# **Ionic basis of the caesium-induced depolarisation in rat supraoptic nucleus neurones**

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- 1. The effects of external  $Cs<sup>+</sup>$  on magnocellular neurosecretory cells were studied during intracellular recordings from 93 supraoptic nucleus neurones in superfused explants of rat hypothalamus.
- 2. Bath application of  $3-5$  mM Cs<sup>+</sup> provoked reversible membrane depolarisation and increased firing rate in all of the neurones tested. Voltage–current analysis revealed an increase in membrane resistance between  $-120$  and  $-55$  mV. The increase in resistance was greater below  $-85$  mV than at more positive potentials.
- 3. Voltage-clamp analysis showed that external  $Cs<sup>+</sup>$  blocked the hyperpolarisation-activated inward current,  $I_{\rm H}$ . Under current clamp, application of ZD 7288, a selective blocker of  $I_{\rm H}$ , caused an increase in membrane resistance at voltages  $\leq -65$  mV. Voltage–current analysis further revealed that blockade of  $I_{\rm H}$  caused hyperpolarisation when the initial voltage was  $<-60$  mV but had no effect at more positive values.
- 4. Current- and voltage-clamp analysis of the effects of  $Cs<sup>+</sup>$  in the presence of ZD 7288, or ZD 7288 and tetraethyl ammonium (TEA), revealed an increase in membrane resistance throughout the range of voltages tested  $(-120 \text{ to } -45 \text{ mV})$ . The current blocked by Cs<sup>+</sup> in the absence of  $I_{\rm H}$  was essentially voltage independent and reversed at  $-100$  mV. The reversal potential shifted by  $+22.7$  mV when external [K<sup>+</sup>] was increased from 3 to 9 mM. We conclude that, in addition to blocking  $I_{\rm H}$ , external Cs<sup>+</sup> blocks a leakage K<sup>+</sup> current that contributes significantly to the resting potential of rat magnocellular neurosecretory cells.

Hypothalamic magnocellular neurosecretory cells (MNCs) are responsible for the release of either vasopressin or oxytocin into the blood (Poulain & Wakerley, 1982). Following their synthesis in MNC somata, these peptides are packaged in vesicles and transported to axon terminals in the neurohypophysis (Brownstein *et al.* 1980) where secretion is triggered by the arrival of action potentials (Dreifuss *et al.* 1971). Previous work has established that different rates and patterns of firing evoked by stimuli affecting MNCs have a profound impact on excitation–secretion coupling (e.g. Dutton & Dyball, 1979; Bicknell & Leng, 1981; Bicknell, 1988). Changes in firing rate and pattern, therefore, are both important features of the response of these neurones to physiological and pathological conditions. In both types of MNC, action potentials are initiated at the soma as a result of interactions between afferent synaptic signals and intrinsic membrane properties (Renaud & Bourque, 1991). Since the integrative properties of the membrane are governed by the complement, density and distribution of ion channels within the somato-dendritic compartment, increasing attention is being placed on the identification

and characterisation of the membrane channels expressed in MNCs (for review see Hatton & Li, 1998).

Experiments on hypothalamic slices have previously shown the presence of the hyperpolarisation-activated inward current  $(I_H)$  in guinea-pig MNCs (Erickson *et al.* 1993). Although this study revealed an involvement of  $I_{\rm H}$  in the control of burst firing, the presence and significance of  $I_H$  in rat MNCs remained somewhat controversial for two primary reasons. First, experiments using rat hypothalamic explants showed that application of external  $\text{Cs}^+$ , a well-known blocker of  $I_{\text{H}}$  (Halliwell & Adams, 1982; Pape, 1996), depolarises MNCs (Stern & Armstrong, 1997; Ghamari-Langroudi & Bourque, 1998). Since blockade of active  $I_{\rm H}$  normally leads to hyperpolarisation (e.g. Maccaferri & McBain, 1996), the occurrence of a depolarising effect of  $Cs<sup>+</sup>$  in rat MNCs seemed inconsistent with a presumed presence of  $I_{\text{H}}$ . Second, the hallmark features of  $I_{\rm H}$  are (i) that it activates slowly during hyperpolarising steps (hundreds of milliseconds) and (ii) that its amplitude increases during steps to more negative potentials (Halliwell & Adams,

1982). Thus, in current-clamped cells expressing  $I_{\rm H}$ , voltage responses to hyperpolarising current pulses typically feature slow depolarising sags whose amplitudes increase with hyperpolarisation (McCormick & Pape, 1990). However, experiments on rat MNCs held at voltages near threshold have revealed depolarising sags whose amplitudes first increase and then decrease in response to hyperpolarising pulses of increasing amplitude (e.g. Stern & Armstrong, 1996, 1997). Despite these seemingly paradoxical observations, we revealed recently that rat MNCs do express a significant density of  $I_{\rm H}$  and that this current plays an excitatory role in the regulation of electrical activity (Ghamari-Langroudi & Bourque, 2000). Our analysis further suggested that the unusual behaviour of depolarising sags in rat MNCs held near threshold is probably due to the fact that the sags reflect not only the progressive activation of  $I_{\rm H}$ , but also the deactivation of the  $K^+$  current responsible for sustained outward rectification (Stern & Armstrong, 1997).

Although these observations have highlighted the complex interactions that result from the presence of overlapping voltage-sensitive currents at sub- and nearthreshold potentials, the basis for the depolarising effects of Cs+ in rat MNCs remains unknown. Moreover, the robust depolarising effect of external  $Cs<sup>+</sup>$  suggests that the presence of a conductance distinct from  $I_{\rm H}$  can also significantly modulate the electrical activity of these neurosecretory neurones. In this study, therefore, we investigated the ionic basis for the depolarising effects of external  $Cs<sup>+</sup>$  in rat MNCs. Our results indicate that  $Cs<sup>+</sup>$ blocks both  $I_{\rm H}$  and a leakage  $K^+$  current that contributes significantly to the resting potential.

## **METHODS**

#### **Preparation of superfused explants**

Hypothalamic explants were prepared as described previously (Ghamari-Langroudi & Bourque, 1998, 2000). Briefly, male Long–Evans rats  $(150-300 \text{ g})$  were briefly  $(5-10 \text{ s})$  restrained in a soft, disposable, plastic cone (Harvard Apparatus Canada, Saint-Laurent, QC, Canada) and killed by decapitation using a small rodent guillotine (model 51330; Stoelting Company, Wood Dale, IL, USA). This tissue-harvesting protocol was approved by the Animal Care Committee of McGill University. The brain was then rapidly removed from the cranial vault. A block of tissue  $({\sim}8~\text{mm} \times 8~\text{mm} \times$ 2 mm) comprising the basal hypothalamus was excised using razor blades and pinned, ventral side up, to the slanted  $(\sim 30 \text{ deg})$  Sylgard base of a temperature-controlled (32–34 °C) superfusion chamber. Within 2–3 min of decapitation, explants were being superfused  $(0.5-1 \text{ ml min}^{-1})$  with an oxygenated  $(95\% \text{ O}_2-5\% \text{ CO}_2)$  artificial cerebrospinal fluid (ACSF; see below) delivered via a Tygon tube placed over the medial tuberal region. The arachnoid membranes covering the ventral surface of the supraoptic nucleus were removed using fine forceps and a cotton wick was placed at the rostral tip of the explant to facilitate drainage of ACSF.

## **Solutions and drugs**

The ACSF (pH 7.4; 295  $\pm$  1 mosmol kg<sup>-1</sup>) was composed of (mM): NaCl, 121; MgCl<sub>2</sub>, 1.3; KCl, 3; NaHCO<sub>3</sub>, 26; glucose, 10; CaCl<sub>2</sub>, 2.5 (all from Fisher Scientific Company, Pittsburgh, PA, USA). The ACSF was supplemented, where indicated in the text, with  $0.3-0.6 \mu M$  tetrodotoxin (TTX; Sigma Chemical Co., St Louis, MO, USA). The effects of  $Cs^+$  were examined by diluting a 1 M stock of CsCl (in  $H_2O$ ) into ACSF stored in a secondary reservoir (50 ml), and by switching the supply to the delivery tube between the control and secondary reservoir. The  $I_{\rm H}$  blocker ZD 7288 (from Tocris Cookson Inc., Ballwin,  $MO$ , USA) was prepared as a 30 mM stock solution (in  $H_2O$ ) and stored at 4 °C. The effects of ZD 7288 were examined by bath application of ACSF containing a dilution of the stock solution as described above. In experiments examining the effects of  $Cs<sup>+</sup>$  in the presence of different concentrations of extracellular K<sup>+</sup>, several accessory reservoirs were used to allow the control and  $\text{Cs}^+$ -containing solutions to carry the same concentration of K+ .

#### **Electrophysiology**

Intracellular recordings were obtained using sharp micropipettes prepared from glass capillary tubes (1.2 mm o.d.; A. M. Systems Inc., Everett, WA, USA) pulled on a P87 Flaming-Brown puller (Sutter Instruments Co., Novato, CA, USA). Pipettes were filled with 2 M potassium acetate, yielding a DC resistance of 70–150 MΩ relative to a Ag–AgCl wire electrode immersed in ACSF. Recordings of membrane voltage (DC to 5 kHz) and current (DC to 0.3 kHz) were obtained through an Axoclamp 2A amplifier (Axon Instruments Inc., Foster City, CA, USA). Voltage recordings were performed in continuous current-clamp ('bridge') mode whereas current recordings were performed using the discontinuous single-electrode voltageclamp (dSEVC) mode. Switching frequencies in dSEVC mode were adjusted (2–3.5 kHz) to ensure that a complete decay of the electrode potential was achieved between periods of current injection. Signals acquired during each experiment were displayed on a chart recorder and digitised (44 kHz; Neurodata Instruments Co., Delaware Water Gap, PA, USA) for back-up storage onto videotape. Current and voltage pulses were delivered through an external pulse generator, or via a Digidata 1200-B interface driven by Clampex 8.0 software (Axon Instruments Inc.) running on a Pentium III computer. All signals were digitised online at 10 kHz and stored on the computer's hard drive. Averaging of current traces and digital subtraction were performed offline using Clampfit 8.0 software.

#### **Statistics**

Throughout the paper, group data are reported as means plus or minus the standard error of the mean  $(+s.\text{E.M.})$ . Differences between mean values recorded under control and test conditions were evaluated using Student's paired *t* test and differences were considered significant when  $P < 0.05$ .

#### Analysis of  $I_{\text{H}}$  and  $G_{\text{H}}$  (hyperpolarisation-activated **conductance)**

The properties of  $I_{\rm H}$  evoked by steps to different voltages were derived from an analysis of the Cs<sup>+</sup>-sensitive, time-dependent currents revealed by subtracting current traces recorded in the presence of Cs+ from traces recorded in control conditions. The timedependent current was fitted using a monoexponential function using Clampfit 8 software. The amplitude of the evoked  $I_{\rm H}$  was defined as the difference between starting and steady-state current values and the time constant of activation was derived directly from the best fit through the data points.

The voltage dependency of  $G<sub>H</sub>$  can be assessed experimentally by plotting the relative amplitude of current tails evoked at a fixed potential following pulses delivered to different conditioning potentials (e.g. McCormick & Pape, 1990). In MNCs, however, relatively large currents are activated or deactivated upon the termination of voltage steps to negative potentials. In particular, the amplitude of the transient  $K^+$  current measured at  $-50$  mV following a prepulse to  $-120$  mV can exceed 1.5 nA (Bourque, 1988), a value about one hundred times greater than that of the  $I_{\rm H}$  tail expected to result from the same protocol. We therefore estimated  $G<sub>H</sub>$  as done previously

(Ghamari-Langroudi & Bourque, 2000), from the amplitude of  $I_{\rm H}$ measured at various test voltages *(V)* divided by the driving force  $(V - E_H)$ , where  $E_H$  is the reversal potential of  $I_H$ . Moreover, since we could not measure  $E_{\rm H}$  directly, the value was arbitrarily set at  $-35$  mV, a value reflecting the median  $E_{\rm H}$  reported during sharp electrode voltage-clamp studies in a variety of cell types (Pape, 1996). Normalised values of  $G_H$  (% $G_H$ ) were calculated as: % $G_H = 100 \times$  $G_{\text{H}(V)}/G_{\text{H(max)}}$ , where  $G_{\text{H}(V)}$  is the value of  $G_{\text{H}}$  at voltage *V*, and where  $G_{\text{H(max)}}$  is defined as the value of  $G_{\text{H}}$  at  $-120$  mV. The data were fitted using a Boltzmann equation:

$$
\% G_{\rm H} = 100/(1 + \exp(V - V_{\nu_{\rm B}}/k)),
$$

using Sigmaplot 5 software (Jandel Scientific, San Rafael, CA, USA), where  $V_{\psi}$  is the half-maximal voltage and *k* is the slope factor characterising the relationship.

#### **RESULTS**

The data presented below were obtained during intracellular recordings made from 93 supraoptic nucleus neurones impaled with sharp microelectrodes in superfused explants of rat hypothalamus. These cells had resting membrane potentials more negative than  $-50$  mV, input resistances exceeding 150  $\text{M}\Omega$ , and fired action potentials whose amplitudes were greater than 60 mV when measured from baseline. Each of these cells also displayed frequency-dependent spike broadening (Andrew & Dudek, 1985; Bourque & Renaud, 1985) and transient outward rectification (Bourque, 1988) when examined from initial membrane potentials below, i.e. negative to,  $-75$  mV. These combined characteristics have been shown to be specific to hypothalamic magnocellular neurosecretory neurones, but not to neighbouring nonneuroendocrine cells, during intracellular recordings *in*

## **Figure 1. Effects of extracellular Cs+ on membrane potential and firing rate**

*A*, chart recording of a spontaneously active supraoptic MNC impaled in a superfused explant of rat hypothalamus. Bath application of  $3 \text{ mm Cs}^+$  (bar) provoked a reversible membrane depolarisation accompanied by an increase in firing rate. *B*, chart recordings from another supraoptic neurone in which the initial membrane potential was adjusted to  $-61$  mV by continuous injection of hyperpolarising current  $(-30 \text{ pA})$ . Bath application of 3 mM Cs<sup>+</sup> in ACSF (bar; upper trace) provoked a reversible membrane depolarisation and the appearance of action potential discharge. Action potentials are truncated in this panel. The lower trace shows the response of the same cell in ACSF containing  $0.6 \mu$ M tetrodotoxin (TTX) to block Na<sup>+</sup> -dependent action potentials.

*vitro* (Renaud & Bourque, 1991; Tasker & Dudek, 1991) and *in vivo* (Bourque & Renaud, 1991; Dyball *et al.* 1991).

## **Effects of external Cs+ on membrane potential and spike discharge**

Previous studies have shown that during current-clamp recordings from MNCs held at membrane potentials below the threshold for action potential discharge, bath application of  $2-5 \text{ mM Cs}^+$  causes a steady-state membrane depolarisation (Stern & Armstrong, 1997; Ghamari-Langroudi & Bourque, 1998). We therefore investigated whether external  $Cs<sup>+</sup>$  could significantly affect the firing of action potentials. When tested on 15 spontaneously active MNCs, bath application of  $3-5$  mM external  $Cs<sup>+</sup>$ caused a consistent and reversible membrane depolarisation with a mean amplitude of  $5.6 + 0.6$  mV (e.g. Fig. 1A). This effect was accompanied by a 5.4-fold increase in the mean  $(\pm$ S.E.M.) frequency of firing from 0.34  $\pm$  0.07 Hz in the control to  $1.84 \pm 0.35$  Hz in the presence of  $Cs^{+}$  ( $P < 0.05$ ). As illustrated in Fig. 1*B*, bath application of  $Cs^+(3-5 \text{ mM})$ to silent MNCs (membrane potential  $(V_m)$  between  $-55$ and  $-70$  mV) also consistently and reversibly depolarised all of the cells tested  $(6.1 \pm 0.4 \text{ mV}, n = 39)$ . In 27 (69%) of these cells the Cs+ -induced membrane depolarisation reached the threshold for action potential discharge and the mean peak firing rate achieved in the presence of  $Cs<sup>+</sup>$ was  $1.3 \pm 0.2$  Hz. In six MNCs the depolarising effects of  $Cs<sup>+</sup>$  were examined both in the control solution and in ACSF containing TTX  $(0.3-0.6 \mu)$  to block Na<sup>+</sup>dependent action potentials. As illustrated in Fig. 1*B*, the depolarising effects of  $Cs<sup>+</sup>$  persisted in the presence of TTX. Moreover, response amplitudes were not significantly



different in the two conditions ( $P < 0.05$ ; paired *t* test), indicating that the Cs+ -evoked depolarisation resulted from an effect on the postsynaptic cell membrane rather than from a presynaptic action.

## **Effects of Cs+ on steady-state voltage–current properties**

As suggested, the intrinsic membrane properties of MNCs appear to be the target of action of  $Cs<sup>+</sup>$ . The possible effects of  $Cs<sup>+</sup>$  on membrane conductance, therefore, were examined by studying voltage–current *(V–I)* relationships in control and in the presence of  $3-5$  mm  $Cs<sup>+</sup>$  in five MNCs. For this purpose, the initial membrane potential of each cell was current clamped 5–15 mV below the threshold for spike discharge and steady-state *V*–*I* relationships were obtained by injecting square current pulses  $(1-3 s)$  of varying amplitude  $(+50 t - 400 pA)$  at intervals  $\geq 8$  s (e.g. Fig. 2*A*). The absolute voltage at the end of each pulse was then measured and plotted as a function of the corresponding absolute current. As illustrated in Fig. 2*B*, steady-state  $V-I$  relationships recorded in the presence of  $Cs<sup>+</sup>$  exhibited an increase in slope compared to the control, and a reversal potential near  $-85$  mV. The increase in slope (i.e. membrane resistance) suggested that the main effect of  $Cs<sup>+</sup>$  was to inhibit one or more of the conductances that were active under steady-state conditions. Since changes in steadystate voltage provoked by  $Cs^+(\Delta V_{Cs})$  at any initial voltage

 $(V<sub>i</sub>)$  partly reflected the magnitude of the  $Cs<sup>+</sup>$ -sensitive conductance at that particular voltage, a qualitative description of the voltage dependency of the affected currents was derived by examining the relationship between  $\Delta V_{\text{Cs}}$  and  $V_i$ . As shown in Fig. 2C, the mean  $(\pm s.E.M.)$  values of  $\Delta V_{Cs}$  were not linearly related to  $V_i$  but were more pronounced at voltages below  $-85$  mV. The depolarising effects of  $Cs<sup>+</sup>$ , therefore, were due either to the inhibition of an inwardly rectifying current reversing near  $-85$  mV, or to the suppression of multiple currents whose different voltage sensitivities and reversal potentials combined to generate the plot shown in Fig. 2*C*. When examined below  $-100$  mV, where the relationship between  $\Delta V_{Cs}$  and  $V_i$  was essentially linear, the slope of the steady-state *V*–*I* relationships obtained from MNCs showed a 71% increase from  $177 \pm 29 \text{ M}\Omega$  in control to 304 ± 49 MΩ in the presence of Cs+ (*P <* 0.05; *n =* 5).

## $\text{Cs}^+$  **blocks**  $I_{\text{H}}$

One of the inwardly rectifying currents that is known to be blocked by external  $Cs<sup>+</sup>$  is the hyperpolarisationactivated inward current,  $I_H$  (Halliwell & Adams, 1982; Pape, 1996). However, as explained earlier, blockade of the inward current mediated by  $I_{\rm H}$  normally leads to hyperpolarisation (e.g. Maccaferri & McBain, 1996), not depolarisation. Moreover, the graphs in Fig. 2*B* and *C* implied that the depolarising effect of  $Cs^+$  above  $-85$  mV was due to the blockade of an active outward current,



**Figure 2. Voltage–current analysis of the Cs+ -mediated depolarisation**

*A*, voltage responses (lower traces) to current pulses (upper) applied in the absence (control) and presence of 3 mM Cs<sup>+</sup>. All samples digitised in the last 100 ms of the voltage responses to each current pulse were averaged to determine steady-state voltage. *B*, plot of the steady-state voltage, measured in the absence and presence of  $Cs<sup>+</sup>$ , as a function of the absolute current being injected into the cell. Using graphs such as these, changes in voltage provoked by  $Cs<sup>+</sup>(\Delta V<sub>Cs</sub>)$  at different initial voltages  $(V<sub>i</sub>)$  were measured as the vertical difference (in mV) between the control and  $Cs<sup>+</sup>$  curves, at 5 mV increments. *C*, a plot of mean  $(+s.E.M.) \Delta V_{Cs} - V_i$  data measured in 5 different MNCs. Note that the plot shows strong inward rectification and a reversal of polarity near  $-85$  mV.

rather than to the recruitment of an additional inward current. We therefore hypothesised that  $Cs<sup>+</sup>$  was blocking a current other than  $I_{\rm H}$  to cause depolarisation and that its effects on steady-state *V*–*I* relationships reflected actions on both conductances. An alternative hypothesis was that  $Cs<sup>+</sup>$  did not block  $I<sub>H</sub>$  in rat MNCs and that the graph in Fig. 2*C* entirely reflected the properties of the conductance that was blocked by  $Cs<sup>+</sup>$  to provoke depolarisation. We thus sought to determine whether  $Cs<sup>+</sup>$ effectively blocked  $I_{\rm H}$  in rat MNCs using voltage-clamp analysis. Current–voltage *(I–V)*relationhips were obtained by delivering a series of prolonged (2–3 s) voltage steps to values between  $-115$  and  $-40$  mV at sufficiently low frequency  $(\leq 0.05 \text{ Hz})$  to allow complete inter-pulse deactivation of voltage-sensitive currents. Such trials were performed at constant intervals before, during and after application of Cs+ (Fig. 3*A*). Digital subtraction of current traces recorded in the presence of  $Cs<sup>+</sup>$  from those obtained in control yielded a family of current traces reflecting the time-dependent, Cs<sup>+</sup>-sensitive current (i.e. putative  $I_{\rm H}$ ) recorded at each voltage (Fig. 3*B*). These traces were fitted with a single exponential function and the time constant of activation  $(\tau_H)$  and steady-state

amplitudes were measured at each of the voltages tested. These data were averaged across the population of cells tested  $(n = 5)$  and plotted as a function of voltage. The mean current–voltage *(I–V)* relationship obtained indicated that the time- and voltage-dependent inward current blocked by  $Cs<sup>+</sup>$  had an apparent activation threshold near \_60 mV (Fig. 3*C)*. For each cell the conductance–voltage *(G–V)* relationship was obtained and the average *G*–*V* data were fitted with a Boltzmann equation (see Methods for details). As shown in Fig. 3*D*, the mean *G*–*V* relationship revealed a half-activation voltage  $(V_{\psi})$  of  $-76$  mV, and a slope factor  $(k)$  of 12. These values are in excellent agreement with results obtained in a previous study of MNCs  $(V_{\nu} = -78 \text{ mV}, k = 12;$ Ghamari-Langroudi & Bourque, 2000) using ZD 7288, a selective blocker of  $I_H$  (Harris & Constanti, 1995). Moreover, when examined between  $-120$  and  $-80$  mV,  $\tau_{\rm H}$  of the time-dependent, Cs<sup>+</sup>-sensitive current varied from  $125 \pm 17$  to  $470 \pm 123$  ms, respectively (Fig. 3*E*). These values are also similar to those reported for  $I_{\rm H}$ blocked by ZD 7288 (123  $\pm$  25 to 563  $\pm$  87 ms; Ghamari-Langroudi & Bourque, 2000). External Cs<sup>+</sup>, therefore, is an effective blocker of  $I_{\rm H}$  in rat MNCs.

#### $Cs<sup>1</sup>$ control wash 45 mV  $-125$  $0.2 nA$  $\overline{1}$  s B C\_-120 -80 -41  $-80$ -40  $V_m$  (mV)  $\tau_{\rm H}$  (ms)  $-55$ - 622 Ŧ, -65 513  $\widetilde{\mathcal{E}}$ 000000  $-75$ 468  $-100$ -85 100 D 424  $\Omega_{\pm}$ -95  $(xeuu\%$ 303 DECESSOR  $-105$ m٧  $\mathbf 0$  $-120$ -80  $-40$ 245 E 600  $-115$  $\mathbf{r}_{\mathbf{H}}$ 173  $(s\omega)$  $-125$ 50 pA  $1<sub>s</sub>$ m۷ 0 140  $-120$  $-80$  $-40$

## **Figure 3.**  $Cs^+$  blocks  $I_H$  in rat MNCs

Voltage-clamp recordings were made of the membrane currents evoked by voltage steps lasting 2.5 s to a variety of potentials  $(V_m)$  between  $-50$  and  $-125$  mV. Multiple trials were evoked in the absence and presence of  $\mathrm{Cs}^+$ , and current traces recorded in each condition were averaged to reduce noise. *A*, superimposed current traces (lower) resulting from voltage steps (upper) applied before (control), during  $(Cs<sup>+</sup>)$  and after (wash) application of 5 mm  $Cs<sup>+</sup>$ . All traces were obtained by averaging 3 individual trials in each condition.  $B$ , the  $Cs<sup>+</sup>$ -sensitive, timedependent currents recorded at each  $V_m$  (noted at the left of each trace) obtained by digitally subtracting traces recorded in the presence of  $Cs<sup>+</sup>$  from control traces. Superimposed on each trace is a monoexponential fit of the data points (thin line extending beyond the trace). The amplitude (difference between the end of the fit and the dotted line) and time constant of activation of the timedependent current  $(\tau_H;$  noted at the right of each trace) were derived from the fits.  $C$ , mean  $(\pm s.E.M.)$ current-voltage relationship of the Cs<sup>+</sup>-sensitive, time-dependent current  $(I_H)$  recorded from 4 MNCs.  $D$ , mean conductance  $(G_H)$ –voltage relationship derived from the data shown in *C*. The points are superimposed by a Boltzmann distribution (see Methods for details).  $E$ , plot of the mean  $(\pm s.E.M.)$ values of  $\tau_H$  as a function of voltage. Note that activation kinetics accelerate with hyperpolarisation.

## $\operatorname{Contribution\ of\ } I_{\text{H}}$  to the  $\operatorname{Cs}^+$ -induced inward **rectification**

Since  $Cs^+$  effectively blocked  $I_H$  in rat MNCs, then part or all of the apparent inward rectification observed in Fig.  $2C$  could be due to the blockade of  $I_{\rm H}$ . We therefore examined the specific effects of blocking  $I_{\rm H}$  on the *V–I* properties of five MNCs. As illustrated in Fig. 4*A*, bath application of the selective blocker of  $I_H$  ZD 7288



Figure 4. Voltage–current analysis of  $I_H$  blockade

*A*, voltage responses (lower) to 2.5 s current pulses (upper) were recorded from a MNC in the absence (control) and presence of 66  $\mu$ M ZD 7288. Holding potential was  $-65$  mV. *B*, the plots show the absolute steady-state voltages achieved in response to current pulses delivered in each of the conditions. *C*, the graph shows the mean  $(\pm s.E.M.)$  amplitude of voltage changes evoked by ZD 7288 ( $\Delta V_{\text{ZD}}$ ) as a function of control steady-state voltage *(V*<sup>i</sup> ) in 5 MNCs. Note the strong inward rectification at negative voltages and that application of  $ZD$  7288 has no effect on  $V_i$  at potentials  $\geq -60$  mV.



## **Figure 5. Effects of Cs+ on membrane potential in** the absence of  $I_{\rm H}$

*A*, superimposed voltage responses (lower) to 3 s current pulses (upper) applied in the absence (control) and presence of 3 mM  $\text{Cs}^+$  (66  $\mu$ M ZD 7288 was present throughout). *B*, the plots show the absolute steadystate voltages achieved in response to current pulses delivered in each of the conditions shown in *A*.  $C$ , filled circles plot the mean  $(\pm s.E.M.)$  amplitude of the voltage changes evoked by  $Cs^+(\Delta V_{Cs})$ , as a function of control voltage  $(V_i)$ , in the presence of ZD 7288. The data were measured in 9 different MNCs as described in Fig. 2*B*, from plots such as that shown in *B*. Note that the plot shows much less rectification in the presence than in the absence of ZD 7288 (open circles, data from Fig. 2*C)*, and that the reversal potential of the effect lies near  $-100$  mV.  $(30-70 \mu M;$  Harris & Constanti, 1995) caused a significant increase in the amplitude of voltage responses to prolonged (2–3 s) hyperpolarising current pulses. When examined below  $-100$  mV, the slope of the *V*–*I*relationship obtained from MNCs showed a 28 % increase from  $172 \pm 19 \text{ M}\Omega$  in control to  $221 \pm 27 \text{ M}\Omega$  in the presence of ZD 7288 ( $P = 0.024$ ;  $n = 5$ ; e.g. Fig. 4*B*). An examination of the mean voltage changes evoked by ZD 7288 ( $\Delta V_{ZD}$ ) at each *V*<sup>i</sup> led to the plot shown in Fig. 4*C*. The graph suggests that the blockade of  $I_{\rm H}$  provoked a hyperpolarising effect of increasing magnitude at values of  $V_i$  below  $-60$  mV, but that it had no effect above  $-60$  mV. These effects were consistent with the voltage dependency of  $I_{\text{H}}$  shown in Fig. 3*C*, and they revealed that the inhibition of  $I_{\rm H}$  by  $Cs<sup>+</sup>$  could have contributed part or all of the apparent inward rectification observed in Fig. 2*C*.

## Effects of  $Cs<sup>+</sup>$  on  $V-I$  properties in the absence of  $I<sub>H</sub>$

The voltage dependency of the Cs<sup>+</sup>-sensitive current responsible for membrane depolarisation was examined by performing steady-state *V*–*I* analysis in the absence of  $I_{H}$ . When tested on nine MNCs pre-exposed to  $30-70 \mu \text{m}$  ZD 7288, application of 3–5 mm Cs<sup>+</sup> caused a mean depolarisation of  $6.3 \pm 0.9$  mV. Under these conditions, steady-state *V*–*I* relationships measured in the absence and presence of  $Cs<sup>+</sup>$  revealed an increase in slope resistance that appeared to prevail across the entire range of voltages tested (Fig. 5*A* and *B)*. Indeed, in the presence of ZD 7288 mean values of  $\Delta V_{\text{Cs}}$  varied as a more linear function of  $V_i$  than in the absence of the drug (Fig. 5*C)*. When measured at values negative to  $-100$  mV and in the continuous presence of ZD 7288, the slope resistance of MNCs increased from  $204 + 11$  to 281 ± 20 MΩ upon applying Cs<sup>+</sup> (*P <* 0.05; *n =* 9). Finally, the polarity of the voltage changes evoked by  $Cs<sup>+</sup>$  in the absence of  $I<sub>H</sub>$  reversed around  $-99$  mV, a value approximating the equilibrium potential for  $K^+$  ions observed previously under similar recording conditions (Kirkpatrick & Bourque, 1996).



**Figure 6.** *I***–***V* **analysis of the ZD 7288-resistant, Cs+ -sensitive current**

*A*, chart recording of membrane current (lower) and voltage (upper) in a voltage-clamped MNC. Vertical deflections are changes in voltage and current associated with repeated series of computer-generated steps for  $I-V$  analysis. The cell was continuously exposed to 66  $\mu$ M ZD 7288 and the effects of bath application of  $3 \text{ mm Cs}^+$  (bar) were examined. Note the steady-state inward current evoked in the presence of Cs<sup>+</sup>, and that a gap (5 min) is inserted in the recording. B, superimposed current responses to 1.5 s steps to voltages between  $-120$  and  $-45$  mV (holding potential,  $-50$  mV) before (1, control), during (2, Cs<sup>+</sup>) and after (3, wash) the application of  $3 \text{ mM } Cs^+$ . The traces shown represent the digitally averaged  $(n = 3$ trials) responses recorded during the corresponding periods numerically identified in *A*. *C*, steady-state *I*–*V* relationships obtained from the data in *B*. *D*, graph showing the mean (±S.E.M.) difference current plots measured in 4 MNCs. The points represent the voltage dependency of the steady-state current blocked by  $\text{Cs}^+$  in the presence of ZD 7288.

## **Analysis of the Cs+ -sensitive current in the absence of**  $I_{\rm H}$

The results described previously suggested that the depolarising effects of extracellular Cs<sup>+</sup> in MNCs might be due to the inhibition of a steady-state  $K^+$  current lacking significant voltage dependency. We therefore examined the properties of the Cs<sup>+</sup>-sensitive, ZD 7288-resistant current under voltage clamp. As illustrated in Fig. 6*A*, bath application of  $3-5$  mM Cs<sup>+</sup> to MNCs pre-exposed to ZD 7288 and TTX evoked a reversible inward current when cells were clamped at a voltage near \_50 mV. Steadystate *I*–*V* analysis was performed in the absence and presence of  $3-5 \text{ mM Cs}^+$  by examining averaged current responses to clamp steps  $(1-2 \text{ s long}; \text{every } 8-16 \text{ s})$  to a variety of voltages (e.g. Fig. 6*B)*. The amplitude of the current responses evoked in each condition was plotted as a function of command potential (Fig. 6*C)* and for each cell a graph of the difference current amplitude was constructed. As illustrated in Fig.  $6D$ , the mean  $(n = 4)$ Cs<sup>+</sup>-sensitive (i.e. difference) current underlying the depolarisation of MNCs showed weak inward rectification and a reversal potential near  $-100$  mV. In three cells, measurements of Cs+ -evoked responses were repeated while explants were superfused with ACSF containing alternately 3 and 9 mm  $\rm[K^+]_{\scriptscriptstyle o}$  (as well as 35  $\mu$ m ZD 7288). In the presence of  $9 \text{ }\mathrm{mm}\,\mathrm{[K^+]}_0$  the reversal potential of the Cs<sup>+</sup>-evoked depolarisation was shifted by  $+22.7 \pm 2.6$  mV compared to values recorded in  $3 \text{ mM } [\text{K}^+]_0$  ( $P < 0.05$ ; data not shown). This value was reasonably close to the shift predicted by the Nernst equation  $(+28.7 \text{ mV})$  for a K<sup>+</sup>selective membrane at our recording temperature, confirming that the Cs<sup>+</sup>-evoked depolarisation is mediated by the suppression of a  $K^+$  conductance.

Externally applied TEA is well known to block the outward currents flowing through a variety of  $K^+$ channels. We therefore examined whether the depolarising effects of  $Cs<sup>+</sup>$  could be occluded by TEA. Intracellular recordings were obtained from MNCs in explants superfused with solutions containing 3–5 mM TEA. In each of three MNCs tested under current clamp, bath application of  $3 \text{ mM Cs}^+$  still provoked reversible membrane depolarisations (e.g. Fig. 7*A*). When tested under voltage clamp  $(n=5)$ , application of  $Cs^+$  in the presence of  $3-5$  mM TEA and  $60-70 \mu$ M ZD 7288 induced an inward current at voltages near rest, and a decrease in slope conductance (Fig. 7*B* and *C)*. In the presence of TEA and in absence of  $I_{H}$ , the mean decrease in slope conductance observed in five MNCs was  $0.80 \pm 0.16$  nS, a value not significantly different from that observed without TEA  $(P > 0.05)$ . The K<sup>+</sup> channels responsible for the depolarising effects of Cs<sup>+</sup> , therefore, are not sensitive to  $3-5$  mM TEA.

## **DISCUSSION**

It has been known for some time that hyperpolarising steps activate a slow inward current in MNCs (Bourque, 1987). Although this behaviour might have been surmised to reflect the presence of  $I_{\rm H}$ , recently published observations appeared to contradict the possible presence of this current in rat MNCs. Indeed, current-clamp experiments in hypothalamic explants (Stern & Armstrong, 1997; Ghamari-Langroudi & Bourque, 1998) revealed that application of external  $Cs<sup>+</sup>$ , a well-known blocker of  $I_H$  (Halliwell & Adams, 1982; Mayer & Westbrook, 1983), causes a membrane depolarisation instead of the hyperpolarisation expected to result



## **Figure 7. The Cs+ -sensitive leak current in MNCs is not blocked by TEA**

 $A$ , the effect of 3 mm  $Cs<sup>+</sup>$  on a MNC recorded in the presence of 5 mM TEA. *B*, current–voltage plots recorded from another cell in the absence (control) and presence of  $3 \text{ mm Cs}^+$ . ZD 7288  $(70 \mu M)$ , TTX  $(0.5 \mu M)$  and TEA  $(5 \mu M)$  were present throughout. The graph in *C* shows the voltage dependency of the Cs<sup>+</sup>-sensitive current obtained by subtraction of the  $Cs<sup>+</sup>$  from the control plot in *B*.

from blockade of  $I_{\rm H}$  (e.g. Maccaferri & McBain, 1996). Moreover, in MNCs held at voltages near threshold, depolarising sags associated with voltage responses to hyperpolarising current pulses of increasing magnitude become progressively smaller as they approach  $-100$  mV (e.g. Stern & Armstrong, 1997). This behaviour is opposite to that normally observed in cells expressing  $I_{\rm H}$ (e.g. McCormick & Pape, 1990). Despite these observations, however, the existence of a functionally significant  $I_{\rm H}$  in rat MNCs was recently established through the combined use of voltage-clamp analysis and bath application of ZD 7288, a selective blocker of  $I_H$  (Ghamari-Langroudi & Bourque, 2000). The analysis further suggested that the progressive reduction in sag amplitude observed during hyperpolarising responses in control conditions might be due to the deactivation of a steady-state voltagedependent  $K^+$  current; perhaps that which underlies slow outward rectification in rat MNCs (Stern & Armstrong, 1997). Although these studies have begun to clarify our understanding of the membrane currents expressed at near- and subthreshold voltages in rat MNCs, the ionic basis for the Cs<sup>+</sup>-induced depolarisation remained unexplained.

## **Cs<sup>+</sup> blocks superimposed inward and outward currents in MNCs**

Our voltage–current analysis indicated that the depolarising effects of bath-applied  $Cs<sup>+</sup>$  were associated with an increase in slope resistance, implying that ion channels are blocked, or otherwise closed, in the presence of the cation. Since membrane depolarisation results necessarily from the generation of a relative inward current, the active current being blocked by  $Cs^+$  must be flowing in the outward direction under steady-state conditions. Analysis of the effects of  $Cs<sup>+</sup>$  at different steady-state voltages initially suggested the involvement of an inwardly rectifying current reversing near \_85 mV. The curve plotted in Fig. 2*C*, however, reflected not only the properties of the  $\bar{Cs}^+$ -sensitive conductance responsible for depolarisation, but also the effects of blocking the superimposed  $I<sub>H</sub>$ . Indeed, our experiments showed that  $Cs^+$  effectively blocks  $I_H$  in rat MNCs (Fig. 3) and that suppression of  $I_{\rm H}$  alone imparts a strong inward rectification to the apparent voltage-dependent properties of the total Cs+ -sensitive conductance derived from *V*–*I* analysis (Figs 2*C* and 4).

## **The depolarising effects of Cs+ are due to blockade of a leakage K<sup>+</sup> current**

Having established that  $I_{\rm H}$  blockade could confound the analysis of the conductance responsible for the production of depolarising responses, we proceeded to examine the effects of  $Cs^+$  in the absence of  $I_H$ , by recording from MNCs in the continuous presence of  $30-70 \mu M$  ZD 7288. Our results indicated that the depolarising effects of  $Cs<sup>+</sup>$ recorded in the absence of  $I_{\rm H}$  were due to the suppression of a K<sup>+</sup> -selective current showing only weak inward rectification, and which reversed polarity near  $-100$  mV (Figs 6*D* and 7*C)*. These observations confirm that much of the inward rectification, and the relatively positive reversal potential  $(-85 \text{ mV})$ , characterising the effects of  $Cs<sup>+</sup>$  in control solutions is due to the simultaneous blockade of  $I_H$  (which reverses near  $-35$  mV; Pape, 1996) and of a relatively linear leakage  $K^+$  current  $(I_{KL})$ . Although previous studies have shown that externally applied  $Cs<sup>+</sup>$  can cause membrane depolarisation through the suppression of inwardly rectifying  $K^+$  currents (Williams *et al.* 1988; Jarolimek *et al.* 1994), the present results suggest that the  $\mathrm{Cs}^+$  -sensitive  $\mathrm{K}^+$  conductance in MNCs is relatively insensitive to membrane potential over the range of voltages examined  $(-40 \text{ to } -120 \text{ mV})$ .

## **Relative contributions of**  $I_{KL}$  **and**  $I_H$  **at subthreshold voltages**

Our results thus show that two distinct Cs<sup>+</sup>-sensitive conductances are present in MNCs of the rat supraoptic nucleus. Since the conductance underlying  $I_{KL}$  ( $G_{KL}$ ) does not show significant voltage dependency, and since  $G_{\text{H}}$ becomes fully active at voltages below about  $-100$  mV (Fig. 3*D)*, the relative maximal contributions of these conductances can be approximated by an inspection of voltage–current relationships below \_100 mV. In control solutions, application of  $Cs<sup>+</sup>$  increased membrane resistance below  $-100$  mV from 177 to 304 MΩ. This represents the suppression of 2.36 nS of the membrane conductance active at this range of voltages. Over the same range of potentials, application of ZD 7288 was found to increase resistance from  $172$  to  $221$  M $\Omega$ . This suggests that  $G_{\rm H}$  alone accounts for 1.32 nS of the steadystate membrane conductance below  $-100$  mV. This value is similar to the maximal value of  $G_H$  derived from voltage-clamp experiments using ZD 7288 (1.1 nS; Ghamari-Langroudi & Bourque, 2000) and from an analysis of the time-dependent current blocked by  $Cs<sup>+</sup>$ under voltage clamp  $(1.2 \text{ nS}; \text{ Fig. 3}).$  The  $\text{Cs}^+$ -sensitive conductance responsible for membrane depolarisation therefore represents about 1 nS of the total membrane conductance. This estimate is supported by the finding that in the presence of  $ZD 7288$ ,  $Cs<sup>+</sup>$  increased membrane resistance from 204 to 281 M $\Omega$ , a change corresponding to the blockade of 1.34 nS. One can surmise, therefore, that  $G_{\text{KL}}$  represents about 20% of the input conductance of rat MNCs impaled in hypothalamic explants.

## $\mathbf{Physiological\ role\ of\ the\ Cs^+}\text{-sensitive}\ \mathnormal{G}_{\mathrm{KL}}$

Although the MNCs from which we recorded in this study were not specifically identified as either vasopressin or oxytocin containing, it is worth emphasising that all 93 of the cells tested in our experiments displayed a Cs<sup>+</sup>induced response consistent with the expression of  $G_{KL}$  at rest. Since quantitative immunocytochemical studies have shown that approximately 69% of the MNCs present in the supraoptic nucleus of Long-Evans rats (Rhodes *et al.* 1981) synthesise vasopressin while the remainder express oxytocin, it is likely that both types of neurone were sampled in our study. We can surmise,

therefore, that both types of MNC express  $G_{KL}$  at rest and that modulation of this conductance could ultimately regulate the secretion of both hormones from the neurohypophysis.

The most obvious physiological role for  $G_{KL}$  is a contribution to the resting potential. The analysis presented in the previous paragraph suggests that approximately 20 % of the input conductance of MNCs at rest is provided by steady-state  $Cs^+$ -sensitive  $G_{KL}$ . The inhibition of this resting  $G_{\text{KL}}$  by 2–5 mM Cs<sup>+</sup> provoked a mean membrane depolarisation of about 6 mV. Given that MNCs typically rest at voltages within 10 mV of the threshold for action potential discharge (e.g. Mason, 1983; Stern & Armstrong, 1996), the outward current normally flowing through  $G_{KL}$  channels evidently provides an important source of tonic intrinsic inhibition. Interestingly, a previous study has shown that the depolarising effects of noradrenaline are in part due to the suppression of a K+ current that is active at the resting potential (Randle *et al.* 1985). Although it is not yet known whether these effects can be occluded by external  $Cs<sup>+</sup>$ , the observation suggests that modulation of the inhibitory effect of  $G_{KL}$ by neurotransmitters might represent an effective mechanism for the regulation of firing and hormone release. It should also be noted that Li & Hatton (1996) previously showed that activation of histamine receptors can depolarise MNCs through the G-protein-mediated suppression of a resting voltage-independent leakage conductance. Whether histamine modulates a residual  $\text{Cs}^+$ -sensitive  $G_{\text{KL}}$  or a distinct  $\text{K}^+$  conductance remains to be determined.

## **Involvement of** *G***KL in depolarising after-potentials**

Another possible function for  $G_{KL}$  might be a contribution to the generation of depolarising after-potentials (DAPs). Previous studies in MNCs have shown that action potentials are followed by DAPs (Andrew & Dudek, 1983; Bourque, 1986) and that the summation of consecutive DAPs contributes to the establishment of the plateau potential that sustains firing during phasic bursts (Bourque *et al.* 1998; Ghamari-Langroudi & Bourque, 1998). Modulation of the DAP, therefore, represents a powerful mechanism for the control of firing pattern by neurotransmitters in MNCs (e.g. Papas & Bourque, 1997; Brown *et al.* 1999). A previous study has provided evidence that the DAP may result from the transient suppression of a baseline  $K^+$  current following each action potential (Li & Hatton, 1997*b*). Although  $Ca^{2+}$  influx appears to be involved in the induction of DAPs (Bourque, 1986; Li & Hatton, 1997*a*), the nature of the putative  $K^+$  conductance being modulated is unknown. In a previous study (Ghamari-Langroudi & Bourque, 1998) we reported that external  $Cs<sup>+</sup>$  effectively blocks DAPs at low millimolar concentrations. Interestingly, the blockade of DAPs during exposure to  $Cs<sup>+</sup>$  is accompanied by membrane depolarisation and the time course of these two effects is similar during both onset and recovery (e.g. Fig. 1 in Ghamari-Langroudi & Bourque, 1998). Since the results of the present study indicate that the Cs<sup>+</sup>evoked depolarisation is due to the blockade of  $G_{KL}$ , DAPs may in fact be due to a transient, action potential-evoked suppression of  $G_{KL}$  and the loss of DAPs in the presence of  $Cs<sup>+</sup>$  may be due to the occlusion of the basal  $K<sup>+</sup>$  current required for its expression. Additional studies will be required to test this hypothesis.

#### Molecular nature of the channels underlying  $G_{KL}$

The participation of specialised 'leakage'  $K^+$  channels in the regulation of resting potential has long been inferred from the observation that many transmitters regulate membrane potential through the modulation of voltageinsensitive, as opposed to voltage-activated or inwardly rectifying,  $K^+$  conductances (Brown, 2000; North, 2000). The molecular nature of putative leakage  $K^+$  channels has recently emerged from the cloning of several members of the two-pore  $K^+$  channel family, such as TASK-1 (Duprat *et al.* 1997), TWIK-1 (Lesage *et al.* 1996), TREK-1 (Fink *et al.* 1996) and TRAAK (Fink *et al.* 1998). Could any of these channels underlie the  $G_{KL}$  of MNCs? Although many additional experiments will be required to provide an answer to this question, it is interesting to note that channels encoded by TASK are insensitive to TEA, but are blocked by external Cs+ (Czirjak *et al.* 2000), as was the  $G_{KL}$  recorded in supraoptic MNCs. Moreover, TASK-1 mRNA is present in a variety of central neurones (Talley *et al.* 2000), including those in the supraoptic nucleus (E. M. Talley & D. A. Bayliss, personal communication). Whether TASK-1 or analogous channels mediate the depolarising effects of  $Cs<sup>+</sup>$  in MNCs remains to be determined.

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