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Light Scattering and Ultrastructure of Healed Penetrating Corneal Wounds

Russell L. McCally^{1,2}, David E. Freund¹, Andrew Zorn³, Jennifer Bonney-Ray¹, Rhonda Grebe², Zenaida de la Cruz², and W. Richard Green²

1 The Johns Hopkins University Applied Physics Laboratory Laurel, MD 20723

2 The Wilmer Ophthalmological Institute The Johns Hopkins Medical Institutions Baltimore, MD 21205

3 The Johns Hopkins University Whiting School of Engineering Baltimore, MD 21218

Abstract

Purpose— To investigate quantitatively for the first time the relationship between light scattering and ultrastructure of semi-transparent scars resulting from penetrating wounds in rabbit cornea.

Methods— 2 mm diameter penetrating wounds were made in the central cornea and allowed to heal for 3.6 to 4.5 years at which time the rabbits were sacrificed. The scar and cornea thickness outside the scar were measured using ultrasonic pachymetry. Corneas were excised immediately and their transmissivity was measured from 400 to 700 nm. The tissue was then prepared for TEM. TEM were analyzed to determine fibril positions and radii. Scattering was calculated using the direct summation of fields (DSF) method.

Results— Scar thickness averaged 0.26 ± 0.04 mm and the scars were flat. Thickness outside the scars averaged 0.40 ± 0.04 mm. Three scars were moderately transparent, five were less transparent, and one was much less transparent. The wavelength dependence of the measured total scattering cross-section was indicative of the presence of voids (lakes) in the collagen fibril distribution and lakes were evident in the TEM. The TEM showed enlarged fibrils and some showed bimodal distributions of fibril diameters. Calculated scattering was characteristic of that expected from regions containing lakes - a finding consistent with the scattering measurements.

Conclusions— Despite the long healing time these scars remained highly scattering. A combination of lakes, disordered fibril distributions, and a significant population of enlarged fibrils can explain the scattering. A possible cellular contribution cannot be ruled out.

INTRODUCTION

Charles Cintron and his co-workers adapted a model originally developed by Heydenreich¹ to investigate healing of penetrating corneal wounds.^{2–4} This highly successful model consisted of removing a centrally located 2 mm diameter full thickness button. Such wounds heal by first forming a fibrin clot, after which cells invade and ultimately produce an avascular network of collagen fibrils.^{2–5} The model had advantages of simplicity, reproducibility of healing with minimal complications, and the production of scar tissue suitable for analysis.⁴ Interestingly,

Corresponding author: Russell L. McCally, The Johns Hopkins University Applied Physics Laboratory, 11000 Johns Hopkins Road, Laurel, MD 20723-6099. russell.mccally@jhuapl.edu

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it was reported that the initially opaque scars ultimately became "transparent." ², ³ This and other biochemical observations led to the initial suggestion that rabbit cornea exhibited a capability for true regeneration of stromal tissue.³ However more extensive investigations showed that the regeneration was incomplete at the biochemical and ultrastructural level² and the statement that the wounds became transparent was qualified to their becoming less opaque and sometimes transparent.⁵

Two primary physical factors lead to the transparency of normal corneas (a third factor, namely the thinness of the normal cornea, is less applicable to the scars which are the subject of the present investigation.).⁶ First, collagen fibrils in the stroma are weak scatterers because their radius is much smaller than the wavelength of visible light and their refractive index is close to that of their surroundings. Second, destructive interference among the fields scattered by different fibrils reduces the scattering from that which would occur if the fibrils scattered independently of one another. A quantity called the total scattering cross-section incorporates both of these effects.^{7–9} It can be determined (within a multiplicative constant) from measurements of transmissivity.⁶, 10, 11 The total scattering cross-section is proportional to a quantity known as the interference factor, which shows how much the scattering is reduced. It depends on the spatial distribution of the collagen fibrils about one another and, to a lesser extent, on the spatial distribution of fibril scattering strengths.⁹ The interference factor is central to most transparency theories.⁷, 12-16 It can be calculated from the fibrillar structures shown in transmission electron micrographs (TEM) and it provides a quantitative measure of the degree (or lack thereof) of fibrillar ordering.⁹ Alterations in the spatial distribution of fibrils that reduce destructive interference, or changes that cause the fibrils to scatter more efficiently, reduce transparency and, if severe, lead to corneal opacity. In normal cornea, keratocytes do not contribute significantly to light scattering except under the special condition of specular scattering, which occurs in reflective confocal microscope images.^{6, 11, 17, 18} Because at present there is no quantitative theory to describe scattering from corneal cells¹⁸ it is not possible to separate their possible contribution to reduced transmissivity in corneal wounds from that of the disrupted fibrillar matrix.

This paper, for the first time, relates experimental measurements of light scattering from healed penetrating wounds to the ultrastructure depicted in TEM of the wounds via calculations of the total scattering cross-section and the interference factor.

MATERIALS AND METHODS

i. Animals

In conducting the experiments we adhered to the ARVO Resolution on the Use of Animals in Research. Two millimeter diameter penetrating wounds in the central corneas of 5 New Zealand white rabbits were produced at different times in Dr. Charles Cintron's laboratory at the Schepens Eye Research Institute as described previously $^{2-4}$ The rabbits were ultimately delivered to the Wilmer Eye Institute where they were housed until the time of the light scattering experiments (3.5–4.5 years after wounding).

Prior to being killed for light scattering measurements, rabbits were anesthetized with an intramuscular injection of xylazine and ketamine hydrochloride in the proportions: 60% of 20 mg/ml xylazine to 40% of 100 mg/ml ketamine by volume. They were killed by an overdose of Buthanasia administered in an ear vein and the eyes enucleated. The scar thickness and the thickness of the adjacent unwounded cornea were measured with an ultrasonic pachymeter and the scars were photographed. Corneas were then excised with a scleral ring and mounted as described previously.¹⁰ Corneas were bathed in BSS during the entire excision and mounting procedure. After mounting, the hydrostatic pressure across the cornea was raised to 18 mmHg and the mounted cornea was inserted in the scattering apparatus where it also was bathed in

BSS. Transmission at 500 nm was monitored for 10 minutes before taking measurements to ensure the stability of the preparation. 10

ii. Transmission measurements

Transmission measurements were made in an apparatus similar to that described in Farrell et al.¹⁰ except for several improvements (primarily in the electronics). The improvements have not altered the optics which ensure that the incident and transmitted light beams have very low angular divergences and that the collection optics minimize the collection of forward scattered light.

The transmissivity, F_T , at a wavelength λ is given by

$$F_T(\lambda) = \frac{I_t}{I_0} = \exp\left(-\rho\sigma_t(\lambda)t\right),\tag{1}$$

where I_t is the irradiance of the light transmitted through the cornea, I_0 is the irradiance of the light incident on the cornea, t is the thickness of the cornea, ρ is the number density of fibrils in a typical lamella, and $\sigma_t(\lambda)$ is the total scattering cross-section. In order to account for cornea–to–cornea variations in scar thicknesses, we average the quantity $-\ln (F_T(\lambda)/t = \rho \sigma_t(\lambda))$. 6, 10, 11 In Results it is shown that the values of $\rho \sigma_t(\lambda)$ vary from scar to scar, but that they can be broken into three groups to facilitate discussion.

iii. Electron microscopy

Immediately following the scattering measurements, the mounted corneas were removed from the apparatus, and while still under pressure, phosphate buffered 1% glutaraldehyde 4% formaldehyde fixative was dripped on the anterior surface for seven minutes. At this point the cornea was sufficiently rigid to retain its shape when it was removed from the mount. Fixation was continued overnight followed by post-fixation in 1% phosphate buffered osmium tetroxide for two hours. Corneas were then dehydrated in a graded series of ethanol, stained with uranyl acetate in 100% ethanol, passed through two changes of propylene oxide and infiltrated in a 1:1 propylene oxide and resin mixture [modified Luft's medium (based on LX 112; Ladd Research Industries, Burlington, VT and Dodecenylsuccinic anhydride (DDSA)] overnight. The infiltration process was continued by immersing the samples in 100% resin mixture under vacuum for two hours. Samples were then polymerized overnight at 37°C and then at 58–66° C for two days. Ultra-thin sections were were doubly stained with tannic acid-uranyl acetate and lead citrate solutions. Tannic acid is an improved stain for elastin and collagen fibers.¹⁹ Samples were examined and photographed with a JEM 100B Transmission Electron Microscope.

iv. Image analysis and scattering calculations

Micrographs (magnification 64,800X) were scanned at 300 dpi on a UMAX PowerLook 1000 scanner. This resolution ensured that the smallest fibrils would be at least 8 pixels in diameter. ²⁰ Image analysis was performed on a Macintosh computer using an augmented version of the public domain NIH Image program (developed at the U.S. National Institutes of Health and available on the Internet at http://rsb.info.nih.gov/nih-image/).²⁰ The general outline of the image-processing algorithm that was used to obtain the fibril positions and their radii was described previously.²⁰

Light scattering was calculated from fibril positions and radii using the direct summation-offields (DSF) method.²¹ This method enables one to estimate the ensemble average of the total scattering cross-section and the interference factor using a single TEM and it can be applied to assemblies of fibrils that have arbitrary spatial distributions and arbitrary distributions of

diameters, such as those encountered in scarred corneas. The calculations herein account for the individual fibril diameters as described by Freund et al. 9

Estimates for the refractive indices of the fibrils and extrafibrillar matrix were obtained using the Gladstone-Dale law of mixtures and the refractive increment law as described previously. 8, 9, 14, 22 Calculations of refractive indices require assumptions of collagen concentration, how it is distributed in the fibrils, and the concentration of components in the extrafibrillar matrix. These quantities are unknown for these scars, thus for the purposes of this study we have assumed that they are the same as for normal cornea.²³ We also have assumed that the fibrils contain a fixed percentage of collagen; thus all fibrils have the same index of refraction – independent of their diameters.⁹

RESULTS

Despite the long healing time these wounds remained highly scattering. Figure 1 shows examples from the three scattering groups that were identified as described later. It is apparent that scattering in the scars is not uniform. The narrow slit photographs also show that the scarred region is flat and thinner than the adjacent unwounded cornea. The average thicknesses of the scars and adjacent cornea were respectively 0.26 ± 0.04 mm and 0.40 ± 0.04 mm. Others have obtained similar results for scar thickness in healed penetrating wounds.²⁴

Table 1 lists values of $\rho\sigma_t(\lambda)$ obtained from transmission measurements on the nine scarred corneas that were investigated (The cornea from the left eye of rabbit C4 was damaged and could not be used). These results show that the values of $\rho\sigma_t(\lambda)$ from the identically prepared scars are in three distinct groups as delineated in the Table. Group 1 has the lowest scattering. The average value of $\rho\sigma_t(\lambda)$ for Group 2 is 1.74 times greater than the Group 1 average at 400 nm and 1.56 times greater at 700 nm. The Wilcoxon-Mann-Whitney test showed that the differences in average values are significant at all wavelengths (*P*<0.04). The single cornea in Group 3 had the highest values of $\rho\sigma_t(\lambda)$ - ranging from 1.5 times greater than the Group 2 average at 400 nm.

Figure 2 shows the average fraction of light transmitted, $\langle F_T \rangle$, as a function of wavelength for the three groups. The values of $\langle F_T \rangle$ were determined from the averages of $\rho \sigma_t(\lambda)$ in Table 1, assuming that the scars had the average thickness noted above (i.e., t = 0.26 mm). The effect on transmissivity resulting from the differences in scattering in the three groups is obvious.

Figure 3 shows that the quantity $\lambda^3 \sigma_t(\lambda)$ has a linear dependence on λ (i.e., it can be expressed in the form $A + B\lambda$, which is a straight line of slope *B*). This implies that the total scattering cross-section has the functional form

$$\sigma_t(\lambda) = \frac{A}{\lambda^3} + \frac{B}{\lambda^2},\tag{2}$$

where *A* and *B* are constants. The first term in equation 2 has the same wavelength dependence as the scattering cross-section of an isolated fibril and is characteristic of fibrils having either short-ranged order or homogeneous disorder in their positions.⁷, 10, 12, 15, 16, 25, 26 The existence of the second term in equation 2 has been associated with the presence of regions in the stromal lamellae that are devoid of fibrils.¹⁰ Such regions have often been called "lakes." 10, 12 Their effect is to increase scattering.⁹, 10, 12, 25–28 Moreover, the fact that the slopes (i.e., the values of *B*) become progressively greater from Group 1 to Group 3 suggests that void regions become more prevalent in the corneas having greater scattering.

Transmission electron micrographs from corneas in the three groups were analyzed to characterize the fibril distributions and to determine fibril positions and diameters. As

discussed later, these data (from suitable TEM¹) were used to calculate the scattering that would be expected from the depicted structures. Figure 4 shows four of the TEM that were analyzed. Fibril positions and diameters were obtained from the entire depicted regions; the rectangles outline the regions that were subsequently used in the DSF light scattering calculations. Figure 5 shows the distributions of fibril diameters from the TEM in Figure 4. Figure 4a from the mid-stromal region of a cornea in Group 1 shows several small lakes, which confirms the prediction of their existence based on the scattering measurements. Interestingly, the distribution of fibril diameters obtained from this micrograph (cf., Figure 5a) is bimodal; i.e., it has two distinct peaks. To our knowledge bimodal distributions of fibril sizes have not been previously observed in either scarred or normal cornea. Figures 4b and 4c, respectively from the anterior and posterior regions of corneas in Group 2, contain several lakes and 4b has a population of very large fibrils. The diameter distributions from these micrographs are unimodal (c.f., Figures 5b and 5c). Figure 4d is from the anterior stroma of the cornea in Group 3. This region contains large lakes, a broad distribution of fibril diameters 5d), and disorganized fibrils.

Table 2 summarizes the fibril statistics from all of the TEM that were analyzed. It shows that several had bimodal diameter distributions. The average position of the distribution peaks for the micrographs in Group 1 having bimodal distributions are at 325A and 400A, whereas the average position of the peaks for those in Group 2 are at 330A and 470A.

It is important to note the TEM chosen for fibril analysis and from which scattering could be calculated have more fibrillar order than was typical throughout the scars. Many regions in the scars had a much less orderly arrangement of fibrils. Such distributions are not amenable to DSF scattering calculations, but are consistent with the high level of scattering found in the measurements. Although not presented here, lower magnification micrographs from all three groups show regions of disorganized lamellar structures, deposits of granular material, and some partially degenerated keratocytes, some of which contain vacuoles filled with granular material. In general all of these characteristics are more prevalent in the corneas in Groups 2 and 3 and are consistent with the higher scattering observed in these groups.

The DSF method ^{9, 21} was used to calculate the total scattering cross-section for rectangular regions in the micrographs listed in Table 2 which contained ~1000 or more fibrils. The calculated total scattering cross-sections and the fibril number densities obtained from the TEM were used in Equation 1 to determine the fraction of light that would be transmitted through corneas assumed to have the structures depicted in the rectangular regions that were analyzed. The calculations also assumed a corneal thickness of 0.26 mm (i.e., the average thickness of the scars). The average values of the product of the fibril number density and the computed total scattering cross-section, $\langle \rho \sigma_t (\lambda) \rangle$, were determined for the micrographs in each of the three groups and were used to compute $\langle F_T \rangle$ for each of the groups. These results are shown in Figure 6. The Wilcoxon-Mann-Whitney test showed that the differences in $\langle \rho \sigma_t(\lambda) \rangle$ for the micrographs in Groups 1 and 2 are significant at all wavelengths (P<0.035). Because only two micrographs from the cornea in Group 3 were analyzed, no statistical test could be done; however $\langle \rho \sigma_t(\lambda) \rangle$ for the Group 3 cornea is ≥ 5 times the Group 2 average at all wavelengths.

Figure 7 displays the wavelength dependence of the average total scattering cross-section for the three groups. In agreement with the experimental results in Figure 3, the calculated values of $(\lambda/550)^3 \langle \rho \sigma_t(\lambda) \rangle$ for each group also have a linear dependence on wavelength and the slopes become progressively greater from Group 1 to Group 2 to Group 3. This result is consistent with TEM from the three groups.

¹The DSF method, which is used to calculate scattering, requires that there be a rectangular region containing ~1000 or more parallel fibrils (in cross-section) in order to obtain adequate statistics.

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DISCUSSION

Cintron et al^{2, 5} reported that scars from 2 mm penetrating corneal wounds became less opaque and in some cases appeared transparent after healing 1.5 to 2 years. This paper presents the first quantitative measurements of light scattering and their relationship to ultrastructure of such scars. Transmissivity in normal rabbit cornea ranges from 0.88 at 400 nm to 0.97 at 700 nm.¹⁰ The scars from the lowest scattering group had transmissivities ranging from 0.39 at 400 nm to 0.69 at 700 nm (cf., Figure 2), values comparable to those of rabbit corneas swollen to twice their normal thickness.¹⁰ Therefore, while these scars were less opaque than in their early stage of healing (photographs not shown), even those in the lowest scattering group did not approach what could reasonably be described as being transparent. It is important to note however, that measurements of transmission average over the diameter of the incident beam (~1 mm) and through the entire depth of the cornea. Slit-lamp views of the scars (cf., Figure 1) show that the scars were not homogeneous. They had some regions that scattered much less than others. It also is noteworthy that scars from opposite eyes of the same rabbit were in different scattering groups for three of the four pairs of corneas (cf., Table 2). Moreover, some scars in Group 2 had healed for the same time as the scars in Group 1. This suggests that the healing process and partial recovery of transparency is probably influenced primarily by slight (unintended) variations in the initial wounding.

The total scattering cross-section of normal transparent rabbit cornea varies as $(1/\lambda^3)$.⁶, 10, 11, 15, 25, 26 This dependence, which is the same as that of an isolated fibril, is indicative of short-ranged ordering in the fibril positions.⁹, 15, 16, 25 15 However, if fibril positions were to become disordered such that the fibrils scattered independently of one another, scattering would increase because there would be less destructive interference, but the scattering cross-section would still vary as $(1/\lambda^3)$.⁶, 14, 15, 25 The scattering data in Figure 3 are inconsistent with this mechanism for explaining the high level of scattering from the scars; because, if this were the case, the product $(\lambda/550)^3 \langle \rho \sigma_t (\lambda) \rangle$ would have been a constant, independent of wavelength. As discussed previously, the results in Figure 3 show that the measured total scattering cross-sections for the scars have the functional form $A/\lambda^3 + B/\lambda^2$. This form is indicative that the scars contain intralamellar fibril-free regions (vis., "lakes") surrounded by regions of disordered fibrils.⁶, 10, 25, 27 Other investigators have reported that that lakes persist in scar tissue up to 9 months, but they stated that they did not appear to be of sufficient size and number to affect transparency. No scattering measurements or analyses were performed to assess the validity of this assertion.²⁹ Our study shows that lakes persist up to 4.5 years (cf., Figure 4) and the scattering data in Figure 3 suggests that, in fact, lakes are a significant factor producing the high level of scattering in the scars.

The wavelength dependence of the calculated scattering is consistent with the wavelength dependence determined from the measurements (cf., Figures 3 and 7). Both the calculations and measurements show that the quantity $(\lambda/550)^3 \langle \rho \sigma_t(\lambda) \rangle$ depends linearly on λ , and that the magnitude and slopes of the linear fits increase from Group 1 to Group 2 to Group 3. The calculated values of $\sigma_t(\lambda)$ are however, much lower than the measured ones. Consequently, the calculated transmissivities are much greater than the measured values (cf., Figures 2 and 6). This is because the measurements probe the entire depth of tissue, including all of the highly disorganized regions noted previously. Thus the lack of quantitative agreement would be expected because, as noted previously, the DSF method can only be applied to regions that contain a sufficient number (~ 1000 or more) of parallel fibrils.²¹ The DSF method implicitly assumes that the fibrils shown in cross section, which are used in the calculations, are arranged parallel to one another over a reasonable distance (at least several wavelengths). Light waves scattered from parallel fibrils can interfere, whereas waves scattered from non-parallel or tangled fibrils cannot. As noted in the Introduction, destructive interference among the waves scattered from the parallel fibrils in the stromal lamellae is a major factor underlying corneal

transparency. Non-parallel or tangled fibrils scatter independently of one another and greatly increase the overall scattering. We are unaware of any theory for calculating scattering from tangled fibrils.

The total scattering cross-section $\sigma_t(\lambda)$ can be written in the form

$$\sigma_t(\lambda) = \left\langle \sigma_{0,t}(\lambda) \right\rangle \sigma_{tN}(\lambda) \tag{3}$$

where $\langle \sigma_{0t} (\lambda) \rangle$ is the total scattering cross-section of an average isolated collagen fibril and $\sigma_{tN}(\lambda)$ is the interference factor which measures the degree of order in the spatial arrangement of fibril axes and fibril sizes.⁹, ²¹ The magnitude of $\sigma_{tN}(\lambda)$ provides a *quantitative* measure of the degree of fibrillar ordering.⁶, ⁹, ³⁰ As employed herein it also accounts for variability in scattering strengths that result from differences in fibril sizes.⁹ The interference factor can have values between zero (corresponding to crystalline order) and one (corresponding to random order).³⁰ In normal rabbit cornea $\sigma_{tN}(\lambda)$ is ~0.11 in the anterior stroma and ~0.085 in the posterior stroma, and it is essentially independent of wavelength in both regions.⁹ These values are indicative of a substantial degree of fibrillar ordering in the normal transparent cornea. In contrast, the average values of σ_{tN} (550) obtained from the DSF calculations are respectively 0.18±0.13, 0.38±0.22, and 0.80±0.33 for Groups 1, 2 and 3. Thus all groups have less fibular order than normal cornea. The progressive increases in the average values of σ_{tN} (550) from Groups 1 to 3 indicate that disorder in the spatial distribution of fibrils, and the resulting decrease in destructive interference in the scattered fields, also is an important factor in the increased scattering.

The *mean* fibril diameters in Table 2 are in close agreement with the larger of the reported values for normal rabbit cornea⁹, ³¹, ³² (however the 476 A diameter for C25RA15675 is an exception in that it is larger). The mean diameters are also in reasonable agreement with earlier measurements obtained from TEM of 1.5 year old scars² and with x-ray diffraction measurements of 21 month old scars.²⁹ In general, however, the diameter distributions in the scars are much broader than those for normal rabbit cornea as evidenced both by their standard deviations and by their ranges. The standard deviations in Table 2 range from 27A (which was for a cornea in Group 1) to 115A, whereas those for normal cornea range from 20A to 28A.⁹ Rawe et al reported a standard deviation of 70A based on their x-ray diffraction measurements of a 21 month old scars. (viz., 100 to 500 A),² perhaps because we analyzed more micrographs and larger regions. Table 2 shows that the majority of the TEM we analyzed contained populations of fibrils with diameters greater than 500A.

Bimodal distributions of fibril diameters have not been previously reported for scars of this type. Fibrils having diameters near the low peak of the bimodal distributions found in this study have diameters within the range reported for normal cornea. However the diameters of fibrils near the high diameter peak for the Group 2 corneas are greater than any that have been found in normal cornea. Moreover the number of fibrils near the high peak in both Groups 1 and 2 is greater than those near the low peak for every cornea except one. Age is unlikely to be a factor in explaining either the broad diameter distributions or the bimodal diameter distributions. Kanai and Kaufman found that collagen fibril diameters in aged (80 year) human corneas are essentially the same as in younger corneas;³³ whereas Daxer et al found fibril diameters of 308±10 A and 322±10A for human corneas that were respectively less than and greater than 65 years of age.³⁴ Both of these values are within the range of those reported for normal human cornea and both have small standard deviations.

X-ray data suggest that the fibrils in older scars produced in an identical manner as this work have hydration levels close to normal. This led Rawe et al to conclude that the larger fibrils

either have a different molecular arrangement or a larger number of collagen molecules per fibril.²⁹ Either of these conclusions would justify our assumption that the fibrils have a fixed percentage of collagen, and therefore that their refractive indices would be independent of their diameters. Because the single fibril scattering cross-section, $\langle \sigma_{0t} (\lambda) \rangle$, is proportional to the fourth power of fibril diameter,⁹, ²¹ the population of large fibrils would be expected to contribute significantly to the scattering.

The effects of the significant populations of large fibrils in both the unimodal and bimodal distributions are manifested in the isolated fibril scattering cross-sections. The mean values of $\langle \sigma_{0t} (550) \rangle$ calculated from the Group 1 and 2 TEM are respectively (1.63±0.59) X 10² A and (2.59±0.53) X 10² A. The Wilcoxon-Mann-Whitney test showed that the differences in the averages are significant (*P*=0.02). The mean value $\langle of \sigma_{0t} (550) \rangle$ calculated from the two TEM from the Group 3 cornea is 8.6 X 10² A. Although no statistical test was performed, it is obviously much greater.

Although it is possible that keratocytes may contribute to the increased scattering, their potential contribution is difficult to assess because at present there is no theory that can quantitatively predict cellular scattering in the cornea.¹⁸ The TEM that were examined did contain some degenerating keratocytes. It is possible that their refractive indices may differ from normal keratocytes,^{35, 36} but even the refractive indices of normal keratocytes are unknown.¹⁸ The way that non-specular angular scattering from normal cornea scales with scattering angle and wavelength indicates that scattering is due to the fibrillar matrix and keratocytes do not contribute significantly.¹⁷ In any event it is unlikely that scattering from keratocytes or other corneal cells such as myofibroblasts, has the same wavelength dependence as the scattering from these scars.³⁷ It is clear however that a comprehensive theory for cellular scattering is needed and that it would contribute not only to a better understanding of light scattering from wounded cornea, but also to scattering from normal cornea.¹⁸, 38

In summary, scars from 2 mm diameter penetrating wounds remained highly scattering even after 3.6–4.5 years of healing. Analysis of light scattering measurements suggested that "lakes" or voids in the collagen fibril distribution were a significant factor contributing to the high level of scattering. TEM from the scars confirmed the presence of the lakes and some also showed bimodal distributions of fibril diameters. Analysis of the TEM using the DSF method showed that, in addition to scattering from lakes, disorder in fibrillar organization and the effects of significant populations of greatly enlarged fibrils on the isolated fibril scattering cross-sections are other significant factors contributing to the increased scattering.

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Figure 1.

Slit-lamp photographs of scars show that there are substantial variations in scattering intensity across the wound with some regions appearing less opaque than others. The photographs with narrow slit illumination show that the scarred region is flat and thinner than the unwounded cornea. There are variations of scattering intensity with depth, with some regions appearing relatively transparent. (a) Cornea C8L from Group 1 (see Table 1 and Figure 2). (b) Cornea C8R from Group 2. (c) Cornea C25L – the single cornea in Group 3.



Figure 2.

Average transmissivity measured for the scars in Groups 1–3. The data are normalized to a thickness of 0.26 mm, which is the average thickness of the scars.



Figure 3.

Wavelength dependence of the total scattering cross-section for the scars in Groups 1–3. Multiplication by $(\lambda/550)^3$ removes the wavelength dependence of the average single fibril scattering cross-section. The straight lines are least squares fits to the data. According to the extension of Benedek's lake theory, positive slopes indicate the presence of lakes in the distribution of collagen fibrils.^{10, 12} The slopes are respectively 2.27, 3.87 and 12.8 for Groups 1, 2, and 3 and the R values of the fits are respectively 0.994, 0.997, and 0.997. Increasing values of the slopes represent greater contributions from lakes.



Figure 4.

Selected TEM of scarred regions that were analyzed to determine fibril sizes and positions. The rectangles outline the regions that were used in the DSF calculations of light scattering. (a) Mid-stroma region of cornea C8LM15983 from Group 1. The structure appears nearly normal, except that a small lake is visible near the center. (b) Mid-stroma region of cornea C8RM16270 from Group 2. Several lakes are visible and there are many enlarged fibrils. (c) Region from the posterior stroma of cornea C11LP16264 from Group 2. Several lakes are visible, but the fibrils appear to have more uniform diameters than those in (c). (d) Region from the anterior stroma of cornea C25LA16077 which was the single cornea in Group 3. Several large lakes are evident and there is a wide distribution of fibril diameters.



Figure 5.

Fibril diameter distributions from the micrographs in Figures 4 a–d. The distribution from cornea C8LM15983 in (a) is bimodal with distinct peaks at 320A and 390A. Bimodal diameter distributions have not been reported previously for such scars.



Figure 6.

Average transmission calculated from TEM. The DSF computations assume that the entire cornea has the structure depicted in the TEM. The transmission is normalized to a thickness of 0.26 mm, which is the average thickness of the scars.



Figure 7.

Wavelength dependence of the average total scattering cross-section calculated from the TEM. Multiplication of $\langle \rho \sigma_t (\lambda) \rangle$ by $(\lambda/550)^3$ removes the wavelength dependence of the average single fibril scattering cross-section. In agreement with the experimental data in Figure 3, plots of the calculated values of $(\lambda/550)^3 \langle \rho \sigma_t (\lambda) \rangle$ are also straight lines whose slopes progressively increase from Group 1 to 3.

Cornea Scar age (yrs)	C4R 4.5	C8L 4.5	C11R 4.0	C8R 4.5	C11L 4.0	C18R 3.6	C18L 3.6	C25R 3.6	C25L 3.6
Group	1	1	1	2	2	2	2	2	3
λ (nm)	$\rho\sigma_t(\mathbf{m})^{-1}$ $\rho\sigma_t(\mathbf{m})^{-1}$								
400	2.841	3.188	4.759	5.197	6.567	6.814	6.183	6.612	9.382
450	2.120	2.617	3.802	4.219	5.127	5.357	4.851	5.093	8.817
500	1.658	2.266	3.104	3.533	4.090	4.315	3.886	4.127	7.594
550	1.397	2.006	2.611	3.056	3.353	3.578	3.218	3.505	6.757
600	1.198	1.812	2.233	2.694	2.803	3.032	2.690	3.020	6.123
650	1.054	1.656	1.951	2.430	2.397	2.623	2.322	2.660	5.624
700	0.941	1.542	1.727	2.217	2.047	2.304	2.045	2.370	5.222

Table 1

Fibril Statistics

Table 2

Cornea/ EM	Scar Age (yrs)	Group	Region	N _{fib}	d _{avg} (A)	Range (A)	Distribution type
C4RM/	4.5	1	Middle	2144	402±40	232-479	Bimodal
15993 C4RP/	4.5	1	Posterior	3141	369±43	168-488	Bimodal
15995 C8LM/	4.5	1	Middle	2763	409±33	194–637	Unimodal
15565 C8LM/ 15082*	4.5	1	Middle	1755	364±40	209-451	Bimodal
C8LP/ 15560*	4.5	1	Posterior	2850	341±27	206-424	Unimodal
C11RP/ 16305 [*]	4.0	1	Posterior	1909	347±46	155–494	Unimodal
C8RA/	4.5	2	Anterior	1420	349±68	149–739	Unimodal
16270 C8RM/	4.5	2	Middle	1054	378±55	239–529	Bimodal
16275 C11LM/ 16257	4.0	2	Middle	1658	373±99	97–644	Bimodal
C11LP/ 16264*	4.0	2	Posterior	2224	342±49	148–489	Unimodal
C18RM/ 15169	3.6	2	Middle	808	428 ±115	119–723	Bimodal
C25RA/ 15675	3.6	2	Anterior	1337	476 ±108	152–742	Bimodal
C25RA/ 15630	3.6	2	Anterior	4352	370±40	160-555	Unimodal
C25RA/ 15632	3.6	2	Anterior	1294	415±92	122-708	Unimodal
15637 [*]	3.6	2	Middle	1/14	455±43	169-605	Unimodal
C25RM/ 15638 C25RM/	3.6	2	Middle	902 3035	456±39	192-593	Unimodal
15680 [*] C25RM/ 15683 [*]	3.6	2	Middle	3029	427±70	163–692	Bimodal
C25LA/	3.6	3	Anterior	2032	392±82	139-651	Unimodal
16077 [*] C25LA/ 16079 [*]	3.6	3	Anterior	3147	366±68	146-625	Unimodal

* TEM used for DSF calculations.