# Compensating increases in protein backbone flexibility occur when the Dead ringer AT-rich interaction domain (ARID) binds DNA: A nitrogen-15 relaxation study

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#### Abstract

AT-rich interaction domains (ARIDs) are found in a large number of eukaryotic transcription factors that regulate cell proliferation, differentiation, and development. Previously we elucidated how ARIDs recognize DNA by determining the solution structure of the *Drosophila melanogaster* Dead ringer protein in both its DNA-free and -bound states. In order to quantitatively determine how ARIDs alter their mobility to recognize DNA, we have measured the relaxation parameters of the backbone nitrogen-15 nuclei of Dead ringer in its free and bound forms, and interpreted these data using the model-free approach. We show that Dead ringer undergoes significant changes in its mobility upon binding, with residues in the loop connecting helices H5 and H6 becoming immobilized in the major groove and contacts to the minor groove slowing down the motion of residues at the C terminus. A DNA-induced rotation and displacement of the N-terminal subdomain of the protein increases the mobility of helix H1 located distal to the DNA interface and may partially negate the entropic cost of immobilizing interfacial residues. Elevated motions on the micro- to millisecond timescale in the N-terminal domain prior to DNA binding appear to foreshadow the DNA-induced conformation change.

**Keywords:** AT-rich interaction domain; ARID; protein–DNA interaction; NMR; backbone dynamics; <sup>15</sup>N-relaxation; Dead ringer protein

Protein–nucleic acid recognition plays a central role in the transcription, replication, repair, and recombination of DNA. Most DNA-binding transcription factors have evolved to recognize specific nucleotide sequences with dissociation constants ranging from  $\sim 10^{-9}$  to  $10^{-12}$  M (Jen-Jacobson 1997). This level of affinity enables their revers-

ible binding in vivo, but is strong enough to facilitate high occupancy and selective binding to regulatory elements within the genome. The three-dimensional structures of numerous protein–DNA complexes and thermodynamic studies have revealed highly complementary binding surfaces that are primarily stabilized by enthalpic effects associated with hydrogen bonding and salt-bridge formation, and the entropically favorable release of water and counterions (Patikoglou and Burley 1997; Nadassy et al. 1999; Garvie and Wolberger 2001). Counterbalancing these interactions is the entropic cost associated with complex formation (the loss of six degrees of freedom, three translational and three rotational), and any unfavorable enthalpic effects (for example, DNA or protein distortions that occur in the complex). Beginning with the pioneering work of Pabo and Sauer on the

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 $\lambda$  repressor, it was realized that flexible regions within proteins could undergo local folding upon DNA binding (Pabo et al. 1982). Although entropically unfavorable, immobilization contributes to binding specificity by enabling proteins to adaptively recognize their DNA ligands (Spolar and Record 1994; Wright and Dyson 1999).

In this work we studied how the AT-rich interaction domain (ARID) from the Dead ringer protein adapts its backbone dynamics to recognize DNA. The ARID is a novel DNA-binding module found in a large number of eukaryotic transcription factors that regulate cell proliferation, differentiation, and development (for review, see Kortschak et al. 2000 and Wilsker et al. 2002). ARIDs were originally named based on the propensity of the ARID containing Bright, Dead ringer and Mrf-2 proteins to sequence-specifically bind to AT-rich DNA (Herrscher et al. 1995; Gregory et al. 1996; Valentine et al. 1998; Whitson et al. 1999), but more recently ARID containing proteins such as Osa and Swi1 have been shown to bind DNA nonsequence specifically (Collins et al. 1999; Dallas et al. 2000). Previous NMR studies of the Mrf-2, Swi1, and Dead ringer proteins have revealed that ARIDs adopt a unique  $\alpha$ -helical fold, but that they are structurally diverse (Yuan et al. 1998; Iwahara and Clubb 1999b; Iwahara et al. 2002; Kim et al. 2004). All ARIDs contain an ~80-amino-acid conserved core region, but can be partitioned into three distinct structural subfamilies because of differences outside the core region: (1) minimal-ARIDs like Mrf-2 consist of a conserved core domain formed by six  $\alpha$ -helices (H2–H7 in the ARID) (Yuan et al. 1998); (2) extended-ARIDs, like the Dead-ringer protein described in this study, contain the core domain and additional α-helices at their N and C termini (Iwahara and Clubb 1999b); and (3) N-terminal extended-ARIDs, which supplement the core domain with an amino-terminal  $\alpha$ -helix. This subfamily was originally predicted based on primary sequence homology (Iwahara and Clubb 1999b) and has recently been visualized in the structure of the Swi1 protein (Kim et al. 2004). Both the conserved core and nonconserved C-terminal regions of ARIDs interact with DNA, while the role of the N-terminal helix specific to proteins in subfamilies 2 and 3 remains unknown. In addition to their structural heterogeneity, ARIDs appear to adaptively recognize their binding sites, since nitrogen-15 relaxation studies of the Mrf-2 protein indicates that its mobility is quenched upon DNA binding (Zhu et al. 2001) and the structure of the Dead ringer-DNA complex has highlighted significant DNA-induced structural changes in the protein (Iwahara et al. 2002). In this paper we have used nitrogen-15 relaxation measurements to quantitatively determine how the Dead ringer ARID alters its mobility to recognize DNA. Because the Dead ringer protein belongs to the extended-ARID subfamily and it is the only ARID for which both the structures of the free and DNA complex are known, this work provides valuable insights in the structural and dynamic mechanisms

through which ARID transcription factors recognize DNA using both well-conserved and less-conserved protein components.

### Results

# Backbone <sup>15</sup>N relaxation measurements of the free and DNA-bound forms of ARID<sup>Phe355Leu</sup>

We studied the NMR relaxation properties of a single amino acid mutant of the Dead ringer ARID (ARID<sup>Phe355Leu</sup>, amino acids Gly262-Gly398 of Dead ringer with leucine substituted for phenylalanine at position 355). This mutant protein was used because the NMR structure of its complex with DNA has been determined previously (the ARID<sup>Phe355Leu</sup>–DNA complex containing ARID<sup>Phe355Leu</sup> bound to the sequence 5'-CCTGTATTGATGTGG-3'; 5'-CCACATCAATACAGG-3') (Iwahara et al. 2002). The amide  ${}^{15}N$  T<sub>1</sub>, T<sub>1</sub>, and  ${}^{1}H$ - ${}^{15}N$  NOE values of nitrogen-15 enriched ARID<sup>Phe355Leu</sup> in the DNA-free state and the ARID<sup>Phe355Leu</sup>–DNA complex were measured at 500 and 600 MHz. The spectra of ARID<sup>Phe355Leu</sup> and the ARID<sup>Phe355Leu</sup>-DNA complex are well resolved (Iwahara and Clubb 1999a), enabling the measurement of relaxation parameters for 111 and 105 residues, respectively. Relaxation data were acquired at different temperatures, 303 K and 310 K for the free ARID<sup>Phe355Leu</sup> and the ARID<sup>Phe355Leu</sup>–DNA complex, respectively. This was required because the DNA free protein is unstable at 310 K and the 26 kDa ARID<sup>Phe355Leu</sup>–DNA complex exhibits increased line broadening at 303 K.

Consistent with its larger molecular size, the protein when bound to DNA exhibits longer  $T_1$  and shorter  $T_2$  values. Figure 1 shows representative  $\{{}^{1}\text{H}-\}{}^{15}\text{N}$  NOE,  $T_1$  and  $T_2$  data for the protein backbone amide nitrogen-15 atoms in the complex as a function of residue number. The median values of the  $T_1$ ,  $T_2$ , and  $\{{}^{1}\text{H}-\}{}^{15}\text{N}$  NOEs for residues located in helices H1 to H8 in the ARID<sup>Phe355Leu</sup>–DNA complex are respectively,  $713 \pm 49$  msec,  $62.4 \pm 6.6$  msec, and  $0.75 \pm 0.09$  at 500 MHz and  $941 \pm 61$  msec,  $58.1 \pm 6.3$  msec, and  $0.79 \pm 0.07$  at 600 MHz. On the other hand, in the free protein the average values are  $612 \pm 37$  msec,  $73.7 \pm 9.1$  msec, and  $0.76 \pm 0.10$  at 500 MHz and  $800 \pm 50$  msec,  $67.7 \pm 8.5$  msec, and  $0.74 \pm 0.09$  at 600 MHz, respectively.

#### Determining the anisotropy of rotational diffusion

In order to accurately extract the order parameters and correlation times of internal motions, we first determined the rotational diffusion anisotropy of the ARID<sup>Phe355Leu</sup> and the ARID<sup>Phe355Leu</sup>–DNA complex using the  $T_1/T_2$  ratios and the NMR solution structures. This was done because erroneous values for internal dynamics parameters can be ob-



**Figure 1.** Representative <sup>15</sup>N relaxation data of the Dead ringer ARID<sup>Phe355Leu</sup>-DNA complex collected at 600 MHz. (*A*) {<sup>1</sup>H-}<sup>15</sup>N NOE, (*B*) <sup>15</sup>N R<sub>1</sub>, and (*C*) <sup>15</sup>N R<sub>2</sub> values as a function of residue number. Sample temperatures were 310 K and a schematic of the secondary structure of the protein is shown. The R<sub>2</sub> values were calculated from T<sub>1</sub> and T<sub>1p</sub> values as described in the text. Errors are based on the standard deviations of the noise level present in the NMR spectra. Residues Gly262, Ser396, Tyr397, and Gly398 are not shown in panel *A* because they have {<sup>1</sup>H-}<sup>15</sup>N NOE values less than zero.

tained even if the anisotropy is as small as  $D_{\parallel}/D_{\perp} \sim 1.2$  (Tjandra et al. 1995; Luginbuhl et al. 1997; Osborne et al. 2001). To obtain a rough estimate of the rotational diffusion anisotropy, the inertia tensors were calculated with the 20-structure ensembles of the free protein (Iwahara and Clubb 1999b) (PDB 1C20) and the ARID<sup>Phe355Leu</sup>–DNA complex (Iwahara et al. 2002) (PDB 1KQQ). The ratio of inertia moments around the principal axes of the inertia tensors ( $I_{xx}$ : $I_{yy}$ : $I_{zz}$ ) are 0.55 ± 0.04:0.84 ± 0.03:1.00 for the free pro-

tein and  $0.61 \pm 0.01:0.72 \pm 0.01:1.00$  for the complex. In both structures, two of three moments are relatively close to one another, indicating that their rotational diffusion is axially symmetric. Moreover, judging from the ratios of the inertia moments, ellipsoids of the diffusion tensors are expected to be "prolate" ( $I_{xx} < I_{yy} \approx I_{zz}$ ;  $D_{zz} > D_{yy} \approx D_{xx}$ ;  $D_{yy}$  $D_{\perp} > 1.0$ ) for the complex and "oblate" ( $I_{xx} \approx I_{yy} < I_{zz}$ ;  $D_{zz} \approx D_{yy} > D_{xx}$ ;  $D_{\parallel}/D_{\perp} < 1.0$ ) for the free protein. Using empirical relationship,  $D_{\parallel}/D_{\parallel} \approx (I_{\parallel}/I_{\parallel})^{(1/\sqrt{2})}$  (Copie et al. 1998), the diffusion anisotropy values of the free protein and the complex are estimated to be  $\sim 1.4$  and  $\sim 0.8$ , respectively. However, it should be noted that the calculated inertial axes from the static structures do not necessarily have to agree well with the principle axes of the actual rotational diffusion tensor, due to lack of information on the hydrodynamics of the macromolecule. We therefore used the experimental data to gain insights into the diffusion tensor (see below).

The diagonal elements and the principle axes of the rotational diffusion tensor were determined by nonlinear least square fitting of the experimental  $T_1/T_2$  ratios determined at 500 and 600 MHz (Tjandra et al. 1995). In this procedure, data from residues that exhibited significant exchange contribution to  $T_2$  in our subsequent analyses, and/or  $\{^{1}H-\}^{15}N$ NOE values smaller than 0.6 were excluded. In the final analysis, only ~70% of the available data was used (73 and 75 residues for the free and complexed proteins, respectively). Because the rotational diffusion anisotropy characterized by  $D_{\parallel}/D_{\parallel}$  was estimated to be relatively small (see above), the effect of noncollinearity of the <sup>15</sup>N-<sup>1</sup>H dipolar and <sup>15</sup>N CSA tensors on the relaxation data was neglected (Fushman and Cowburn 1999). As pointed out previously (Blackledge et al. 1998), the  $\chi^2$  function defined as  $\chi^2 = \Sigma [(T_1/T_2)_{calc} - (T_1/T_2)_{obs}]^2 / \sigma^2$ , could show two minima with similar magnitudes: one at  $D_{\parallel}/D_{\perp} > 1.0$  (prolate ellipsoid) and the other at  $D_{\parallel}/D_{\perp} < 1.0$  (oblate ellipsoid). Since the inertia tensors of the free and bound proteins had previously been determined from their overall molecular shapes, the nonlinear least square fitting procedure was performed constraining  $D_{\parallel}/D_{\perp}$  ( $D_{\parallel}/D_{\perp} > 1.0$  for the free protein and  $D_{\parallel}/D_{\perp}$  $D_{\perp} < 1.0$  for the complex). This fitting procedure produced  $D_{\parallel}/D_{\perp}$  values of 1.41 for the free protein and 0.84 for the complex, which are in close agreement to the values predicted using the inertia moments. Note that the  $D_{\parallel}$  symmetry axis of the rotational diffusion tensor should theoretically correspond to the X-axis of the inertia tensor for the prolate ellipsoid ( $I_{xx} < I_{yy} = I_{zz}$ ) and the Z-axis for the oblate ellipsoid model ( $I_{xx} = I_{yy} < I_{zz}$ ). In fact, despite the lack of information on the hydrodynamics of the free protein and the complex, the symmetry axes of the diffusion tensors were roughly parallel with the corresponding inertial axes calculated from the static structures; the angle differences were 19° for the free protein and 7° for the complex. Residuals of the  $\chi^2$  functions divided by the number of data

points ( $\chi^2/N$ ) were 16.7 and 3.4 for the free protein and the complex, respectively. The higher  $\chi^2/N$  value for the free protein is mainly due to the smaller experimental errors in  $T_1/T_2$  and may also be due to lower quality of the structure of the free protein; residual dipolar coupling constraints were used for the complex, but not for the free protein, and a larger number of long-range NOE restraints define the structure of the protein in the complex (Iwahara and Clubb 1999b; Iwahara et al. 2002). A fully asymmetric diffusion tensor  $(D_{zz} > D_{xx} > D_{yy})$  was also examined. Although the number of fitted parameters is larger (six as opposed to four for the axially symmetric model), calculations with the fully asymmetric diffusion tensor resulted in only a slight reduction of  $\chi^2/N$  (the values were 16.7 for free and 2.9 for the complex). This suggests that it is meaningless to use the fully asymmetric model for both systems. Thus, an axially symmetric diffusion model was employed in the subsequent model-free analysis, using  $D_{\parallel}/D_{\perp} = 1.43$  for the free protein and 0.83 for the complex.

#### Effective overall correlation times

The above analyses predicted effective correlation times of 10.4 nsec for the 16 kDa ARID<sup>Phe355Leu</sup> protein at 303 K, and 12.7 nsec for the 26 kDa ARID<sup>Phe355Leu</sup>–DNA complex at 310 K. In order to cross-validate these numbers we used nitrogen-15 relaxation data to determine the effective correlation time of the complex at the same temperature as that used for the free protein (data not shown). The sensitivity of the NMR data acquired on the complex at 303 K was significantly lower, and only 43 residues could be used for the same analysis of the rotational diffusion tensor. The complex has an effective correlation time of 14.6 nsec at 303 K and the ratio of correlation times of the complex to the free protein is 1.4 ( $\tau_{c,eff}^{complex}/\tau_{c,eff}^{free} = 14.6/10.4$ ). This value is in good agreement with the structural data, which predict a ratio of 1.36; the rotational correlation time is proportional to the cubic of the radius of gyration, which is  $16.6 (\pm 0.4)$ Å and 18.4 ( $\pm 0.2$ ) Å in the structures of the free protein and the complex, respectively. Because of the generally higher quality of the data collected at 310 K, it was used to characterize the dynamics of the complex.

#### Model-free analysis

The model-free approach was used to determine the generalized order parameter ( $S^2$ ) and internal correlation time ( $\tau_e$ ) of backbone amide motion, and the exchange–contribution ( $R_{ex}$ ) to <sup>15</sup>N transverse relaxation rates. Five model functions were used to fit the relaxation data; the functions contained either (1)  $S^2$ ; (2)  $S^2$  and  $R_{ex}$ ; (3)  $S^2$  and  $\tau_e$ ; (4)  $S^2$ ,  $S_f^2$ ,  $\tau_e$ ; or (5)  $S^2$ ,  $\tau_e$ ,  $R_{ex}$  (Clore et al. 1990; Mandel et al. 1995). For each residue the method proposed by Ishima and colleagues was used to select the best model function to fit the relaxation data acquired at the two magnetic fields (Ishima and Nagayama 1995; Ishima et al. 1995). The values of  $S^2$ ,  $\tau_e,$  and  $R_{ex}$  for individual residues in the free protein and complex are shown in Figures 2 and 3, respectively. The mean values of  $S^2$  value for residues Trp263–Glu391 is 0.87 and 0.89, for the free ARIDPhe355Leu and the ARID<sup>Phe355Leu</sup>–DNA complex, respectively. As observed in other proteins, residues within regular secondary structure exhibit larger order parameters, while residues within loops have smaller values (Goodman et al. 2000); residues in helices in the free and the bound proteins have mean  $S^2$ values 0.94 and 0.93, respectively. In both binding states, the largest degree of mobility in regular secondary structural elements is observed in the N-terminal portion of helix H1, the  $\beta$ -hairpin, and the C-terminal portion of helix H8. Large internal motions in the pico- to nanosecond timescale are also observed in all protein loops and 15 residues in the free protein and seven residues in the complex exhibit Rex values larger than 1 Hz. Of these residues, only Asn279 and Ile323 in both the free and DNA-bound states, and residue Val272 in the free protein, exhibit Rex values larger than 3 Hz. A detailed discussion of the relaxation data is provided below.

#### Discussion

The Dead ringer ARID can be envisioned as consisting of two subdomains (Fig. 4). A C-terminal helical subdomain is formed by helices H3 to H8, and interacts with the phosphodiester backbone and the major and minor grooves, while an N-terminal subdomain consisting of helices H1, H2 and a short  $\beta$ -hairpin, rotates upon DNA binding to form stabilizing interactions with the phosphodiester backbone of the duplex. In order to aid in the discussion of the dynamics data, the interface can be divided into four regions: (1) the loop connecting helices H5 and H6 (the H5/H6 loop; residues Leu344-Ser352), which is inserted into the major groove (Fig. 6A, below); (2) the  $\beta$ -hairpin region that contacts the phosphodiester backbone on the 5' side of the major groove interface (Fig. 6A, below); (3) the concave surface formed by residues within helices H4 to H6 and H8, which contacts the phosphodiester backbone on the 3' side of the major groove interface (Fig. 6B, below); and (4) the C-terminal tail that contacts the minor groove (Fig. 6B, below). A comparison of the backbone atom coordinate precision in the DNA free and bound forms of the protein suggests that the H5/H6 loop, the  $\beta$ -hairpin, and the C terminus undergo a disordered to ordered transition upon binding DNA, since the coordinates of atoms in these regions are more precisely defined in the complex (Fig. 5). However, this structural argument for changes in protein dynamics is weak, since the interfacial amino acids protrude into the solvent in the DNA free protein (e.g., the H5/H6 loop, β-hairpin, and C terminus) and are thus defined by fewer



**Figure 2.** Model-free parameters for the Dead ringer protein (ARID<sup>Phe355Leu</sup>) in the DNA-free state. (*A*) Order parameters  $S^2$  as function of residue number. (*B*) Correlation times of internal motions ( $\tau_e$ ) as a function of residue number. (*C*) Exchange contributions ( $R_{ex}$ ) to the <sup>15</sup>N relaxation data as a function of residue number. The secondary structure of ARID<sup>Phe355Leu</sup> is drawn above each panel and the standard deviations of each parameter are represented by error bars.

NOE restraints as compared to the complex. In order to determine the magnitude and timescales of motion, we therefore performed a nitrogen-15 relaxation study.

# *The flexible turn in the helix-turn-helix motif becomes partially immobilized*

Binding specificity is achieved by inserting the H5/H6 loop into the major groove. The loop can be considered to be the "turn" of a noncanonical helix-turn-helix (HTH) motif and contributes hydroxyl-bearing side chains that make basespecific hydrogen bonds. Figure 7 shows a plot of the differences between the  $S^2$  values of the DNA-free and -bound proteins. The data indicate that motions in the loop are quenched upon DNA binding, since the average value of the available order parameters in this region increases from 0.70 to 0.82 upon DNA-binding (His345, Ser349, Ile350, and Thr351 in the free protein, and Leu346, Ser348, Ser349, and Ser352 in the DNA-bound state could be measured). Because of distinct patterns of resonance overlap, the backbone amide of Ser349 is the only residue whose relaxation could be analyzed in both states and it experiences a 0.09 increase in its order parameter. The mobility of the loop in the complex is reduced because several residues form new hydrogen bonds and hydrophobic interactions with the duplex (Fig. 6A). For example, in the complex the hydrophobic side chains of Thr351 and Ile350 pack against the methyl group of Thy21 and the H8 atom of Gua9, respectively. These hydrophobic contacts position the side chains of Thr351 and Ser352 above the Ade<sup>10</sup>-Thy<sup>11</sup> base-step for a series of intermolecular hydrogen bonds. In particular, in 100% and 85% of the 20 conformers in the ensemble, the hydroxyl groups of Thr351 and Ser352 accept a hydrogen bond from the Ade20 N6 amine, respectively. Further anchoring the loop in the groove are an array of less frequently observed interfacial direct and water mediated hydrogen bonds from the hydroxyls of Thr351 and Ser352 to bases near Ade20 (Iwahara et al. 2002).

## *Micro- to millisecond timescale motions foreshadow a DNA-dependent conformational change*

A conformational change in the protein is coupled to DNA binding and helps to immobilize the H5/H6 loop. As shown in Figure 4B, DNA binding drives the rotation of the N-terminal subdomain (composed of helices H1 and H2 and the  $\beta$ -hairpin) relative to the remainder of the polypeptide. This results in an ~10 Å shift of the hairpin towards the H5/H6 loop enabling the side chains of Met308 and Ala309 within the hairpin to insert into the hydrophobic pocket formed by the side chains of Thr356 and Ile357 within helix H6 and Pro347 and Ile350 within the H5/H6 loop (Fig. 6A). Interestingly, in addition to stabilizing the H5/H6 loop, DNA-binding induced rotation also immobilizes residues within the  $\beta$ -hairpin, since in the free protein the average  $S^2$  for these residues is 0.80, which increases to 0.86 upon complex formation. The most pronounced DNA-binding



**Figure 3.** Model-free parameters for the Dead ringer ARID<sup>Phe355Leu</sup>-DNA complex. (*A*) Order parameters  $S^2$  as function of residue number. (*B*) Correlation times of internal motions ( $\tau_e$ ) as a function of residue number. (*C*) Exchange contributions ( $R_{ex}$ ) to the <sup>15</sup>N relaxation data as a function of residue number. The secondary structure of ARID<sup>Phe355Leu</sup> is drawn *above* each panel and the standard deviations of each parameter are represented by error bars.

dependent increases in  $S^2$  occur in residues Lys310 and Ser311, which increase by 0.09 (Fig. 7). These changes would seem to result from contacts to the 5' portion of the

major groove interface, since the backbone amide atoms of Ala309 and Lys310 donate hydrogen bonds to the phosphate oxygens of Thy408 in over 90% of the conformers and salt-bridges to the adjacent phosphate groups are formed through the side chains of Arg304 and Lys310.

An inspection of the R<sub>ex</sub> data reveals that the DNA-free ARID is more dynamic on the micro- to millisecond timescale as compared to the DNA-bound protein (cf. Figs. 2C, 3C). Excluding the termini, most of the residues in the free protein exhibiting dynamics on this timescale either reside in the N-terminal subdomain (helices H1, H2, or the  $\beta$ -hairpin) or the surface on the C-terminal domain contacted by the N-terminal subdomain (helix H3) (Fig. 2C). This suggests that slow conformational dynamics that occur prior to binding may facilitate the DNA dependent rotation of the N-terminal subdomain. It is unclear why residues within the H5/H6 do not exhibit micro- to millisecond time scale motions, since this region of the protein forms a contact surface with β-hairpin after rotation of the N-terminal domain. Interestingly, motions in the H5/H6 loop slow down upon binding DNA, increasing from ~300 psec to 1.1 nsec. Although we did not directly measure Rex, it should be mentioned that slow dynamics on the micro-to-millisecond time scale can be characterized more accurately by applying experiments that can directly determine the Rex values (Kroenke et al. 1998; Loria et al. 1999).

### The C-terminal tail is partially immobilized in the complex

An analysis of the relaxation data indicates that the C-terminal tail is partially immobilized in the complex (Fig. 7). Residues in the C terminus are very mobile on the picosecond timescale in the absence of DNA. For example, residues Asn388, Arg389, and Arg390 have  $S^2$  values 0.55, 0.33, and 0.26, respectively. These residues are only partially immobilized in the complex, since the  $S^2$  values of Asn388 and Arg389 in the complex are 0.70 and 0.52, respectively (Fig. 3A). The decrease in mobility can be attributed to the formation of a salt-bridge between the side chain of Arg389 and the phosphate of Thy21, which helps to position the side chain of Arg390 for an array of contacts to the minor groove (Fig. 6B). Arg390 is poised for a variety of favorable contacts to the  $\mathrm{O}_2$  atoms of Thy13 and Cyt19, and the adjacent phosphodiester backbone. Further limiting the mobility of the tail are hydrogen-bonding contacts between Arg358 and Asn388, which link the HTH unit to helix H8. These residues also favorably contact the phosphate group of Ade20. In contrast to the C-terminal tail, the pocket formed by residues in helices H4 to H6 and H8 appears to act as a rigid scaffold to engulf the phosphodiester backbone of nucleotides Ade18 to Thy21 (Fig. 6B), since their order parameters change only slightly. Interestingly, nitrogen-15 relaxation studies of the free and bound forms of the ARID



**Figure 4.** Structure of the ARID<sup>Phe355Leu</sup>-DNA complex. (*A*) Ribbon schematic of the Dead ringer ARID<sup>Phe355Leu</sup>-DNA complex. The figure displays the lowest energy structure (PDB 1KQQ) (Iwahara et al. 2002). The N and C termini of the protein and the 5' and 3' ends of the DNA are labeled. (*B*) Overlay of the backbone structure of ARID<sup>Phe355Leu</sup> in the DNA-free and -bound states (Iwahara and Clubb 1999b; Iwahara et al. 2002). The overlay was constructed by superimposing helices H3 to H8 of each structure (rmsd for the backbone atoms of residues 315–388 is 1.5 Å). For clarity, helices H1 and H2 are represented as cylinders and the  $\beta$ -hairpin is represented with arrows. DNA binding causes a large rotation of the two subdomains with respect to one another; in the energy minimized structures, helices H1 and H2 of the N-terminal subdomain are shifted by ~31° and ~17° relative to the C-terminal subdomain upon binding DNA, respectively.

from the Mrf protein reveal that the loop analogous to the H5/H6 loop in Dead ringer is immobilized in its complex with DNA (Zhu et al. 2001). Moreover, Mrf replaces helix H8 in Dead ringer with a disordered C-terminal that is also substantially immobilized. Although the structure of the Mrf–DNA complex has not yet been determined, these data suggest that all members of the structurally diverse ARID

family undergo significant changes in their mobility upon binding DNA.

#### The interface retains flexibility after binding DNA

Although residues in the interfacial H5/H6 loop,  $\beta$ -hairpin, and C terminus are immobilized upon binding, in the complex they remain more mobile than residues located in regions of regular secondary structure (Fig. 3A). Elevated mobility within the major groove is compatible with the structure, since the H5/H6 loop-DNA interaction surface is largely exposed to the solvent and thus presumably not restricted in its motion. This suggests that in addition to the primary hydrogen bonding contacts to the Ade20 N6 amine, flexibility may enable the formation of additional hydrogen bonds to nearby bases from the hydroxyls of Thr351 and Ser352; these contacts are present in the ensemble of structures, but occur less frequently than contacts to the Ade20 N6 amine. The C-terminal tail is especially mobile in the complex, with the DNA contacting residues of Asn388 and Arg389 exhibiting order parameters of 0.55 and 0.33, respectively. Compatible with its elevated mobility, few NOEs define the positioning of these residues and the structure of the tail is poorly defined (Fig. 5).

# The effect of the Phe355Leu mutation on protein dynamics

Previously we demonstrated that the NMR spectrum of the DNA in the ARID<sup>Phe355Leu</sup>–DNA complex has narrower line widths as compared to the DNA spectrum of the wild-type protein–DNA complex, a phenomenon that was critical



**Figure 5.** Coordinate precision in the DNA and free bound forms of the Dead ringer protein. Plot of the root mean square deviation (rmsd) of the backbone atom coordinates in the ensemble of conformers to the mean structure as a function of residues number for the DNA free (PDB 1C20; Iwahara and Clubb 1999b) (open circles) and bound structures (PDB 1KQQ; Iwahara et al. 2002) (black diamonds). A schematic of the secondary structure of the protein is drawn *above* the figure.



**Figure 6.** Isolated views of the protein–DNA interface. (*A*) View of the major groove and adjacent  $\beta$ -hairpin interface. Contacts between residues in the H5/H6 loop and  $\beta$ -hairpin are indicated. (*B*) View of the minor groove interface containing residues from the C-terminal tail. The figure also shows residues in helices H4 to H6 and H8, which form a pocket that contacts the phosphodiester backbone. Details of the contacts in the complex have been described elsewhere (Iwahara et al. 2002). In both panels the surface of the DNA is displayed with the oxygen atoms of the backbone shaded light gray.

for solving the NMR structure of the complex. In order to gain insights into the beneficial properties of the Phe355Leu mutation we measured the <sup>15</sup>N T<sub>1</sub>, T<sub>2</sub>, and {<sup>1</sup>H-}<sup>15</sup>N NOEs of the wild-type ARID–DNA complex and interpreted these data using the model free formalism. A comparison of the T<sub>2</sub> values of the wild-type and ARID<sup>Phe355Leu</sup>–DNA complexes reveals no major differences, consistent with our previous finding that the mutation primarily affects the linewidths of the DNA. Intriguingly, a comparison of the S<sup>2</sup> parameters of the mutant and wild-type complexes reveals that the backbone amide groups at the C-terminal interface are more rigid in the wild-type complex (Asn388, 0.85 [wild-type] and 0.70 [mutant]; Arg389, 0.73 [wild-type] and 0.52 [mutant]). This finding is compatible with the three-dimensional structure of the Dead ringer ARID<sup>Phe355Leu</sup>–

DNA complex, since the side chain of Leu355 makes hydrophobic contacts with a thymine base near the C-terminal interface (Iwahara et al. 2002). It is possible that the mutation introduces unfavorable contacts at this thymine, which makes the C-tail less stable; a conclusion that is consistent with the lower affinity of the mutant protein for the target site as compared to the wild-type protein (Iwahara et al. 2002).

#### DNA binding increases the mobility of helix H1

As shown in Figure 7, the  $S^2$  parameters of the N-terminal portion of the helix H1 (residues Phe265, Glu266, Glu267, Lys270) decrease upon DNA-binding, indicating that it becomes more mobile in the DNA-bound state on the picosecond timescale. In principle, the decrease might be caused by the higher temperature used to collect the relaxation data of the complex (310 K for the complex; 303 K for the free protein). However, this seems unlikely, since judging from the recent publications on the temperature-dependence of  $S^2$ (Mandel et al. 1996), the temperature difference of 7 K can not fully account for the observed  $S^2$  changes that are as large as  $\sim 0.2$  in this region (Fig. 7). These changes likely result from the conformational change in the protein that occurs when it binds DNA (Fig. 4B) (Iwahara et al. 2002). Presumably, the DNA-induced rotation and displacement of the N-terminal subdomain breaks protein-protein contacts between residues in helix H1 and the body of the protein, thereby increasing helix H1 mobility. This would seem to have a favorable effect of the energetics of binding, partially negating the entropic cost of immobilizing residues at the interface.



**Figure 7.** Plot of the differences between the  $S^2$  values of the DNA-free and -bound forms of the protein.  $S^2$  (Dead ringer ARID<sup>Phe355Leu</sup>-DNA complex)- $S^2$  (free Dead ringer ARID<sup>Phe355Leu</sup>) is plotted as a function of residue number. A schematic of the secondary structure of the protein is drawn *above* the figure.

It is also conceivable that helix H1 serves a regulatory role in DNA binding that is analogous to the inhibitory helix of the Ets DNA-binding domain. In Ets, the energetic cost associated with unfolding its inhibitory helix reduces the affinity of the domain for DNA (Petersen et al. 1995). Interestingly, the helix serves a regulatory role, since binding of GABPB displaces it from an inhibitory position, thereby increasing the binding affinity of Ets (Batchelor et al. 1998). Helix H1 may play a similar role, since modeling studies using the structure of the protein solved in the absence of DNA indicate that the N-terminal portion of helix H1 must be moved through the aforementioned conformational change in order for Dead ringer to bind B-form DNA. Helices analogous to H1 are found in the extended and Nterminally extended members of the ARID family, and similar to the Ets domain, could potentially be sites of interaction with other proteins that stabilize protein conformations better suited for DNA binding.

#### Material and methods

### Backbone <sup>15</sup>N relaxation measurements

The NMR samples of free proteins of <sup>15</sup>N-labeled Dead ringer ARID<sup>Phe355Leu</sup> and its complex with the 15-bp DNA fragment (dCCTGTATTGATGTGG/dCCACATCAATACAGG; unlabeled) were prepared in the same method as described previously (Iwahara and Clubb 1999a,b; Iwahara et al. 2002). The NMR buffer contained 20 mM Tris HCl (pH 6.7), 2 mM DTT, 0.5 mM EDTA, 0.01% NaN<sub>3</sub>, 7% D<sub>2</sub>O. For the sample of the free protein, 100 mM NaCl was also included for the sample stability. The sample concentrations of the free protein and the protein–DNA complex were 1.5 mM. The DNA fragment in the complex was unlabeled.

NMR measurements were performed with Bruker DRX-500 and DRX-600 spectrometers operated at <sup>1</sup>H frequencies of 500.09 MHz and 600.13 MHz, respectively. Backbone amide <sup>15</sup>N T<sub>1</sub>, T<sub>2</sub>, and {<sup>1</sup>H-}<sup>15</sup>N NOE were measured at both magnetic fields. The temperatures during NMR measurements were 303 K for the free protein and 310 K for the complex. The temperature at the probe head was determined using the temperature-dependence of the difference between the methyl and hydroxyl <sup>1</sup>H chemical shifts of methanol. Backbone <sup>15</sup>N T<sub>1</sub> relaxation times were measured with 2D 1H-15N experiments according to Farrow and colleagues (Farrow et al. 1994). To gain sensitivity and better resolution, the "semi-constant time" technique (Grzesiek and Bax 1993) was implemented in the <sup>15</sup>N dimension. Typically, seven time points in the range of 0.01–1.8 sec for T<sub>1</sub> relaxation were recorded. Because it has been noted that off-resonance effects in CPMG experiments can introduce errors in the measured T<sub>2</sub> values (Ross et al. 1997; Korzhnev et al. 2000), the T<sub>2</sub> relaxation times were calculated using the equation

$$T_{10}^{-1} = T_1^{-1} \cos^2 \lambda + T_2^{-1} \sin^2 \lambda,$$

where the tilt angle  $\lambda$  of the effective field is calculated for each resonance based on the r.f. strength of the spin-lock and <sup>15</sup>N chemical shift (Davis et al. 1994; Tjandra et al. 1996b). T<sub>1p</sub> relaxation times under a 2 kHz <sup>15</sup>N spin-lock field were measured with 2D <sup>1</sup>H-<sup>15</sup>N experiments (Palmer 2001; Peng et al. 2001).

During the <sup>15</sup>N-spin-lock, <sup>1</sup>H 180 pulses were applied with an interval of 4 msec to prevent cross-correlation between dipolarinteraction and CSA relaxation mechanisms (Kay et al. 1992; Palmer et al. 1992; Palmer 2001). Typically, seven different spinlock lengths from 4 msec to 152 msec were used.  $T_1$  and  $T_{10}$ relaxation times were determined with two-parameter fitting to an exponential function using peak heights. Errors in the  $T_1$  and  $T_{10}$ values were estimated from the noise levels in the spectra {1H-}15N NOE values were measured with the water-flip back 2D <sup>1</sup>H-<sup>15</sup>N experiment (Grzesiek and Bax 1993). The errors of the NOE values were estimated from the noise levels of the spectra and were set to no less than 0.08 because of the relatively low accuracy of the standard {1H-}15N NOE measurements (Idiyatullin et al. 2001). Spectra were processed using NMRPipe. The program NMRView (version 4.1.3) was utilized for peak intensity quantification and to calculate the  $T_1$  and  $T_{1\rho}$  values.

### Calculation of inertia tensor and radius of gyration of the free protein and the complex

The inertia tensor and radius of gyration of the free protein and the complex were calculated from the structures of the free protein and protein–DNA complex (Iwahara and Clubb 1999b; Iwahara et al. 2002). The inertia tensors were calculated according to Browner et al. (1992). The inertia tensor  $\mathbf{I}$  is defined as:

$$I_{xx} = \sum m_i (r_i^2 - x_i^2), I_{yy} = \sum m_i (r_i^2 - y_i^2), I_{zz} = \sum m_i (r_i^2 - z_i^2)$$
  
$$I_{xy} = I_{yx} = \sum m_i x_i y_i, I_{xz} = I_{zx} = \sum m_i x_i z_i, I_{yz} = I_{zy} = \sum m_i y_i z_i$$

where  $r_i$  is the position vector of the *i* th atom from the center of the mass of the molecule;  $x_i$ ,  $y_i$ ,  $z_i$ , are the x, y, and z components of the vector  $r_i$ ;  $m_i$  is the mass of the *i* th atom. The eigenvectors and eigenvalues of **I** were calculated, which correspond to principle axis of inertia and inertial moments around them, respectively. The mass of each nitrogen atom in the protein was assumed to be 15, since the NMR samples were labeled with <sup>15</sup>N. The radius of gyration of the free protein and the complex were calculated as an atomic RMS difference between the centroid and individual atoms (Kuszewski et al. 1999):

$$R_{gyration} = \sqrt{\sum |\mathbf{r}_i - \mathbf{r}_c|^2 / N}$$

where  $\mathbf{r}_{c}$  is the position vector of the centroid and N is the total number of atoms. Disordered regions at the N and C termini were included in the calculation, because they also affect the inertia moments and the radius of gyration (Clubb et al. 1995).

#### Calculation of rotational diffusion tensor

Calculations of the rotational anisotropy were performed by nonlinear least square fitting of the  $\chi^2$  function defined as:

$$\chi^{2} = \sum \frac{\{(T_{1}/T_{2})_{calc} - (T_{1}/T_{2})_{obs}\}^{2}}{\sigma^{2}}$$

where the  $\sigma$  values are experimental errors in  $T_1/T_2$ . Two anisotropy models were used, full asymmetric diffusion ( $D_{zz} > D_{xx} > D_{yy}$ ) and axially symmetric diffusion ( $D_{zz} > D_{xx} = D_{yy}$  or  $D_{zz} = D_{xx} > D_{yy}$ ). Nonlinear least square fitting was performed using the Levenberg-Marquardt algorithm. The Jacobian matrix in the fitting procedure was calculated with numerical differentiation rather than the derivative of the fitting function. In-house programs were used to calculate the inertia tensor, radius of gyration, and the diffusion tensor.

# Calculation of order parameters, internal correlation times, and exchange contributions

The program Modelfree version 4.0 (A.G. Palmer 3rd, Columbia University) was used to calculate the order parameters  $S^2$ , the internal correlation times  $\tau_e$ , and the exchange contributions to the transverse relaxation rates. Calculations were performed assuming axially symmetric rotational diffusion with  $D_{\parallel}/D_{\perp}$ , the effective correlation time  $\tau_{c,eff}$ , and the polar angles about the symmetric axis obtained using the methods described above. Five model functions were used: (1)  $S^2$ ; (2)  $S^2$  and  $R_{ex}$ ; (3)  $S^2$  and  $\tau_e$ ; (4)  $S^2$ ,  $S_f^2$ ,  $\tau_e$ ; and (5)  $S^2$ ,  $\tau_e$ ,  $R_{ex}$  (Clore et al. 1990; Mandel et al. 1995). Model selection was based on an analysis of the J( $\omega$ ) profiles according to Ishima et al. (Ishima and Nagayama 1995). Uniform <sup>15</sup>N CSA values of -172 ppm and  $r_{NH} = 1.02$  Å were assumed (Tjandra et al. 1996a; Palmer 2001).

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