

Packing, tiling, and covering with tetrahedra

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It is well known that three-dimensional Euclidean space cannot be tiled by regular tetrahedra. But how well can we do? In this work, we give several constructions that may answer the various senses of this question. In so doing, we provide some solutions to packing, tiling, and covering problems of tetrahedra. Our results suggest that the regular tetrahedron may not be able to pack as densely as the sphere, which would contradict a conjecture of Ulam. The regular tetrahedron might even be the convex body having the smallest possible packing density.

tessellations | polyhedra

The problem of how densely given solid objects can pack in space has been a source of fascination since the dawn of civilization. Dense packing of convex objects is intimately related to the arrangement of molecules in condensed states of matter (1) and to the best way to transmit encoded messages over a noisy channel (2). Three-dimensional Euclidean space \mathbb{R}^3 (3-space) already provides many challenging open problems. It was only recently that Kepler's conjecture, which postulated that the densest packings of congruent spheres in 3-space have packing density (fraction of space covered by the spheres) $\Delta = \pi/\sqrt{18} = 74.048\dots\%$, realized by variants of the face-centered cubic (FCC) lattice packing, was proved (3). Much less is known about the packing characteristics of other congruent convex objects that do not tile 3-space. For example, an ellipsoid is simply obtained by an affine (linear) transformation of a sphere, and yet the densest packing of ellipsoids is an open problem. The rotational degrees of freedom of an ellipsoid (absent in a sphere) enables such packings to achieve densities greater than $\pi/\sqrt{18}$, the densest sphere packing density (4–6). There is a family of periodic arrangements of nearly spherically shaped ellipsoids that surpass the density of the optimal sphere packing and that has a maximal density of $\Delta = 0.7707\dots\%$ (6), which is the highest known density for any ellipsoid packing.

The evidence below suggests that the regular tetrahedron is a counterexample to Ulam's conjecture (Martin Gardner, private communication; see also ref. 7), which states that the optimal density for packing congruent spheres is smaller than that for any other convex body. Indeed, it suggests that perhaps the regular tetrahedron achieves this minimum. However, our interest in tetrahedra in this work goes beyond their packing characteristics. Tetrahedra have interesting connections to sphere packings, certain tilings of space (including foams), liquids, and glasses, and complex alloy structures. It is well known that the maximum number of spheres in 3-space that can be locally packed such that each sphere contacts the others is four. The polyhedron that results by taking the sphere centers as vertices is the regular tetrahedron, but such a tetrahedron cannot tile space because its dihedral angle $[\cos^{-1}(1/3) \approx 70.53^\circ]$ is not a submultiple of 360° . Interestingly, the ratio of the volume of the portion of this tetrahedron covered by the spheres to the volume of the tetrahedron leads to the Rogers upper bound of $77.96\dots\%$ on the sphere packing density (8).

It was Frank and Kasper (9, 10) who proposed that the underlying "polytetrahedral" network of sphere packings can serve to explain the crystalline structure of complex alloys, particularly those of transition metals. These arrangements are now known as "Frank–Kasper" phases. If 5 regular tetrahedra are packed around a common edge, there remains a small gap of 7.36° , and if 20 regular tetrahedra are packed around a common vertex, the gaps amount

to a solid angle of 1.54 steradians (see Fig. 1). By closing these gaps by a slight deformation, we get a regular icosahedron, which corresponds to a 12-coordinated sphere whose vertices correspond to the vertices of the icosahedron. A 12-fold coordination is an important case of the Frank–Kasper phases; others include 14-, 15-, and 16-fold coordinations (13 is not possible), corresponding to triangular-faced tetrahedra constructed from tetrahedra sharing a single vertex. The Frank–Kasper phases thus consist of various tilings of space by "almost-regular" tetrahedra with atoms at their vertices. It also has been shown that the structure of atomic liquids and glasses has significant polytetrahedral character (11). The dual Voronoi regions of the vertices of the Frank–Kasper structures (obtained by joining the centers of every pair of tetrahedra in face contact) together fill space and consist of polyhedra with 12, 14, 15, and 16 faces. It is well known that periodic structures with atoms at the vertices of these dual tilings occur in clathrate hydrates (12). This type of polyhedral tiling inspired Weaire and Phelan (13) in their discovery of a minimal area foam with a smaller average surface area per cell ("bubble") than Kelvin's best solution (14).

The fact that congruent regular tetrahedra cannot be used to tile 3-space gives rise to several mathematical questions.

- (i) What are the "closest-to-regular" tetrahedra that will tile 3-space?
- (ii) What is the least covering density for congruent regular tetrahedra in 3-space?
- (iii) What is the greatest packing density for such tetrahedra?

Of course question (i), which we discuss first, is itself several different problems, according to how we interpret closest-to-regular. We will show that our answers to these questions will introduce us to many of the polyhedral tilings of space discussed immediately above. The best solution we offer for question (iii) leads to the possibility that the regular tetrahedron contradicts Ulam's conjecture and might itself have the least packing density of any convex body.

Scottish, Irish, and Welsh Configurations

In 1887, Lord Kelvin conjectured that a certain system of equal-volume bubbles, which we shall call the "Scottish bubbles," was optimal in the sense that it minimized the mean surface per bubble (14). This conjecture was disproved by the "Irish bubbles" found by Weaire and Phelan in 1994 (13). We define both systems in this work, along with a third system, the "Welsh bubbles." Many of the other configurations we need are related to these three bubble-systems and so are also appropriately described as Scottish, Irish, and Welsh.

For example, the Scottish, Irish, and Welsh

- "Nuclei" are the centers of the bubbles;
- "Vocells" are the cells of the Voronoi tessellation that they determine;

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Abbreviations: 3-space, three-dimensional Euclidean space \mathbb{R}^3 ; BCC, body-centered cubic.

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Table 3. Typical Irish tetrahedra

Points	Irish high: 8 symmetries		Irish medial: 3 symmetries		Irish low: 2 symmetries	
	Coordinates	Type	Coordinates	Type	Coordinates	Type
Vertices	(-1/4, 1/2, 0)	+	(0, 0, 0)	$\mathbf{0}$	(0, 0, 0)	+
	(1/4, 1/2, 0)	+	(0, 1/4, 1/2)	+	(-1/4, 1/2, 0)	+
	(0, 1/4, 1/2)	+	(1/2, 0, 1/4)	+	(1/4, 1/2, 0)	+
	(0, 3/4, 1/2)	+	(1/4, 1/2, 0)	+	(0, 1/4, 1/2)	+
Node	(0, 1/2, 1/4)	-	(5/24, 5/24, 5/24)	$\mathbf{0}_3$	(-5/32, 0, 10/32)	$-\mathbf{0}$

Dodecahedral cells have nuclei $\mathbf{0}$, $\mathbf{1}$, $\mathbf{2}$, or $\mathbf{3}$ and volume 125/1024. Dodecahedral ones have nuclei + and volume 129/1024. Here $\mathbf{0}_3 = (7/12)\mathbf{0} + (5/12)\mathbf{3}$ and $-\mathbf{0} = (5/8)\mathbf{-} + (3/8)\mathbf{0}$. Nodes are all points $-\mathbf{0}$, $\mathbf{1}$, $\mathbf{2}$, $\mathbf{3}$, $\mathbf{0}_3$, $\mathbf{1}_0$, $\mathbf{2}_1$, $\mathbf{3}_2$, $-\mathbf{0}$, $-\mathbf{1}$, $-\mathbf{2}$, or $-\mathbf{3}$ with above proportions.

divided into identical regions F called “fundamental cells,” each of which contains the center of just one body. We denote by $\text{Vol}(C)$ and $\text{Vol}(F)$ the volumes of the convex body and fundamental cell, respectively. The packing density of a lattice packing P_L is therefore given by

$$\Delta = \frac{\text{Vol}(C)}{\text{Vol}(F)}. \quad [2]$$

A “periodic packing” P_P of congruent copies of a convex particle C is obtained by placing a fixed nonoverlapping configuration of n particles (where $n \geq 1$) in each fundamental cell of a lattice. Thus, the packing is still periodic under translations by a lattice vector, but the n particles can be positioned anywhere in the fundamental cell with arbitrary orientation subject to the nonoverlap condition. The density of such a periodic packing is given by

$$\Delta = \frac{n\text{Vol}(C)}{\text{Vol}(F)}. \quad [3]$$

There are more general types of packings, but in this work we will restrict ourselves to periodic packings.

We now turn to the packing of congruent regular tetrahedra. Our first remark is that the method of (Bravais) lattice packing, which produces good packings for many other solids (including an optimal sphere packing; see ref. 2), is of no use here. The optimal lattice packing for any tetrahedron, found by Hoylman (16), has density $\Delta = 18/49 = 36.73 \dots \%$, and each tetrahedron meets others at 14 points (see Fig. 3). [It is noteworthy that there is a lattice packing of tetrahedra with much smaller density ($\Delta = 1/3$) but in which each tetrahedron is in contact with 18 others (17), a rather counterintuitive result.] Clearly, denser packings can be achieved by orienting the tetrahedra in different directions. Fig. 4 shows a simple packing

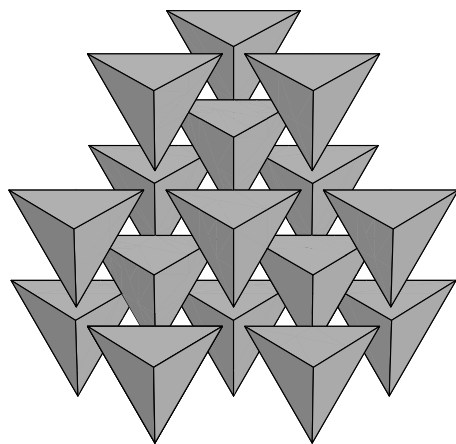


Fig. 3. A portion of the densest (Bravais) lattice packing of regular tetrahedra (16). It has density $\Delta = 18/49 = 36.73 \dots \%$, and each tetrahedron is in contact with 14 others.

that achieves density $\Delta = 2/3 = 66.666 \dots \%$. This is the best density we have been able to achieve with a “uniform packing,” i.e., one in which the tetrahedra are embedded in the same way, meaning that there is a symmetry of the packing that takes any one tetrahedron to any other.

Scottish and Welsh Regulars. Another idea is to insert regular tetrahedra into one of our three systems of irregulars. From Irish irregulars, this idea produces a very low density, and from Scottish ones, a density Δ of $1/2 = 50\%$ is produced (the regular tetrahedra in that case being obtained by shrinking the long edges of the Scottish irregulars by a factor of $\sqrt{1/2}$). However, a more interesting packing can be obtained from the Scottish Vocells as follows. There is a well known way [see, e.g., Coxeter (18)] to inscribe an icosahedron in an octahedron. In fact, the icosahedron fits entirely inside the corresponding truncated octahedron (see Fig. 5) and occupies 8/9 the volume of the truncated octahedron. We obtain “Scottish icosahedra” by inscribing icosahedra in this way in all of the Scottish truncated octahedra. Each icosahedron touches eight others as in Fig. 6, the contact spots of that figure (which form the vertices of a cube) being at the centers of eight faces.

The tetrahedra of the Scottish regular packing (or just the “Scottish regulars”) are obtained by packing 20 tetrahedra in each icosahedron (see *Appendix*). This tetrahedron packing has density $\Delta = 45/64 = 70.3125\%$. We shall see later that the Scottish regulars can be displaced slightly to increase their density to $>71.655\%$.

Another good packing is obtained from the Welsh irregulars. We retain the high ones that are already regular and put regular tetrahedra of the same size as these into each of the medial and low ones. The density of the resulting Welsh regulars is $\Delta = 17/24 = 70.8333 \dots \%$. Once again, there exist displacements that increase the density.

Displaced and Reformed Regulars. We have remarked that the density of both Scottish and Welsh regulars can be improved by suitably repositioning them. The tetrahedra fall into clumps, and we

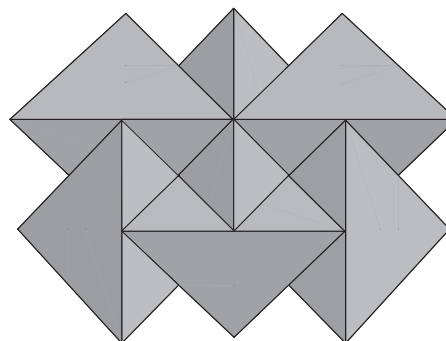


Fig. 4. A portion of the densest uniform packing of regular tetrahedra that we have been able to find. It has density $\Delta = 2/3 = 66.666 \dots \%$.

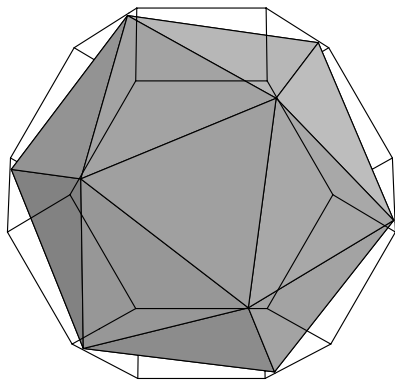


Fig. 5. A regular icosahedron inscribed in a truncated octahedron.

can distinguish between “displacements,” which treat the clumps bodily, and “reformations,” which then move individual tetrahedra.

The Scottish regulars do not form an optimal packing because the icosahedra are related by the translations of the BCC lattice, and this lattice can be slightly deformed in a way that increases the density without causing the icosahedra to overlap. The reason, in brief, is that any icosahedron I (centered at zero) touches only eight others, namely, $I \pm v_1, \dots, I \pm v_4$, where v_1, \dots, v_4 are generators of the BCC lattice that satisfy $v_1 + \dots + v_4 = 0$. They can be replaced by four nearby vectors w_1, \dots, w_4 , provided that $w_1/2, \dots, w_4/2$ lie in the same faces of the icosahedron and add to zero. Since this requirement imposes only four conditions on the lattice, whereas six parameters are needed to specify the shape of a lattice, it can be varied with two degrees of freedom, and it turns out that some values increase the density. The optimal (Bravais) lattice packing of icosahedra (the “displaced Scottish icosahedra”) was found by Betke and Henk (19); it has the density $\Delta = 83.63574 \dots \%$ and yields the displaced Scottish packing of regular tetrahedra that has density $\Delta = 71.65598 \dots \%$ (see Fig. 7).

However, this density is still not optimal, because the clumps may be reformed by adjusting the individual tetrahedra so as to increase the density still further. To see this, observe that none of the “contact spots” S of Fig. 7 lies at the center C of its face.

Now consider two tetrahedra from clumps centered at P and P' , whose contact spots S and S' presently coincide with each other, but not with the centers C and C' of their corresponding faces. Then they can be rotated about axes through P and P' perpendicular to the line CC' so as to take S and S' further away from the initial positions of C and C' .

Do likewise for all tetrahedra of every clump. After this rotation, if it is through a sufficiently small angle, the tetrahedra still do not overlap and now touch only at the centers of their clumps (see Fig.

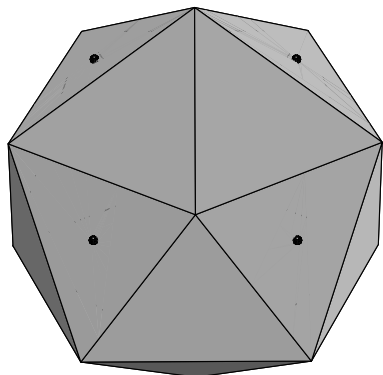


Fig. 6. Each Scottish icosahedron is placed so that its eight contact spots coincide with those of its neighbors. It has density $\Delta = 82.13 \dots \%$.

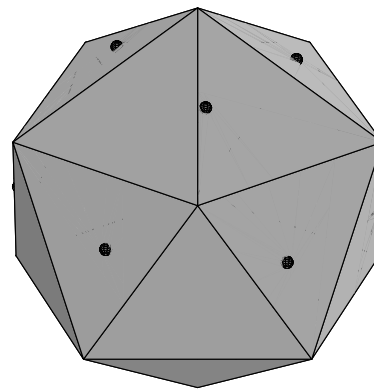


Fig. 7. For Betke and Henk’s displaced Scottish packing of icosahedra, the number of contact spots increases to 12. It has density $\Delta = 83.63574 \dots \%$ and leads to the displaced Scottish regulars that have density $\Delta = 71.65598 \dots \%$.

8). The density therefore can be increased by bringing the clumps closer together.

It is very difficult to say exactly how dense such “reformed Scottish” packings can be (especially because they will involve slight changes to the lattice), but we suspect it will be $\approx 72\%$.

Welsh regulars also fail to be optimal, because these tetrahedra fall into “clumps” of 17, and each clump touches others only at four points. If we replace these points by universal joints, the resulting structure is not rigid, and suitable displacements increase density. One such displaced Welsh packing has density $\Delta = 71.7455\%$, the highest we have yet explicitly achieved. It is obtained by rotating the clumps alternately through ± 0.1131 radians about their vertical (dyad) axes. Before this rotation, each low tetrahedron is hinged as far as possible about the edge it shares with a high one, either “centrifugally” outwards if that edge is horizontal and otherwise in the “lagging” direction (thus, each low tetrahedron goes to the place it would if the displacement of the other tetrahedra were impulsive). Again, still denser reformed Welsh packings might be obtainable by allowing individual tetrahedra to move more freely.

There are various systems of Irish regulars, obtained by repositioning (some of) the Irish irregulars. So far, we have not found any that are as dense as our Scottish or Welsh ones, but we cannot rule out this possibility. We therefore suspect that the optimal packing of regular tetrahedra will be some displacement of either the Scottish or Welsh regulars, but we do not know which! However, it appears unlikely that the density of the optimal packing of regular tetrahedra will exceed the optimal density of $74.048 \dots \%$ for congruent spheres.

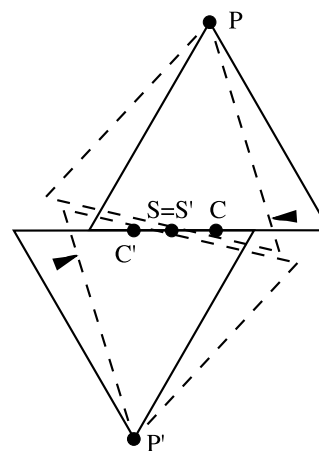


Fig. 8. Adjusting the Scottish regulars (a two-dimensional schematic).

