Effect of β -amyloid block of the fast-inactivating K⁺ channel on intracellular Ca²⁺ and excitability in a modeled neuron

(Alzheimer disease/mathematical model)

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ABSTRACT β -Amyloid peptide (A β), one of the primary protein components of senile plaques found in Alzheimer disease, is believed to be toxic to neurons by a mechanism that may involve loss of intracellular calcium regulation. We have previously shown that A β blocks the fast-inactivating potassium (A) current. In this work, we show, through the use of a mathematical model, that the A β -mediated block of the A current could result in increased intracellular calcium levels and increased membrane excitability, both of which have been observed *in vitro* upon acute exposure to A β . Simulation results are compared with experimental data from the literature; the simulations quantitatively capture the observed concentration dependence of the neuronal response and the level of increase in intracellular calcium.

Intracellular Ca^{2+} is an important second messenger in the central nervous system. Increases in intracellular Ca^{2+} concentration ($[Ca^{2+}]_i$) trigger neurotransmitter release, increase enzyme activity, and up-regulate gene expression, and they are believed to play a role in synaptic plasticity (1). Loss of Ca^{2+} homeostasis is therefore very detrimental to cell function and may precede certain types of neurotoxicity (2).

 β -Amyloid (A β) is a major component of senile plaques, one of the defining pathological features of Alzheimer disease (AD), and is believed by many to play an important role in the onset and/or progression of the disease. A β is a short peptide, between 38 and 43 amino acids long, produced by proteolytic cleavage of the amyloid precursor protein (APP) (3). Evidence supporting the role of A β in neurotoxicity in AD includes the following: mutations in APP lead to early-onset AD (3), transgenic mice overexpressing one of the mutant forms of APP have AD-like pathology and cognitive impairment (4), and A β is toxic to cultured neurons (5–15).

Loss of Ca^{2+} regulation has been implicated in the observed neurotoxicity of $A\beta$ (6, 16). In *in vitro* experiments, addition of $A\beta$ resulted in abnormally large increases in $[Ca^{2+}]_i$ upon depolarization of the neuron (17). $A\beta$ increased spontaneous Ca^{2+} fluctuations in neurons, in both frequency and magnitude (18). Anticonvulsants (19), Ca^{2+} channel blockers (8), glutamate receptor antagonists, and Na⁺ channel blockers (18) reportedly reduce or prevent $A\beta$ -mediated Ca^{2+} rise, neurotoxicity, or increased excitability.

The mechanism by which $A\beta$ induces loss of Ca^{2+} regulation has yet to be firmly established. We endeavored, through the use of a mathematical simulation of a hippocampal neuron, to demonstrate the feasibility of one mechanism by which $A\beta$ could cause elevations in $[Ca^{2+}]_i$, specifically, by blocking the fast-inactivating potassium (A) channel. Previously we showed that $A\beta$ blocks the A current, that the block is independent of membrane potential, and that the kinetics of channel opening and deactivation are not altered (20). From our data we postulated that $A\beta$ bound to a closed channel and prevented



FIG. 1. Schematic of components of neuron model. Ca^{2+} buffers were uniformly distributed within the cytosol. The membrane contained a Na⁺ channel; the A, DR, C, and AHP K⁺ channels; the L, N, and T Ca²⁺ channels; and a Ca²⁺ pump. Ca²⁺ concentration was computed as a function of time *t* and radial distance *r* inside the sphere.

it from opening. In results reported here, we simulated the changes in ion fluxes and intracellular Ca^{2+} in a hippocampal neuron in response to the A β block. We demonstrate that A β block of the A current would increase $[Ca^{2+}]_i$ and neuronal excitability, quantitatively consistent with literature data.

METHODS

Structure of Model. We modeled calcium dynamics in a hippocampal neuron, treating the neuron as a "wrinkled" sphere with radius R and surface area S. The spherical geometry, although incapable of describing signal integration in dendrites, has been used successfully to describe calcium dynamics and properties related to neuron excitability (21–23). The components of the model neuron are shown in Fig. 1. The only source or sink of calcium included in the cytosol was immobile calcium buffers. Calcium release from intracellular stores was ignored because it occurs on a longer time scale than that of the present simulations (24, 25). Unless otherwise stated, parameters were taken from the literature on hippocampal neurons and were not adjusted. Parameter values are listed in Table 1.

The governing equation for intracellular Ca^{2+} concentration, $[Ca^{2+}]_{i}$, everywhere inside the neuron is

$$\frac{\partial [\operatorname{Ca}^{2+}]_i}{\partial t} = D\nabla^2 [\operatorname{Ca}^{2+}]_i - r_{\text{buffer}},$$
[1]

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Abbreviations: $[Ca^{2+}]_i$, intracellular Ca²⁺ concentration; A β , β -amyloid; AD, Alzheimer disease; AMPA, α -amino-3-hydroxy-5-methyl-4-isoxazolepropionic acid.

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where *t* is time, *D* is the diffusion coefficient of Ca^{2+} inside the cell, and r_{buffer} is the rate of Ca^{2+} uptake by buffers. Ca^{2+} buffering is described as a reversible saturable process, with a single binding site for Ca^{2+} on the buffer. r_{buffer} is expressed as

$$r_{\text{buffer}} = k_{b1} [\text{Ca}^{2+}]_i [\text{B}] - k_{b2} [\text{Ca}^{2+}\text{B}],$$
 [2]

where k_{b1} and k_{b2} are the forward and reverse rate constants for the buffering reaction, respectively, [B] is the concentration of the nondiffusing Ca²⁺ buffering proteins, and [Ca²⁺B] is the concentration of buffering proteins with Ca²⁺ bound (21, 27). The value of k_{b1} is taken from the literature (21, 24). k_{b2} was adjusted to be consistent with experimental estimates of the equilibrium dissociation constant of calcium from the buffers (24–26, 28). Everywhere inside the sphere

$$\frac{\partial[\mathbf{B}]}{\partial t} = -\frac{\partial[\mathbf{Ca}^{2+}\mathbf{B}]}{\partial t} = -r_{\text{buffer}}.$$
 [3]

Diffusive terms were not included in Eq. 3, as buffers were assumed to be large and immobile at short times (28).

The boundary conditions are

$$\frac{\partial [\operatorname{Ca}^{2+}]_{i}}{\partial r} = 0 \text{ at } r = 0$$
 [4]

$$\frac{\sum I_{Ca^{2+}}}{2F} + \left(D\frac{\partial [Ca^{2+}]_i}{\partial r}\right)S + r_{pump} = 0 \text{ at } r = R, \quad [5]$$

where *r* is the distance from the center of the neuron in the radial direction, *F* is Faraday's constant, $I_{Ca^{2+}}$ is the calcium current through the membrane, and r_{pump} is the rate that Ca^{2+} is being pumped out of the cell. A positive current is defined as a cation flowing out of the cell. L, N, and T calcium currents were included. The ion currents and the rate of pumping are defined for the entire membrane; therefore, the diffusive flux of calcium is multiplied by the surface area. The Ca²⁺ pump is described by a Michaelis–Menten type formulation with

$$r_{\text{pump}} = \frac{k_{\text{p}} P_0 [\text{Ca}^{2+}]_{r=R}}{\frac{1}{K_{\text{p}}} + [\text{Ca}^{2+}]_{r=R}} - r_{\text{leak}},$$
 [6]

where k_p is the rate that Ca²⁺ is transported out of the cell, P_0 is the total surface concentration of the calcium pump, $[Ca^{2+}]_{r=R}$ represents Ca²⁺ concentration at the inside surface of the cell, K_p is the equilibrium constant for Ca²⁺ binding to the pump, and r_{leak} is the basal rate at which Ca²⁺ leaks into the cell at resting potentials. Although there are at least two Ca²⁺ transport mechanisms out of the cell, a Na⁺⁻ and an ATP-dependent pump, only the first mechanism was included (22, 26). r_{leak} was calculated such that r_{pump} equals zero when $[Ca^{2+}]_i$ equals $[Ca^{2+}]_0$. The rate of calcium extrusion predicted from our simulations is consistent with experimental values of calcium extrusion rates from two types of neurons (29).

Table 1. Model parameter va

Value	Ref(s).	
12	23	
4×10^{-5}	23	
1	26, 27	
$6 imes 10^{-6}$	21, 22	
10^{8}	21, 24	
600	24-26, 28	
125,000	22	
4×10^{-18}	22, 26	
2.5×10^{-20}	22, 26	
	$\begin{tabular}{c} \hline Value \\ \hline 12 \\ 4 \times 10^{-5} \\ 1 \\ 6 \times 10^{-6} \\ 10^8 \\ 600 \\ 125,000 \\ 4 \times 10^{-18} \\ 2.5 \times 10^{-20} \end{tabular}$	ValueRef(s).1223 4×10^{-5} 23126, 27 6×10^{-6} 21, 22 10^8 21, 24 600 24–26, 28125,00022 4×10^{-18} 22, 26 2.5×10^{-20} 22, 26

Initial conditions at t = 0 for all r are defined for $[Ca^{2+}]$, [B], and $[Ca^{2+}B]$ such that their distribution is homogeneous throughout the cell. The homogeneous distribution of buffers within the cytosol is not uniformly assumed in the literature (21, 27), but experimental data supporting a nonuniform buffer distribution are lacking. $[Ca^{2+}]_0$ was fixed at 50 nM. Reported values for buffer concentrations in various cell types range from 0.003 to 1 mM (21, 26, 28). For most simulations, [B]₀ was set at 0.225 mM. $[Ca^{2+}B]_0$ was fixed such that at the start of the simulation, the equilibrium condition was satisfied. Table 2 gives a complete list of the initial conditions used in the model.

The membrane potential on the surface of the soma is (21)

$$C\frac{\partial V}{\partial t} + I_{\text{Na}} + I_{\text{Ca}(\text{L})} + I_{\text{Ca}(\text{N})} + I_{\text{Ca}(\text{T})} + I_{\text{K}(\text{DR})} + I_{\text{K}(\text{A})}$$
$$+ I_{\text{K}(\text{C})} + I_{\text{K}(\text{AHP})} + I_{\text{inject}} + I_{\text{AMPA}} + I_{\text{leak}} = 0, \quad [7]$$

where *C* is the membrane capacitance; *V* is the membrane potential; I_{Na} is the sodium current; $I_{Ca(L)}$, $I_{Ca(N)}$, and $I_{Ca(T)}$ are the L, N, and T type calcium currents, respectively; $I_{K(DR)}$, $I_{K(A)}$, $I_{K(C)}$, and $I_{K(AHP)}$ are the delayed rectifier, fastinactivating, calcium-dependent, and calcium-dependent–long afterhyperpolarizing potassium currents, respectively; I_{inject} is the current stimulus injected into the cell; I_{AMPA} is the synaptic current corresponding to the α -amino-3-hydroxy-5-methyl-4isoxazolepropionic acid (AMPA) receptor activation; and I_{leak} is the current leaking into or out of the cell. It was assumed that there were no spatial gradients of voltage or current on the surface of the membrane, implying that all ion channels and pumps were uniformly distributed. In general, the *j*th ion current is described by

$$I_j = \bar{g}_j m^a h^b (V - V_j), \qquad [8]$$

where \bar{g}_j is the maximum conductance through the *j*th channel multiplied by the number of *j* ion channels, *m* and *h* are voltage- and time-dependent parameters describing the probability of a channel being open, *a* is an empirically determined parameter related to the number of closed but active states the channel has, *b* has a value of 1 if the channel inactivates and a value of 0 if it is noninactivating, and V_j is the Nernst potential for the *j*th ion. The Nernst potentials for Na⁺ and K⁺ were assumed to be constant, as ion fluxes were unlikely to significantly change intracellular and extracellular ion concentrations under *in vitro* conditions. For the calcium currents, the Goldman–Hodgkin–Katz potential was used (30).

For each ion current, m and h are computed by solving

$$\frac{\partial x}{\partial t} = \alpha_x (1-x) - \beta_x x, \qquad [9]$$

where x represents either m or h, and α_m , α_h , β_m , and β_h are voltage-dependent rate constants for the state variables m and h. The α and β functions for the ion currents were taken from Traub *et al.* (23), with the exception of the L, N, and T currents, which were taken from Jaffe *et al.* (26), and $\alpha_{mK(AHP)}$ and the Ca²⁺ dependence of $I_{K(C)}$, which were modified to adjust for differences in numerical methods used in determining $[Ca^{2+}]_i$.

Table 2.Initial conditions

Variable	Initial value
[Ca ²⁺]0, nM	50
[B] ₀ , mM	0.225
$[Ca^{2+}B]_0, \mu M$	1.87
All <i>m</i>	0
All h	1

and

$$I_{\rm K(C)} = \bar{g}_{\rm K(C)} \frac{100[\rm Ca^{2+}]_{r=R}}{2.5 \times 10^{-6} + [\rm Ca^{2+}]_{r=R}} (V - V_{\rm K}).$$
 [11]

Functional forms and parameters for I_{AMPA} (26) and I_{leak} (23) were taken from the literature. Ca²⁺ current through the *N*-methyl-D-aspartate (NMDA) receptor was not included in the model because it is not a significant source of Ca²⁺ entry except under conditions of high-frequency stimulation such as occur in long-term potentiation (31). Other currents not incorporated into the model were $I_{K(M)}$ and $I_{Ca(P)}$. Membrane excitability and calcium dynamics have been accurately modeled without the inclusion of either current (23, 26). Table 3 lists all channel conductances.

Effect of A β . Previously, we showed that A β blocks the A current, probably by binding to closed A channels and preventing them from opening, without altering the kinetics of channel opening, kinetics of deactivation, or voltagedependent properties of the channel (20). We therefore assumed that A β has no effect on the state variables m and h. The effect of $A\beta$ on the current was immediate and reversible, hence, we assumed that $A\beta$ binding to the channel is at equilibrium. On the basis of these data, we propose that $A\beta$ blocks a fraction of the channels $x_b = [A\beta]/(K_I + [A\beta])$, where $K_{\rm I}$ is the inhibition constant for A β blocking the channel and has a value of 1 μ M. The value of K_I determined previously (10 μ M) was based on a concentration of peptide in the delivery pipette and thus represented an upper limit of the true value of the inhibition constant. On the basis of experience, we assumed that a 10-fold dilution of peptide occurred between the delivery pipette and the neuronal membrane, yielding the lower value for K_{I} . The maximum conductance through all of the fast-inactivating potassium channels was expressed as the maximum conductance through all channels in the absence of A β , $\bar{g}_{K(A)_0}$, multiplied by the fraction of channels still open, 1 $-x_{\rm b}$. $I_{\rm leak}$ was fixed to keep the neuron membrane quiescent when unstimulated and roughly balanced the magnitude of I_A at low potentials (-65 mV). The size of I_{leak} was therefore modified for different [A β] to reflect the change in I_A .

Model Computations. The equations were solved using PDASAC (32). The solution of the system of equations yields $[Ca^{2+}]_i$ as a function of radial distance and time. Values of $[Ca^{2+}]_i$ reported were averaged over the entire volume of the neuron. For some simulations, $[Ca^{2+}]_i$ was averaged over the time of the simulation as well.

Variable	Conductance, mS/cm ²
$\bar{g}_{ m Na}$	135*
\bar{g} Ca(L)	2.5
$\bar{g}_{Ca(N)}$	2.5
$\bar{g}_{Ca(T)}$	0.25
$\bar{g}_{\rm K(DR)}$	32
$\bar{g}_{\mathrm{K}(\mathrm{A})_{0}}$	5
Γ ^g K(C)	0.8
$\bar{g}_{\rm K(AHP)}$	0.4
\bar{q}_{laak}	0 - 0.05

Channel conductances were taken from the literature (27). ${}^{*}\bar{g}_{Na}$ was increased by a factor of 3 to yield an action potential of the appropriate half-width.

RESULTS

The goal of this work was to demonstrate that $A\beta$ block of the fast-inactivating potassium channel, which we previously identified (20), is a plausible molecular explanation for observed changes in cellular function (principally increased [Ca²⁺]_i and increased excitability) in response to $A\beta$ (16–19, 33, 34). We therefore constructed a model which described the neuron in terms of elementary reactions and diffusion. The model was evaluated under several conditions, first, to ensure that it captured essential features of normal neuronal function, and second, to evaluate the effect of $A\beta$ on neuronal excitability and [Ca²⁺]_i.

Case 1: Single Action Potential. We simulated the response of the neuron to an injected current pulse I_{inject} of adjustable magnitude but with a fixed duration of 2 ms. The action potential triggered by an injected current of 0.4 nA in the absence of A β has a half-width of approximately 1 ms, and it is followed by a period during which the membrane is hyperpolarized (Fig. 2). The time between the peak membrane potential and the peak $[Ca^{2+}]_i$ is less than 2 ms. The time constant for $[Ca^{2+}]_i$ decay is on the order of 500 ms. These time constants and the shapes of the curves are generally characteristic of experimental data and model simulation of hippocampal neurons (22, 26, 35).

Excitability of the model neuron was assessed as a function of $[A\beta]$ by determining the minimum I_{inject} needed to elicit a response typical of an action potential (Fig. 3*A*). The threshold current decreased substantially in the presence of A β , indicating that A β increases the excitability of the neuron at micromolar concentrations.

Peak $[Ca^{2+}]_i$ levels were computed at a number of concentrations of A β when the stimulus applied to the neuron was 0.08-nA injected current (Fig. 3B). This level is subthreshold in



FIG. 2. Single action potential in the absence of $A\beta$. (*A*) Current pulse injected into the model neuron to evoke the action potential. (*B* and *C*) Computed average $[Ca^{2+}]_i(B)$ and voltage trajectory (*C*) of the model neuron in response to the injected current, typical of an action potential in a hippocampal neuron.



FIG. 3. Response of neuron to single action potential in the presence of $A\beta$. (A) Threshold current needed to evoke an action potential. (B) Peak $[Ca^{2+}]_i$ during an injected current pulse as a function of the extracellular $[A\beta]$. Simulation results (lines) and experimental data (\bullet ; ref. 18) are shown.

the absence of $A\beta$. The simulation demonstrates that $A\beta$ increases the peak $[Ca^{2+}]_i$ during a single action potential. Experimental data (18) of peak $[Ca^{2+}]_i$ during spontaneous action potentials in cultured hippocampal neurons in the presence of $A\beta$ -(25–35) were compared with the simulations (Fig. 3*B*). Both model simulations and data represent acute responses to $A\beta$. The simulations reasonably approximate the experimental data in terms of the minimum concentration of $A\beta$ needed to elicit an increase in peak $[Ca^{2+}]_i$ and the asymptotic value for $[Ca^{2+}]_i$ at high $A\beta$ concentrations.

Case 2: Repetitive AMPA Stimulation. Previous reports have shown that $A\beta$ increases neuronal sensitivity to glutamate (6, 16). We therefore simulated glutamatergic stimulation of the neuron in the presence of $A\beta$. A repetitive stimulus was applied to the neuron that approximated the current entering a glutamatergic neuron through AMPA-gated channels (26).

Fig. 4 shows the membrane potential and $[Ca^{2+}]_i$ in response to a stimulating AMPA current when no $A\beta$ is present. To assess the excitability of the model neuron in the presence of $A\beta$, the frequency of response (or action potential generation) was monitored during repetitive AMPA stimulation at different $A\beta$ concentrations. As $A\beta$ concentration increased, the frequency of the response exceeded the frequency of the stimulus (Fig. 5A).

The increase in $[Ca^{2+}]_i$ upon repetitive stimulation as a function of $A\beta$ was simulated. $[Ca^{2+}]_i$ was averaged over the 100-ms simulation and compared with experimental data from the literature (18). As shown in Fig. 5*B*, $A\beta$ increased the average change in $[Ca^{2+}]_i$ upon AMPA stimulation. The simulation was able to reproduce the experimental data in terms of both concentration dependence and level of response. It should be noted that there are some differences between simulation and experiment: experimentally, $[Ca^{2+}]_i$ was obtained by averaging intracellular concentrations over a 1-min



FIG. 4. Response to repetitive stimulation at the AMPA receptor in the absence of $A\beta$. (*A*) The AMPA current, applied at 30 Hz, used to elicit a repetitive response in the model neuron. (*B*) The computed average $[Ca^{2+}]_i$ and (*C*) voltage trajectory in response to AMPA stimulation.

period, and spontaneous activity as opposed to a defined stimulus was used to evoke a cellular response.

Correlation with Cell Viability. Data from the literature (5–15) on cell viability as a function of concentration of $A\beta$ were plotted along with model simulations of the average $[Ca^{2+}]_i$ during 100 ms of repetitive AMPA stimulation (Fig. 6). Only studies which reported viability of cortical or hippocampal neurons with $A\beta$ -(25–35) or aged (aggregated) $A\beta$ -(1–40) and $A\beta$ -(1–42) were included. The concentration of $A\beta$ needed to induce increased $[Ca^{2+}]_i$ (by simulation) corresponds to the concentration needed to induce toxicity.

DISCUSSION

Brorson et al. (18) showed that acute (\approx 1-min) exposure of hippocampal neurons to A β caused repeatable, reversible, and sustainable increases in intracellular Ca²⁺ and in membrane excitability, and they suggested that $A\beta$ somehow moved neurons closer to the threshold for firing. Other investigators (17, 33) have shown similar acute responses experimentally, but did not provide a mechanism by which A β caused these responses. In this work, we used a mathematical model of a hippocampal neuron to causally link these reported effects to our recent experimental observation that acute application of A β blocks the fast-inactivating K⁺ channel. We compared our simulations of Ca²⁺ response to A β directly with experimental data from the literature (18). When a subthreshold I_{iniect} is used to stimulate the neuron, the magnitude of the simulated changes in $[Ca^{2+}]_i$ are similar to the experimentally observed changes in $[Ca^{2+}]_i$ (Fig. 3B). When a repetitive AMPA stimulation is used and $[Ca^{2+}]_i$ is averaged over time, again, reasonable agreement between simulation and experiment is observed with respect to the minimum $A\beta$ concentration



FIG. 5. Response to repetitive stimulation at the AMPA receptor in the presence of $A\beta$. (A and B) Voltage trajectory (A) and average increase in $[Ca^{2+}]_i$ (B) during repetitive (30-Hz) AMPA stimulation as a function of extracellular $[A\beta]$. The average $[Ca^{2+}]_i$ increase without application of $A\beta$ was subtracted from the average $[Ca^{2+}]_i$ increase with application of $A\beta$ and plotted. Simulation results (lines) and experimental data (\bullet ; ref. 18) are shown.

required to elicit an increase in $[Ca^{2+}]_i$ and with respect to the level of the increase in $[Ca^{2+}]_i$ (Fig. 5B). Several factors must be considered in comparing simulation with experiment. First, simulation results are sensitive to parameter choices. Wherever possible we used well-established values for our simulations. Second, the experimental paradigm could not be simu-



FIG. 6. Neuronal viability and Ca²⁺ response as a function of extracellular [A β]. Experimental data (\bullet) of viability were taken from the literature (5–15). Error bars represent the SD of reported values at each [A β]. Simulations (broken line) were performed as in Fig. 5B. The left and right scales refer to the experimental data and simulation results, respectively.

lated exactly: the simulations investigated response of a single neuron to a controlled stimulus, whereas the experimental paradigm most likely represents the additive effects of a neuron being more responsive in the presence of $A\beta$ as well as the neuron receiving more stimulus from the surrounding population, which is again more highly excited in the presence of $A\beta$. In spite of these caveats, the agreement between data and simulation on both the dosage required for response and the level of response is supportive of our hypothesis.

Our mechanism is consistent with a number of in vitro and in vivo observations. For example, y-aminobutyric acid (GABA)ergic and calretinin-immunoreactive neurons are relatively spared from A β toxicity, possibly due to high Ca²⁺ buffering capacity (9, 12, 36), and Ca2+ channel density is lower in areas of the AD brain that are less affected by the disease (37, 38). We modified the Ca^{2+} buffering capacity or Ca²⁺ channel density in our simulations and showed that neurons with increased Ca2+ buffering ability or fewer Ca2+ channels are less responsive to $A\beta$ block of the A current. For example, a 3-fold increase in the Ca²⁺ buffer concentration or a 5-fold decrease in Ca²⁺ channel density leads to 20% lower $[Ca^{2+}]_i$ after A β exposure. In other studies, increased neuronal excitability has been observed in individuals at risk for developing AD (39). Anticonvulsants, which decrease excitability of the neuron, also decrease A β neurotoxicity in vitro (19). Our simulations show that $A\beta$ increases the excitability of the neuron, both by decreasing the amplitude of the stimulus necessary to excite the neuron (Fig. 3A) and by increasing the frequency of response to the same stimulus (Fig. 5A).

As seen in Fig. 6, the transition between nontoxic and toxic concentrations of $A\beta$ seen experimentally overlaps with the transition between concentrations of $A\beta$ incapable and capable of elevating Ca²⁺ levels determined from our simulations. The agreement in concentration dependence between experiment and simulation is consistent with our hypothesis that $A\beta$ block of the A current is an early critical step in the molecular mechanism by which $A\beta$ kills neurons. Toxicity is generally assayed after long-term (hours to days) exposure of neuronal cultures to $A\beta$ and is presumably the cumulative result of a number of events. Interactions of $A\beta$ with a number of other cell membrane components have been reported (for example, see refs. 40–42); these interactions could have additional effects on neuronal toxicity that are not accounted for in our model.

Our model of the action of $A\beta$ on a hippocampal neuron demonstrates that the block of the fast-inactivating potassium channel by $A\beta$ is sufficient to initiate a cascade of events membrane depolarization, calcium influx, and increased excitability—that have been observed *in vitro*. Other events induced by $A\beta$, possibly subsequent to and/or caused by Ca^{2+} entry, such as tyrosine phosphorylation (43) and free radical damage (44), will also determine the response of neurons to chronic exposure to $A\beta$.

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