

α -Helix nucleation by a calcium-binding peptide loop

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ABSTRACT A 12-residue peptide AcDKDGDGY-
ISAAENH₂ analogous to the third calcium-binding loop of
calmodulin strongly coordinates lanthanide ions ($K = 10^5$
 M^{-1}). When metal saturated, the peptide adopts a very rigid
structure, the same as in the native protein, with three last
residues AAE fixed in the α -helical conformation. Therefore,
the peptide provides an ideal helix nucleation site for peptide
segments attached to its C terminus. NMR and CD investi-
gations of peptide AcDKDGDGYISAAEAAAQNH₂ presented
in this paper show that residues A13–Q16 form an α -helix of
very high stability when the La³⁺ ion is bound to the D1–E12
loop. In fact, the lowest estimates of the helix content in this
segment give values of at least 80% at 1°C and 70% at 25°C.
This finding is not compatible with existing helix-coil tran-
sition theories and helix propagation parameters, s , reported
in the literature. We conclude, therefore, that the initial steps
of helix propagation are characterized by much larger s values,
whereas helix nucleation is even more unfavorable than is
believed. In light of our findings, thermodynamics of the
nascent α -helices is discussed. The problem of CD spectra of
very short α -helices is also addressed.

As follows from the Zimm–Bragg theory (1) of the helix-coil transition in polypeptides, the statistical weight of an α -helical segment formed by n amino acid residues with φ and ψ torsional angles close to -57° and -47° , respectively, can be described by the following simple expression (2, 3):

$$k_n = \sigma s^{n-2}, \quad [1]$$

where σ and s are the helix nucleation and propagation parameters, respectively, and $n-2$ is equal to the number of intrahelical hydrogen bonds between CO and NH groups of amino acid residues $i, i+4$.

The propagation parameters s in water solutions have been determined for all 20 amino acid residues from studies of various model peptides (2). For helix makers, $s > 1$ at low (1°C) or even room (25°C) temperatures and long helical segments ($n > 15$) built of such residues are substantially populated. But short helices ($n < 10$) are always very unstable because of the small value of the nucleation parameter σ . Estimated at 0.0013, this value is believed to be the same for all amino acid residues (4).

A straightforward way to stabilize the α -helical conformation in short peptides is, therefore, to introduce an artificial helix nucleation site composed of a few residues fixed in the helical conformation ($\sigma = 1$), or at least with reduced conformational entropy ($0.0013 \ll \sigma < 1$).

Several such systems have been devised and used successfully to nucleate the helix by covalent (5–12) or metal-induced (13, 14) side chain–side chain bridges. All of them reduce the entropy change related to the helix nucleation and, consequently, lead to a large increase of σ parameter, but none of them provides an

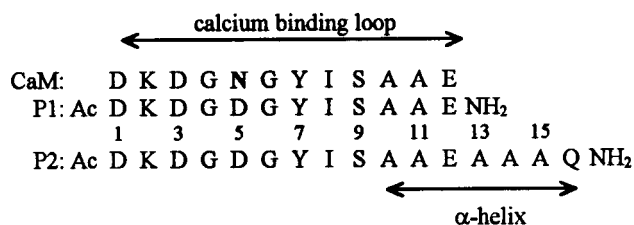


FIG. 1. Sequences of calcium binding loop III of native rat testis calmodulin (CaM) and its synthetic analogs, P1 and P2 peptides.

ideal nucleation site composed of a few residues fixed in the rigid α -helical structure. Moreover, helix geometry at such nucleation sites is often distorted, which may have some destabilizing effect on the helical conformation (15).

The ideal helix-nucleation sites have been created by nature in calcium-binding loops of EF-hand proteins. A typical loop is composed of 12 residues providing six ligands that bind the metal ions directly or through a water molecule. The coordinated metal induces a very rigid structure of the loop in which the last three residues, 10, 11, and 12, are fixed in the α -helical conformation that initiates helix structure flanking the C terminus of the loop. Isolated loops preserve, to a degree, their metal-binding abilities and, therefore, can be used to nucleate α -helix formation in model peptide systems.

Peptide P1, whose sequence is analogous to the third calcium-binding loop of calmodulin (Fig. 1) weakly coordinates calcium with the binding constant of about $10^2 M^{-1}$. Consequently, the calcium-saturated loop still preserves some conformational flexibility (16). Nevertheless, lanthanide ions, known as excellent substitutes for calcium (17, 18), bind to the loop quite strongly, with binding constants of about $10^5 M^{-1}$ (19), and fix it in a very rigid conformation close to that known from the x-ray structure of calmodulin (16).

The main object of our studies was a 16-residue synthetic polypeptide P2 with the sequence shown in Fig. 1, composed of the calcium-binding loop (residues Asp-1–Glu-12) and the sequence of four residues, Ala-Ala-Ala-Gln-NH₂, attached to its C terminus. Glutamine is used here to increase peptide solubility and prevent association of peptide helical segments. Peptides built of Ala₄Gln blocks are well soluble in water and do not show any signs of aggregation up to a concentration of 120 μM (20). Such peptides were used to study helix-coil transitions in water (20, 21) as well as in water–trifluoroethanol mixtures (4). Because alanine is known to be the best helix maker, the Ala-13–Gln-16 segment of the peptide should readily assume the helical structure when it is nucleated by the La³⁺-saturated loop Asp-1–Glu-12.

This paper was submitted directly (Track II) to the *Proceedings* office. Abbreviations: NOE, nuclear Overhauser effect; NOESY, NOE spectroscopy; rmsd, root-mean-square deviation; TFE, trifluoroethanol. Data deposition: The atomic coordinates have been deposited in the Protein Data Bank, Biology Department, Brookhaven National Laboratory, Upton, NY 11973 (PDB ID code 1NKF).

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MATERIALS AND METHODS

Peptide Synthesis and Purification. Peptides P1 and P2, synthesized by using fluorenylmethoxycarbonyl chemistry, were generous gifts of P. S. Kim. The peptides were cleaved by using standard procedures, desalted on a G-10 Sephadex column in 5% acetic acid, and purified by reverse-phase HPLC chromatography on a Vydac C18 semipreparative column, by using a 10–30% acetonitrile gradient in 0.1% trifluoroacetic acid. The purified peptides were checked for purity on an analytic Vydac C18 HPLC column to monitor possible transamidation products, as was described elsewhere (16).

NMR Measurements. P2 peptide was dissolved in H₂O/D₂O (9:1) containing 100 mM NaCl and 20 mM LaCl₃. pH was adjusted to 6.0 (direct pH meter reading) by microliter additions of 0.1 M HCl or NaOH. The final peptide concentration was 4 mM. It was shown (16) that peptide P1 saturated with lanthanide ions dimerizes, forming a native-like two-loop structure stabilized, most probably, by hydrogen bonds between CO and NH groups of isoleucine residues of two peptide molecules. Peptide P2, containing the same calcium-binding loop, should dimerize in the same way. Nevertheless, the dimerization constant is low. According to the data given in ref. 16, its value at 275.2 K is 2.5 mol⁻¹. At 4 mM concentration, the fraction of dimerized peptide molecules equals only 2% and cannot have any noticeable effect on NMR measurements.

¹H NMR measurements were performed on a Bruker AMX-600 (Bruker Analytik, Rheinstetten, Germany) spectrometer at a frequency of 600.13 MHz and a temperature of 275.2 K. One-dimensional ¹H NMR spectra were measured with the jump-and-return method ($\Delta = 116 \mu\text{s}$) (22) to suppress the strong water signal without affecting exchangeable solute protons. Two-dimensional clean total correlation spectroscopy (23) and nuclear Overhauser enhancement spectroscopy (NOESY) (24) spectra, however, were measured with water presaturation during relaxation delay (1.8 s) and mixing time (NOESY). Mixing times in total correlation spectroscopy and NOESY spectra were equal to 80 ms and 200 ms, respectively. Correlation of ¹H and ¹³C chemical shifts was measured on a Bruker AMX-400 spectrometer in ²H₂O solution with heteronuclear multiple quantum coherence sequence (25) at natural ¹³C abundance.

¹H and ¹³C chemical shifts were calibrated relative to the fictitious 2,2-dimethyl-2-silapentane sulfonate at 298 K. Such an approach is required to set a reference point identical to that used for the chemical shift index (26, 27).

Structure Calculations. Nuclear Overhauser effect (NOE) crosspeaks were divided into three categories and assigned to distance ranges according to their intensity: strong, 0.18–0.27

nm; medium, 0.18–0.40 nm; weak, 0.18–0.55 nm (28). Peak intensities were estimated from the number of contours in the NOESY spectrum. Harmonic restraints for the La³⁺ ion were deduced from the position of the corresponding Ca²⁺ ion in the crystal structure of calmodulin (Protein Data Bank code 1cdm). A total of six harmonic distance restraints was included to fix the distance and the octahedral arrangement of the six ligands relative to the La³⁺ ion, assuming the same coordination as for the Ca²⁺ ion in the calmodulin crystal structure.

The structure calculations used the *ab initio* simulated annealing (SA.INP) and refinement (REFINE.INP) protocols from the X-PLOR program package (Yale University, New Haven, CT) (29). The calculations started from an extended template with randomized backbone torsion angles followed by 50 cycles of energy minimization to remove close nonbonded contacts. The high-temperature phase comprised 50 ps of dynamics at 1,000 K; the final 16 ps had an increased weight on covalent geometry restraints and the NOE derived distance restraints. In the next phase, the system was slowly cooled from 1,000 K to 100 K in 30 ps, followed by 200 steps of energy minimization. For the NOE effective energy term representing the interproton distances, a soft square-well potential was applied (30). The refinement protocol consisted of slow cooling from 1,000 K to 100 K within 45 ps. A force constant of 200 kcal mol⁻¹rad⁻¹ was used for the dihedral angle restraints, while the NOE derived distance restraints and harmonic restraints were represented by a square-well potential function with a force constant of 50 kcal mol⁻¹Å⁻².

Of 200 resulting structures, these 30 structures that showed the lowest energy and the least violation of the experimental data were selected for further characterization. Geometry of the structures and elements of secondary structure were analyzed by using PROCHECK (31) and DSSP (32).

CD Measurements. The spectra were recorded on an Aviv DS62 spectropolarimeter (Aviv Associates, Lakewood, NJ) equipped with an HP Model 89100A temperature controller. The wavelength was calibrated with epi-androsterone (*trans*-androsterone) and with 4% holmium oxide in 10% perchloric acid. All measurements were carried out in 0.1 M NaCl/5 mM cacodylic buffer, pH 6.9, at 1°C, 25°C, and 50°C, with peptide concentration of *ca.* 10 μM in a 10-mm path length cell. Samples were prepared by dilution of concentrated stock solutions of the peptides. Peptide concentration in the stock solutions was determined with 2% accuracy by fitting tyrosine absorption within the range 230–280 nm (33) to sample spectra measured on a Cary 3E spectrometer by using our own [trifluoroethanol (TFE)] least-squares fitting procedure.

Samples for trifluoroethanol titration experiments were prepared by mixing 10-μM stock solutions of P2 peptide in a

Table 1. ¹H chemical shifts in P2 peptide

Residue	$\delta(\text{H}_\text{N})/{}^3\text{J}_{\text{HN}\alpha}$	$\delta(\text{H}_\alpha)$	$\delta(\text{H}_\beta)$	Others
Asp-1	8.71/6.0	4.48	2.63, 2.37	N-term CH ₃ 1.92
Lys-2	8.39/7.5	4.41	1.98, 1.86	H _γ 1.58; H _δ 1.68; H _ε 3.01; H _ζ 7.48
Asp-3	8.07/9.1	4.75	3.13, 2.61	
Gly-4	7.99	3.89, 3.79		
Asp-5	7.81/8.8	4.70	3.22, 2.67	
Gly-6	9.04	4.00, 3.52		
Tyr-7	8.50/8.4	4.88	2.95, 2.70	H _δ 7.03; H _ε 6.80
Ile-8	9.53/9.7	4.47	1.83	γCH ₂ 1.57; γCH ₃ 0.86; H _δ 0.86
Ser-9	9.78/6.9	4.51	4.44, 4.10	
Ala-10	9.09/3.0	4.19	1.49	
Ala-11	8.53/4.6	4.20	1.42	
Glu-12	8.04/5.8	4.02	2.40, 2.13	H _γ 2.57, 2.31
Ala-13	8.79/4.4	4.11	1.45	
Ala-14	7.89/5.1	4.20	1.49	
Ala-15	7.58/5.8	4.28	1.49	
Gln-16	7.75/7.1	4.21	2.19, 2.13	H _γ 2.50, 2.42; H _ε 7.36, 7.33
NH ₂	7.66, 6.90			

Table 2. Chemical shifts of protonated carbons in P2 peptide

Residue	$\delta(C_\alpha)$	$\delta(C_\beta)$	Others
Asp-1	53.95	40.67	N-term CH ₃ 23.93
Lys-2	56.37	33.43	γ CH ₂ 24.47; δ CH ₂ 28.40; ϵ CH ₂ 41.62
Asp-3	52.12	39.57	
Gly-4	47.33		
Asp-5	53.40	40.51	
Gly-6	44.95		
Tyr-7	56.53	41.59	δ CH 133.65; ϵ CH 117.69
Ile-8	60.82	37.59	γ CH ₂ 28.08; γ CH ₃ 17.75; δ CH ₃ 13.09
Ser-9	57.04	65.30	
Ala-10	54.77	*	
Ala-11	55.54	18.05	
Glu-12	58.37	28.76	γ CH ₂ 38.15
Ala-13	54.37	17.71	
Ala-14	53.73	*	
Ala-15	53.07	*	
Gln-16	56.27	29.41	γ CH ₂ 33.85

*Ambiguity because of the degeneracy of ¹H chemical shifts, $\delta(C_\beta)$: 17.37, 17.91, 18.47.

buffer containing 0.1 M NaCl/5 mM sodium cacodylate, pH 6.9/65% vol/vol mixture of TFE, and the buffer.

The scans were recorded with 0.2-nm step, 1 s averaging time, except for the measurements of P1(apo) and P2(apo) solutions, when the scans were recorded with 2-nm step, 10 s averaging time.

RESULTS

NMR Experiments. One-dimensional ¹H NMR spectra of P2 showed good dispersion for the amide proton resonances stretched from 6.9 to 9.8 ppm, so that the sequence-specific resonance assignments could be performed with a standard approach based on two-dimensional total correlation spectroscopy and NOESY spectra (34). The complete ¹H resonance assignments are given in Table 1. These assignments allowed us to obtain almost all ¹³C chemical shifts for proton-bearing carbons from the two-dimensional ¹H/¹³C heteronuclear multiple

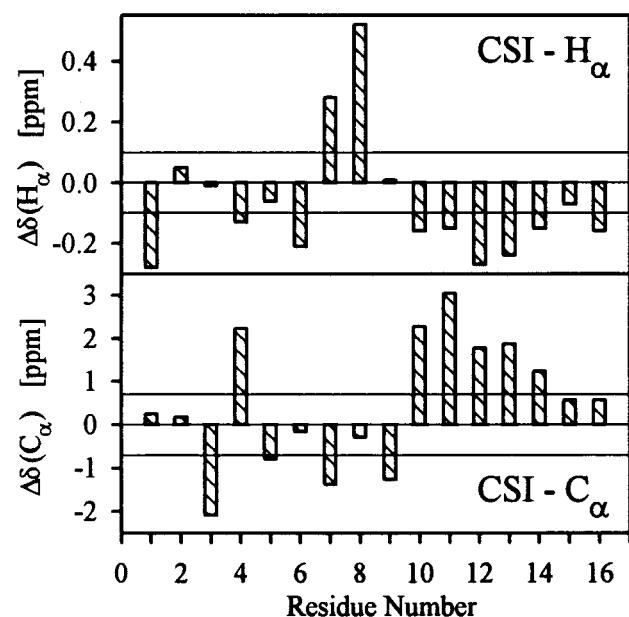


FIG. 2. Chemical shift index of H_α and C_α nuclei in P2 peptide. Chemical shift differences $\Delta\delta = \delta_i - \delta_{i,\text{ref}}$ and appropriate limits determining chemical shift indices are displayed rather than their discrete values -1, 0, 1. Stretches of negative $\Delta\delta(H_\alpha)$ and positive $\Delta\delta(C_\alpha)$ values are characteristic of the α -helical structure.

quantum coherence spectrum of P2 (Table 2). Only β -carbons of Ala-10, Ala-14, and Ala-15 showed ambiguity because of degeneracy of the corresponding ¹H chemical shifts (Table 1).

H_α and C_α chemical shift data available from our experiments were used to perform a secondary structure estimate according to the chemical shift index strategy (26, 27). Accuracy of this essentially statistical method increases with the number of types of involved nuclei. The chemical shift index method relies on calculation of chemical shift deviations for a given type of nuclei in all amino acid residues, if random coil chemical shifts are used as references. Consecutive deviations large enough to be meaningful indicate either helical structures or β -strands, depending on their sign. Negative deviations of H_α chemical shifts and positive deviations of C_α shifts correspond to local α -helical structures, whereas the opposite is true for local β -sheet structures. For P2 chemical shift indices of H_α nuclei indicate an α -helical region extending from Ala-10 to Gln-16 (Fig. 2). Although chemical shift indices of C_α nuclei for two N-terminal residues are equal to zero, suggesting that the α -helical region stretches from Ala-10 to Ala-14, their chemical shift differences are equal to +0.57 ppm and, therefore, are close to the "quantization" limit of +0.7 ppm.

This finding is supported by the values of vicinal ³J_{H_Nα} coupling constants. Regular secondary structures are characterized by the following values for ³J_{H_Nα}: helices, J ≈ 4 Hz and β -sheets, J ≈ 9 Hz. For helices, the uniqueness of structure grows with the number of subsequent amino acid residues showing J < 6 Hz (87% for five subsequent residues) (34). Residues 10 to 15 are the only ones exhibiting coupling constants smaller than 6 Hz (Table 1). The value of 7.1 Hz observed for Gln-16 reflects some end effects, either partial unwinding of the helix or distortion of its geometry.

The calculation of the final structures was based on 83 NOE distance restraints (including 14 medium-range and 10 long-range NOEs), 14 φ angle restraints and 6 harmonic restraints

Table 3. Structural statistics for the family of 30 P2 structures

Measurement	Value
rmsd from experimental distance and torsion angle restraints*	
NOE restraints (83), Å	0.022
La ³⁺ ion restraints (6), Å	0.022
φ angles (14), °	0.014
rmsd from idealized covalent geometry	
Bond lengths, Å	0.001
Angles, °	0.239
Improper, °	0.206
X-PLOR potential energies (±SD), kcal·mol ⁻¹ †	
E _{total}	11.6 (±0.7)
E _{repel}	2.2 (±0.2)
E _{NOE}	1.5 (±0.4)
E _{bond}	0.5 (±0.0)
E _{angle}	3.3 (±0.4)
E _{improper}	0.7 (±0.1)
E _{dihedral}	0.0 (±0.0)
E _{harm}	3.5 (±0.2)
rmsd of cartesian coordinates‡, Å	
Backbone N, C _α and C (residues 1–16)	0.30
All heavy atoms (residues 1–16)	0.84
Backbone N, C _α and C (residues 8–16)	0.10
All heavy atoms (residues 8–16)	0.25

*The total number of restraints is given in parentheses. None of the 30 structures showed distance violations of more than 0.20 Å or dihedral angle violations of more than 2.0°. No distance or dihedral angle restraints were consistently violated by more than 0.10 Å or 1.0°, respectively.

†E_{tot}, total energy, E_{repel}, repulsive energy term, E_{NOE}, effective NOE energy term. The harmonic energy term (E_{harm}) was used to maintain the geometry of the six La³⁺ ion ligands.

‡rmsd values result from a superposition on the N, C_α, and C atoms of the protein backbone and were calculated by averaging the individual rmsds between the average structure and each member of the family.

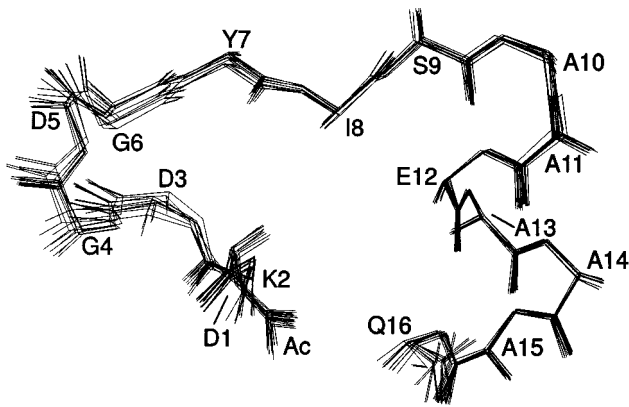


FIG. 3. Overlay of a family of ten P2 peptide structures. The structures were accepted by the criterion of the best energy function after restrained molecular dynamics calculations.

for the ligands of the La^{3+} ion. The coordinates for the ten lowest energy structures have been deposited in the Protein Data Bank (PDB ID code 1NKF).

Analysis of the energies in Table 3 shows that both the NOE and harmonic energy term are close to zero, confirming that the assumption of an octahedral ligation of the La^{3+} ion (identical to that observed for the Ca^{2+} ion in calmodulin) is in total agreement with the experimental data.

The structure is overall well defined, showing root-mean-square deviation (rmsd) values of less than 0.4 \AA for most parts of the peptide chain and particularly low rmsd values for the C-terminal half (Fig. 3; Table 3). All backbone torsion angles of P2 are found in the energetically most favored (85%) or allowed (15%) regions of the Ramachandran plot.

The conformation of the N-terminal part of P2 closely resembles that of the third calcium binding loop of calmodulin, showing a pairwise backbone rmsd of $0.68 (\pm 0.07) \text{ \AA}$ between residues 93–104 of calmodulin and residues 1–12 of P2.

As already expected from the presence of helix-typical $\{i\}\text{C}_\alpha\text{H}-\{i+3\}\text{NH}$, $\{i\}\text{C}_\alpha\text{H}-\{i+3\}\text{C}_\beta\text{H}$, and $\{i\}\text{C}_\alpha\text{H}-\{i+4\}\text{NH}$ NOEs and small $^3J_{\text{HN}\alpha}$ coupling constants, the calculated structures show a helical secondary structure from Ser-9 to Gln-16 with a regular $i, i+4$ hydrogen bonding pattern. A greater structural variability is observed only for the very C-terminal part of this helix, as evidenced by the presence of alternative $i, i+3$ hydrogen bonds for the amide protons of residues Ala-15 and Glu-16 in approximately 30% of the calculated structures. This finding, which indicates a higher conformational flexibility, is in good agreement with the $7.1 \text{ Hz } ^3J_{\text{HN}\alpha}$ coupling constant observed for Gln-16.

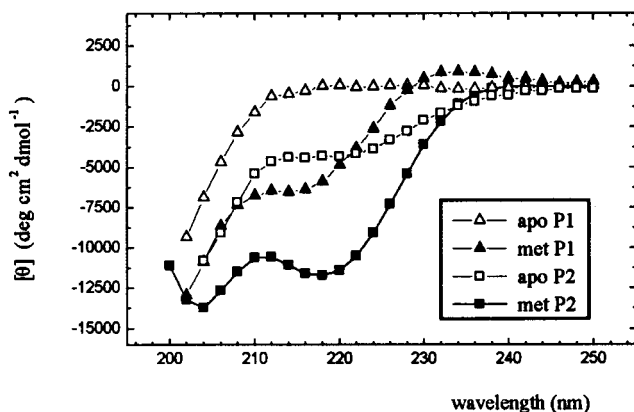


FIG. 4. The CD spectra of P1 and P2 peptides at 1°C in 5 mM cacodylic buffer/0.1 M NaCl, pH 6.9. In the met state, the peptides were saturated by a 50-fold excess of lanthanum ions.

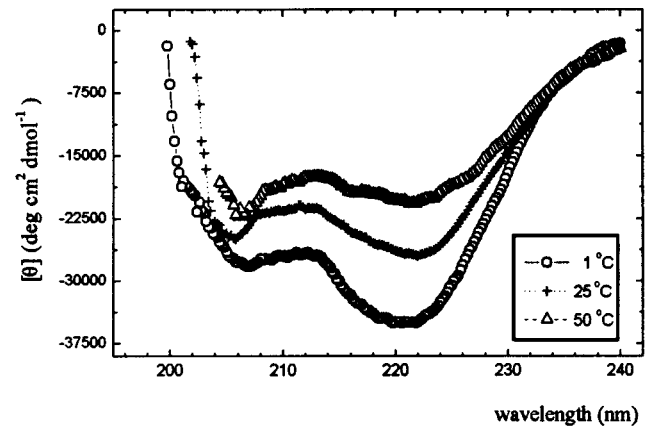


FIG. 5. The temperature dependence of CD difference spectra of P2 minus P1 peptides in 5 mM cacodylic buffer, 0.1 M NaCl/pH 6.9.

CD Measurements. The CD spectrum of free P1 peptide is characteristic of the coil conformation (Fig. 4). That La^{3+} -saturated peptide looks like a mixture of the coil component and the spectrum with a clearly visible minimum at 214 nm. Some small contribution from the α -helical conformation is visible in the spectrum of free P2 peptide. This contribution increases dramatically on La^{3+} ion binding.

As we have already reported (16), the spectrum of La^{3+} -saturated P1 peptide is temperature independent and, within error, remains the same with a temperature increase from 1°C to 80°C , indicating that the loop structure is very rigid. Therefore, the difference spectra of metal-saturated P2 and the loop alone (P1) can be expected (see *Discussion*) to provide direct specific information about the conformation of the P2 Ala-13—Gln-16 segment alone. Such difference spectra, measured at 1, 25, and 50°C , are shown in Fig. 5. Their general shape, with two minima at 206 and 222 nm, is characteristic for the pure α -helical structure.

Within the temperature range for which the measurements have been performed (1 – 50°C), the CD spectrum of La^{3+} -saturated P1 peptide does not change noticeably with TFE concentration. The P2–P1 difference values of mean residue ellipticity θ_{222} remain constant, within the error, at 1 and 25°C , and increase slightly with TFE concentration at 50°C (Fig. 6).

DISCUSSION

Calculation of the Helix Content. The α -helical conformation of residues Ala-13—Gln-16 is not stabilized or by any specific interactions between the residue side chains or charge–helix dipole interactions. Only the C-cap effects (35) should be intro-

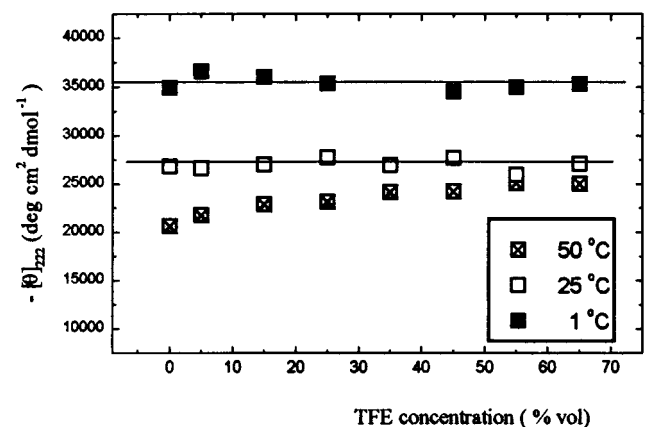


FIG. 6. CD difference signal P2 minus P1 peptides at 222 nm as the function of TFE concentration.

duced into the original Zimm–Bragg equation (Eq. 1) to calculate the helix content. Because the helix is fully nucleated, $\sigma = 1$ and the statistical weight of an n -residue-long helix is given by the following expression:

$$k_n = c_{13+n} \prod_{i=13}^{12+n} s_i, \quad [2]$$

where n changes from 1 to 4, s_i denotes the propagation parameters of residues number i , and c_{13+n} is the C-cap parameter of consecutive cap residues, including the terminal NH_2 group.

Several sets of s parameters have been determined by using various model peptide systems (2). Because the amino acid composition of the peptide fragment studied by us is very close to that of the systems used by Baldwin and his collaborators (2), we will use the s and c values reported by them here (4, 35). The calculated helix contents in water and 40% TFE are shown in Table 4.

Estimation of the Helix Content from CD Measurements. Determination of the helix content in very short peptides from CD measurements is not a trivial task. Long ago the chain-length dependence of CD spectra of the α -helical structure was predicted theoretically (36) and confirmed experimentally by analysis of protein CD spectra (37). The following empirical expression was proposed to describe the molar ellipticity, at any wavelength, of an α -helix built of r peptide units:

$$\Theta^{rh} = \theta^\infty(r - k), \quad [3]$$

where θ^∞ is the molar peptide unit ellipticity of infinite α -helix and k a wavelength-dependent constant. Or, in terms of molar peptide unit ellipticities (θ^{rh}):

$$\theta^{rh} = \theta^\infty(1 - k/r). \quad [4]$$

Both types of values, Θ and θ , are given in the same units: $\text{deg}\cdot\text{cm}^2\cdot\text{dmol}^{-1}$. The only difference is that mol refers in one case to peptide concentration and in the other case to concentration of peptide units. Therefore, $\theta^{rh} = \Theta^{rh}/r$ (compare Eqs. 3 and 4).

Numerous theoretical and experimental studies have been done to determine the precise values of θ^∞ and k , particularly at 222 nm because the signal at this wavelength is used as a measure of helix fraction in peptides. The latest data obtained from studies of a large number of peptides in water and water–TFE mixtures have been published by Luo and Baldwin (38). According to them, $k = 4$ (their x parameter, equal to 3, refers to the number of residues in helical conformation and has to be increased by one, if the peptide groups are counted) and

$$\theta_{222}^\infty = -44,000 + 250t, \quad [5]$$

where t is the temperature in $^\circ\text{C}$.

The CD signal at 222 nm of La^{3+} -saturated peptide P1 (Θ_{222}^1) arises from the helical structure of the four C-terminal peptide groups beginning with that formed between Ser-9 and Ala-10 (Θ_{222}^{4h}) and from the rest of the loop (Θ_{222}^l):

$$\Theta_{222}^1 = \Theta_{222}^{4h} + \Theta_{222}^l. \quad [6]$$

The contribution of the helical structure is small, if any, as evidenced by the spectrum shown in Fig. 4. At 230 nm, θ is close to zero, whereas in the CD spectrum of the α -helix it attains approximately half of the θ_{222} value. Therefore, the CD signal from short helices is indeed very small and $k \approx 4$. In P2 peptide, the helix is longer and contains eight peptide groups.

$$\Theta_{222}^2 = \Theta_{222}^{8h} + \Theta_{222}^l \quad [7]$$

Therefore, the difference signal

$$\Delta\Theta_{222}^h = \Theta_{222}^2 - \Theta_{222}^1 = \theta_{222}^\infty(8 - k) - \theta_{222}^\infty(4 - k) = 4\theta_{222}^\infty \quad [8]$$

and the helix fraction f_H in the Ala-13–Gln-16 segment should be calculated (38) from the equation:

$$f_H = (\theta_{222}^{bs} - \theta_{222}^c)/(\theta_{222}^h - \theta_{222}^c), \quad [9]$$

where θ_{222}^{bs} and θ_{222}^c are observed and coil mean peptide group ellipticities, respectively, and $\theta_{222}^h = \theta_{222}^\infty$. Nevertheless, if $k > 4$, as reported by Gans *et al.* (39), $\Theta_{222}^1 = 0$ and from Eq. 8 we obtain:

$$\Delta\Theta_{222}^h = \theta_{222}^\infty(8 - k). \quad [10]$$

In this case $\theta_{222}^h = \theta_{222}^\infty(2 - k/4)$.

Assuming that $k \leq 4$, we can estimate the lowest limit of the helix fractions. They are shown in Table 4. In water, they are considerably higher than those calculated from Eq. 2, particularly at low temperature. The s parameters given by Rohl *et al.* (4) have to be multiplied by the factors of 2.1, 1.4, and 1.2 to obtain the helix contents determined at 1, 25, and 50 $^\circ\text{C}$, respectively. At 1 $^\circ\text{C}$, it corresponds to additional favorable contribution $\Delta\Delta G_s$ of -400 cal/mol to ΔG_s of each step of the helix propagation, and the total contribution to the Ala-13–Gln-16 helix stability can be evaluated at -1.6 kcal/mol.

The high helix content at 1 and 25 $^\circ\text{C}$ is additionally confirmed by the fact that it does not increase with TFE concentration. The helix fractions determined in 40% TFE, assuming that $k \leq 4$, are somewhat smaller than the calculated ones, which seems unreasonable. Even if $\Delta\Delta G_s$ tends to zero in concentrated TFE solutions, it cannot be expected to exceed this value. In other words, there is no reason to believe that in our system TFE can really destabilize the helix. The helix fractions determined from the experiment are equal to, or greater than, the calculated ones for $k = 4.5$ (compare columns 3 and 7 in Table 4). In this case, though, the helix content at 1 $^\circ\text{C}$ may exceed 90% and the helix would be additionally stabilized by $\Delta\Delta G_s$ of as much as -1 kcal/mol. Further studies of peptides with the calcium-binding loop and helical segments of various length are needed to definitely elucidate the problem of CD spectra of very short helices and make it possible to determine more precisely the helix content in these peptides.

Unexpectedly High Stability of Prenucleated Helices. Even with the lowest estimates of the helix content, it is evident that the helical conformation of the Ala-13–Gln-16 segment of P2 peptide is considerably more stable than predicted by the

Table 4. Helix fractions in A-13–Gln-16 segment of P2 peptide saturated with La^{3+} ions

Temp., $^\circ\text{C}$	Calculated*		From CD, $k \leq 4^\dagger$		From CD, $k = 4.5^\ddagger$	
	H_2O	40% TFE	H_2O	40% TFE	H_2O	40% TFE
1	0.68	0.82	0.81	0.79	0.93	0.92
25	0.64	0.80	0.72	0.69	0.82	0.80
50	0.60	0.78	0.65	0.73	0.75	0.85

*Calculated from Eq. 2. s values at 273 K taken from Rohl *et al.* (4) and calculated for 1, 25, and 50 $^\circ\text{C}$ using ΔG values reported therein and ΔH equal -1 kcal/mol and -800 cal/mol in water (41) and in concentrated TFE solutions (3), respectively.

† Calculated from Eqs. 9 and 5 with θ_{222}^∞ values taken from Luo and Baldwin (38).

helix-coil transition theory and helix propagation parameters determined from studies of long polypeptide chains. This phenomenon cannot be accounted for by any specific interactions between the loop and the helix. The NMR structure of the peptide indicates that there are no contacts between the loop residues and the helix. Electrostatic interactions of La^{3+} ion and its ligands with the dipoles of C-terminal peptide groups beginning from Glu-12—Ala-13 are negligible. Calculations made by using the DELPHI program (40) have shown that their energy does not exceed -200 cal/mol (K. Pawłowski, unpublished results). Therefore, we have to conclude that what we observe is an intrinsic property of a nascent helix. Namely, that the s parameters describing the very first steps of helix propagation are greater than those characteristic for long helices. Apparently this effect is enthalpy driven, because $-\Delta\Delta G_S$ drops with temperature increase.

What can the physical nature of this phenomenon be? We believe that it is related to hydration of the helix ends. The results of TFE titration experiments suggesting that $-\Delta\Delta G_S$ drops to zero in concentrated TFE solutions point in this direction. The hypothesis explaining $-\Delta\Delta G_S$ by helix end hydration effects has been presented by us in a previous paper (3). In short, the nonhydrogen bonded CO and NH groups at C- and N-termini of the α -helix, respectively, are perhaps better solvated when they are separated from each other in long helices than when they merge with each other in the helix nucleus. Therefore, the helix nucleation is related to an additional increase of hydration energy compensated during the first steps of the helix propagation that lead to separation of the helix ends. Unfortunately, we failed to verify this hypothesis by molecular dynamics calculations of helix end hydration energies. Even after long runs of simulations, the statistical errors of the calculated values were too large to allow for any conclusions. We hope that further studies will make it possible to confirm or dismiss this hypothesis.

Conclusions. Calcium-binding loops from EF-hand proteins saturated with lanthanide ions nucleate extremely effectively the α -helical structure in peptide segments attached to their C termini. Short peptides containing such nucleation sites provide unique model systems to study the very first steps of helix propagation and to investigate in detail various specific interactions stabilizing or destabilizing the helical conformation, such as hydration of the helix ends or side chain–side chain interactions within the helix.

An important advantage of these systems lies in that they are built of natural amino acid residues. Therefore, they can be produced by using standard methods of molecular engineering. The material can be labeled with ^{15}N and ^{13}C nuclei to make it suitable for sophisticated NMR experiments. The calcium-binding loops attached to N-termini of protein fragments or proteins with N-terminal α -helices, e.g., globins, can be used to significantly increase their structural stability in the presence of lanthanide ions.

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