.68620846515025 -0.38550755085006 -3.89743691175879
H -5.59402029389089 -1.73139093368148 -1.91991355137942
H -5.33872079244902 -0.00206617048913 -2.08720830723050
H -4.04695233890537 -1.02343832869364 -1.48156140319029

```

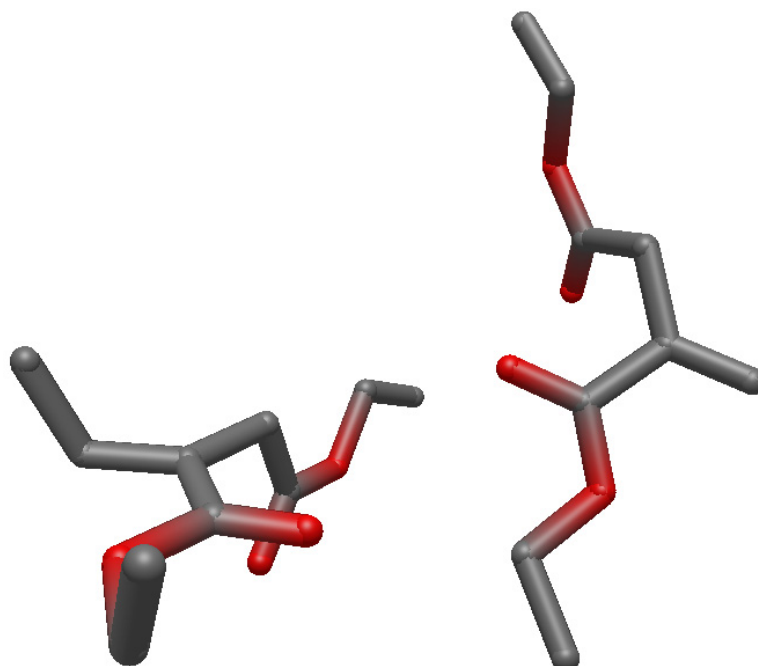


Figure A1. Structure of the reactants on UHF/6-31G(d) level of theory.

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TS UHF/6-31G(d)

```

C -2.84758329469564 -0.60649978656624 -0.35173091904958
C -3.47487656697929 0.67281650602152 -0.69218860128065
C -3.55472220027557 1.66082573684066 0.44012508768508
H -4.15851870941808 2.51469364647842 0.14752602094326
C -2.17076736081956 2.21218446327552 0.76780052421886
H -4.02699926242146 1.20732082219422 1.29793260123681
O -2.37853165100162 -0.81041029727699 0.73762180050786
O -1.41235165557832 2.49882171602563 -0.10039626506317
O -2.84933008160362 -1.48558407651005 -1.33176369020927

```

C -2.19960394673800 -2.74805070080274 -1.12824762851276
O -1.83173239905996 2.43331079819363 2.02868690080066
C -2.55087438021810 1.92566714844963 3.14350614153460
C -3.11843172362171 -3.73239536663273 -0.43066657627182
H -1.28476954398572 -2.59602904116585 -0.58112513653719
H -1.95694673813358 -3.08261012013724 -2.12545612828486
C -1.74871064876934 2.27094685075197 4.38067005218506
H -2.66491389394941 0.85539962284188 3.04179449838902
H -3.53222003290299 2.38406925600524 3.18726109676217
H -2.62772884930282 -4.69814852441827 -0.36306955244555
H -4.04290923096498 -3.85660274198890 -0.98517311104905
H -3.35315580445981 -3.40034850572737 0.57317343240027
H -0.76717665857079 1.81482615852842 4.33393178450831
H -2.25947451406035 1.90665701863723 5.26621241567900
H -1.62511979519397 3.34401516574444 4.46827123098726
C 2.45871053755559 -0.26177323025727 0.56762667061529
C 2.79559014608955 -1.71516497412086 0.87362760554040
C 3.61166865724519 0.64959838060042 0.31640811668741
C 2.36276327964106 -2.64269924413079 -0.25522296185348
H 2.22573018964650 -2.03420907136342 1.73606841408987
H 3.84303814407620 -1.80421504547035 1.10979341341543
C 1.21370556363695 0.18845387067221 0.53453341572562
O 1.22302307625872 -2.95947287516606 -0.37567425946026
O 4.75138414735338 0.28560266220910 0.31494859834329
O 3.25990506702815 1.90066757854423 0.09691365205411
C 4.29380168181435 2.84598959018137 -0.15811209860918
O 3.25092318791082 -3.09808863238834 -1.12612844328230
C 4.65135217212255 -2.86463776788425 -1.03645456782551
C 3.63156516776283 4.19033256304175 -0.37090343771677
H 4.85059252285008 2.53064888670376 -1.02987695753387
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C 5.27640883314746 -3.46428606983801 -2.27856471421145
H 5.03901151894821 -3.34541579398755 -0.14555211646049
H 4.85187306822269 -1.80671745273000 -0.96978488472731
H 2.94832090725509 4.15365783576250 -1.21135921475387
H 4.38626883919165 4.94414590789859 -0.57156183280494
H 3.07299069277300 4.48727374435213 0.50915042355221
H 5.06662336133763 -4.52580732703888 -2.33870880032320
H 6.35211780068496 -3.32282494115468 -2.25478158261550
H 4.88443282638933 -2.98689481786365 -3.16908667469771
C -3.97084269224552 1.00933932810088 -2.06935207777060
H 0.38495916637729 -0.47045358745645 0.71628589321382
H 0.98343794180324 1.21339663085629 0.32059547184296
C -5.47876335273979 0.76682926161363 -2.23434806872066
H -3.43270737829562 0.42171878332917 -2.80110921489571
H -3.74481810647014 2.05270880870212 -2.26764947487017
H -5.71983714072145 -0.27989805219422 -2.08236802290827
H -5.80030204259206 1.04437265979783 -3.23290135138030
H -6.05723300624604 1.35210768974649 -1.52524402102384

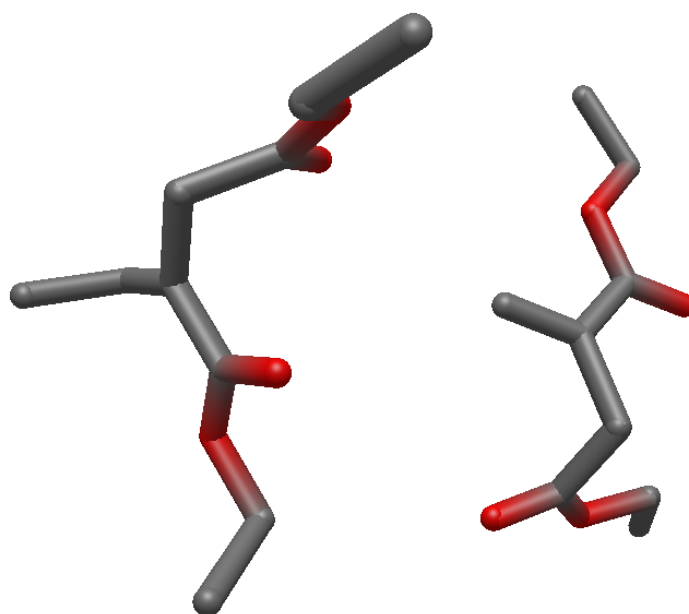


Figure A2. Structure of the transition state on UHF/6-31G(d) level of theory.

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Product UHF/6-31G(d)

C -5.33010044431225 -5.17468805041027 -2.76289465235525
C -4.78315772241794 -3.74187542846288 -2.83238392739415
C -4.91490745292611 -3.16240216496752 -1.40859455783708
H -5.94684554863516 -3.24388163191316 -1.08693595720722
C -4.49533849687114 -1.71842485435190 -1.20853512882611
H -4.36388407521240 -3.78964728582604 -0.72205388163776
O -5.58232282598512 -5.74716981209004 -1.74444107985903
O -4.23777261302297 -0.96449508470472 -2.08821498403281
O -5.47736666409740 -5.70975276825688 -3.95247798831326
C -5.84141129580070 -7.08827845545773 -4.05745733149766
O -4.44462443011922 -1.26384739619001 0.04233381138976
C -4.68964894858728 -2.02009241215792 1.21933639715133
C -7.34751799177768 -7.26186776811452 -4.01977300066089
H -5.35587379264847 -7.63996552105348 -3.26963765395188
H -5.42893426724061 -7.40053252050025 -5.00366130343885
C -3.39021667860755 -2.51346724885203 1.82991718662646
H -5.37280836558159 -2.83407224254007 1.02766841874399
H -5.18154748745600 -1.33074205897361 1.89050822281558
H -7.59816478269545 -8.30664551845971 -4.17670802357371
H -7.82162518208360 -6.67825133235716 -4.80173076367217
H -7.74968835253830 -6.95655894609922 -3.06141754447941
H -2.86885690800861 -3.20089536207003 1.17303179101418
H -3.59294050548452 -3.02702278893814 2.76445031994558
H -2.73133992171486 -1.67742526636710 2.03373806871067
C -2.25638271644363 -4.29802356482885 -2.34672033075670
C -2.14994288261521 -5.75792961322905 -1.95614587436087
C -1.27560572228266 -3.38516847385001 -1.73135212323383
C -1.74501007255527 -6.62518498669894 -3.14442466679776
H -3.11702183256505 -6.12824495221061 -1.64444855005516
H -1.47069622074738 -5.85160883143027 -1.12546239144077
C -3.29627120196457 -3.79678101600397 -3.32477710297106
O -2.55018701171806 -6.97278116004941 -3.94624601760557

O -0.52009825073052 -3.72531508379593 -0.86073931989745
O -1.30784908742213 -2.16235825899936 -2.21815076002935
C -0.45753853908387 -1.17878825589158 -1.63106437802248
O -0.47828520965378 -6.97610937240425 -3.30670509988317
C 0.55342894625031 -6.71420088237352 -2.36305862143452
C -0.72664336509408 0.12928609417126 -2.34257114085629
H 0.57036870257388 -1.49752056363488 -1.73829056138870
H -0.68133578016605 -1.11308157749710 -0.57537710329629
C 1.84663386032022 -7.20646861946591 -2.97771857588070
H 0.34209716384925 -7.24440752572095 -1.44200542427960
H 0.60382294400404 -5.65799792058409 -2.14495476578319
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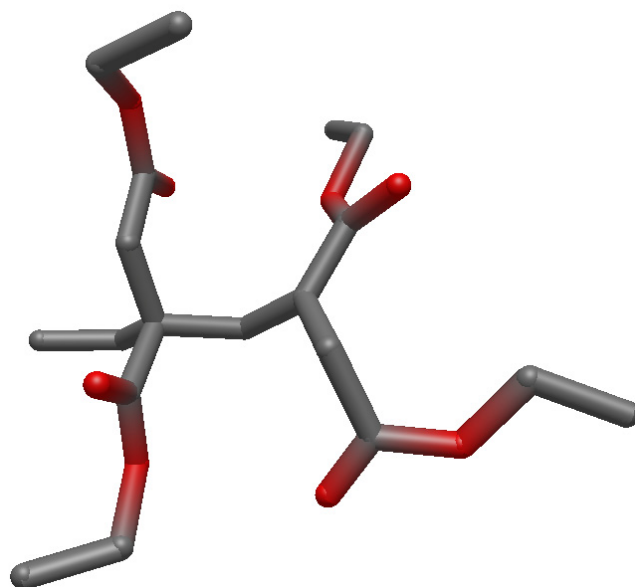


Figure A3. Structure of the product on UHF/6-31G(d) level of theory.

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