The Multifunctionality of Lanthanum−**Strontium Cobaltite Nanopowder: High-Pressure Magnetic Studies and Excellent Electrocatalytic Properties for OER**

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ABSTRACT: Simultaneous study of magnetic and electrocatalytic properties of cobaltites under extreme conditions expands the understanding of physical and chemical processes proceeding in them with the possibility of their further practical application. Therefore, La_{0.6}Sr_{0.4}CoO₃ (LSCO) nanopowders were synthesized at different annealing temperatures *t*_{ann} = 850–900 °C, and their multifunctional properties were studied comprehensively. As *t*ann increases, the rhombohedral perovskite structure of the LSCO becomes more single-phase, whereas the average particle size and dispersion grow. Co^{3+} and Co^{4+} are the major components. It has been found that LSCO-900 shows two main Curie temperatures, T_{C1} and T_{C2} , associated with a particle size distribution. As pressure *P* increases, average $\langle T_{C1} \rangle$ and $\langle T_{C2} \rangle$ increase from 253 and 175 K under ambient pressure to 268 and 180 K under *P* = 0.8 GPa, respectively. The increment of $\langle dT_C/dP \rangle$ for the smaller and bigger particles is sufficiently high and equals 10 and 13 K/GPa, respectively. The magnetocaloric effect in the LSCO-900 nanopowder demonstrates an extremely wide peak $\delta T_{\text{fwhm}} > 50$ K that can be used as one of the composite components, expanding its working temperature window. Moreover, all LSCO samples showed excellent electrocatalytic performance for the oxygen evolution reaction (OER) process (overpotentials of only 265−285 mV at a current density of 10 mA cm⁻²) with minimal $η_{10}$ for LSCO-900. Based on the experimental data, it was concluded that the formation of a dense amorphous layer on the surface of the particles ensures high stability as a catalyst (at least 24 h) during electrolysis in 1 M KOH electrolyte.

KEYWORDS: *cobaltites, nanopowder, high pressure magnetic studies, OER, electrochemistry*

1. INTRODUCTION

Interest in lanthanum-doped strontium La1[−]*x*Sr*x*CoO3 cobaltites constantly increases due to their physical and chemical multifunctionality.¹ They demonstrate unique magnetic and transport properties 2 as well as charge, spin, and orbital ordering;³ spin-state transition;^{[4](#page-13-0)} electromagnetic phase sepa-ration;^{[5](#page-13-0)} and metal−insulator transition^{[6](#page-13-0)} with strong mutual spin, orbital, charge, and lattice coupling. It makes them ferromagnetic (FM), ferroelectric, and electrocatalytically active with high electrode performance and a giant magneto-resistance effect.^{[7](#page-13-0)} Moreover, they show high ionic and electronic conductivity, good catalytic properties, and high stability, δ mainly because of the low mobility of Sr on the

crystal surface and the formation of $Co⁴⁺$ ions and oxygen vacancies. Besides that, the excellent properties of cobaltites can be changed by external perturbations such as pressure and annealing process, i.e., "core−shell" ratio. All these attract great scientific interest to cobaltites and make them promising

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candidates for practical application as electrode materials for high-temperature solid oxide fuel cells, catalysts, current and magnetic field sensors, gas sensors, etc. $9,10$ $9,10$ $9,10$ For studying the electrocatalytic and magnetic properties, applying the nanoparticle materials is a very useful tool, because the influence of the surface-dependent reasons is dominant in both cases. However, there still needs to be clarified and unstudied physical and chemical processes as surface-dependent reasons for obtaining high electrode performance. A simultaneous study of physical and chemical properties is relevant to improving the catalytic properties of such electrodes. The possibilities of changing the size of cobaltite particles and their effect on electrocatalytic and magnetic properties are of great interest. Moreover, the contactless heating and cooling of cobaltites as electrodes, using the magnetocaloric effect (MCE), should influence their reaction rate. 11 In this regard, studying the MCE or the effect of heating under an alternating electromagnetic field is particularly interesting. Additionally, as shown recently in our previous work, 12 the perovskite ceramics, after electrocatalysis, exhibited a decrease in its magnetic moment, making it possible to control the depletion of the electrodes and predict their working time and performance.

Original $LaCoO₃$ has an antiferromagnetic (AFM) order and dilution with divalent Sr^{2+} up to $x = 0.4$ enhance its FM order making this $La_{0.6}Sr_{0.4}CoO_3$ compound fully FM with the highest Curie temperature $T_C \approx 230$ K and the closest to the room temperature among other La_{1−*x*}Sr_{*x*}CoO₃ compositions.^{[6](#page-13-0)} However, the T_C can significantly differ and be 246 K^{[13](#page-13-0)} or even 321 K^{[14](#page-13-0)} for the same $x = 0.4$. Moreover, besides classical FM double-exchange (DE) between Co^{3+} and Co^{4+} ions, the AFM superexchange interactions Co³⁺-Co³⁺ and Co⁴⁺-Co⁴⁺, and phase separation can coexist¹⁵ that makes this $La_{0.6}Sr_{0.4}CoO_3$ system more complicated along with the sophisticated spin states distribution. Additionally, there is no data about particle size distribution's influence on the magnetic properties, Curie temperature, and MCE of the $La_{0.6}Sr_{0.4}CoO_3$ nanopowders. The effect of high hydrostatic pressure on the functional properties is of particular interest. External pressure should modify the Co−O distance and Co−O−Co angle and, as a result, change exchange interactions, phase transition temperatures, degree of magnetic uniformity, and MCE parameters. All of these issues and unspecified moments need to be clarified.

Among the various functional properties of perovskite-based materials containing 3d-metals, a good electrocatalytic activity for oxygen release should also be noted.¹⁶ Such electrocatalysts are extremely important for energy conversion and storage, ecologically clean fuel (hydrogen) generation, and highly efficient current sources. However, a significant problem remains, such as the large overpotential and sluggish kinetics in the electrochemical reactions of the oxygen evolution reaction (OER) and oxygen reduction reaction, which slow the adoption of new technologies and commercial developments. From this perspective, one of the critical issues is the replacement of existing industrially expensive Ir/Ru-based catalysts with cheaper catalysts with similar characteristics. The leading search for such catalysts is quite successfully carried out among many compounds of transition metals or their combinations (borides, carbides, nitrides, phosphides, chalcogenides, oxides, or hydroxides, etc.). In addition, these multicomponent systems and composite materials containing carbon nanosheets or nanotubes have been actively studied in recent years as electrocatalysts for Zn-air batteries and water splitting.^{17,[18](#page-14-0)} It fully applies to Ln- and Co-containing complex oxides of the *ABO*₃-type, where for perovskites, *A*-positions usually correspond to a rare-earth or alkaline earth element, and *B* is commonly a transition metal or a combination thereof.[19](#page-14-0)−[28](#page-14-0) In general, it is considered that both types of metal influence the oxygen-catalytic activity of perovskites in the A - and B -sites.^{[29](#page-14-0)} However, the high electrocatalytic activity of perovskites for OER in an alkaline medium is primarily the result of the redox behavior of the hybridization of transition metal 3d and oxygen 2p orbits. 30 Numerical studies of perovskites as electrocatalysts revealed that, in addition to the chemical composition, activity in oxygen electrocatalysis reactions is also affected by a significant number of factors, namely, method and conditions of preparation of the material (annealing at different temperatures and in a controlled atmosphere can cause different types of defects on the surface); $22,31$ particle size, porosity, and shape (affects active surface area);^{20–[22,32](#page-14-0)} doping;^{[19,23,25](#page-14-0),[33,34](#page-14-0)} and amorphization of the surface. 35

It should be noted that, as demonstrated by Mefford et al., 36 for LSCO catalysts, oxygen vacancy defects are a crucial parameter in improving the electrocatalysis of oxygen on metal oxide surfaces, due to which they can control the physical parameters of the ion diffusion rate and reflect the basic electronic structure of the catalyst. Moreover, due to electrochemical processes or an aggressive alkaline electrolyte, complex oxides, including perovskites, $37,38$ $37,38$ can be transformed *in situ* to other compounds,^{[11,12](#page-13-0),[39](#page-14-0)} directly affecting the electrocatalytic activity during long-term electrolysis. Interestingly, although many perovskites remain crystalline during OER in alkaline media, some highly active compounds become amorphous.^{[38](#page-14-0)} As a result, it has been suggested that dynamic surface self-reconstruction is the key to highly active perovskite electrocatalysts for water splitting, 40 and later atomistic insights into activation and degradation of LSCO electrocatalysts under oxygen evolution conditions were considered. 41 For these reasons, when studying perovskites as electrocatalysts for the OER, evaluating their primary activity and conducting research after electrochemical testing is crucial.

Additionally, studying electrochemical activity and physical properties directly on the same samples is very important since the obtained data supplement each other and allow us to make more correct conclusions. Moreover, knowledge of the behavior of other physical properties, such as temperature changes of magnetic phase transitions and magnetoactive phenomena (MCE or magnetoresistance effect) under changing internal and external factors is essential as well when it is necessary to change the conditions of electrocatalytic applications.

Thus, in this paper, the cobaltite $La_{0.6}Sr_{0.4}CoO_3$ nanopowders have been obtained under different annealing temperatures $t_{\text{ann}} = 850$, 875, and 900 °C and their structure, morphology, particle size distribution, and magnetic, magnetocaloric, and electrocatalytic properties have been studied comprehensively. The most single-phase $La_{0.6}Sr_{0.4}CoO_3$ nanopowder with t_{ann} = 900 °C has been selected to study the phase transitions and magnetic properties under high hydrostatic pressure up to $P \approx 0.8$ GPa. The same sample showed excellent electrocatalytic activity and high stability for the OER in alkaline electrolytes.

Figure 1. XRD patterns of the LSCO measured at room temperature and fitted by the Rietveld method. The experimental (black cycles) and calculated (red up curve) values and a difference curve (blue bottom line) normalized to a statistical error are presented. Vertical bars are the calculated positions of diffraction peaks corresponding to the main rhombohedral *R*3̅*c* perovskite structure and rhombohedral *R*32 impurity $SrCo_{0.78}O_{2.48}$ in (a) LSCO-850, (b) LSCO-875, and (c) LSCO-900.

2. EXPERIMENTAL SECTION

2.1. Sample Preparation. Cobaltite La_{0.6}Sr_{0.4}CoO₃ (LSCO) nanopowders were obtained using the sol−gel method from the initial La(NO3)3·6H2O (99.9% trace metals basis, Sigma−Aldrich), Sr- (NO_3) ₂ (ACS reagent, ≥ 99.0%, Sigma–Aldrich) and Co (NO_3) ₂. $6H₂O$ (reagent grade, 98%, Sigma-Aldrich) reagents. The calculated amounts of nitrates and citric acid (the ratio of the molar sum of metals and $H_3Cit·H_2O$ was 1:1) were dissolved in deionized water. The solution was evaporated to form a homogeneous gel, dehydrated, and heated from 200 to 500 °C (the heating rate equals 50 °C h⁻¹). The obtained powder was ground in an agate mortar and annealed at temperatures of $t_{\text{ann}} = 850$, 875, and 900 °C for 10 h at each temperature. Finally, three samples obtained at different t_{ann} values were named LSCO-850, LSCO-875, and LSCO-900 for further investigation.

2.2. Characterization. The phase composition and size of coherent scattering regions of the samples were examined using the X-ray diffraction (XRD) method on a Shimadzu LabX XRD-6000 diffractometer in Cu_{K*a*1}-radiation (λ = 0.15406 nm) before electrocatalysis and X-ray diffractometer MicroMax-007 HF (Rigaku, Japan) in Mo- K_a radiation ($\lambda = 0.071146$ nm) after electrocatalysis, both at room temperature. The JCPDS PDF-2 database was used for defining the phase composition. The refinement of crystal structures was carried out with Rietveld analysis^{[42](#page-14-0)} using FullProf software.^{[43](#page-14-0)}

The morphology and size of particles were determined using the transmitting electron microscopy (TEM) method on a JEM-2200FS Transmission Electron Microscope and scanning electron microscopy (SEM) method on a FEI MAGELLAN 400 and Bruker XFlash 6|60 in the Hitachi Regulus 8100. The distribution function of nanoparticle size was obtained by analyzing SEM and TEM images using Nano

Measure 1.2.5 software. 11 Additionally, the chemical composition of the samples was clarified using energy-dispersive X-ray spectroscopy (EDS) mode.

X-ray photoelectron spectroscopy (XPS) measurements were performed on an ESCALAB 250 X-ray photoelectron spectrometer. The XPS spectra were fitted and analyzed by using OriginPro 2018 software. The background of the X-ray photoelectronic lines was cut by using the Shirley method. The spectra were excited by using monochromatized Al−K*α* radiation. The state of the surfaces was monitored through the C 1s line, allowing calibration of the energy scales for all spectra. The C 1s line binding energy was approximately 285 eV.

Magnetic measurements were carried out using a Quantum Design SQUID MPMS 3 instrument in a temperature range from 2 to 300 K and in a magnetic field up to 30 kOe. The temperature dependences of magnetization $M(T)$ under $H = 50$ Oe were carried out in two regimes: zero-field cooling (ZFC) and field cooling (FC). The heating and cooling of the samples were carried out with a constant rate of 1 K min[−]¹ . The MCE was determined using magnetization isotherms $M(H)$ near the Curie temperature T_C with a $\Delta T = 2$ K step. Each isotherm *M*(*H*) was measured with an increase in the magnetic field *H* from 0 to 10 kOe with a step of Δ*H* = 100 Oe. Before each measurement of the isotherms, the magnetic nanopowder was demagnetized. Additionally, magnetic measurements under high pressure up to $P \approx 0.8$ GPa were performed using a piston-type pressure cell made of Ni−Cr−Al alloy.^{44,45} Silicon oil with a low viscosity was used as a pressure-transmitting medium. The pressure inside was measured by using the pressure dependence of the superconducting transition temperature of high-purity lead.

2.3. Electrode Fabrication and Electrochemical Measurements. To fabricate the working electrodes, mixtures of LSCO samples, Super P (conductive carbon black), and PVDF (poly- (vinylidene fluoride)) were prepared in a weight ratio of 85:10:5, respectively. Next, NMP (*N*-methylpyrrolidone) was added to the mixture. After that, the liquid suspension was applied onto a carbon fiber (HCP331N) with a working area of 1 cm² and dried at 60 °C for 24 h in a vacuum oven. Each electrode contained approximately 30 mg of electroactive material. The benchmark $RuO₂/C$ electrode for investigating OER was made using a similar preparation procedure.

The electrochemical tests were conducted using an electrochemical workstation (CHI 760E, Shanghai, China). A typical three-electrode setup included an LSCO working electrode, a commercial Hg/HgO reference electrode, and a platinum foil counter electrode in 1 M KOH (pH 14). The linear sweep voltammetry (LSV) curves were collected at a scan rate of 1 mV s^{-1} with *iR*-compensation. Electrochemical impedance spectroscopy (EIS) was obtained in the frequency range from 0.1 to 100 kHz with an amplitude of 5 mV at the open circuit potential. Chronopotentiometry (CP) was carried out under a constant current density of 10 mA cm^{-2} for 24 h. All the potentials versus Hg/HgO were converted to versus reversible hydrogen electrode (RHE): $(E_{\text{RHE}} = E_{\text{Hg/HgO}} + 0.098 + 0.059 \text{pH}).$

3. RESULTS AND DISCUSSION

3.1. Structural Properties. XRD patterns of the LSCO nanopowders are presented in [Figure](#page-2-0) 1. The main crystalline phase is a perovskite crystal structure with a rhombohedral *R*3̅*c* space group. Some of the second phase was also detected as $SrCo_{0.78}O_{2.48}$, which decreases with increasing t_{ann} from 17.4% for LSCO-850 to 14.8% for LSCO-875 and 8% for LSCO-900. It should be noted that this impurity with increased content of cationic $Co_{0.78}$ and anionic $O_{2.48}$ vacancies does not influence the magnetic properties of the main LSCO phase near its high Curie temperature $T_C \approx 250$ K (see the following [Section](#page-4-0) 3.2.) because the closest $SrCoO_{2.5}$ composition to our impurity is AFM with high Neel temperature T_N = 570 K.^{[46](#page-14-0)} Moreover, the less content of oxygen and cobalt, the much lower T_C compared with our $T_C \approx 250$ K and much weaker FM interactions.[3,](#page-13-0)[47](#page-14-0),[48](#page-14-0) The main lattice parameters are listed in Table 1 and are consistent with other literature data.^{[13](#page-13-0)[,49](#page-14-0)} With increasing t_{ann} , a unit cell volume *V* is increased nonmonotonically, indicating the successive replacement of smaller La^{3+} cation ($r_{\rm La}^{3+}$ = 0.136 nm) by bigger $\rm Sr^{2+}$ cation ($r_{\rm Sr}^{2+}$ = 0.144 nm) in 12-fold oxygen coordination.^{[50](#page-15-0)} At the same time, the average size of coherent scattering regions D_{XRD} , determined via the Scherrer equation, 51 increases with increasing t_{ann} from 48 nm for the LSCO-850 to 54 nm for the LSCO-875 and 60 nm for the LSCO-900 (Table 2 and Supporting [Information](https://pubs.acs.org/doi/suppl/10.1021/acsami.3c06413/suppl_file/am3c06413_si_001.pdf) SI [1](https://pubs.acs.org/doi/suppl/10.1021/acsami.3c06413/suppl_file/am3c06413_si_001.pdf)).

The morphology and particle size distribution of the LSCO were studied using SEM and TEM data (see [Figures](#page-4-0) 2 and [3](#page-5-0)).

Figure 2. SEM and EDS images of the LSCO compositions obtained at $t_{ann} = (a, b) 850$, (c, d) 875, and (e, f) 900 °C. Insets show particles' distribution with an average particle size D_{SEM} .

All LSCO nanopowders demonstrate a well-developed structural nature with a spherical-like shape of nanoparticles. An average particle size *D* and their distribution (see [Table](#page-3-0) 2) were defined using a detailed approach described in our previous work.^{[11](#page-13-0)} With increasing t_{ann} , the particle size is increased from D_{SEM} = 74 nm and D_{TEM} = 39 nm for LSCO-850 to 75 and 44 nm for LSCO-875 and 79 and 47 nm for LSCO-900, which correlates with the XRD data. The differences between SEM and TEM data are possible because of insufficient data sets by size and a more accurate determination approach of the average particle size based on

the TEM method using a smaller 100−200 nm scale and having less error with a higher coefficient of determination *R*² (see [Table](#page-3-0) 2). It should be noted that dispersion σ is relatively high, indicating the existence of particles with a big difference in their sizes. Additionally, the chemical composition of all LSCO is approximately confirmed as 0.6:0.4:1 according to the EDS data (see Figure 2b, d, f).

3.2. Magnetic Properties. The field dependences of magnetization $M(H)$ for LSCO-900 at 2, 77, and 300 K are presented in [Figure](#page-5-0) 4. The LSCO-900 is in FM at 2 and 77 K and PM at 300 K states. At 2 K, the spontaneous

Figure 3. TEM images of the LSCO nanopowders obtained at $t_{\text{ann}} = (a) 850, (b) 875,$ and (c) 900 °C. Insets show particles' distribution with an average particle size D_{TEM} .

Figure 4. Field dependences of magnetization *M*(*H*) at 2, 77, and 300 K for the LSCO-900. The inset shows an enlarged area of the *M*(*H*).

magnetization M_S achieves 10.5 emu/g, and there is no saturation even under magnetic field 3 T that indicates competing between FM DE $Co³⁺-Co⁴⁺$ and AFM superexchange $Co^{3+} - Co^{3+}$ and $Co^{4+} - Co^{4+}$ interactions.^{[52](#page-15-0)} At the same time, Co ions can be in high-spin (HS), low-spin (LS), and/or even intermediate-spin (IS) states.⁵² Based on the

Table 3. Main Magnetic Parameters: Spontaneous Magnetization *M*_S, Residual Magnetization *M*_r, Coercivity *H*_C, Characteristic Curie Temperatures T_{C1} , T_{C2} , T_{C3} , T_{C4} , AFM Transition Temperature T_v and Blocking Temperature T_B under Different Pressures *P* for LSCO-900

	FM						near phase transitions					
	M_s (emu/g)		$M_{\rm r}$ (emu/g)		$H_C(kOe)$		FC	ZFC	FC	ZFC		
P(GPa)	2 K	77 K	2 K	77 K	2 K	77 K	$T_{C1}(K)$	$T_{C2}(K)$	$T_{C3}(K)$	$T_{C4}(K)$	$T_{AFM}(K)$	$T_{\rm B}({\rm K})$
$\mathbf{0}$	10.5	8.7	7.2	4.9	4.3	1.4	251	255	163	187	145	169
0.16						-	259	261	161	186	150	164
0.43							260	263	162	190	153	168
0.64	-						265	267	165	191	156	170
0.76							266	269	166	194	158	170

Figure 5. Temperature dependences of the magnetization (a) $M_{\text{FC}}(T)$ and (b) $M_{\text{ZFC}}(T)$ for LSCO-900 under different pressures *P* and in the field of *H* = 50 Oe. The insets show the defined characteristic Curie temperatures (a) T_{C1} and T_{C3} and (b) T_{C2} and T_{C4} , as well as (b) AFM transition temperature T_{AFM} and blocking temperature T_{B} .

experimentally defined FM moment $\mu_{\text{FM}}^{\text{exp}} = 0.42 \mu_{\text{B}}$ and other $data$, 2,6,47 2,6,47 2,6,47 2,6,47 2,6,47 2,6,47 2,6,47 it can be concluded that the most probable spin state for Co^{3+} and Co^{4+} ions are HS $(t_{2g}^4e_g^2)$ with $S = 4/2$ and LS $(t_{2g}^5e_g^0)$ with $S = 1/2$ states, respectively. Theoretical calculations give μ_{FM}^{theor} = 2.8 μ_B . The big difference between theoretical and experimental data may be associated with competing for different magnetic interactions and the presence of oxygen vacancies[.44,](#page-14-0)[53](#page-15-0),[54](#page-15-0) The main magnetic parameters are given in Table 3. It should be noted that LSCO-900 is a ferromagnet with a strong coercivity $H_C = 4.3$ kOe and residual magnetization $M_r = 7.2$ emu/g at 2 K (see the inset of [Figure](#page-5-0) [4](#page-5-0)).

The temperature dependences of magnetization $M_{\text{ZFC}}(T)$ and $M_{\text{FC}}(T)$ for the LSCO-900 under different pressures P and in the magnetic field $H = 50$ Oe (Figure 5) correspond to a typical behavior of the nanoparticles with FM and AFM inclusions, as well as with randomly oriented uniaxial magnetic anisotropy.[55](#page-15-0) From the temperature dependences of the extremum of the derivative of magnetic susceptibility d(*M*/ *H*)/d*T*, where $\chi = M/H$ (see the insets in Figure 5), several characteristic phase transition temperatures were determined (Table 3). Noteworthy, four characteristic Curie temperatures

 T_{C1} = 251 K and T_{C3} = 163 K from FC and T_{C2} = 255 K and T_{C4} = 187 K from ZFC curves were defined as an example for the LSCO-900 without pressure $P = 0$ (see the insets in Figure 5). The observed difference between ZFC and FC Curie temperatures is associated with "pinning" magnetic moments on the magnetic inhomogeneities and structural defects, as well as the competing different magnetic exchange interactions.[10](#page-13-0),[12,](#page-13-0)[53,56](#page-15-0) The obtained average ZFC and FC Curie temperatures 175 and 253 K indicate the presence of smaller and bigger nanoparticles in the LSCO-900. $11,57$ $11,57$ The highest Curie temperature, 253 K (*P* = 0), complies with other literature data 230 and 246 K for the same ceramic bulk composition.[6](#page-13-0) An additional contribution to the ZFC and FC curves at *T* < 187 K is due to the magnetization from smaller particles. As shown,^{[11](#page-13-0),[57](#page-15-0)} a change in the Curie temperature T_C value is caused by the particle size distribution. In our case, the change of the average Curie temperatures ∼ 1.4 times corresponds to the existence of bigger $D_{\text{XRD}} \approx 60$ nm (see [Table](#page-3-0) 2) with $T_C = 253$ K and smaller $D_{\text{XRD}} \approx 43$ nm with T_C $= 175$ K nanoparticles. Moreover, the blocking temperature T_B = 169 K, determined from the maximum temperature on the $M_{\text{ZFC}}(T)$ curve, is observed for LSCO-900 ($P = 0$). At $T < T_{\text{B}}$,

Figure 6. Pressure dependences of characteristic Curie temperatures (a) T_{C1} and T_{C2} and (b) T_{C3} and T_{C4} defined from the ZFC-FC curves for the LSCO-900. The dashed lines show linear behavior $T_{\rm C}$ versus *P*.

Figure 7. (a) Isotherms of magnetization $M(H)$ and (b) Arrott's plots $M^2(H/M)$ within 200–300 K with a $\Delta T = 2$ K step for the LSCO-900.

Figure 8. Temperature dependences of magnetic entropy change $-\Delta S_M(T)$ (a) without pressure, *P* = 0, and (b) under high pressure, *P* ≈ 0.8 GPa in the magnetic field Δ*H* up to 10 kOe for the LSCO-900.

the magnetic nanoparticles are in a blocked state. At $T_B \le T \le T$ T_{C2} , the thermal fluctuations of the magnetization exceed the difference in the spin states between two energy minima, and the magnetic nanoparticles are in an unblocked state. At *T* > *T*_{C2}, the magnetic nanoparticles pass into the PM state. Additionally, an abnormal decreasing magnetization $M_{ZFC}(T)$ at an AFM transition temperature T_{AFM} is observed [\(Figure](#page-6-0) [5](#page-6-0)b) that is caused by AFM contribution of nanoparticles' "dead" skin.⁵³

As pressure *P* increases, the FM and AFM subsystems become more magnetically homogeneous. The characteristic Curie temperatures T_{C1} and T_{C2} for bigger and T_{C3} and T_{C4} for

Figure 9. Electrochemical testing for the LSCO electrodes in 1 M KOH: (a) Nyquist plots; (b, c) LSV curves and Tafel plots before the long-term stability test for the LSCO catalysts and RuO_2 at a scan rate of 1 mV s⁻¹ (*iR*-compensation 90%); (d) time dependence of potential under a constant current density of 10 mA cm^{−2}; (e, f) LSV curves and Tafel plots after the long-term stability test at a scan rate of 1 mV s^{−1}.

smaller nanoparticles increase with a pressure rise [\(Table](#page-6-0) 3), indicating the strengthening of the FM DE $Co^{3+} - O^{2-} - Co^{4+}$. At the same time, the T_{AFM} also increases, demonstrating the improvement of the AFM superexchanges $Co^{3+}-O^{2-}-Co^{3+}$ and $Co^{4+}-O^{2-}-Co^{4+}$, whereas the blocking temperature T_B increases slightly ([Table](#page-6-0) 3). The characteristic Curie temperatures show linear-like behavior versus pressure [\(Figure](#page-7-0) 6). It allows us to define an average $\langle dT_C/dP \rangle$ ratio for smaller and bigger nanoparticles, 10 and 13 K/GPa, respectively. The obtained values show that larger nanoparticles are more sensitive to the external pressure *P*, and their magnetic parameters can be easily controlled.

The isothermal magnetization curves $M(H)$ around T_C in a wide temperature range from 200 to 300 K with a $\Delta T = 2$ K step for LSCO-900 are shown in [Figure](#page-7-0) 7a. At $T > T_{C}$, the magnetization increases linearly with an applied magnetic field, indicating the PM state of the sample. At $T < T_C$, the linear isothermal magnetization behavior turns to a curve, which indicates passing to the FM state. At the same time, the magnetization cannot be saturated even at lower temperatures and higher magnetic fields, probably because of AFM phase contribution.⁵⁸ According to the Banerjee criterion,^{[59](#page-15-0)} the positive slope of Arrott's curves *M*² (*H*/*M*) (see [Figure](#page-7-0) 7b) near the Curie temperature is observed, which indicates a second-order magnetic phase transition for the LSCO-900.

The temperature dependences of the magnetic entropy change $-\Delta S_M(T, H)$ near phase transition(s) [\(Figure](#page-7-0) 8) were plotted using the isotherms of magnetization *M*(*H*) and the numerical integration method of the Maxwell relation:^{[60](#page-15-0)}

$$
\Delta S_{\rm M}(T, \mu_0 H) = S_{\rm M}(T, \mu_0 H) - S_{\rm M}(T, 0)
$$

=
$$
\int_0^{\mu_0 H} (\partial M / \partial T)_{\mu_0 H} d(\mu_0 H)
$$
 (1)

The magnetic entropy change $-\Delta S_M(T, H)$ shows a vast indistinct peak with relatively low values of MCE. Upon application of the external high hydrostatic pressure $P \approx 0.8$ GPa and with increasing the magnetic field Δ*H* up to 10 kOe, the MCE peak position moves toward higher temperatures from 250 to 256 K, indicating the domination of the FM DE over AFM superexchange interactions, whereas without pressure $P = 0$ we have the opposite situation with temperature reducing from 247 to 233 K ([Figure](#page-7-0) 8). Moreover, the δT_{fwhm} for the $-\Delta S_M^{\text{max}}$ is far greater than 50 K, making the studied LSCO-900 a promising component for creating magnetic refrigerant composites in a wide temperature range.

3.3. Electrocatalytic Properties. As mentioned in the introduction, interest in Ln- and Co-containing perovskites as electrocatalysts for OER in an alkaline environment has increased significantly recently.[19](#page-14-0),[21](#page-14-0)−[27](#page-14-0),[32](#page-14-0)−[35](#page-14-0) In order to find out the OER activity for water splitting of the prepared LSCO samples, electrochemical tests were carried out in 1 M KOH electrolyte (pH 14) with the following sequence: multiple current−voltage (CV) scans (until stable electrode operation is achieved), alternative current (AC) impedance, recording of LSV curves with *iR* correction, long-term chronopotentiometry tests (at a current density of 10 mA·cm[−]²) and reregistration of LSV curves. It should be emphasized that Co- and Nicontaining complex oxides can be entirely or partially *in situ* transformed into oxides/hydroxides due to electrochemical processes in an alkaline medium.^{61,[62](#page-15-0)} Accordingly, after a long electrochemical treatment, we investigated the electrode materials to verify the stability of the LSCO catalysts.

Before the electrochemical test, the manufactured electrodes were placed in an electrolyte environment for 1 h to establish equilibrium at the electrode−electrolyte interface. As shown in [Figure](https://pubs.acs.org/doi/suppl/10.1021/acsami.3c06413/suppl_file/am3c06413_si_001.pdf) S2, for all samples the shape of the CV curves in the range of 0.9−1.8 V is similar at the corresponding scanning

Figure 10. XPS spectra for the LSCO-900 samples before and after electrocatalysis: (a) survey spectra; (b) La 3d spectra with A and B features from charge transfer satellites, plasmon lines, and MNN Auger lines; (c) Sr 3d spectra with spin–orbit splitting Sr $3d_{5/2}$ and Sr $3d_{3/2}$; (d) O 1s spectra with O1 lattice oxygen and O2 adsorbed oxygen on the surface components; and (e) Co 2p spectra (the dashed and dotted lines indicate the energy position of the $Co₃O₄$ and $Co²⁺$ satellite peaks, respectively).

rate (from 1 to 50 mV s^{-1}). Only the peak around 1.3 V (anodic scans) is clearly observed, corresponding to the $Co⁴⁺$ \rightarrow Co³⁺ transformation.^{[39](#page-14-0)} As AC impedance measurements show [\(Figure](#page-8-0) 9(a)), the LSCO-900 sample has the best conductivity.

Nyquist plots correspond to the equivalent circuit and fitting parameters using ZView software, which includes series resistance *R_s*, the double layer capacitance at the active material/electrolyte interface C_{db} electron transfer resistance

 R_{ct} , and the diffusion impedance Warburg impedance [\(Figure](https://pubs.acs.org/doi/suppl/10.1021/acsami.3c06413/suppl_file/am3c06413_si_001.pdf) [S3](https://pubs.acs.org/doi/suppl/10.1021/acsami.3c06413/suppl_file/am3c06413_si_001.pdf)). The results of the fitting are presented in [Table](https://pubs.acs.org/doi/suppl/10.1021/acsami.3c06413/suppl_file/am3c06413_si_001.pdf) S2. The impedance spectra include two parts, a line in a low-frequency region and a semicircle in a high-frequency region. The numerical value of the diameter of the semicircle on the Z_0 axis is approximately equal to the charge transfer resistance $R_{\rm ct}$, and the sloping line corresponds to Warburg impedance, which is associated with electrolyte diffusion in electrode materials. As known, the Helmholtz double layer capacitance $C_{\rm dl}$ appears

Table 4. Energy positions of the La 3d, Sr 3d, and O 1s X-ray photoelectron lines in the LSCO-900 samples before and after electrocatalysis

due to the porosity of the active material and the uneven distribution of the current on the electrode surface.^{[17,18](#page-14-0),[63](#page-15-0)} The double-layer capacitance C_{dl} was estimated by eq [S3](https://pubs.acs.org/doi/suppl/10.1021/acsami.3c06413/suppl_file/am3c06413_si_001.pdf) plotting the Δj versus RHE as a function of the scan rate [\(Figure](https://pubs.acs.org/doi/suppl/10.1021/acsami.3c06413/suppl_file/am3c06413_si_001.pdf) S4e− h). [Figure](https://pubs.acs.org/doi/suppl/10.1021/acsami.3c06413/suppl_file/am3c06413_si_001.pdf) S4 shows CV curves recorded in the potential windows from 0 to 0.1 V vs Hg/HgO at a scan rate of 1−10 mV s⁻¹. The electrochemically active surface area (ECSA) can be calculated from the C_d using eq S4 [\(see](https://pubs.acs.org/doi/suppl/10.1021/acsami.3c06413/suppl_file/am3c06413_si_001.pdf) SI 2). In this case, CC-based with the Super P electrode can be accepted as a standard, showing the measured capacitance of 36.59 mF cm[−]² . The ECSA values of LSCO-850, LSCO-875, and LSCO-900 are $15.97\;\mathrm{cm^2}$, 17.26 $\mathrm{cm^2}$, and 24.53 $\mathrm{cm^2}$, respectively. It is known that the ECSA value corresponds to the number of potential active sites for electrocatalysis and R_{ct} is inversely proportional to electrochemical activity. A low charge transfer resistance and high ECSA indicate good electronic conductivity and a high rate of OH ion transfer across the boundaries between the electrolyte and active electrode materials. The above AC impedance results show that the LSCO-900 has the attributes of an efficient electrocatalyst for OER in an alkaline environment.

According to onset LSV measurements, all LSCO samples exhibited highly efficient electrocatalytic activity for alkaline OER and a slow boost of current density with an increase of applied potential ([Figure](#page-8-0) 9b). The obtained overpotential values' intervals were only 265−285 and 360−380 mV when reaching a current density of 10 and 100 mA cm $^{-2}$, respectively (more details are given in [Table](https://pubs.acs.org/doi/suppl/10.1021/acsami.3c06413/suppl_file/am3c06413_si_001.pdf) S3). In addition, the intrinsic per-site activity of the electrocatalysts is studied by the turnover frequency (TOF), which accurately describes the activity of electrode materials. Therefore, higher values correspond to a large electrolyte-electrode interface and good OH diffusion and indicate excellent electrochemical perform-ance.^{[18](#page-14-0)} In this way, values for the LSCO samples have been calculated (details in the SI and [Figures](https://pubs.acs.org/doi/suppl/10.1021/acsami.3c06413/suppl_file/am3c06413_si_001.pdf) S5 and S6). The LSCO-900 material demonstrates a high intrinsic activity of about 0.11–0.34 s⁻¹ at 300–350 mV (1.53–1.58 V vs RHE), corresponding to the high intrinsic activity compared to other samples.

The LSV curves for the LSCO-900 and LSCO-875 electrodes are almost identical, and the overpotential of the OER for the LSCO-850 is slightly higher. It is essential that the commercial electrocatalyst of noble metal oxide $(RuO₂)$, used for comparison, showed worse activity. These results confirm the calculations of the Tafel slope [\(Figure](#page-8-0) 9c): for the synthesized LSCO samples, the range of values is 74−82 mV dec^{-1} , and for commercial $RuO₂$ it is slightly more than 92 mV dec[−]¹ . The subsequent long-term test of electrodes during continuous electrolysis (current density of 10 mA cm[−]²) is presented in [Figure](#page-8-0) 9d. The continuous formation of oxygen bubbles on the surface of the electrodes was observed when the stability tests were carried out. For all samples, an increase in overpotential OER is observed during the first 4 h of the chronopotentiometry test, and the values of η_{10} reach 300, 320,

and 335 mV for LSCO-900, LSCO-875, and LSCO-850, respectively. However, the activity of the electrocatalysts decays much more slowly during the next 20 h of electrolysis, especially in the case of the LSCO-900 sample (the magnitude of the overpotential increases by only 12 mV); these results for all electrodes are presented in [Table](https://pubs.acs.org/doi/suppl/10.1021/acsami.3c06413/suppl_file/am3c06413_si_001.pdf) S4. The slight increase in voltage after 4 h of the test proves the final stabilization of all electrodes and the high efficiency of the LSCO materials as OER catalysts. It should also be considered that carbon fiber (the base of the electrode used in this study) is gradually destroyed in an alkaline environment under OER conditions.^{[62](#page-15-0)} It leads to breakdowns of the interface between the catalyst and collector electrode and is an additional factor in the growth of overpotential OER during long-term tests. LSV curves measured after long-term stability tests [\(Figure](#page-8-0) 9e) are in good agreement with CP results: the values of overpotential *η*¹⁰ for the corresponding samples are identical [\(Tables](https://pubs.acs.org/doi/suppl/10.1021/acsami.3c06413/suppl_file/am3c06413_si_001.pdf) S3 and [S4](https://pubs.acs.org/doi/suppl/10.1021/acsami.3c06413/suppl_file/am3c06413_si_001.pdf)). Reregistration LSV measurements revealed that in the range of current densities from 10 to 100, the OER overvoltage increased by only 14−17% for the LSCO-900 electrode, while for the LSCO-875 and LSCO-850, the changes were 22−24 and 27−44%, respectively [\(Table](https://pubs.acs.org/doi/suppl/10.1021/acsami.3c06413/suppl_file/am3c06413_si_001.pdf) S3). Calculation of the Tafel slope after long-term electrolysis [\(Figure](#page-8-0) 9f) also indicates the preservation of the high catalytic activity of the LSCO-900 material, in contrast to the LSCO-875 and especially LSCO-850 samples.

Electrode materials were also studied in detail after the CP tests. According to XRD, all LSCO samples retained their crystallinity. As shown in [Figure](https://pubs.acs.org/doi/suppl/10.1021/acsami.3c06413/suppl_file/am3c06413_si_001.pdf) S7, the diffraction patterns show reflections from both crystalline phases $(La_{0.6}Sr_{0.4}CoO_3)$ and $SrCo_{0.78}O_{2.48}$). It proves the high stability of Ln and Cocontaining complex oxides in an alkaline environment during long-term electrolysis, in contrast to similar Mn-containing compounds.^{[11](#page-13-0)}

The survey spectra of the LSCO-900 samples before and after electrocatalysis show the existence of La, Sr, Co, and O peaks [\(Figure](#page-9-0) 10a). Nevertheless, the XPS spectra revealed some differences. All peak positions for the La, Sr, and O ions are listed in Table 4.

In the XPS spectra of the La 3d, two major peaks, corresponding spin doublet La $3d_{5/2}$ and La $3d_{3/2}$, are observed ([Figure](#page-9-0) 10b). However, the spectra have a more complex profile due to charge transfer satellites, plasmon lines, and MNN Auger lines (marked A and B in [Figure](#page-9-0) 10b). Similar results are also observed for the XPS spectrum of La_2O_3 .^{[64](#page-15-0)} The binding energy positions for the LSCO-900 before electrocatalysis are 832.9 eV (La $3d_{5/2}$) and 849.7 eV (La $3d_{3/2}$), with a difference of 16.8 eV, which agrees with the standard deviation.^{[64](#page-15-0)} Noteworthy, the XPS spectrum for LSCO-900 after electrocatalysis is transformed to singlepeaks with binding energies of 833.3 eV (La $3d_{5/2}$) and 850.9 eV (La $3d_{3/2}$). It is associated with the amorphization of the LSCO-900 surface and the appearance of chemical defects due to a positive core-level shift of $0.4-1.2$ eV.^{[35](#page-14-0)}

[Figure](#page-9-0) 10c shows Sr 3d spectra, which exhibit two main components from the perovskite lattice-bound Sr_{lat} at binding energies of 131.6 eV (Sr $3d_{5/2}$) and 133.4 eV (Sr $3d_{3/2}$) for LSCO-900 before electrocatalysis. In addition, a shoulder with higher binding energies of 135.2 eV from surface Sr_{surf} is also presented, which is usually associated with surface strontium-rich phases often observed for LSCO materials.^{[41](#page-14-0)} After electrocatalysis, there is a redistribution of intensity between the main Sr_{lat} spectral lines, small peaks position shifting, and hiding and/or disappearing of Sr_{surf} component ([Table](#page-10-0) 4).

As for O 1s XPS before electrochemical tests [\(Figure](#page-9-0) 10d), two components with energy positions of about 528.4 eV (O1) and 531.2 eV (O2) can be clearly distinguished in the spectrum [\(Table](#page-10-0) 4), which may be attributed to the lattice oxygen and adsorbed oxygen on the surface, respectively.^{[35](#page-14-0)} After electrochemical tests, the O 1s profile is modified significantly to a single broad peak with an energy position of about 530.3 eV. Such a redistribution in the intensities of the peaks, as well as variation in their shape and positions, indicates the transformation of the LSCO-900 surface.

Changes are also observed for the XPS spectra of Co 2p ([Figure](#page-9-0) 10e): exhibit two prominent peaks at around 780.1 and 795.6 eV for the electrode before and at around 778.8 and 794.0 eV after electrocatalysis, corresponding to the Co $2p_{3/2}$ and Co $2p_{1/2}$ states, respectively. It should be noted that for LSCO-type materials, the position of the Co 2p peaks has a metallic character and is practically unchanged for compounds with different strontium content. 65 After electrocatalysis, the surface layer of LSCO is significantly transformed [\(Figure](#page-9-0) [10e](#page-9-0)): (1) the Co 2p peaks are shifted toward lower energy \sim 1.3−1.6 eV; (2) the simultaneous reducing $Co²⁺$ satellite peak around ∼786.0 eV and, as a result, its content and increasing $Co₃O₄$ satellite peak ∼789.0 eV, consisting of the mixed $Co²⁺$ and $Co³⁺$ states, are observed, which gives evidence about the formation of a new surface; 41 (3) the full-width at halfmaximum (fwhm) of the Co 2 $p_{3/2}$ main peak reduces ~2.0 eV, suggesting more $Co³⁺$ and fewer $Co²⁺$ ions on the surface.^{[66](#page-15-0)}

Such significant changes in La 3d, Sr 3d, O 1s, and Co 2p XPS spectra indicate the complex nature of transformations on the surface of LSCO-900 during OER in an alkaline environment and is an indirect confirmation of the amorphization of the catalyst surface as recently shown in work.^{[41](#page-14-0)}

SEM and elemental mapping of the electrode surface after prolonged electrolysis [\(Figures](https://pubs.acs.org/doi/suppl/10.1021/acsami.3c06413/suppl_file/am3c06413_si_001.pdf) S8−10) mostly show nanoparticles as in the initial LSCO material, which agrees with the elemental mapping ([Figure](https://pubs.acs.org/doi/suppl/10.1021/acsami.3c06413/suppl_file/am3c06413_si_001.pdf) S9). This morphology is characteristic of cobalt hydroxide, CoO(OH), which can be produced *in situ* during the electrochemical destruction of complex oxides in an alkaline environment.³⁹ According to the work,⁴¹ the mechanism of destruction of LSCO material involves the formation of mixed surface phases: the formation of the selfassembled active $CoO(OH)$ phase, the relative enrichment of passivating La $(OH)_{3}$, and the leaching of Sr $(OH)_{2}$. In our case, a similar way of transformation is implemented only partially. First, the surface areas of CoO(OH) formation are only a small number of fragments, and second, according to XPS, the peak from Srsurf still remains after electrocatalysis. TEM results for the LSCO-900 nanoparticles before and after electrolysis are shown in [Figures](https://pubs.acs.org/doi/suppl/10.1021/acsami.3c06413/suppl_file/am3c06413_si_001.pdf) S11 and S12. It should be noted that the TEM images demonstrate the formation of an additional thin amorphized layer on the surface of the particles after electrolysis. SAED images confirm the crystalline state of the catalyst [\(Figure](https://pubs.acs.org/doi/suppl/10.1021/acsami.3c06413/suppl_file/am3c06413_si_001.pdf) S11). In addition, the ring diffraction pattern of the LSCO-900 sample after prolonged electrolysis additionally confirms the presence of an amorphized phase (see [SI](https://pubs.acs.org/doi/suppl/10.1021/acsami.3c06413/suppl_file/am3c06413_si_001.pdf) for more details). Accordingly, the consequence of electrochemical transformations is the formation of an amorphized layer on the surface of LSCO, which should be described as forming a stable core−shell system.

[Table](#page-11-0) 5 lists the comparative characteristics of the OER for electrocatalysts based on La, Co-containing perovskites in an alkaline electrolyte. Among them, the LSCO material described here occupies a leading position. It should be noted that very recently, it was reported that the amorphization of crystallized nanometer perovskite $LaCoO₃$ to a highly disordered state leads to improved electrocatalytic properties of OER.^{[35](#page-14-0)} In our case, a highly active amorphized layer is formed on the surface of the LSCO nanoparticles during electrolysis. The catalytic activity increases with increasing proportion of this crystalline phase for a number of LSCO-850, LSCO-875, and LSCO-900 samples. However, the role of the second phase, $SrCo_{0.78}O_{2.48}$, can also be very important. It is known that less active perovskite catalysts often remain crystalline during OER, while more active ones undergo amorphization.[38](#page-14-0) It can be assumed that the destruction of $SrCo_{0.78}O_{2.48}$ is slower than that of LSCO, which leads to the stabilization of the amorphized layer. It can prevent further degradation of the nanoparticles, protect the initial morphology and provide a significant electroactive surface area for a long time. However, this assumption needs further experimental verification, including using similar multiphase samples. Finally, forming *in situ* core−shell nanostructures based on perovskites is a potentially successful strategy for creating highly active and stable electrocatalysts.

Finally, based on the above-mentioned results, the interconnection between magnetic and electrocatalytic properties of the LSCO nanoparticles can be established, including additionally the magnetic properties of the LSCO-900 sample before and after long-term electrolysis (see SI 2 and [Figure](https://pubs.acs.org/doi/suppl/10.1021/acsami.3c06413/suppl_file/am3c06413_si_001.pdf) S13 for more details). As it turned out, the material demonstrates the same behavior of magnetization *M* with almost unchanged phase transition temperatures (Curie and blocking temperatures). Moreover, the change in the sample surface and its modification is confirmed, manifesting in the relative change of the magnetization $\Delta M/M_{\text{max}}$ by 35% under FC and 10% under ZFC regimes at 50 K. Similar behavior was also observed for La_{0.9}Mn_{1.1}O_{3−δ} manganites annealed under different temperatures, but in the K_2SO_4 (0.5 M) media.^{[12](#page-13-0)} Thus, despite testing the LSCO-900 sample in rigid media and further transforming its surface and structure, this material still demonstrates its magnetic properties, indicating its multifunctionality. In other words, the cobaltite nanoparticles retain their practical potential to be used simultaneously for both magnetic and electrocatalytic purposes.

4. CONCLUSIONS

The structure, morphology, particle size distribution, and magnetic phase transitions, as well as magnetic, magnetocaloric, and electrochemical properties of the LSCO nanopowders obtained under different annealing temperatures t_{ann} = 850, 875, and 900 °C, have been studied comprehensively. All samples indicate a rhombohedral *R3c* perovskite structure, improving its single-phase nature with increasing t_{ann} . At the same time, it leads to increasing unit cell volume *V*, an average particle size *D*, and its dispersion *σ*. The LSCO-900 composition mainly consists of the $Co³⁺$ and $Co⁴⁺$ ions, as well as minor Co^{2+} and Co_3O_4 traces that play a crucial role in forming magnetic and electrochemical properties. The main magnetic parameters such as spontaneous magnetization M_S , coercivity H_C , residual magnetization M_r , phase transition temperatures, and magnetic entropy change $-\Delta S_M$ have been defined. The LSCO-900 sample is a ferromagnet with a high

 H_C of 4.3 kOe. It has been found that LSCO-900 exhibits two characteristic Curie temperatures, T_{C1} and T_{C2} , which are associated with the presence of nanoparticles of different sizes. Additionally, FM and AFM subsystems are stabilized and become more magnetically homogeneous because of applied external hydrostatic high-pressure *P*. As pressure increases, the average $\langle T_{C1} \rangle$ and $\langle T_{C2} \rangle$ increase from 253 and 175 K under ambient pressure to 268 and 180 K under $P \approx 0.8$ GPa, respectively. At the same time, the AFM transition temperature T_{AFM} and blocking temperature T_{B} also increase from 145 and 169 K to 158 and 170 K, respectively. The rise of $\langle dT_{C1}/dP \rangle$ for bigger particles and $\langle dT_{C2}/dP \rangle$ for smaller particles differs and is 13 and 10 K/GPa, respectively. The MCE in the LSCO-900 nanopowder is near the Curie temperature and has a second-order phase transition. The values of $-\Delta S_{\rm M}$ are quite low and equals 0.05 J/kg K without pressure and 0.017 J/kg K under $P \approx 0.8$ GPa in the magnetic field 1 T, but with extremely wide peak $\delta T_{\text{fwhm}} > 50$ K.

The obtained series of the LSCO nanomaterials showed excellent characteristics as electrocatalysts for the OER process in a 1 M KOH electrolyte. The values of the initial overpotential of oxygen generation were within 265−285 mV (a current density of 10 mA/cm²), and the sample with the highest content of the $La_{0.6}Sr_{0.4}CoO_3$ phase showed the best performance. The LSCO materials described here retain high catalytic activity during long-term electrolysis, which should be attributed to the formation of a stable amorphization layer on the surface of nanoparticles. Accordingly, a significant area of the catalytically active surface, resulting from the emergence of stable core−shell nanosystems, can ensure a high efficiency of practical use of the material.

The parallel study of nanopowders' physical and electrocatalytic properties in this work made it possible to establish the physical and chemical basis of the various samples obtained at different *t*ann values and find ways to obtain materials for practical application in electrocatalysis. For instance, the research of magnetic and magnetocaloric properties of the samples opens up the possibility of changing the conditions of catalysis by contactless heating or cooling the temperature of electrodes and electrolytes using the MCE. Moreover, using the alternating field can increase the temperature of electrodes.

■ **ASSOCIATED CONTENT**

\bullet Supporting Information

The Supporting Information is available free of charge at [https://pubs.acs.org/doi/10.1021/acsami.3c06413](https://pubs.acs.org/doi/10.1021/acsami.3c06413?goto=supporting-info).

> SI1, determination of the size of the coherent scattering region in the $La_{0.6}Sr_{0.4}CoO_3$ nanopowders; SI2, electrochemical properties of the LSCO nanopowders ([PDF\)](https://pubs.acs.org/doi/suppl/10.1021/acsami.3c06413/suppl_file/am3c06413_si_001.pdf)

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Author Contributions

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Notes

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