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# Evaluating the feasibility of creatine-weighted CEST MRI in human brain at 7 T using a Z-spectral fitting approach

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# Abstract

The current study aims to evaluate the feasibility of creatine (Cr) chemical exchange saturation transfer (CEST)-weighted MRI at 7 T in the human brain by optimizing the saturation pulse parameters and computing contrast using a Z-spectral fitting approach. The Cr-weighted (Cr-w) CEST contrast was computed from phantoms data. Simulations were carried out to obtain the optimum saturation parameters for Cr-w CEST with lower contribution from other brain metabolites. CEST-w images were acquired from the brains of four human subjects at different saturation parameters. The Cr-w CEST contrast was computed using both asymmetry analysis and Z-spectra fitting approaches (models 1 and 2, respectively) based on Lorentzian functions. For broad magnetization transfer (MT) effect, Gaussian and Super-Lorentzian line shapes were also evaluated. In the phantom study, the Cr-w CEST contrast showed a linear dependence on concentration in physiological range and a nonlinear dependence on saturation parameters. The in vivo Cr-w CEST map generated using asymmetry analysis from the brain represents mixed contrast with contribution from other metabolites as well and relayed nuclear Overhauser effect (rNOE). Simulations provided an estimate for the optimum range of saturation parameters to be used for acquiring brain CEST data. The optimum saturation parameters for Cr-w CEST to be used for brain data were around  $B_{1rms} = 1.45 \ \mu T$  and duration = 2 seconds. The Z-spectral fitting approach enabled computation of individual components. This also resulted in mitigating the contribution from MT and rNOE to Cr-w CEST contrast, which is a major source of underestimation in asymmetry analysis. The proposed modified z-spectra fitting approach (model 2) is more stable to noise compared with model 1. Cr-w CEST contrast obtained using fitting was

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CONFLICTS OF INTEREST

The authors declare no financial interests.

SUPPORTING INFORMATION

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 $6.98 \pm 0.31\%$  in gray matter and  $5.45 \pm 0.16\%$  in white matter. Optimal saturation parameters reduced the contribution from other CEST effects to Cr-w CEST contrast, and the proposed Z-spectral fitting approach enabled computation of individual components in Z-spectra of the brain. Therefore, it is feasible to compute Cr-w CEST contrast with a lower contribution from other CEST and rNOE.

#### **Keywords**

chemical exchange saturation transfer; creatine; creatine-weighted chemical exchange saturation transfer; Lorentzian; magnetic resonance imaging; Z-spectra

# 1 | INTRODUCTION

The chemical-exchange-saturation-transfer (CEST) MRI technique is used for the highresolution mapping of metabolites/molecules/macromolecules possessing protons exhibiting exchange phenomena with protons of bulk water.<sup>1,2</sup> In a typical CEST experiment, exchangeable protons of the solute pool are saturated using a long radio-frequency (RF) pulse, which results in a reduction of bulk water proton signal due to the transfer of saturated magnetization from solute to bulk water. CEST-MRI techniques are being developed for in vivo mapping of labile molecules and metabolites.<sup>3–5</sup>

Creatine (Cr) is a significant metabolite in the brain and its alteration has been linked to various diseases.<sup>6–9</sup> It has been shown that Cr concentration ([Cr]) in the brain decreases with aging as well as in diseases including Alzheimer's disease (AD) and Huntington disease, and in brain tumors.<sup>7,8,10</sup> Using both invasive and noninvasive methods, the brain's Cr level is reported to be in the range of 6–14mM.<sup>11–16</sup> Proton magnetic resonance spectroscopy (<sup>1</sup>H MRS) provides a noninvasive approach to detect the total Cr (tCr) (phosphocreatine [PCr] + Cr) level in the brain. In the brain, PCr concentration is smaller compared with Cr (the Cr/PCr ratio is ~ 2:1).<sup>12,17</sup> However, using <sup>1</sup>H MRS, the free Cr level in the brain cannot be quantified and, moreover, it is unable to provide higher resolution mapping of Cr. The CEST-based method has been recently developed to map the Cr at a higher resolution by exploiting the exchangeable amine protons resonating at a frequency offset of ~ 1.8 ppm downfield from water resonance.<sup>18,19</sup> Cr-weighted (Cr-w) CEST imaging in human calf muscle has been used to probe the exercised induced changes in Cr level on 7<sup>20</sup> and 3 T human scanners.<sup>21</sup> Further, Cr-w CEST mapping is used to detect the failing bioenergetics in the infarcted myocardium.<sup>22</sup> Recently, Cr-w CEST imaging mapped the spatial distribution of Cr in rat brain on a 9.4 T animal scanner.<sup>23</sup> Another study reported that the molecular origins of CEST signals at 2 ppm in rat brain is predominantly due to Cr.24

The exchange rate of amine protons of Cr has been reported in the range of 500–1190 Hz.<sup>4,19,25</sup> The exchange rate of amine protons (1.8 ppm downfield from water resonance) and amide protons (2.5 ppm downfield from water resonance) of PCr were reported as ~ 120 and ~ 140 Hz, respectively, in a phantom study.<sup>19</sup> Due to the intermediate exchange rate, the CEST effect from Cr is dominant over PCr.<sup>19</sup> The CEST effect from PCr (for physiological range) for a particular saturation parameter was reported to be negligible in the phantom

study at 3 T.<sup>19</sup> However, by using low saturation power with a very long saturation duration, the CEST effect from PCr can be enhanced.

In vivo mapping of the CEST effect from amide proton transfer (APT-w CEST),<sup>3,26–29</sup> myoinositol (MI-w CEST),<sup>30</sup> glutamate (Glu-w CEST)<sup>4,31,32</sup> and glucose (Gluco-w CEST)<sup>33</sup> has been performed in the human brain. Various clinical applications of these CEST techniques are emerging for investigating human brain diseases. As such, amide protons of APT resonate at 3.5 ppm, amine protons of glutamate at 3.0 ppm, and hydroxyl protons of MI at ~ 0.6 ppm downfield of the bulk water resonance; however, the line shapes of their CEST effect are usually broad, and therefore they may contribute to the CEST effect from Cr at 1.8 ppm. Therefore, it is imperative to evaluate the contribution from various metabolites/ molecules to Cr-w CEST.

In general, computation of CEST contrast is challenging due to contribution from direct saturation (DS), magnetization transfer (MT) and other metabolites showing a broad CEST effect. Although the conventional CEST asymmetry (CEST<sub>asy</sub>) approach enables a substantial reduction of DS and MT effects, this approach is highly sensitive to static field (B<sub>0</sub>) inhomogeneity and requires accurate correction before interpretation of CEST contrast. CEST<sub>asy</sub> is also known as magnetization transfer ratio (MTR<sub>asy</sub>). Another issue with CEST<sub>asy</sub> analysis is the negative contribution to CEST contrast from relayed nuclear Overhauser effect (rNOE) (centered at 3.5 ppm upfield to the water resonance), which reduces the sensitivity of measurement of CEST<sub>asy</sub> contrast. Alternative approaches like Z-spectral fitting are used to mitigate these effects.<sup>23,24</sup> Several studies have used a combination of Lorentzian functions<sup>22–24,34,35</sup> for computing CEST contrast. Recently, superposition of Lorentzian function was used to calculate CEST contrast at 2 ppm in rat brain at 9.4 T<sup>23</sup>. Two studies<sup>36,37</sup> have also used partial fitting of Z-spectra corresponding to wide-offset frequencies using a two-pool Henkelman's model for the computation of APT and rNOE at 3 and 4.7 T. It was reported that with steady state condition, super-Lorentzian line shapes performed better than Lorentzian and Gaussian in computing the broad MT effect from the Z-spectra of tissue.<sup>38</sup> However, in another study, for nonsteady state saturation, Gaussian line shape was reported as better than super-Lorentzian.<sup>39,40</sup> The choice of line shape for fitting MT component depends on several factors like shape of the saturation pulse, saturation parameters (amplitude and duration), and steady state versus nonsteady state conditions. Therefore, depending upon the type of Z-spectra data, an appropriate line shape should be used. Another study used an apparent exchange-dependent relaxation method (AREXfit)<sup>41</sup> for computing the CEST effect at 2 ppm.<sup>24</sup>

The CEST effect is highly dependent on saturation parameters, field strength and tissue type. Therefore, a systematic study to investigate the feasibility of Cr-w CEST mapping of the human brain with optimal saturation pulse parameters and evaluation of possible contributions from different metabolites is required.

The objectives of the current study were to carry out optimization of saturation parameters using phantom and simulations; to evaluate contributions from other sources to Cr-w CEST contrast using simulations; and to propose an improved Z-spectral fitting approach for better computation of individual Z-spectral components for in vivo human brain data from 7

T. Reproducibility of Cr-w CEST contrast was also evaluated. The effect of Lorentzian, Gaussian and super-Lorentzian line shapes for a broad MT pool on Z-spectra fitting was also investigated.

# 2 | MATERIALS AND METHODS

MR imaging of phantoms and healthy human brain was performed using a vendor-supplied 32-channel receive and single-channel transmit RF coil on a Siemens 7 T human research scanner (Siemens Medical Solutions, Malvern, PA, USA).

# 2.1 | CEST MRI pulse sequence

The pulse sequence used in the current study consisted of a frequency selective saturation pulse train (user-selected saturation offset  $[\Delta\omega]$ , saturation duration and saturation power) followed by a segmented RF-spoiled gradient echo (GRE) readout acquisition with centric phase encoding order.<sup>42</sup> The saturation pulse train consisted of a series of Hanning-windowed rectangular pulses each of 100 ms duration separated by a 200 µs delay (duty cycle = 99.80%). The excitation bandwidth of this saturation pulse train was 5 Hz for a 1 second saturation duration with a 1% bandwidth of 20 Hz. At the end of the GRE acquisition segments, a variable delay was added to provide T<sub>1</sub> recovery and to keep the sequence under the permitted RF safety limits on the scanner. This sequence was flexible enough to be used for water saturation shift referencing (WASSR)<sup>43</sup> data acquisition and CEST imaging.

### 2.2 | Phantom preparation

To study the dependence of CrCEST contrast on concentration, Cr solution phantoms with varying concentrations (0, 2, 4, 6, 8 and 10mM) (Sigma Aldrich, USA) at pH 7.0 were prepared in phosphate buffered saline (PBS). These samples were added to small test tubes (of 10 mm diameter) and immersed inside a large cylindrical container with PBS. This cylindrical container was placed inside a thermally isolated rectangular box for temperature maintenance. The temperature of PBS inside the container was maintained at  $37 \pm 1^{\circ}$ C during the MRI experiment.

#### 2.3 | Acquisition protocol for phantom MR imaging

The MRI study protocol consisted of the following steps: a three-plane localizer,  $B_0$  map,  $B_1$  map, Z-spectra or CEST data, and a control image without saturation. The CEST MRI imaging parameters were: slice thickness = 10 mm, GRE flip angle (FA) = 10°, GRE readout repetition time (TR)/echo time (TE) = 5.6/2.7 ms, field of view (FOV) =  $120 \times 120$  mm<sup>2</sup>, matrix size =  $192 \times 192$ , and a single shot acquired every 15 seconds. Z-spectrum data were acquired using a saturation pulse of a root mean square  $B_1$  ( $B_{1rms}$ ) of 2.2 µT and 1 second duration at frequencies offset ranging from -4 to +4 ppm in steps of 0.2 ppm. The total time for the Z-spectra data acquisition was 10.5 minutes. To evaluate the  $B_1$  dependence of Cr CEST, CEST-weighted images were acquired at varying  $B_{1rms}$  (1.4, 2.1, 2.9 and 3.6 µT) with 1 second saturation duration over the limited offset frequency range ( $\pm 1.4$  to  $\pm 2.2$  ppm with a step-size of 0.2 ppm). For studying the saturation duration dependence of Cr-w CEST, CEST-weighted images at  $B_{1rms}$  of 2.2 µT and durations of 1, 2 and 3 seconds were acquired

over the limited offset frequency range ( $\pm 1.4$  to  $\pm 2.2$  ppm with a step-size of 0.2 ppm). Both saturation power and duration dependence data were acquired in 17.5 minutes.

Data for generating  $B_0$  and  $B_1$  maps from the same imaging slice were obtained for removing the  $B_0$  and  $B_1$  field inhomogeneities effect on CEST data. For the  $B_0$  field map, WASSR data were acquired using a  $B_{1rms}$  of 0.4 µT with a duration of 200 ms and offset frequency range of (±0 to ±1) ppm with a step-size of 0.1 ppm. The total time for WASSR data acquisition was 2.2 minutes. Data for the  $B_1$  field map were obtained using a single slice fast spin echo readout sequence with TR/TE = 6 s/12 ms. Two images were obtained using preparation square pulses with FAs of 30° and 60° (pulse duration = 0.3 ms) followed by a spoiler gradient. It took 12 seconds to acquire  $B_1$  map data.

#### 2.4 | Bloch-McConnell numerical simulations

Bloch-McConnell equation solvers were written in in-house developed code using MATLAB R2014b-R2017a incorporating relaxation and chemical exchange for the experimental conditions used in the current study.<sup>4,44,45</sup> We simulated the magnetization available at the end of the saturation pulse train and calculated Z-spectra, CEST<sub>asy</sub> plots and CEST<sub>asy</sub> values. Table 1 shows the details of parameters used for simulations, mimicking in vivo situations. Specifically, we used three sites (free and bound pools of water and CEST pool at 1.8 ppm) for optimizing B<sub>1rms</sub> (0.37 to 7.4  $\mu$ T with a step-size of 0.37  $\mu$ T) and a duration of 0.2 to 4 seconds (with a step-size of 0.2 seconds). Surface plots of CEST<sub>asy</sub> contrast with respect to B<sub>1rms</sub> and duration were generated.

A multi-pool exchange model mimicking in vivo brain simulations<sup>4</sup> was considered to estimate the amount of overlap to Cr-w CEST contrast from other brain metabolites. The multi-pool model consists of DS pool (free pool of water) centered at 0 ppm, MT (bound pools of water) centered at -2.4 ppm, MI-w pool centered at 0.6 ppm, Glu-weighted pool centered at 3 ppm and APT-w pool centered at 3.5 ppm (Table 1). Simulations were carried out for different B<sub>1rms</sub> (0.37 to 7.4  $\mu$ T with a step-size of 0.37  $\mu$ T) and duration (0.2 to 4 seconds with a step-size of 0.2 seconds) values. The surface plot of CEST<sub>asy</sub> contrast from various metabolites (except Cr at 1.8 ppm) was generated for different B<sub>1rms</sub> and duration values.

Contribution from individual components to Cr-w CEST contrast was also estimated by simulating the CEST effect from the corresponding component. For obtaining estimates of bound water, which provide the MT effect, to Cr-w CEST contrast, the concentrations of amide protons, Cr, myo-inositol and glutamate metabolites were set to zero in simulations. For estimating the contribution of a specific metabolite to Cr-w CEST contrast, the concentrations of all other metabolites were set to zero keeping the bound water concentration fixed. The dependence of Cr-w CEST on exchange rate was also evaluated.

#### 2.5 | In vivo human brain MR imaging

All the human scans were performed according to the approved protocol of the Institutional Review Board committee of the University of Pennsylvania, USA. Four healthy volunteers aged 28–40 years were recruited, and informed consent was obtained from each volunteer after the study protocol had been explained to them.

The MRI study protocol consisted of the following steps: a localizer, T<sub>1</sub>-weighted anatomical three-dimensional (3D) magnetization-prepared rapid gradient echo (MPRAGE) data, B<sub>0</sub>/WASSR map data, B<sub>1</sub> map data, CEST or Z-spectra data, and a control image without saturation. Whole-brain MPRAGE protocol consisted of 176 axial slices, TR/TE/ inversion time (TI) = 2800/4.4/1500 ms, FA = 7°, resolution =  $0.8 \times 0.8 \times 0.8 \text{ mm}^3$ , iPAT = 2. The scan time for MPRAGE was 4.4 minutes.

Image without saturation pulse ( $M_0$ ) was acquired using imaging readout parameters: FOV =  $240 \times 240 \text{ mm}^2$ , matrix =  $192 \times 192$ , slice thickness = 8 mm, FA =  $10^\circ$ , TR/TE = 5.4/2.6 ms and a single shot. For the same slice, CEST data were acquired at  $B_{1rms}$  of 0.75, 1.1, 1.45, 2.2 and 2.9 µT and a duration of 2 seconds. Offset frequency range used in this protocol were  $\pm 0$  to  $\pm 3$  ppm in steps of 0.3 ppm,  $\pm 3.5$  to  $\pm 6$  ppm in steps of 0.5 ppm, and  $\pm 8$  to  $\pm 20$  ppm in steps of 4 ppm,  $\pm 40$  and  $\pm 100$  ppm. Time per shot including delay and duration of CEST pulse was 8 seconds. The number of averages used for CEST data was one. The acquisition time for the entire CEST data (corresponding to five  $B_{1rms}$  and 46 offset frequencies) was 31 minutes. For the same slice, WASSR data and  $B_1$  map data were also acquired using the parameters described in section 2.3. In the current study, a total of four subjects were scanned for CEST data. Two of these subjects were scanned twice to obtain data for reproducibility. Proper cushions and pads were used for minimizing head movement during MRI.

#### 2.6 | CEST contrast computation

CEST contrast computed using asymmetry analysis is given by:

$$CEST_{asy}(\Delta w) = 100 \times \frac{(M_{sat}(-\Delta \omega) - M_{sat}(+\Delta \omega))}{M_0}$$
(1)

where  $M_{sat}$  ( $\pm \Delta \omega$ ) are the magnetizations obtained with saturation at '+' and '-'  $\Delta \omega$  offset (= 1.8 ppm for Cr) of the water resonance and M<sub>0</sub> is the magnetization without saturation. Pixelwise Z-spectra were obtained from CEST-weighted images by plotting the normalized image intensity as a function of the frequency offset of the saturation pulse for each sample.

$$Z(\bigtriangleup \omega) = 100 \times \frac{M_{sat}(\bigtriangleup \omega)}{M_0}$$
(2)

In this study, for fitting we have used scaled Z-spectra, which is also known as the magnetization transfer ratio (MTR), given by the following equation:

$$\operatorname{Zsc}(\bigtriangleup \omega) = 100 - \operatorname{Z}(\bigtriangleup \omega)$$

(3)

### 2.7 | Z-spectra fitting methods

For estimating CEST effects from individual components, scaled Z-spectrum (Zsc) data were fitted using superposition of Lorentzian functions corresponding to multiple pools.<sup>23</sup>

$$F(\Delta\omega) = 100 \times \sum_{n=1}^{N} \frac{A_n}{1 + 4\left(\frac{\Delta\omega - \Delta\omega_n}{\sigma_n}\right)^2}$$
(4)

where  $\Delta \omega$  is the frequency offset from the water resonance,  $A_n$ ,  $\Delta \omega_n$  and  $\sigma_n$  are the amplitude, frequency offset and line width of the nth proton pool, respectively. In this study, two Z-spectra fitting approaches (models 1 and 2) were used. In model 1, the scaled Z-spectra (Zsc) were fitted with the sum of five Lorentzian functions. For the in vivo brain data, the following five pools were considered: DS (at 0 ppm), MT (at -2.4 ppm), rNOE (at -3.5 ppm), Cr-w CEST (at 1.8 ppm) and CEST@3.5 ppm (at 3.5 ppm). In the current study, CEST@3.5 ppm corresponds to both APT-w and Glu-w contrasts. Resonances of these exchanging pools are well reported; therefore, in this study the position of exchanging pools was fixed with respect to the center of the DS component. This approach enabled reduction of fitting parameters. For MT effect, super-Lorentzian and Gaussian line shapes were also evaluated. Lorentzian function corresponding to MT line shape in Equation 4 was replaced by super-Lorentzian and Gaussian line shapes. Equations for these line shapes have been described previously.<sup>38</sup>

Model 2 is a modified version of model 1. In model 2, instead of fitting Zsc directly, we initially used limited data from  $\pm 12$  to  $\pm 40$  ppm to fit and compute the MT component. It was assumed that, in this range, there is only MT contribution and negligible contribution from DS, rNOE or any other CEST effect. In the current study, different line shapes, namely, Lorentzian, Gaussian and super-Lorentzian, were used to fit broad MT component. After fitting broad MT component, the MT spectrum was interpolated over the whole offset frequency range. This fitted MT component was removed from the Zsc and the remaining part was fitted ( $\pm 0$  to  $\pm 12$  ppm) using superposition of 4-pool Lorentizian function corresponding to DS (at 0 ppm), rNOE (at -3.5 ppm), Cr-w CEST (at 1.8 ppm) and CEST@3.5 ppm (at 3.5 ppm).

#### 2.8 | Data analysis/postprocessing

All image processing and data analysis were performed using in-house developed programs written in MATLAB (R2009b to R2017a) (Natick, MA, USA). For motion correction, all images (WASSR data, B<sub>0</sub> map, B<sub>1</sub> map, Z-spectra or CEST) were registered to the control image without saturation using affine transformation followed by manual desculping. Pixelwise Z-spectra data were interpolated using inbuilt cubic spline function in MATLAB for generating an offset frequency range of  $\pm 0$  to  $\pm 3$  ppm in steps of 0.2 ppm,  $\pm 3.5$  to  $\pm 6$  ppm in steps of 0.5 ppm,  $\pm 8$  to  $\pm 20$  ppm in steps of 1 ppm, and  $\pm 25$  to  $\pm 40$  ppm in steps of 5 ppm.

Z-spectra fitting was performed using an in-house written program which used the inbuilt nonlinear fitting "lsqcurvefit" routine of MATLAB. Parameters used as initial guess and bound constraints are listed in the next paragraph. Before fitting Z-spectra, frequency offset corresponding to minimum intensity value was estimated and Z-spectra were automatically shifted accordingly for mitigating  $B_0$  field inhomogeneity. This shifting takes care of large drift in  $B_0$ . Moreover, the resonance frequency of DS is one of the free parameters of the fitting model, which enable further improvement in the accuracy of automatic  $B_0$  correction. Therefore,  $B_0$  or WASSR map data are not required for computing the CEST contrast based on Z-spectra fitting approaches. This hypothesis was tested by simulating the effect of offset frequency shift on z-spectra fitting as well as by asymmetry analysis. In the current study, we only used the WASSR data-based  $B_0$  map for computation of conventional CEST<sub>asy</sub>.

Four circular regions of interest (ROIs) (~ 15 pixels) were drawn in gray matter (GM) and four in white matter (WM) tissues on the MTR map. The MTR map was computed as  $100*(1-M_{sat} [20 \text{ ppm}]/M_0)$ . First of all, fitting was carried out for representative Z-spectra from ROIs in GM and WM tissues. Lower and upper bounds for the amplitude and width of each component were set based on physiological constraints and prior knowledge from simulation and reported literature. Lower bounds of 20, 0, 0, 0 and 0% and upper bounds of 100, 70, 20, 10 and 10% were set for amplitudes of DS, MT, rNOE, Cr-w CEST and CEST@3.5 ppm, respectively. Lower bounds of 0.1, 10, 0.5, 0.5 and 0.5 ppm and upper bounds of 5, 100, 10, 4 and 8 ppm were set for the width of DS, MT, rNOE, Cr-w CEST and CEST@3.5 ppm, respectively. The initial values for the amplitude and width of each component were taken as averages of the lower and upper bounds. The initial value for the center of DS was set as 0 ppm. Lower and upper bounds for the DS center were – 0.2 and + 0.2 ppm, respectively.

The average value of the fitted width parameter in representative GM and WM ROIs was taken as an initial guess for the pixelwise Z-spectra fitting. Lower and upper bounds for width were set as 50 and 150% of the average initial values, respectively. For the amplitude, the initial value and constraints were the same as those described in the previous section. For the center of DS, the initial guess and constraints were the same as those described in the previous section.

Residual errors and the difference between Z-spectra and fitted Z-spectra were calculated for evaluating the accuracy of fit. Average residual error ( $\sigma$ ) was calculated using the following equation:

$$\sigma = \frac{\sum_{i=1}^{N} \sqrt{\left(Y(i) - Y_{fit}(i)\right)^2}}{N}$$

(5)

where Y represents Z-spectra data,  $Y_{\rm fit}$  represents fitted Z-spectra data and N represents the number of data points in Z-spectra. The goodness-of-fit parameter ( $R^2$ ) was also calculated for each fitting.

Moreover, WM and GM masks were segmented using a MTR map followed by morphological operation erosion (kernel size:  $3 \times 3$ ) to obtain pure normal appearing WM (NAWM) and normal appearing GM (NAGM) masks and to remove contribution from partial volume effects.

#### 2.9 | Monte-Carlo simulations

Z-spectra from the GM data obtained at a  $B_{1rms}$  of 1.45  $\mu$ T and duration of 2 seconds was subjected to Gaussian noise for testing the stability of models 1 and 2 for five pools. Different levels of noise, obtained by changing standard deviation (SD), were used for simulations. For each noise level, 500 noisy Z-spectra data were generated and were fitted using models 1 and 2. Original Z-spectra without additional added noise were also fitted using both models. Absolute relative percentage error (ARPE) between fitted parameters corresponding to noisy Z-spectra and fitted parameters corresponding to Z-spectra without additional noise were computed. Average and SD for each of the fitted parameters over 500 iterations were computed for both models.

Similarly, the noise stability of model 2 corresponding to six pools was evaluated. For the sixth pool, the following pools were considered, one by one: rNOE at -1.6 ppm (rNOE@2), CEST pool at 1 ppm (CEST@1 ppm) and CEST pool at 3 ppm (CEST@3 ppm).

#### 2.10 | B<sub>0</sub> inhomogeneity correction for CEST<sub>asy</sub> contrast

Because calculation of  $\text{CEST}_{asy}$  is highly sensitive to  $B_0$  inhomogeneity, it requires an accurate  $B_0$  map. An offset frequency map generated by finding a minimum of Z-spectra is less accurate and is not suitable for the  $\text{CEST}_{asy}$  map. Therefore, a  $B_0$ map  $\delta\omega_0$  was generated from WASSR data followed by the correction of CEST data for  $B_0$  inhomogeneity using the previously described method.<sup>4,45</sup> Briefly, CEST data were smoothed and interpolated using a cubic spline to generate data with a step-size of 0.01 ppm. For  $B_0$  inhomogeneity correction, each voxel data value at offset frequency  $\Delta\omega$  ppm was replaced by the interpolated data value from ( $\Delta\omega - \delta\omega 0$ ) ppm.  $B_0$  corrected CEST<sub>asy</sub> contrast maps were generated using Equation 1.

# 2.11 | B<sub>1</sub> inhomogeneity corrections

The transmit  $B_1$  field map ( $B_{1rel}$ ) was generated using the double angle method described previously.<sup>4</sup> Amplitude and width of all the maps as well as the CEST<sub>asy</sub> map were corrected using the  $B_1$  calibration-based method described previously<sup>49</sup> considering multiple ROIs (section 2.8).

#### 2.12 | Statistical analysis

Studen s t-test with two tails was used to compare the contrast between GM and WM. The difference was significant if P < 0.05. For reproducibility, CEST maps of the same volunteers at two time points were used to measure the coefficient of variation (COV) in %. The COVs were computed for multiple ROIs (section 2.8) in GM and WM tissues as well as in NAWM and NAGM masks.

# 3 | RESULTS

Z-spectra from Cr solution phantom (10mM) showed a clear dip at ~ 1.8 ppm downfield to the water resonance while no CEST effect was observed from PBS (Figure 1A). Figure 1B shows the plot of Cr-w CEST contrast as a function of Cr concentration. Cr-w CEST contrast linearly correlated with Cr concentration with a detection sensitivity of ~ 0.65% per mM Cr concentration for a  $B_{1rms}$  of 2.2  $\mu$ T and duration of 1 second. Cr-w CEST contrast showed a nonlinear dependence on  $B_{1rms}$  (Figure 1C). For a fixed  $B_{1rms}$ , Cr-w CEST contrast increased with increase in saturation duration (Figure 1D).

Figure 2 shows surface plots demonstrating the dependence of simulated Cr-w CEST contrast on B<sub>1rms</sub> and saturation duration corresponding to an exchange rate of 500 Hz. CEST contrast showed a nonlinear dependence on saturation parameters. The surface plot in Figure 2A shows that Cr-w CEST contrast peaks at long durations for low saturation B<sub>1rms</sub>. The peak value of Cr-w CEST contrast shifts towards short duration with further increase in B1rms. The optimum range of B1rms and saturation duration provide CEST contrast in the region highlighted in yellow-red color. For in vivo experiments, this optimum range is further limited by the hardware set limits and restrictions on scan time. The MR scanner used in the current study allowed a maximum  $B_{1rms}$  of 2.9  $\mu$ T at a duration of 2 seconds. For a 1 second saturation duration, the scanner was allowed to acquire data up to  $B_{1rms}$ of 3.6 µT. Figure 2B shows the contribution from other brain metabolites to Cr-w CEST, illustrating the CEST effect from brain metabolites (myo-inositol, glutamate and amide protons) at 1.8 ppm as a function of the saturation parameters. The dotted ellipse on the surface plot (Figure 2A) highlights an optimum region for selection of saturation parameters for Cr-w CEST in terms of maximum Cr-w CEST contrast, less overlap from other brain metabolites, and hardware constraints. Simulations showed that an optimum  $B_{1rms}$  of 1.45 µT and duration of 2 seconds resulted in small contribution (~ 30%) from other known metabolites to Cr-w CEST. Similar surface plots were obtained for Cr exchange rates of 900 and 1190 Hz (Figure S1). There was a small shift in the region of maximum CEST value compared with an exchange rate of 500 Hz.

Figure 3 shows Z-spectra (Figure 3A) and asymmetry plots (Figure 3B) from an ROI in GM tissue of human brain at 7 T with  $B_{1rms}$  of 0.7, 1.1, 1.45, 2.2 and 2.9 µT and duration of 2 seconds. At low  $B_{1rms}$ , the dip around 3.5 ppm upfield to the water resonance corresponds mainly to rNOE, which overlaps with DS and MT effects with increase in  $B_{1rms}$ . Similarly, at low  $B_{1rms}$ , the dip at 3.5 ppm downfield to the water resonance corresponds to APT-w contrast. For other resonances, the CEST effect seems to be broad, overlapped, and is not clearly visible on the Z-spectra or MTR<sub>asy</sub> analysis for the saturation parameters and sampling used in the current study. Asymmetry plots shown in Figure 3b were normalized with  $M_0$  signal. Asymmetry plots show mixed contrast from CEST and rNOE. At low  $B_{1rms}$ , contrast is negative due to high rNOE and MT effect compared with the CEST effect. Higher CEST contrast and lower rNOE were observed with increased  $B_{1rms}$  (Figures 3).

Figure S2 shows unsaturated image (a), CEST-weighted image at 1.8 ppm (b),  $B_0$  field map (c), relative  $B_1$  field map (d) and Cr-w CEST<sub>asy</sub> map (e). A substantial amount of inhomogeneity was present in the  $B_0$  and  $B_1$  field maps obtained from human brain

at 7 T. Cr-w CEST<sub>asy</sub> maps obtained without correction of these field inhomogeneities were heterogeneous even in the same type of tissue. After correction for the field inhomogeneities, maps become uniform, particularly in the same type of tissues (e). For a B<sub>1rms</sub> of 1.45  $\mu$ T and duration of 2 seconds, Cr-w CEST<sub>asy</sub> contrast in GM was 4.1 ± 0.91%, and in WM it was -0.31 ± 29%.

Plots in Figure 4 show an example of fitting in vivo Z-spectra from an ROI in GM tissue  $(B_{1rms} = 1.45 \,\mu\text{T}, \text{duration} = 2 \text{ seconds})$  using model 1. The Z-spectra fitted well with  $R^2 = 0.99$  and  $\sigma = 0.055$  (Figure 4C). The pointwise residual plot (Figure 4C) also shows that error was within  $\pm 2\%$ , except at water resonance where it was ~ 3%. Individual fitted components are also shown (Figure 4D).

Plots in Figure 5 show an example of fitting in vivo Z-spectra from ROIs in GM tissue  $(B_{1rms} = 1.45 \,\mu\text{T}, \text{duration} = 2 \text{ seconds})$  using model 2. The Z-spectra fitted well with R<sup>2</sup> = 0.99 and  $\sigma$  = 0.059. Figure 5A,B shows original Z-spectra (Z1) and scaled or inverted Z-spectra (Z2) along with the fitted MT component and Z-spectra after removal of MT (Z3). Figure 5C shows four-pool model fitting of Z3 along with various components like DS, rNOE, Cr-w CEST and CEST@3.5 ppm. This is due to the broad spectrum of Glu-w CEST, which overlaps with the APT-w effect and the use of suboptimial parameters for both of these effects. Individual CEST and rNOE components are also shown in the zoomed pan (Figure 5D). Widths of Cr-w, CEST@3.5 ppm and rNOE pools were small for model 2 compared with those obtained using model 1. The amplitude of the rNOE pool decreased in model 2. The amplitude of Cr-w CEST and CEST@3.5 ppm was similar in both models. Width and amplitude of DS and MT pools remained similar for both models.

Figure 6A–C shows fitting of Z-spectra using model 2 at  $B_{1rms}$  of 0.7, 1.45 and 2.1  $\mu$ T and duration of 2 seconds. For all of the cases, fitting was reasonable ( $\sigma \sim 0.06-0.07$ ) and residual errors of fitting at all the offset frequencies remained within ±2%. Figure 6D–F shows the variation in amplitude of rNOE, Cr-w CEST and CEST@3.5 ppm pool with different  $B_{1rms}$ . rNOE decreases with increase in  $B_{1rms}$ . Cr-w CEST and CEST@3.5 ppm increase then decrease with increase in  $B_{1rms}$ .

Figure 7 shows amplitude and width maps of DS, MT, rNOE, Cr-w CEST and CEST@3.5 ppm computed by pixelwise fitting of Z-spectra data from a representative healthy human brain using model 1 (rows 1 and 3) and model 2 (rows 2 and 4). Model 2 maps show a better distribution of rNOE, Cr-w CEST contrast and CEST@3.5 ppm contrast between GM and WM than model 1. DS and MT distributions are similar for both models in GM and WM. The Cr-w CEST contrast (Figure 7D,I) was higher in GM compared with WM and was significantly different (P < 0.01) for both models. The Cr-w CEST contrast ratio in GM/WM was 1.28 in model 2 and 1.43 in model 1. For width maps, the distribution of contrasts was similar in GM and WM. The white arrows in Figure 7S point to representative subarachnoid space of brain, which show high Cr width and low MT amplitude (Figure 7B,G) compared with other regions. Z-spectra of these regions showed some artifacts. Thus, Z-spectra fitting and fitted values of various parameters such as rNOE, Cr-w CEST and CEST@3.5 ppm are not reliable in these regions. Therefore, these voxels should be removed from final maps, otherwise this can lead to misinterpretation. Further studies should be carried out to obtain

good quality Z-spectra in these regions for reliable parameter computation. Table 2 shows mean  $\pm$  SD of individual components of Z-spectra using model 2 averaged over all healthy volunteers taken from multiple ROIs on GM and WM. Preliminary results on two subjects show that the maps obtained using the fitting approach are reproducible with COV < 5%.

The shift in resonance frequency based upon minimum of Z-spectra resulted in first order correction of Z-spectra data for  $B_0$  inhomogeneity. Automatic estimation of the position of DS using the Z-spectra fitting approach further improves its accuracy. Figure 8 shows the sensitivity of Cr-w CEST contrast computed using the fitting approach to shift in the resonance frequency. It was observed that the Cr-w CEST contrast map computed using the fitting approach is less sensitive to  $B_0$  inhomogeneity and does not require a separate  $B_0$  map. On the other hand, the Cr-w CEST<sub>asy</sub> contrast is highly sensitive to  $B_0$  inhomogeneity and even a small error in  $B_0$  estimation can affect the contrast value.

Table 3 shows the results of noise stability of model 2 versus model 1 for Z-spectral fitting. ARPE for all the fitted parameters (except amplitude of MT) were lower in model 2 compared with model 1. Simulation results showed that ARPE for amplitudes of DS, rNOE, Cr-w CEST and CEST@3.5 ppm were 26, 14, 12 and 4%, respectively, lower in model 2 compared with model 1. Similarly, ARPE for widths of DS, MT, rNOE, Cr-w CESTand CEST@3.5 ppm were 6, 25, 22, 7 and 17%, respectively, and found to be lower in model 2 compared with model 1. For an amplitude of MT both models showed a similar stability; however, for width, model 2 was more stable to noise compared with model 1.

Figure 9 shows fitting of Z-spectra using model 2 ( $B_{1rms} = 1.45 \mu$ T, duration = 2 seconds) considering multi-pool models. Figure 9A shows fitting using a five-pool model (DS pool at 0 ppm, MT pool at -2.4 ppm, rNOE pool at -3.5 ppm, Cr-w CEST pool at 1.8 ppm and CEST@3.5 ppm). The reliability of fitting is tested by increasing the number of pools. Figure 9B shows fitting considering six pools, where the first five pools are similar to those of Figure 9A and the sixth pool (CEST@3 ppm) is considered at 3 ppm. In Figure 9C, the sixth pool (CEST@1 ppm) is considered at 1 ppm. In Figure 9D, the sixth pool (rNOE-2) is considered at -1.6 ppm. Residual errors of fittings were more than 2% and  $\sigma = \sim 0.05$  for all cases. However, Monte-Carlo simulations showed that the stability and reliability of the fit decreases with increase in the number of pools. Figures S4 and S5 show the fitting of Z-spectra using model 2 for  $B_{1rms} = 0.7$  and 2.1  $\mu$ T, respectively, at a duration of 2 seconds considering multipool models.

Figure S3 shows a comparison of partial fitting of Z-spectra, at far-off resonance where MT effects dominate and other CEST and rNOE pools have a negligible effect, using Lorentzian, Gaussian and super-Lorentzian line shapes for  $B_{1rms}$  of 0.7, 1.45 and 2.1 µT at a saturation duration of 2 seconds. For the current set of saturation parameters, the Lorentzian line shape fitted better to broad MT spectra with the lowest average residual error ( $\sigma = 0.0003$  to 0.0012) compared with other line shapes for all  $B_{1rms}$ . For the super-Lorentzian line shape, interpolated Z-spectra resulted in a pole at –2.4 ppm, which affected the accuracy of the entire Z-spectra fitting.

# 4 | DISCUSSION

In the current study, we have used numerical simulations to obtain optimized saturation parameters for Cr-w CEST contrast with minimal overlap from known metabolites exhibiting CEST contrast. Optimized parameters were used to obtain in vivo Z-spectral data. The improved Z-spectral fitting approach resulted in an estimate of CEST contrast at 1.8 ppm with reduced overlap from other metabolites.

In phantom studies, a linear dependence of Cr-w CEST contrast on a physiological concentration range of Cr was observed. The Cr-w CEST contrast showed a nonlinear dependence on saturation duration and saturation B<sub>1rms</sub>. For in vivo data, the dependence of Cr-w CEST contrast on saturation parameters might be quite different compared with Cr solution phantom data due to substantial direct saturation of bulk water, MT effects from macromolecules, and overlap from other CEST and rNOE. Therefore, in this study, we have performed numerical simulations, mimicking the in vivo situation for DS and MT effects, exchange rates, relaxation times and different experimental conditions. Although it is not possible to completely simulate an in vivo environment, major contributors to in vivo Cr-w CEST contrast have been evaluated in this study. For optimization, we have not simulated contribution from rNOE to Cr-w CEST as we planned to separate out rNOE using a fitting-based approach. Based upon preliminary observations using simulations, it was possible to reduce overlap and obtain optimal Cr-w CEST contrast. However, it was difficult to find saturation parameters to completely remove the overlap from other metabolites. Simulation results show that 10mM of Cr provide ~ 4.3% of CEST effect at 1.8 ppm using a  $B_{1rms}$  of 1.45  $\mu$ T and duration of 2 seconds at 7 T. It was also observed that PCr exhibits a small CEST effect at 1.8 and 2.5 ppm.<sup>19</sup> Due to a slow exchange rate compared with amine protons of Cr, the CEST effect from PCr is negligible (~ 0.2% per mM). Amine protons of GABA also exhibit the CEST effect at ~ 3 ppm, with an exchange rate similar to that of glutamate. The concentration of GABA in human brain is ~ 1-2mM. Therefore, in this study, GABA was considered a part of glutamate for the purposes of simulation.

The superposition of Lorentzian functions has been successfully used to fit in vivo Z-spectra data from rat brain at 9.4 T.<sup>23</sup> In the reported studie,<sup>23</sup> CEST peaks at 3.5 and 1.8 ppm were visible and fitted nicely to the superposition of Lorentzian functions. However, the in vivo Z-spectral data obtained at 7 T using saturation parameters (1.4  $\mu$ T, 2 seconds) showed broad CEST effects, and because of this model 1 resulted in higher values of width compared with model 2. Model 2 fits a broad MT component using partial Z-spectra data with negligible contribution from other effects and removes this component from the Z-spectra before finer fitting. This approach results in lower values of the width of rNOE and CEST pools compared with model 1. The amplitude of rNOE showed a clear decrease for model 2. As such, both models fitted reasonably well to Z-spectra data and there were differences in the width and amplitude. However, it is difficult to compare these two models on the basis of only residual error, and change in width and amplitude. Thus, Monte-Carlo simulations were carried out to check the stability of both models and to determine which one works better. Monte-Carlo simulations showed that ARPE decreases for all of the parameters, except for MT in model 2 compared with model 1. Thus, it can be concluded that model 2 is less sensitive to noise than model 1.

A previous study<sup>50</sup> at 9.4 T compared different fitting quantification methods such as Lorentzian difference, multiple-pool Lorentzian fitting or model 1 and three-point methods and reported that model 1 was better than the other methods. However, model 2 was not included in this study and the saturation parameters were quite different compared with the current study. The fitting model accuracy depends on saturation parameters and field strength, etc. Therefore, one should evaluate the accuracy of different models for their Z-spectra data. Our in vivo human study at 7 T showed that model 2 is slightly better than model 1.

In general, contrast computed using CEST asymmetry analysis is highly sensitive to  $B_0$  field inhomogeneity, and accurate  $B_0$  mapping and correction are required before interpretation of the results. On the other hand, CEST contrast computed using a fitting-based approach is less sensitive to  $B_0$  field inhomogeneity and there is no need to acquire separate WASSR/ $B_0$ map data. An estimate of DS offset frequency based upon the minimum value of Z-spectra followed by an improvement during fitting is good enough for automatic correction of  $B_0$ field inhomogeneity artifacts.

In this study, Cr-w CEST contrast in the GM and WM tissue was  $6.98 \pm 0.31\%$  and  $5.45 \pm 0.16\%$ , respectively. A spectroscopic studies have reported that the Cr ratio in GM/WM is in the range of 1.2 to 1.5.<sup>13</sup> In the current study, the Cr-w CEST contrast ratio in GM/WM was 1.28 and 1.43 using models 2 and 1, respectively. Therefore, Cr-w CEST correlated with spectroscopy-based Cr measurements in GM and WM tissues.

For brain tissue, a ratio of 2:1 is reported for the Cr/PCr ratio. However, for the saturation parameter used in the current study, the PCr contribution is negligible in brain tissue. As such, although we have tried to mitigate overlap from different effects using optimized saturation parameters and a fitting-based approach, the Cr-w CEST contrast might still have some overlap. However, a major contribution from Cr to Cr-w CEST contrast cannot be ruled out. A recently published study<sup>24</sup> on rat brain data has reported that amine protons from side chain proteins also exhibit the CEST effect at ~ 2 ppm, which can also overlap with Cr-w CEST contrast. Another study<sup>51</sup> also showed that CEST effects from side chain protein amines are very broad and can contribute to the CEST signal at ~ 2 ppm. However, due to the use of low saturation power in the current study, contribution from these fast exchanging amine protons should be small. Further studies are required to estimate the overlap of such broad peaks to CrCEST.

The chemical shift  $(\Delta \omega)$  of the exchangeable proton is proportional to field strength while the exchange rate is constant. In the current study, in which we performed the phantom experiments at 7 T, due to the slow intermediate exchange rates of Cr it is possible to observe the Cr-w CEST effect even with low field scanners such as 3 T.<sup>22</sup> The feasibility of Cr-w CEST mapping for calf muscle and myocardium has been demonstrated previously at 3 T.<sup>21,22</sup> The CEST effect is also highly dependent on saturation pulse parameters, field homogeneity and tissue type. Further studies to optimize the parameters for in vivo mapping of Cr-w CEST in human brain on a lower field scanner are required.

In this study, MT components of Z-spectra data were fitted using Lorentzian, Gaussian and super-Lorentzian line shapes. It was observed that all three line shapes fitted reasonably well to far-off resonance MT components of Z-spectra data, having negligible contribution from other effects. However, for model 1, super-Lorentzian line shape resulted in poor fitting close to water resonance due to the presence of a large hump at -2.4 ppm. Similarly, for model 2, the interpolation of fitted super-Lorentzian component created a large hump at -2.4 ppm and hence resulted in poor fitting of Z-spectra, particularly close to water resonance. Due to this pole of super-Lorentzian line shape, a recent study<sup>52</sup> has used symmetric interpolation between  $\pm 20$  ppm using a cubic-spline interpolation to remove the pole that occurs close to its resonance frequency. In the current study, Lorentzian line shape for MT provided the least residual error compared with other line shapes. Several other studies<sup>22,23,34</sup> have also used either Lorentzian line shapes or polynomials. Computation of rNOE and CEST components using a multipool Z-spectra-fitting approach (models 1 and 2) with super-Lorentzian line shape for MT pool is not suitable for the data presented in the current study. It should be noted that when super-Lorentzian line shape for MT pool was proposed for Z-spectra fitting, CEST and rNOE were not considered and the main emphasis was on fitting Z-spectra over far-off resonance frequencies. The previously reported analytical model was valid for two-pool Z-spectra under steady state conditions. Z-spectra data acquired in the current study are nonsteady state, have multiple pools and were acquired using low saturation B<sub>1rms</sub> and were optimized for Cr-w CEST contrast for a limited frequency range of -40 to 40 ppm. Due to the presence of multiple pool and nonsteady state conditions, the derivation of analytical models is not possible. Therefore, for such nonsteady state data, empirical models based upon combinations of different line shapes such as sum of Lorentzians can be used. Thus, for the set of saturation parameters used in this study, Lorentzian line shape is a better approximation to the MT component of Z-spectra and thus we used Lorentzian line shape for removal of the MT component.

In general, the CEST MRI signal depends upon direct saturation of free water pool, semisolid macromolecular concentration (MT effect), solute pool concentration, exchange rate, pH, water longitudinal relaxation time (T<sub>1</sub>), saturation pulse parameters and field strength. Thus, CEST contrast is usually compared with similar saturation parameters and at the same field strength. CEST contrast computed using a conventional CEST<sub>asy</sub> analysis still has some tissue-specific bias due to differences in MT effect and T<sub>1</sub>. Different approaches have been adopted to mitigate these effects on computed CEST contrast. The AREX method<sup>41</sup> and normalization by negative offset frequency<sup>4,49</sup> are two such examples. In the current study, CrCEST contrast between GM and WM tissue also has some bias due to different MT effects and T1 in these tissues. CEST contrast computation using the proposed Z-spectra fitting using both models 1 and 2 mitigates the MT effect and is less sensitive to T<sub>1</sub>; however, complete removal of these effects from true CEST contrast has not been achieved to date. Another study<sup>53</sup> reported that for nonsteady-state acquisition, the effect of T<sub>1</sub> is substantially less. In this study, we have considered only two CEST pools and a single rNOE. For testing purposes, when additional CEST pools were included in the Z-spectra fitting then the stability of the model to noise was reduced. An additional rNOE pool centered at -1.6 ppm was also considered for the Z-spectra fitting, which resulted in negligible amplitude corresponding to  $B_{1rms}$  of 1.4  $\mu$ T. The same additional pool resulted in

a broad line shape for  $B_{1rms}$  of 0.7  $\mu$ T and a narrow line shape for  $B_{1rms}$  of 2.1  $\mu$ T. These results show that the fitting of this additional rNOE pool is not reliable for the saturation parameters and models considered in the current study.

Limitations of the current study are stated below along with directions for future research. The acquisition of Z-spectra data over a wide range of offset frequencies increases the scan time, as do the extra data required for  $B_1$  inhomogeneity correction. We have tried to optimize the data acquisition strategy to minimize the number of data points required. Alternative approaches based upon machine learning can be considered in future studies to reduce the scan time. Cr-w CEST contrast in GM and WM was correlated with a reported spectroscopy-based ratio of [Cr]. In future study, MR spectroscopic data could be collected in representative GM and WM regions to evaluate subject-specific correlation between [Cr] and Cr-w CEST contrast.

# 5 | CONCLUSION

It is feasible to perform Cr-w CEST mapping in human brain at 7 T using optimal saturation parameters and a Z-spectra fitting approach. In the limited, far-off resonance frequency range, the Lorentzian line shape provided a slightly better approximation for the broad MT component of Z-spectra compared with Gaussian and super-Lorentzian functions. Although goodness of fit and residual errors of fit were similar in both Z-spectra fitting models, Monte-Carlo simulations showed that model 2 is more stable to noise than model 1. Imaging protocol and the fitting approach proposed in the current study can be used for acquiring patient data and computing Cr-w CEST contrast.

# Supplementary Material

Refer to Web version on PubMed Central for supplementary material.

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# Abbreviations used:

ARPE	absolute relative percentage error
CEST	chemical exchange saturation transfer
CEST@1 ppm	CEST pool centered at 1 ppm
CEST@3.5 ppm	CEST pool centered at 3.5 ppm
Cr-w	creatine-weighted
GM	gray matter

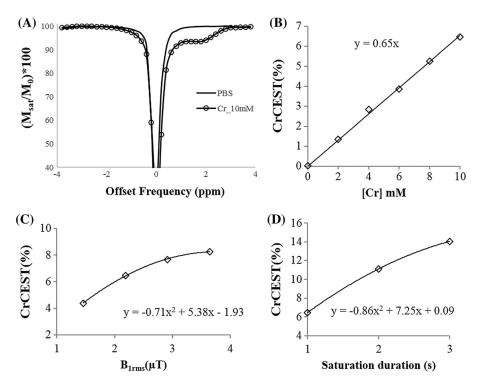
RF	radio frequency
rNOE	relayed nuclear Overhauser effect
WM	white matter

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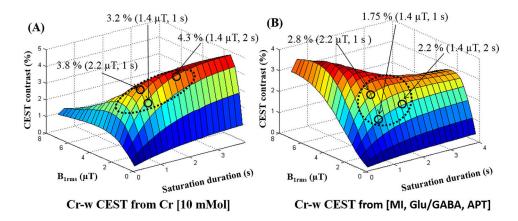
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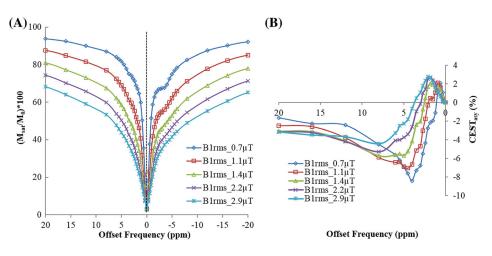
# FIGURE 1.

(A) Z-spectra of Cr (10 mM) and PBS using a saturation  $B_{1rms}$  of 2.2  $\mu$ T and duration of 1 second. (B) Dependence of Cr-w CEST as contrast on Cr concentration ([Cr]) using a  $B_{1rms}$  of 2.2  $\mu$ T and duration of 1 second. (C) Dependence of Cr-w CEST contrast on saturation power for a fixed duration of 1 second and (D) dependence on saturation duration for a fixed saturation  $B_{1rms}$  of 2.2  $\mu$ T. Phantom data was acquired at temperature of  $37 \pm 1^{\circ}$ C and pH of 7.0. Here, y represents CrCEST (%) and x represents [Cr] in (B),  $B_{1rms}$  in (C) and saturation duration in (D)



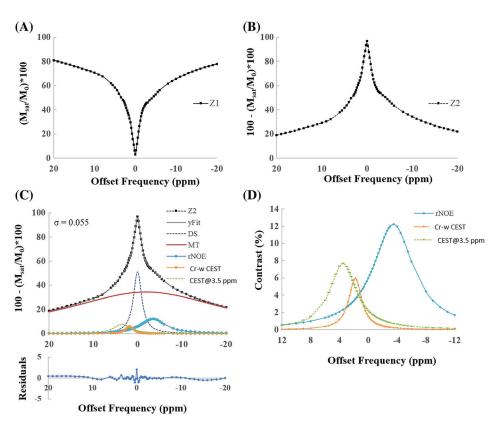
## FIGURE 2.

Surface plots demonstrating dependence of Cr-weighted (w) CEST contrast at 1.8 ppm on  $B_{1rms}$  and saturation duration using numerical simulations at 7 T. (A) CEST contrast from Cr [10mM] only. (B) CEST effect from combined metabolites [MI, Glu/GABA, APT] except Cr at 1.8 ppm. Arrows point to the three different saturation values and corresponding contrast. Circles show regions corresponding to optimum range of saturation parameters. Dotted ellipse highlight the regions for optimal CrCEST contrast with less contamination. APT, amide proton transfer; Glu, glutamate; MI, myo-inositol



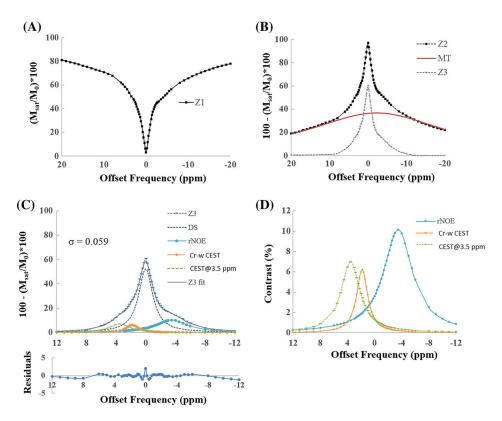
# FIGURE 3.

(A) Z-spectra and (B) asymmetry plot from a region of interest in gray matter tissue of human brain at 7 T with a saturation pulse  $B_{1rms}$  of 0.7, 1.45, 2.2 and 2.9  $\mu$ T and duration of 2 seconds. Asymmetry plots in (B) use normalization by signal without saturation (M<sub>0</sub>)



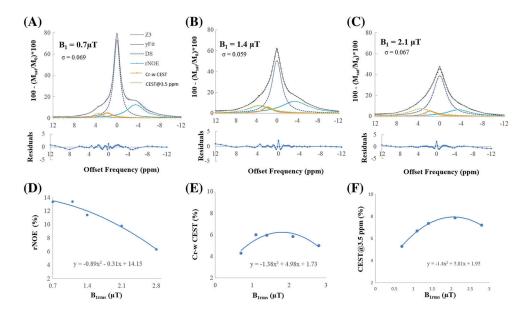
# FIGURE 4.

Plots show an example of fitting in vivo Z-spectra, from a region of interest in gray matter tissue, for a  $B_{1rms}$  of 1.45 µT and duration of 2 seconds using model 1. (A) Original Z-spectra (Z1). (B) Scaled Z-spectra (Z2). (C) Five-pool fitting of scaled Z-spectra (Z2) along with various fitted components DS, MT, rNOE-w, Cr-w CEST and CEST@3.5 ppm contrast.  $R^2 = 0.99$  and residual errors are less than 2%. (D) Individual CEST and rNOE components. Fitting is performed over a range of -40 to +40 ppm and shown for a range of -20 to +20 ppm for better visualization. CEST@3.5 ppm, CEST contrast at 3.5 ppm; Cr-w CEST, CEST contrast at 1.8 ppm; DS, direct saturation; MT, magnetization transfer; rNOE, relayed nuclear Overhauser effect;  $\sigma$ , average residual error;  $R^2$ , goodness of fit



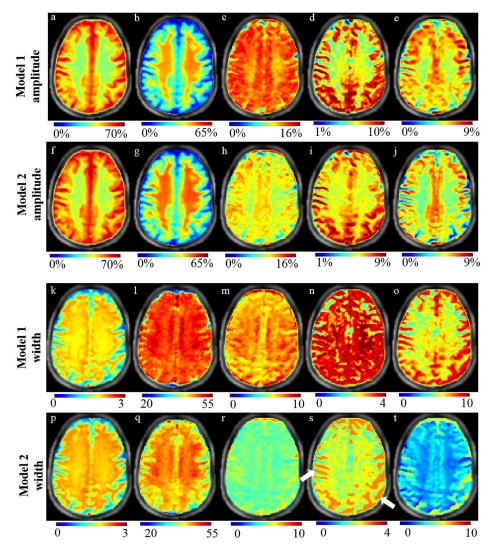
# FIGURE 5.

Plots show an example of fitting in vivo Z-spectra, from region of interests in gray matter tissue, for a  $B_{1rms}$  of 1.45 µT and duration of 2 seconds using model 2. (A) Original Z-spectra (Z1). (B) Scaled Z-spectra (Z2) along with fitted MT component and Z-spectra after removal of MT component (Z3). (C) Four-pool fitting of Z3 along with various fitted components like DS, rNOE, Cr-w CEST and CEST@3.5 ppm contrast.  $R^2 = 0.99$  and residual errors are under 2%. (D) Individual CEST and rNOE components. Fitting of Z2 is performed over a range of -40 to +40 ppm and shown for a range of -20 to +20 ppm for better visualization. CEST@3.5 ppm, CEST contrast at 3.5 ppm; Cr-w CEST, CEST contrast at 1.8 ppm; DS, direct saturation; MT, magnetization transfer; rNOE, relayed nuclear Overhauser effect;  $\sigma$ , average residual error;  $R^2$ , goodness of fit



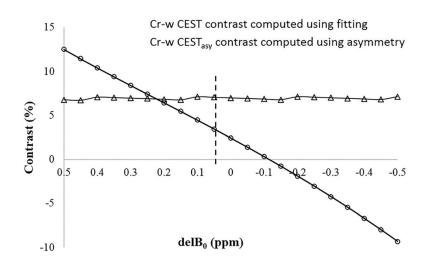
# FIGURE 6.

Fitting of in vivo Z-spectra from gray matter tissue using model 2 at durations of 2 seconds and  $B_{1rms}$  of (A) 0.7, (B) 1.45 and (C) 2.2  $\mu$ T. Residual errors of fitting are less than 2%. rNOE decreases with increase in  $B_{1rms}$ . Cr-w CEST and CEST@3.5 ppm initially increase and then decrease with  $B_{1rms}$ . CEST@3.5 ppm, CEST contrast at 3.5 ppm; Cr-w CEST, CEST contrast at 1.8 ppm; DS, direct saturation; MT, magnetization transfer; rNOE, relayed nuclear Overhauser effect;  $\sigma$ , average residual error; R<sup>2</sup>, goodness of fit



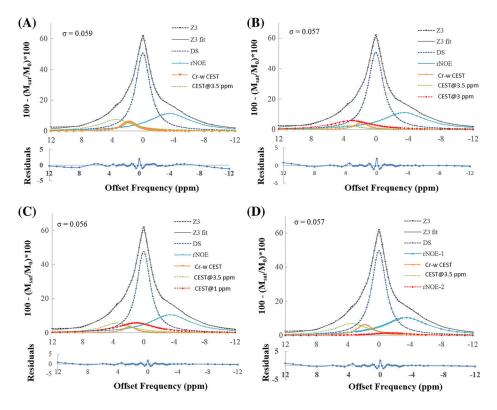
# FIGURE 7.

Amplitude and width maps of various components of Z-spectra from a representative healthy human brain using model 1 (rows 1 and 3) and model 2 (rows 2 and 4). Amplitude maps of DS (A, F), MT (B, G), rNOE (C, H), Cr-w CEST (D, I), CEST@3.5 ppm (E, J). Width maps of DS (K, P), MT (L, Q), rNOE (M, R), Cr-w CEST (N, S), CEST@3.5 ppm (O, T). The scales for width maps are in ppm. White arrows point to representative subarachnoid space of brain. Images are cropped and descalped



#### FIGURE 8.

Plots show the sensitivity of Cr-w CEST contrast computed using fitting approach to shift in resonance frequency (delB<sub>0</sub>). Z-spectra from a region of interest in gray matter tissue was selected followed by fitting. Simulated Z-spectra were generated by varying the center, corresponding to direct saturation of free water obtained using fitting, from -0.5 to 0.5 ppm. These simulated Z-spectra were fitted using model 2 for computing Cr-w CEST as well as for computing asymmetry contrast



# FIGURE 9.

Fitting of Z-spectra using model 2 at a B<sub>1rms</sub> of 1.45  $\mu$ T and duration of 2 seconds using multi-pool models. (A) Fitting using a five-pool model (DS pool at 0 ppm, MT pool at -2.4 ppm, rNOE pool at -3.5 ppm, Cr-w CEST pool at 1.8 ppm and CEST@3.5 ppm). (B) Fitting considering six pools, where the first five pools are similar to those of (A) and the sixth pool (CEST@3 ppm) is considered at 3 ppm. (C) Sixth pool (CEST@1 ppm) is considered at 1 ppm. (D) Sixth pool (rNOE-2) is considered at -1.6 ppm. Residual errors of fittings were less than 2% and  $\sigma = -0.05$  for all cases. CEST@3 ppm, CEST contrast at 3 ppm; CEST@3.5 ppm, CEST contrast at 3.5 ppm; Cr-w CEST, CEST contrast at 1.8 ppm; DS, direct saturation; MT, magnetization transfer; rNOE, relayed nuclear Overhauser effect; rNOE-2, rNOE pool at -1.6 ppm;  $\sigma$ , average residual error; R<sup>2</sup>, goodness of fit

#### TABLE 1

Simulation results on potential contributions to in vivo human brain creatine-weighted chemical exchange saturation transfer contrast at 7 T

Metabolites	$T_{1}\left(s\right)$	$T_{2}\left(s ight)$	ω (ppm)	Concentration (M)	Exchange rate (Hz)
Free water <sup>a</sup>	2	0.115	0.0	81	-
Bound water <sup>b</sup>	1	0.00001	-2.4	7	10
Amide <sup>C</sup>	1	0.033	3.5	0.072	30
Creatine <sup>d</sup>	1	0.010	1.8	0.010	500, 900, 1190
Glutamate <sup>e</sup>	1	0.010	3.0	0.011	2000
$Myo-inositol^f$	1	0.010	0.6	0.006	600

<sup>*a*</sup>For simulation,  $T_1$  and  $T_2$  for free water were taken from reference.<sup>4</sup>

 $^{b}$  For simulation, the bound water parameters were used from reference <sup>51</sup>. This reflects MTR asymmetry contribution.

<sup>c</sup> For simulation, all parameters were taken from reference.<sup>3</sup>

 $^{d}$ For simulation, the concentrations were taken from references 47,48, and exchange rates of 500, 900 and 1190 Hz were taken from references, 4,19 and,<sup>25</sup> respectively. Four amine protons were considered for each creatine molecule.

 $^{e}$ For simulation, the concentrations were taken from references 52,53, and exchange rate measured in vivo from rat brain.<sup>4</sup> Three amine protons were considered for each glutamate molecule.

f For simulation, the concentration and exchange rate were taken from reference<sup>30</sup> and relaxation parameters for -OH were assumed to be the same as creatine. Six -OH protons were considered for each myo-inositol molecule.

# TABLE 2

Fitted amplitude parameter obtained at  $B_{1rms}$  of 1.4  $\mu$ T and duration of 2 seconds using model 2 for multiple regions of interest on gray matter (GM) and white matter (WM) tissue averaged over all healthy human volunteers

	Amplitude (%	6)			
	DS	MT	rNOE	Cr-w CEST	CEST@3.5 ppm
GM	$55.08 \pm 0.52$	$32.77 \pm 1.46$	$10.86\pm0.67$	$6.98 \pm 0.31$	$7.01\pm0.31$
WM	$38.9\pm0.8$	$50.51 \pm 1.01$	$9.74 \pm 0.25$	$5.45\pm0.16$	$4.79\pm0.28$

Abbreviations: CEST@3.5 ppm, CEST contrast at 3.5 ppm; Cr-w CEST, CEST contrast at 1.8 ppm; DS, direct saturation; MT, magnetization transfer; rNOE, relayed nuclear Overhauser effect.

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# TABLE 3

Comparison of model 2 and model 1 fitting to Z-spectra data obtained at B<sub>1rms</sub> of 1.4 µT and duration of 2 seconds from gray matter tissue of human brain using Monte-Carlo simulations

		Amplitude (%)					Width (ppm)				
		DS	TM	rNOE	Cr-w CEST	CEST@3.5 ppm	DS	MT	rNOE	Cr-w CEST	CEST@3.5 ppm
Model 2 Par0	ar0	52.22	36.51	10.38	6.22	7.09	2.31	45.12	5.27	1.99	3.39
m	ean $\pm$ SD	mean $\pm$ SD 51.96 $\pm$ 1.61	$38.38 \pm 2.01$	$11.12\pm1.68$	$6.86 \pm 1.89$	$7.19 \pm 1.73$	$2.30\pm0.12$	$50.12\pm3.65$	$50.12 \pm 3.65$ $5.51 \pm 1.14$	$2.25\pm0.90$	$3.60\pm1.29$
A	ARPE	2.46	6.18	14.68	25.96	19.92	4.27	12.12	17.60	38.74	31.70
Model 1 Par0	ar0	51.07	33.99	12.55	5.96	7.79	2.27	46.72	6.96	2.19	4.58
m	ean $\pm$ SD	mean $\pm$ SD 50.55 $\pm$ 2.02	$34.90\pm2.42$	$34.90 \pm 2.42$ $13.97 \pm 2.15$	$6.64\pm2.16$	$8.04\pm2.00$	$2.25\pm0.13$	$53.75 \pm 5.64$	$53.75 \pm 5.64$ $7.65 \pm 1.79$ $2.46 \pm 1.01$	$2.46 \pm 1.01$	$5.41\pm2.02$
IA	ARPE	3.33	6.06	17.15	29.55	20.79	4.56	16.13	22.66	41.70	37.97

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