

# Conformational gating of the electron transfer reaction $Q_A^-Q_B \rightarrow Q_AQ_B^-$ in bacterial reaction centers of *Rhodobacter sphaeroides* determined by a driving force assay

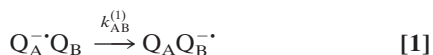
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Contributed by George Feher, July 29, 1998

**ABSTRACT** The mechanism of the electron transfer reaction,  $Q_A^-Q_B \rightarrow Q_AQ_B^-$ , was studied in isolated reaction centers from the photosynthetic bacterium *Rhodobacter sphaeroides* by replacing the native  $Q_{10}$  in the  $Q_A$  binding site with quinones having different redox potentials. These substitutions are expected to change the intrinsic electron transfer rate by changing the redox free energy (i.e., driving force) for electron transfer without affecting other events that may be associated with the electron transfer (e.g., protein dynamics or protonation). The electron transfer from  $Q_A^-$  to  $Q_B$  was measured by three independent methods: a functional assay involving cytochrome  $c_2$  to measure the rate of  $Q_A^-$  oxidation, optical kinetic spectroscopy to measure changes in semiquinone absorption, and kinetic near-IR spectroscopy to measure electrochromic shifts that occur in response to electron transfer. The results show that the rate of the observed electron transfer from  $Q_A^-$  to  $Q_B$  does not change as the redox free energy for electron transfer is varied over a range of 150 meV. The strong temperature dependence of the observed rate rules out the possibility that the reaction is activationless. We conclude, therefore, that the independence of the observed rate on the driving force for electron transfer is due to conformational gating, that is, the rate limiting step is a conformational change required before electron transfer. This change is proposed to be the movement, controlled kinetically either by protein dynamics or intermolecular interactions, of  $Q_B$  by  $\approx 5$  Å as observed in the x-ray studies of Stowell *et al.* [Stowell, M. H. B., McPhillips, T. M., Rees, D. C., Soltis, S. M., Abresch, E. & Feher, G. (1997) *Science* 276, 812–816].

Protein dynamics play an important role in many biological processes (1), including electron transfer reactions (2). When protein changes precede and are slower than the intrinsic rate of electron transfer, protein dynamics will control the rate of reaction. This situation has been called conformational gating (2). In this work, we study the role of protein dynamics in an electron transfer reaction of the bacterial reaction center (RC). The RC is a membrane protein that mediates the light-induced electron and proton transfer reactions in bacterial photosynthesis. (3) After photoexcitation of the primary donor, a bacteriochlorophyll dimer (D), an electron is transferred from the donor via a bacteriopheophytin (BPhe) to the tightly bound primary quinone,  $Q_A$  ( $k \approx 10^{10} \text{ s}^{-1}$ ). Subsequent electron transfer occurs from  $Q_A^-$  to the secondary quinone,  $Q_B$ , with observed rate,  $k_{AB}^{(1)} \approx 10^4 \text{ s}^{-1}$  (4, 5), that is,

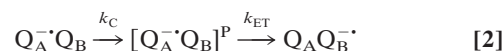


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The observed rate ( $10^4 \text{ s}^{-1}$ ) and activation energy ( $\approx 200$  to  $500$  meV) measured for the electron transfer between the two quinones are within ranges that are common for reactions governed by protein dynamics. Previous studies have suggested that this reaction is accompanied by conformational changes (6–10). In this work, we examine the role of conformational gating in the observed reaction (Eq. 1).

A simplified gating model for  $Q_B$  reduction involves two consecutive steps:



where  $Q_A^-Q_B$  represents the initial state,  $[Q_A^-Q_B]^P$  is the state activated for electron transfer and,  $Q_AQ_B^-$  is the final product state. In this model, the overall (observed) rate of the reaction depends on the values of  $k_{ET}$  and  $k_C$  (2). If  $k_C < k_{ET}$ , the reaction is conformationally gated. If  $k_C > k_{ET}$ , the reaction is rate-limited by electron transfer. To determine which of these two cases prevails, a driving force assay was used (11). In this assay, quinones with different redox potentials (i.e., driving force) are substituted for  $Q_A$ , thereby changing the driving force for electron transfer and, thus, the intrinsic rate constant,  $k_{ET}$ . This will change the overall rate  $k_{AB}^{(1)}$  if electron transfer is rate limiting but will not change  $k_{AB}^{(1)}$  if the reaction is gated.

Naphthoquinones (NQ) with redox potentials different from that of native  $Q_{10}$  were substituted into the  $Q_A$  site whereas  $Q_{10}$  was retained in the  $Q_B$  site. (11–13). The substitution of  $Q_A$  changes the intrinsic rate of electron transfer,  $k_{ET}$ , without affecting  $k_C$  since the  $Q_B$  site is unchanged and the structures of the substituted  $Q_A$  molecules are nearly identical. By using this procedure, the observed rate  $k_{AB}^{(1)}$  was measured as a function of the driving force for electron transfer, from which the mechanism of the  $k_{AB}^{(1)}$  reaction was determined. A preliminary account of this work has been published (14).

## MATERIALS AND METHODS

**Sample Preparation.** RCs from *Rhodobacter sphaeroides* strains R26.1 and 2.4.1 were isolated and purified in lauryldimethylamine-N-oxide (LDAO) as described (15). The ratios of absorbance ( $A^{280\text{nm}}/A^{800\text{nm}}$ ) of the purified RCs were  $< 1.25$ .  $Q_A$  and  $Q_B$  were removed from RCs as described (16). Reduced *R. sphaeroides* cytochrome (cyt)  $c_2$  was purified as described (17).  $Q_{10}$ ,  $MQ_0$ ,  $MQ_P$ , and  $MQ_4$  were obtained from Sigma. 2,3,5-trimethyl-1,4-naphthoquinone ( $Me_3NQ$ ) and 2,3,6,7-tetramethyl-1,4-naphthoquinone ( $Me_4NQ$ ) kindly were provided by J. M. Bruce (Univ. of Manchester, U.K.). They all were solubilized in ethanol before use.  $Q_{10}$  also was solubilized in 1% LDAO and 1%

Abbreviations: RC, bacterial reaction center; NQ, naphthoquinone;  $Me_3NQ$ , 2,3,5-trimethyl-1,4-naphthoquinone;  $Me_4NQ$ , 2,3,6,7-tetramethyl-1,4-naphthoquinone; LDAO, lauryldimethylamine-N-oxide; cyt  $c_2$ , *R. sphaeroides* cytochrome  $c_2$ ; BPhe, bacteriopheophytin;  $MQ_0$ , menadione;  $MQ_P$ , vitamin  $k_3$ ;  $MQ_4$ , menatetronone.

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Triton X-100. Unless otherwise indicated, experiments were performed in BMK buffer, which consisted of 0.04%  $\beta$ -D-dodecyl-maltoside, 10 mM KCl, 0.1 mM EDTA, and 5 mM of each of the following buffers: citrate, mes, pipes, tris, ches, and caps. Reagents were of analytical grade; their sources are given in ref. 11.

A detailed procedure for reconstitution of the  $Q_A$  binding site with naphthoquinone and  $Q_B$  binding site with ubiquinone has been presented (11). The occupancy of the  $Q_A$  and  $Q_B$  sites after reconstitution were determined from the amplitudes of transients associated with charge recombination,  $k_{AD}$  ( $D^+Q_A^- \rightarrow DQ_A$ ), and  $k_{BD}$  ( $D^+Q_AQ_B^- \rightarrow DQ_AQ_B$ ) (18). One-hundred percent  $Q_A$  site occupancies and 75–90%  $Q_B$  site occupancies were achieved at pH 7.2.

**Electron Transfer Measurements.** Absorbance changes were measured at 23°C by using a single beam spectrophotometer of local design (19). Voltage output was recorded on a digital oscilloscope (LeCroy 9310M, Chestnut Ridge, NY) and was digitally filtered (LabView, National Instruments, Austin, TX) and fitted (Peakfit, Tablecurve, or Sigmaplot, Jandel, San Rafael, CA). Actinic illumination was provided by a pulsed dye laser (PhaseR DL2100c, 590 nm,  $\approx 0.2$  J/pulse, 0.4- $\mu$ s full width at half maximum) or by a ruby laser (Optic Technologies Model 130, 695 nm,  $\approx 0.5$  J/pulse,  $\approx 10$ -ns full width at half maximum). Laser light intensity was determined to be saturating ( $>98\%$ ) before and after each set of experiments.

**The Electron Transfer from  $Q_A^-$  to  $Q_B$  Was Measured by Three Independent Methods.** (i) Optical absorbance changes at 412 and 470 nm. The difference in extinction between ubisemiquinone and naphthosemiquinones provides a spectral signature of the electron transfer,  $(NQ)_A^-Q_B \rightarrow (NQ)_AQ_B^-$ , in hybrid RCs. Peaks in the differential extinction occur at wavelengths of 412 and 470 nm (20, 21). In native RCs,  $Q_A$  and  $Q_B$  are the same ( $Q_{10}$ ); there is, however an absorption change at 412 nm caused by an electrochromic shift of the BPhe absorption. No electrochromic shift is observed at 470 nm.

(ii) Optical absorbance changes at 757 nm. At this wavelength, the absorption changes associated with electron transfer,  $Q_A^-Q_B \rightarrow Q_AQ_B^-$ , are caused by an electrochromic shift of the BPhe absorption (9). The observed kinetics were corrected for fluorescent artifacts by subtracting the absorbance changes when electron transfer was blocked by *tert*butyryne.

(iii) Two-flash cyt  $c_2$  assay (22). Optical absorbance changes at 865 nm or at 550 nm were measured after two laser flashes separated by a variable time interval,  $\Delta t$ , from 1 to 10,000  $\mu$ s. The fraction of photoactive RCs,  $f(\Delta t) = 1 - \exp[-(k_{AB}^{(1)}\Delta t)]$ , was calculated from the absorbance change at 865 nm due to oxidation of the primary donor, D, from the relation

$$f(\Delta t) = \Delta A_2^{865}(\Delta t) / \Delta A_{(N_o \text{ cyt})}^{865} \quad [3]$$

where  $\Delta A_2^{865}(\Delta t)$  is the absorbance change induced by the second flash, and  $\Delta A_{(N_o \text{ cyt})}^{865}$  is the absorption change after a single flash in the absence of cyt  $c_2$ . Normalization of the data required correction for several small effects, whose magnitudes were determined independently, i.e., (i) the fraction of RCs without a functional  $Q_B$ , (ii) the fraction of RCs without cyt  $c_2^+$  bound before the first flash, (iii) the thermodynamic equilibrium between the  $Q_A^-Q_B$  and  $Q_AQ_B^-$  states, and (iv) the exchange of cyt  $c_2^+$  for cyt  $c_2^{2+}$  between flashes.

From absorption changes at 550 nm, which result from cyt  $c_2^{2+}$  oxidation, the fraction of photoactive RCs was obtained from

$$f(\Delta t) = \Delta A_2^{550}(\Delta t) / \Delta A_2^{550}(\Delta t = 10 \text{ ms}) \quad [4]$$

where  $\Delta A_2^{550}(\Delta t)$  is the absorbance change induced by the second flash given  $\Delta t$  seconds after the first and  $\Delta A_2^{550}(\Delta t = 10 \text{ ms})$  is the absorbance change induced by the second flash given 10 ms after the first. Because at 10 ms the reaction is complete,  $\Delta A_2^{550}(\Delta t = 10 \text{ ms})$  is the maximum possible change. Normalization to  $\Delta A_2^{550}(\Delta t = 10 \text{ ms})$ , instead of the amplitude change induced by the first flash, avoids the need to apply corrections *i*, *ii*, and *iii*.

## RESULTS

**Hybrid RCs [(NQ)<sub>(A)</sub>, Q<sub>10(B)</sub>].** Hybrid RCs containing different naphthoquinones in the  $Q_A$  site and  $Q_{10}$  in the  $Q_B$  site [(NQ)<sub>(A)</sub>, Q<sub>10(B)</sub>] were prepared and assayed as described in *Materials and Methods*. The fraction of hybrid RCs, [(NQ)<sub>(A)</sub>, Q<sub>10(B)</sub>]/{[(NQ)<sub>(A)</sub>, Q<sub>10(B)</sub>] + [Q<sub>10(A)</sub>, Q<sub>10(B)</sub>]}, ranged from 0.45 to 0.80 for samples containing MQ<sub>P</sub>, MQ<sub>4</sub>, Me<sub>3</sub>NQ, and Me<sub>4</sub>NQ. A lower value of 0.35 was determined for samples with MQ<sub>0</sub>. From measurements of the recombination rates  $k_{AD}$  and  $k_{BD}$ , the *in situ* free-energy changes relative to  $Q_{10}$ ,  $\delta\Delta G^0$  of the substituted quinones were obtained (see Table 1) (11–13, 23, 24).

**Determination of Electron Transfer Rates,  $Q_A^-Q_B \rightarrow Q_AQ_B^-$ , in Hybrid [(NQ)<sub>(A)</sub>, Q<sub>10(B)</sub>] and Native (Q<sub>10(A)</sub>, Q<sub>10(B)</sub>) RCs.** The electron transfer rates  $k_{AB}^{(1)}$  were determined by three

Table 1. Parameters describing the characteristic time of electron transfer,  $Q_A^-Q_B \rightarrow Q_AQ_B^-$ , in hybrid and native RCs

$Q_A$	$\delta\Delta G^0$ , meV	Two-flash assay		Optical absorption change assays			$1/k_{AB}^{(1)}$ , $\mu$ s (avg)
		865 and 550 nm	412 nm	470 nm	757 nm		
$Q_{10}$	0	47 $\mu$ s (37%) 200 $\mu$ s (57%)	120 $\mu$ s (–27%)	NA	27 $\mu$ s (26%) 250 $\mu$ s (74%)	110 $\pm$ 3	
MQ <sub>0</sub>	15 $\pm$ 5	ND	42 $\mu$ s (36%) 340 $\mu$ s (27%)	48 $\mu$ s (25%) 310 $\mu$ s (58%)	76 $\mu$ s (47%) 320 $\mu$ s (53%)	150 $\pm$ 20	
MQ <sub>P</sub>	–35 $\pm$ 9	ND	51 $\mu$ s (27%) 220 $\mu$ s (43%)	32 $\mu$ s (44%) 230 $\mu$ s (45%)	64 $\mu$ s (29%) 260 $\mu$ s (71%)	140 $\pm$ 30	
MQ <sub>4</sub>	–46 $\pm$ 9	ND	47 $\mu$ s (18%) 240 $\mu$ s (42%)	50 $\mu$ s (13%) 210 $\mu$ s (77%)	51 $\mu$ s (28%) 170 $\mu$ s (72%)	150 $\pm$ 20	
Me <sub>3</sub> NQ	–95	ND	52 $\mu$ s (15%) 270 $\mu$ s (48%)	44 $\mu$ s (23%) 340 $\mu$ s (63%)	33 $\mu$ s (35%) 230 $\mu$ s (65%)	190 $\pm$ 30	
Me <sub>4</sub> NQ	–120 $\pm$ 10	46 $\mu$ s (48%) 200 $\mu$ s (49%)	50 $\mu$ s (24%) 320 $\mu$ s (34%)	NA	50 $\mu$ s (42%) 320 $\mu$ s (58%)	140 $\pm$ 20	

Data were fitted with two exponentials. The characteristic times are given in microseconds together with their relative contributions in parenthesis (757 nm). The percent amplitude relative to the steady-state change ( $NQ_A^-$  minus  $Q_{10(13)}$ ) is given in parentheses for 412 and 470 nm.  $\delta\Delta G^0$  is given relative to  $Q_{10}$ , which is defined as 0. The tabulated values in the seventh column represent the average 1/e time of the three methods. Errors were calculated as the SD of the mean. NA, not applicable; ND, not determined.

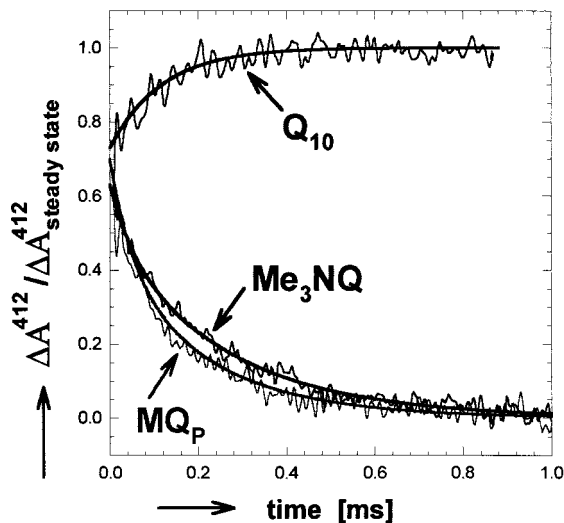


FIG. 1. Optical absorbance changes at 412 nm associated with electron transfer from  $Q_A^-$  to  $Q_B$  in hybrid ( $Me_3NQ^-$  or  $MQP^-$  in the  $Q_A$  site and  $Q_{10}$  in the  $Q_B$  site) and native ( $Q_{10}$  in both  $Q_A$  and  $Q_B$  sites) RCs. In native RCs, the absorbance change is due to an electrochromic shift of the BPhe absorption. In hybrid RCs, the main effect is caused by the difference in extinction coefficient between the naphtho-semiquinones ( $NQ^-$  in the  $Q_A$  site and  $Q_{10}^-$  in the  $Q_B$  site). The observed transients in hybrid RCs are reduced by the electrochromic shift, which was assumed to be the same as that observed in native RCs. The data for the hybrid RCs were normalized to the expected absorption determined from the known difference in extinction coefficients between  $(NQ)_A^-$  and  $Q_{10}^-$  (11, 20, 21). A shallow slope determined by monitoring the recombination reactions ( $D^+Q_A^- \rightarrow DQ_A$  and  $D^+Q_AQ_B^- \rightarrow DQ_AQ_B$ ) was subtracted for clarity. Addition of tertbutryne, which inhibits electron transfer to  $Q_B$ , eliminated the transient kinetic phases. [Conditions: 5  $\mu M$  2.4.1RCs, BMK buffer (pH = 7.2), 0.03% LDAO. Spectral bandwidth, 10 nm].

independent methods described in *Materials and Methods*. The use of several methods helped to guard against possible artifacts caused by transient signals not associated with electron transfer. Of particular importance is the two-flash functional assay that directly measures electron transfer. The decay kinetics were fitted by one or two exponential rate constants (see Table 1). The fitted values should be viewed as empirical parameters and are not intended to imply a mechanism. Biexponential decays were observed even with  $Q_{10}$  in both binding sites and, therefore, are not related to the mixture of hybrid and native RCs. The results of the kinetic measurements are described below.

Transient absorbance changes at 412 nm after a laser flash are shown for hybrid and native RCs in Fig. 1. The solid lines represent fits to a biexponential function. Despite the change in electron driving force, the decay rates of the two hybrid samples were essentially the same. The amplitudes were normalized by using the steady-state semiquinone absorption spectra (with  $\Delta\epsilon_{(NQ)_A^- - Q_{10}^-}^{412\text{ nm}} = 4.3\text{ mM}^{-1}\text{ cm}^{-1}$ ) (11, 20, 21). In native RCs, the absorbance change is caused by an electrochromic shift of the BPhe absorption. The same shift is assumed to be present also in hybrid RCs. To estimate the fraction of hybrid RCs displaying the measured electron transfer, the electrochromic shift observed for native RCs must be added to the absorbance changes of hybrid RCs. This results in an amplitude for the  $(NQ)_A^-$  transient decay that is  $90 \pm 5\%$  of the steady-state (i.e., expected) semiquinone decay.

At 470 nm, no electrochromic shifts are observed, and, therefore, no transient signals are obtained in native RCs. The amplitudes of the observed transients in hybrid RCs caused by  $(NQ)_A^-$  absorption decay amounted to  $89 \pm 5\%$  of the expected steady-state signal (with  $\Delta\epsilon_{(NQ)_A^- - Q_{10}^-}^{470\text{ nm}} = 2.5\text{ mM}^{-1}\text{ cm}^{-1}$  for

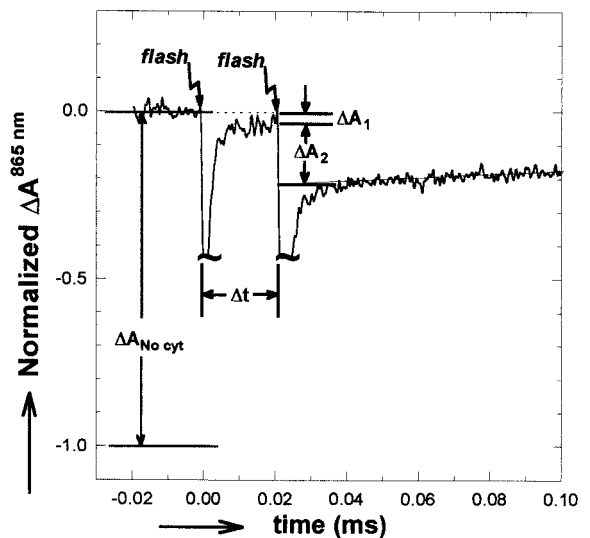


FIG. 2. Optical absorbance changes at 865 nm for native RCs in the presence of  $cyt\ c_2^{2+}$  after two laser flashes spaced  $\Delta t$  apart. Laser flash artifacts are truncated. The ratio of  $\Delta A_2$  to  $\Delta A_{No\ cyt}$  gives the fraction of photoactive RCs at the time of the second flash (see Eq. 3).  $\Delta A_1$  is due to the fraction of RCs that do not have a bound  $cyt\ c_2^{2+}$ . Addition of tertbutryne to inhibit electron transfer to  $Q_B$  eliminated the absorbance change in response to the second laser flash. (Conditions: 3  $\mu M$  2.4.1 RCs with  $Q_{10}$  reconstituted into the  $Q_A$  and  $Q_B$  sites by addition of  $Q_{10}$  in LDAO. After reconstitution, LDAO was 0.03%. 100  $\mu M$   $cyt\ c_2^{2+}$ , 5  $\mu M$  ascorbate, 1 mM Tris, and 0.04% maltoside; spectral bandwidth, 10 nm.)

$MQ_0$ ,  $MQ_P$ , and  $MQ_4$  (20, 21) and  $5.0\text{ mM}^{-1}\text{ cm}^{-1}$  for  $Me_3NQ$ ) (11). The characteristic times associated with the observed transients at 470 nm were, within experimental error, the same as those obtained at 412 nm (see Table 1).

The transient absorption changes at 757 nm that arise from an electrochromic shift associated with  $k_{AB}^{(1)}$  (9, 21, 25, 26) were measured in native as well as in all hybrid RCs. The data were fitted with a biexponential function and were in general agreement with the results obtained by the other methods. Replacement of  $Q_{10}$  with  $Q_2$  and  $Q_0$  in the  $Q_B$  site had no effect on the observed decay rates.

The functional two-flash assay was performed at two wavelengths (865 and 550 nm) for hybrid and for native RCs. The result of an 865 nm assay is shown in Fig. 2. After one laser flash, the charged species,  $D^+Q_A^-Q_B$ , was formed.  $Cyt\ c_2^{2+}$  reduced the primary donor to form  $DQ_A^-Q_B$  ( $\tau = 0.8\ \mu s$  in RCs with a bound  $cyt\ c_2^{2+}$ ). A small residual absorbance ( $\Delta A_1$ ) remained due to RCs without  $cyt\ c_2^{2+}$  bound before the first flash. For RCs in which a second flash was given before the electron left  $Q_A^-$ , no charge separation on the timescale of our experiments and consequently no optical transients were observed. In RCs in which electron transfer had occurred before the second flash, the photoactive state,  $DQ_AQ_B^-$ , could undergo charge separation (to form  $D^+Q_A^-Q_B^-$ ) with a concomitant optical absorbance change observable at 865 nm  $\Delta A_2$ . The absorbance change after the second flash is directly related to the rate of electron transfer (see Eq. 3). Similar measurements of optical absorbance changes at 550 nm (not shown) caused by oxidation of  $cyt\ c_2^{2+}$  by  $D^+$  after a second flash also were made (see Eq. 4).

The applicability of the measurements at 865 nm is limited to the time interval  $1\ \mu s < \Delta t < 100\ \mu s$ . The lower limit is given by the electron transfer time from  $cyt\ c_2^{2+}$  to  $D^+$  ( $\approx 0.8\ \mu s$ ). The upper limit is given by the  $cyt\ c_2^{3+}/cyt\ c_2^{2+}$  exchange time (0.5–0.7 ms): If  $cyt\ c_2^{3+}/cyt\ c_2^{2+}$  exchange takes place between flashes,  $D^+$  will be reduced by  $cyt\ c_2^{2+}$ . Accurate measurements at 550 nm were obtained only for time intervals  $\Delta t > 40\ \mu s$  because of the relatively poor signal-to-noise ratio associated with the small

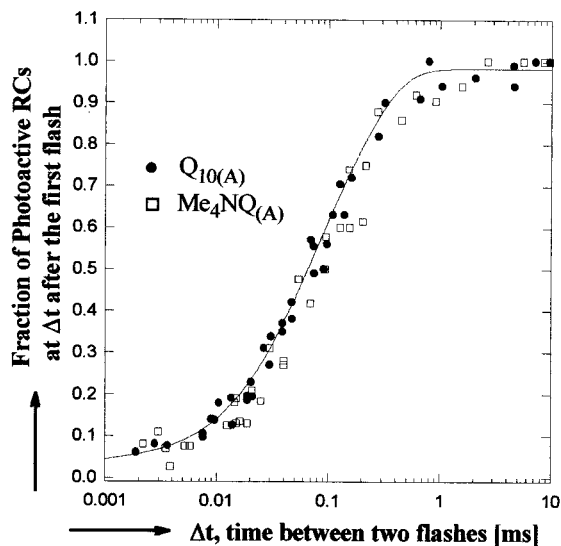


Fig. 3. The fraction of photoactive RCs as a function of the time interval between two laser flashes for native (●), and hybrid (□) RCs. The solid line represents a least squares biexponential fit to the data. A minimum of 20 min dark recovery was allowed between sets of flashes. (865 nm assay conditions given in Fig. 2; at 550 nm, cyt  $c_2^{2+}$  concentration was reduced to  $\approx 30 \mu\text{M}$ ; spectral bandwidth, 4 nm.)

difference in extinction between oxidized and reduced cyt  $c_2$  ( $\approx 21 \text{ mM}^{-1} \text{ cm}^{-1}$ ). For time  $40 \mu\text{s} < \Delta t < 100 \mu\text{s}$ , measurements were performed at both wavelengths; the two sets of measurements agreed within experimental error.

The normalized fraction of photoactive RCs (calculated by using Eqs. 3 and 4) is plotted in Fig. 3 as a function of the time interval between flashes for RCs having either  $Q_{10}$  or  $\text{Me}_4\text{NQ}$  in the  $Q_A$  site. The solid line represents a biexponential fit to the data (see Table 1). The results in Fig. 3 show that the rate of electron transfer is the same for both quinones. RCs from the strain 2.4.1 containing carotenoid were used for the studies shown in Fig. 3 to eliminate a small absorption change at very short times in carotenoid-free RCs (R26.1). Control experiments (not shown) confirmed that 2.4.1 and R26.1 RCs display the same electron transfer rates.

The results determined from the three independent methods are consistent with one another (Table 1). A model independent way to compare the data from samples with different  $Q_A$  molecules is to compare the time required to decay to  $1/e$  of the initial state. The  $1/e$  times, averaged over the three methods, for RCs with various quinones in the  $Q_A$  sites are similar and show no correlation with the redox potential of  $Q_A$  (see Table 1). The rates of each of the two components of the biexponential fit derived from the two-exponential analysis also do not vary systematically with driving force.

**pH Dependence of Electron Transfer in Hybrid ( $\text{Me}_3\text{NQ}_{(A)}$ ,  $Q_{10(B)}$ ) and Native ( $Q_{10(A)}$ ,  $Q_{10(B)}$ ) RCs.** The rate  $k_{AB}^{(1)}$  was measured at 470 and 757 nm over the pH range of 4.7 to 10.3 for hybrid RCs with  $\text{Me}_3\text{NQ}$  in the  $Q_A$  site and for native RCs. The rate was independent of the driving force throughout the entire range. The hybrid RCs retained the pH dependence of native RCs; that is, the rate was pH-independent below pH 8 and decreased by a factor of  $\approx 10$  per pH unit above pH 8 (19, 25, 27).

**Effect of Exogenous Quinone Concentration [ $Q_{10}$ ] on Electron Transfer.** We measured the effect of quinone concentration [ $Q_{10}$ ] on  $k_{AB}^{(1)}$  in native RCs. One phase of the biexponential rate was unaffected by [ $Q_{10}$ ] and remained  $\approx 7,000 \text{ s}^{-1}$  whereas the second rate depended strongly on [ $Q_{10}$ ] (e.g., at  $5 \mu\text{M}$  [ $Q_{10}$ ],  $k_2 = 1,300 \text{ s}^{-1}$  and, at  $85 \mu\text{M}$  [ $Q_{10}$ ],  $k_2 = 25,000 \text{ s}^{-1}$ ). These experiments are still in progress.

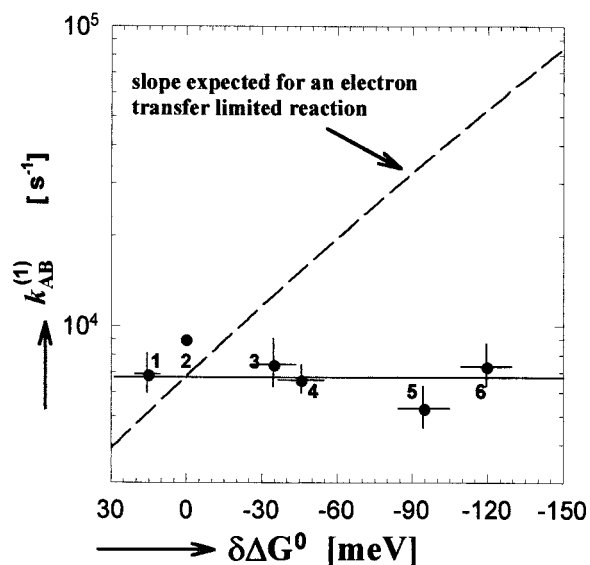


Fig. 4. The electron transfer rate  $k_{AB}^{(1)}$  as a function of redox free energy (driving force) for electron transfer (data from the second and seventh columns of Table 1). The solid line gives a best fit to the data. The dashed line represents the dependence predicted by the Marcus theory for an electron transfer limited reaction using the measured free energy difference between the initial,  $Q_A^-Q_B$ , and final,  $Q_AQ_B^-$ , states of 60 meV (19) and 1.1 eV for the reorganization energy (13, 31). Error bars represent one SD of the mean. Quinones incorporated into the  $Q_A$  site listed in order of decreasing redox potential were: 1,  $\text{MQ}_0$ ; 2,  $Q_{10}$ ; 3,  $\text{MQP}$ ; 4,  $\text{MQ}_4$ ; 5,  $\text{Me}_3\text{NQ}$ ; and 6,  $\text{Me}_4\text{NQ}$ . Redox energies are quoted with respect to  $Q_{10}$  (see ● without error bars).

## DISCUSSION

In this work, we tested whether conformational changes or intrinsic electron transfer determines the kinetics of the observed reaction,  $Q_A^-Q_B \rightarrow Q_AQ_B^-$ , in RCs from *R. sphaeroides*. This was accomplished by substituting naphthoquinone molecules with redox potentials different from that of native  $Q_{10}$  into the  $Q_A$  binding site while retaining  $Q_{10}$  in the  $Q_B$  site. The experimental results showed that the observed reaction rates were independent of the electron transfer driving force.

**Analysis of the Driving Force Assay.** For a typical electron transfer reaction, the observed rate depends on the free energy for electron transfer as given by the Marcus theory (28):

$$k_{\text{ET}} = \frac{2\pi|T_{\text{AB}}|^2}{\hbar\sqrt{4\pi\lambda k_{\text{B}}T}} \exp\left[-\frac{(\Delta G_{\text{ET}}^0 + \lambda)^2}{4\lambda k_{\text{B}}T}\right] \quad [5]$$

where  $k_{\text{ET}}$  is the intrinsic electron transfer rate constant,  $\Delta G_{\text{ET}}^0$  is the free energy for electron transfer (defined as the energy of the final minus the initial state),  $\lambda$  is the reorganization energy,  $T_{\text{AB}}$  is the tunneling matrix element, and  $k_{\text{B}}T$  is the thermal energy. In principle, the independence of  $k_{AB}^{(1)}$  on  $\Delta G_{\text{ET}}^0$  could be caused by the reaction being activationless ( $-\Delta G_{\text{ET}}^0 = \lambda$ ). However, this is not the case for  $k_{AB}^{(1)}$ , for which a large activation energy is observed (9, 19, 29, 30). Values for  $\lambda$  and  $\Delta G_{\text{ET}}^0$  have been determined to be  $\approx 1.1 \text{ eV}$  and  $\approx -0.1 \text{ eV}$ , respectively (13, 19, 31, 32). For these values,  $k_{\text{ET}}$  is expected to vary by a factor of 10 over the range of free energy spanned in this study\* (see dashed line in Fig. 4). This is in contrast to the observed invariance of the rate with electron transfer driving force shown in Fig. 4. A fit of the data to a biexponential function (see Table 1) shows that neither com-

\*We assume that the values of  $\lambda$  and  $T_{\text{AB}}$  remain the same on NQ substitution (11). A driving force dependence similar to the dashed line in Fig. 4 was observed for the second electron reduction of  $Q_B$  (11).

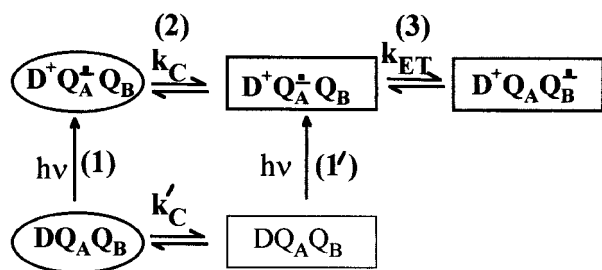


FIG. 5. Proposed model for conformational gating of electron transfer from  $Q_A^-$  to  $Q_B$ . Two conformational states are shown, distal (D; oval), which is inactive, and proximal (P; rectangle), which is active in electron transfer. The main sequence of steps in the reaction are (1) excitation of RCs in the distal state, (2) conversion to the proximal state, and (3) electron transfer. A fraction of RCs in the proximal state before light excitation may give rise to a faster observed rate through reactions 1' and 3.

ponent of the decay depends on electron transfer driving force. These results show that  $k_{AB}^{(1)}$  is not limited by electron transfer.

**Conformational Gating.** The independence of  $k_{AB}^{(1)}$  on driving force can be explained by a conformational gating mechanism (2). Fig. 5 depicts a kinetic scheme for an electron transfer reaction coupled with a conformational change. The protein can exist in two different states indicated by different shapes in Fig. 5: the D (oval) and P (rectangle) states, for Distal and Proximal, as will be discussed below. In the D state,  $Q_B$  is in an inactive conformation, and electron transfer is impaired; in the P state,  $Q_B$  is in the active configuration, and electron transfer can proceed. The rate of equilibration between the D and P states in the dark is  $k_C$ . The gating model requires that, in the dark, the RC is predominantly in the D state. Light excitation results in formation of  $D^+Q_A^-Q_B^-$ , with  $Q_B$  inactive. Conversion to the  $Q_B$  active configuration occurs with the conformational gating rate  $k_C$ . Once the RC is in the  $Q_B$  active configuration, electron transfer can occur with a rate  $k_{ET}$  (given by Eq. 5). Thus, for the conformational gating model the observed  $k_{AB}^{(1)} \approx 10^4 \text{ s}^{-1} = k_C$ . The intrinsic electron transfer rate must be much greater than the observed rate (i.e.,  $k_{AB}^{(1)} \ll k_{ET}$ ). This is consistent with recent calculations suggesting that the intrinsic electron transfer rate for this reaction is  $\gg 10^4 \text{ s}^{-1}$  (33). The model is also consistent with the work of Li *et al.* (30). These authors reported an additional fast phase (in the  $10^5$ – $10^6 \text{ s}^{-1}$  range) in the  $k_{AB}^{(1)}$  kinetics measured at 398 nm and attributed this phase to electron transfer in a population of RCs active for electron transfer in the dark (i.e., our P state in the dark; see Fig. 5).

The gating model accounts for the observation by Kleinfeld *et al.* (6) that electron transfer from  $Q_A^-$  to  $Q_B$  proceeds in RCs cooled to cryogenic temperature under illumination but does not proceed in RCs cooled in the dark. RCs frozen in the dark are in the inactive conformation (D state). At low temperature, there is insufficient thermal energy for conformational interconversion. RCs cooled under illumination are frozen in the active P state, which explains the observed electron transfer in these RCs.

**Molecular Basis of the Gating Mechanism.** *Structural changes responsible for the gating.* Structural changes involved in the gating process were elucidated recently by comparing the x-ray crystal structure of the RC cooled to cryogenic (90 K) temperatures under illumination (light structure,  $D^+Q_A^-Q_B^-$ ) with the structure of the RC cooled in the dark (dark structure,  $DQ_AQ_B$ ) (10). In the light structure,  $Q_B$  had moved closer to  $Q_A$  and had undergone a  $180^\circ$  propeller twist about the isoprenoid chain (see Fig. 6). The movement of  $Q_B$  is expected to affect both the rate and thermodynamic equilibrium of the electron transfer reaction. The more direct pathway (and shorter distance) between  $Q_A^-$  and  $Q_B$  in the light structure increases the rate of electron transfer. The greater number of hydrogen bonds to  $Q_B$ , compared with the dark structure, stabilizes  $Q_AQ_B^-$  with respect to  $Q_A^-Q_B$ , thereby increasing the

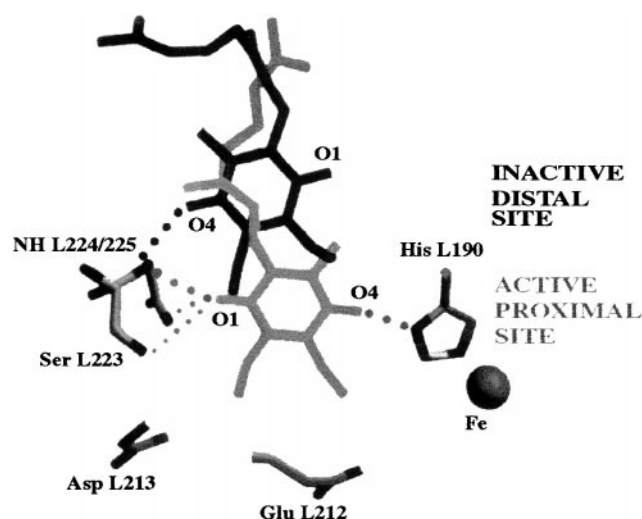


FIG. 6. Comparison of the binding positions for  $Q_B$  determined from the light ( $D^+Q_A^-Q_B^-$ ) and dark ( $DQ_AQ_B$ ) x-ray crystal structures of the RC (10). Movement from an inactive–distal (black) to an active–proximal (gray) binding site is proposed as the major structural change involved with conformational gating of electron transfer from  $Q_A^-$  to  $Q_B$ . Hydrogen bonding partners are connected by dotted lines. Modified from ref. 10.

driving force for electron transfer. The position of  $Q_B$  in the light structure, hydrogen bonded to Ser-L223 and to His-L190, corresponds to the active P (proximal) conformation in Figs. 5 and 6. The position of  $Q_B$  in the dark structure, hydrogen-bonded only to the amide hydrogen of Ile-L224, corresponds to the inactive D (distal) conformation in Figs. 5 and 6. Thus, the conformational gating step involves the movement of  $Q_B$  by  $\approx 5 \text{ \AA}$  from the distal to the proximal site.

*What controls the kinetics of the gating process?* The unhindered movement and rotation of the quinone in solution would proceed on the nanosecond timescale (34, 35), that is, 4–5 orders of magnitude faster than the observed value of  $1/k_{AB}^{(1)}$ . Consequently, the rate of  $Q_B$  movement from the distal to the proximal site must be controlled by a process other than diffusion. What is the rate limiting step for the  $Q_B$  movement? Several possibilities can be envisioned; three of these are discussed below.

The most likely gating process is associated with protein dynamics that allow the quinone to move into the proximal position. Support for gating by protein dynamics comes from experiments that showed (i) light-induced proton binding changes on the timescale of  $k_{AB}^{(1)}$  that occur in the absence of electron transfer from  $Q_A^-$  to  $Q_B$  (36), (ii) light-induced electrogenic transients in the absence of electron transfer from  $Q_A^-$  to  $Q_B$  with the same rate and pH dependence as  $k_{AB}^{(1)}$  in native and EQ(L212) mutant RCs (8), and (iii) the importance of protein fluctuations in electron transfer reactions of the RC (37, 38). Structural changes associated with protein dynamics should, in principle, be observable in the x-ray crystal structures. However, no large changes between protein structures of light and dark adapted RCs were reported (10). The largest change was the loss of electron density for Glu-H173 (close to  $Q_B$ ) in the light structure, which was attributed to greater disorder. The absence of large changes can be explained if the gating is due to protein “breathing” modes not detected at the cryogenic temperatures at which the structures were determined or if the gating motion involves only small movements of many parts of the protein or changes in water structure around the quinone region, not sufficiently well resolved in the x-ray crystal structures.

Another possibility is that conformational changes at the  $Q_B$  site are triggered by  $Q_A$  reduction. In the gating scheme shown in Fig. 5, a triggering mechanism implies that coupling between

the  $Q_A$  and  $Q_B$  sites exists and that  $k_C > k'_C$ . Evidence for coupling comes from experiments which show (i) an altered recombination rate ( $D^+Q_A^- \rightarrow DQ_A$ ) and EPR spectrum of  $Q_A^-$  when inhibitors are present in the  $Q_B$  site (19, 39) and (ii) proton binding by amino acids near  $Q_B$  in response to  $Q_A$  reduction (32, 40, 41). Electrostatic calculations also suggest coupling between the  $Q_A$  and  $Q_B$  binding sites (42, 43).

A third possibility is that the rate-limiting step involves the breaking of the intermolecular interactions between  $Q_B$  in the distal state and the protein. The result reported in this work and by McComb *et al.* (44), that the rates of  $k_{AB}^{(1)}$  for RCs with  $Q_{10}$ ,  $Q_2$ , or  $Q_0$  in the  $Q_B$  site are the same, excludes interactions associated with the isoprenoid tail as rate limiting. However, interactions of the quinone head group, for example,  $\pi$ -stacking with Phe-L216 or hydrogen-bonding between the O-4 carbonyl of  $Q_B$  and the amide hydrogen of Ile-L224, may limit the rate of movement between the D and P states.

**Non-Exponential Rates.** The kinetics of  $k_{AB}^{(1)}$  observed in this study were nonexponential and could be fitted with a biexponential function. Biphasic rates have been reported by several groups (9, 30, 45). In this work one component of the rate was observed to depend on the exogenous quinone concentration. We postulate that the origin of this component of the kinetics is associated with the fraction of RCs that do not have a bound quinone in the  $Q_B$  site. In these RCs, the rate limiting (gating) step is associated with quinone binding. This may explain both the variability and the discrepancy between rates reported by different groups. The discussion of the molecular basis of the gating mechanism presented in the previous section refers to the invariant part of the kinetics associated with those RCs that have a bound quinone.

**Conformational Gating in Other Systems.** Electron transfer reactions gated by ligand binding or redox-state dependent protein changes occur in other biological systems (46, 47). It is likely that some of the electron transfer reactions of energy transducing proteins, including cytochrome oxidase, cyt *bc<sub>1</sub>/bc<sub>1</sub>*, complex I, and others are conformationally gated. For example, crystallographic data of the *bc<sub>1</sub>* complex show an  $\approx 10$  Å movement of the Fe-S subunit between two sites, one for reduction of the Fe-S center by ubiquinol and another for oxidation of the Fe-S center by cytochrome *c* (48, 49). To understand protein-catalyzed electron transfer reactions in general, it is necessary to establish whether the intrinsic reaction step or the associated conformational changes are rate limiting. The technique described in this work, whereby the free energy for electron transfer is varied by cofactor substitution, illustrates a general approach to investigate reaction mechanisms and energetics.

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