Accounting for Near-Normal Glucose Sensitivity in K_{ir}6.2[AAA] Transgenic Mice

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ABSTRACT $K_{ir}6.2[AAA]$ transgenic mouse islets exhibit mosaicism such that ~70% of the β -cells have nonfunctional ATPsensitive potassium (K_{ATP}) channels, whereas the remainder have normal K_{ATP} function. Despite this drastic reduction, the glucose dose-response curve is only shifted by ~2 mM. We use a previously published mathematical model, in which K_{ATP} conductance is increased by rises in cytosolic calcium through indirect effects on metabolism, to investigate how cells could compensate for the loss of K_{ATP} conductance. Compensation is favored by the assumption that only a small fraction of K_{ATP} channels are open during oscillations, which renders it easy to upregulate the open fraction via a modest elevation of calcium. We show further that strong gap-junctional coupling of both membrane potential and calcium is needed to overcome the stark heterogeneity of cell properties in these mosaic islets.

INTRODUCTION

ATP-sensitive potassium channels are key regulators of glucose homeostasis in insulin-secreting β -cells in the pancreatic islets of Langerhans. Their conductance for K⁺ is affected by changes in nucleotide levels, and these channels can thereby couple plasma membrane electrical activity to cell metabolism (1-7). When blood glucose levels are low, glycolysis in β -cells is limited by the low influx of glucose, and the intracellular ADP/ATP ratio is high. Under these conditions, KATP channels in the membrane of β -cells remain open, holding the membrane potential at ~-70 mV. The negative membrane potential keeps voltage-dependent Ca2+ channels shut and the intracellular Ca²⁺ concentration low. With such limited Ca²⁺ entry, no insulin is released. When blood glucose levels are high, more glucose is taken up, glycolysis is increased, and the intracellular ADP/ATP ratio decreases. The decreased intracellular ADP/ATP ratio closes most of the K_{ATP} channels (8–14); >97% of the K_{ATP} channels are closed at physiological stimulatory glucose concentrations. This allows the membrane potential to become more positive, opening the voltage-gated Ca²⁺ channels and raising the intracellular Ca^{2+} concentration inside the cell. The influx of Ca²⁺ both contributes to the ensuing oscillations of electrical activity and stimulates the release of insulin from secretory vesicles (15,16).

Thus, closing of the K_{ATP} channels represents a crucial step in the complex chain of events that links a rise in blood glucose concentration to insulin release (1,16–20). Various mutations that directly or indirectly affect β -cell K_{ATP}

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channel function decouple insulin release from blood glucose concentration in different ways and degrees. Accumulating evidence (1,5,16,17,19-21) links such mutations to a broad range of diseases from hyperinsulinism to development delay, epilepsy, and neonatal diabetes (DEND) and type 2 diabetes. Therefore, experimental and modeling studies of the behavior of islets of Langerhans taken from animals with genetically modified K_{ATP} channel activity can give valuable insight into the mechanisms of regulation of these channels.

The focus of this article is a theoretical investigation of the previously reported behavior of K_{ir}6.2[AAA] transgenic mice, whose islets exhibit mosaicism such that approximately two-thirds of the β -cells have nonfunctional K_{ATP} channels (1,2,21). Despite the substantial reduction in the number of functional K_{ATP} channels, these studies found nearly normal calcium responses and only a small shift of ~2–3 mM in the glucose dose-response curve. We show that a previously published model of pancreatic β -cell electrical activity (3) is able to reproduce the experimental findings, and we use bifurcation analysis to interpret the model's behavior.

In Modeling and Methods, we describe the model, modified from one of the models in Bertram and Sherman (3) that produces fast bursting (period < 1 min) involving oscillations in K_{ATP} conductance. Although similar results can be obtained with a model in which K_{ATP} conductance is constant, we choose this one, since, a priori, it seems more challenging for a model that depends on oscillations in K_{ATP} conductance to compensate for the reduced K_{ATP} conductance.

In Results, we study the behavior of the model in the cases of normal (WT) islet and mutant (AAA) islets. We consider first a simplified case of perfectly coupled islets, in which membrane potential, calcium, and nucleotide concentrations are uniform and can be represented by a single cell with

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a reduction of two-thirds of its functional K_{ATP} channels. We find that the model can compensate for this reduction by increasing the fraction of open KATP channels, mediated by an increase in intracellular calcium concentration. We then fit experimental data (22,23) relating changes in external glucose concentration to the plateau fraction of bursting to calculate the glucose dose-response curve for the model. The results agree with the observed modest shift of ~2-3 mM in K_{ir}6.2[AAA] islets (1,2,21). We apply bifurcation analysis of the simplified model to show geometrically how the compensation is achieved. Finally, we study a model islet consisting of 216 coupled β -cells in a cubic lattice, with two-thirds of the cells randomly assumed to lack K_{ATP} channels. These simulations confirm that electrical coupling among those drastically heterogeneous β -cells can account for the nearly normal response to glucose stimulation. However, the model indicates that coupling through diffusion of calcium through gap junctions may be needed in addition to coupling via membrane potential to equalize Ca^{2+} responses in a mosaic islet. We also contrast the AAA islets with islets that are heterozygous for a different mutation (21) in Kir6.2 or in SUR1, another component of the KATP channel, and have a similar reduction in KATP conductance, but uniformly among the cells, rather than in a mosaic pattern.

MODELING AND METHODS

In the first models for β -cell bursting (24) the slow negative feedback for oscillations was provided by Ca²⁺-activated K⁺ channels. K_{ATP} current was added later as a modulating current to transduce the effects of glucose metabolism (23). There is, however, more recent experimental evidence (18,25–28) that glucose metabolism results in oscillations in the ADP/ATP ratio, which would in turn lead to oscillations in K_{ATP} channel activity and hence membrane potential. K_{ATP} channels would then participate in the pacemaking of bursting β -cells electrical activity together with Ca²⁺-activated K⁺ channels.

We use a modified version of the bursting model with three slow variables described in Bertram and Sherman (3), which is based on the hypothesis that rises in Ca²⁺ increase the ADP/ATP ratio. The parameter values are given in Table 1. A few have been changed from Bertram and Sherman (3), mainly an increase in \overline{g}_{KATP} so that it can be interpreted as the total K_{ATP} conductance in the cell, rather than the residual conductance at stimulatory glucose; this also entails a change in the range of values of *r*. The decrease in s_a was chosen to limit the shift in the glucose-dose response curve. The equations for membrane potential, *V*, delayed rectifier activation, *n*, cytosolic free Ca²⁺ concentration, *c*, the concentration of Ca²⁺ in the endoplasmic reticulum (ER), c_{er} are as follows:

$$C_{\rm m} \frac{dV^{(\rm i)}}{dt} = -\left(I_{\rm Ca}^{(\rm i)} + I_{\rm K}^{(\rm i)} + I_{\rm K_{\rm Ca}}^{(\rm i)} + I_{\rm K_{\rm ATP}}^{(\rm i)} + I_{\rm c}^{(\rm i)}\right), \quad (1)$$

$$\frac{dn^{(i)}}{dt} = \frac{n_{\infty}^{(i)}(V^{(i)}) - n^{(i)}}{\tau_{n}},$$
(2)

$$\frac{dc^{(i)}}{dt} = f_{\rm cyt} \Big(J_{\rm mem}^{((i)} + J_{\rm er}^{(i)} - J_{\rm diff}^{(i)} \Big), \tag{3}$$

TABLE 1 Parameter values for Fig. 1

Param.	Value	Param.	Value	Param.	Value
. *	0.2 μM	s_a^*	0.01 µM	Cm	5 300 fF
f _{cvt} ; f _{er}	0.01	$ au_{ m a}$	300,000 ms	$g_{K_{Ca}}$	300 pS
$V_{\rm cyt}/V_{\rm er}$	5	p_{leak}	0.0005 ms^{-1}	$\nu_{\rm n}$	-16 mV
K _D	0.3 μM	α	4.5×10^{-6}	<i>s</i> _n	5 mV
			$fA^{-1} \mu M ms^{-1}$		
k _{PMCA}	0.2 ms^{-1}	k_{SERCA}	0.4 ms^{-1}	$\nu_{\rm m}$	-20 mV
Vĸ	-75 mV	$ au_{ m n}$	16 ms	s _m	12 mV
V _{Ca}	25 mV	$g_{\rm K}$	3000 pS	р	5
KATP *	30,000 pS	$g_{\rm Ca}$	1200 pS		

The main change is an increase in $\overline{g}_{K_{ATP}}$ so that it can be interpreted as the total K_{ATP} conductance in the cell, rather than the residual conductance at stimulatory glucose; this also entails a change in the range of values of *r*. The decrease in s_a was chosen to limit the shift in the glucose dose response curve (compare Figs. 2 and 4). Figure legends indicate variations from the table.

*Values modified from Bertram and Sherman (3).

$$\frac{dc_{\rm er}^{\rm (i)}}{dt} = -f_{\rm er} \left(V_{\rm cyt} / V_{\rm er} \right) J_{\rm er}^{\rm (i)},\tag{4}$$

where the superscripts (*i*) index each variable for cell number for islet calculations. To carry out bifurcation analysis, we also consider a simplified case of perfectly coupled islets that can be represented by a single cell. In all cases, C_m is the membrane capacitance, τ_n is the activation time constant for the delayed rectifier channel, $n_{\omega}^{(i)}$ is the steady-state function for the activation variable *n*, and $J_{mem}^{(m)}$ and $J_{er}^{(m)}$ are the Ca²⁺ fluxes through the plasma membrane and the ER membrane, respectively. The ionic currents in Eq. 1 are

$$I_{\rm Ca}^{(i)}(V^{(i)}) = g_{\rm Ca} m_{\infty}^{(i)}(V^{(i)}) (V^{(i)} - V_{\rm Ca}),$$
(5)

$$I_{\rm K}^{(i)}(V^{(i)}, n^{(i)}) = g_{\rm K} n^{(i)} (V^{(i)} - V_{\rm K}), \tag{6}$$

$$I_{K_{Ca}}^{(i)}(V^{(i)}, c^{(i)}) = g_{K_{Ca}}\omega(c^{(i)})(V^{(i)} - V_{K}),$$
(7)

$$I_{K_{ATP}}^{(i)}(V^{(i)}) = g_{K_{ATP}}(V^{(i)} - V_K).$$
(8)

The steady-state activation functions are

$$m_{\infty}^{(i)}(V^{(i)}) = \left(1 + \exp\left(\frac{\nu_{\rm m} - V^{(i)}}{s_{\rm m}}\right)\right)^{-1},$$
 (9)

$$n_{\infty}^{(i)}(V^{(i)}) = \left(1 + \exp\left(\frac{\nu_{n} - V^{(i)}}{s_{n}}\right)\right)^{-1}$$
. (10)

The function ω ($c^{(i)}$) represents the fraction of Ca²⁺-sensitive K⁺ (K_{Ca}) channels that are activated by cytosolic Ca²⁺,

$$\omega(c^{(i)}) = \frac{c^{(i)}}{c^{(i)} + k_{\rm D}^{\rm p}},\tag{11}$$

where k_D is the dissociation constant for Ca²⁺ binding to the channel. Ca²⁺ fluxes across the plasma membrane, $J_{mem}^{(i)}$, and the ER membrane, $J_{er}^{(i)}$, are given by

$$J_{\rm mem}^{(i)}(V^{(i)}, c^{(i)}) = -\left(\alpha_{\rm Ca}^{(i)}(V^{(i)}) + k_{\rm PMCA}c^{(i)}\right), \quad (12)$$

$$J_{\rm er}^{(i)}(c^{(i)}, c_{\rm er}^{(i)}) = -p_{\rm leak}(c_{\rm er}^{(i)} - c^{(i)}) - k_{\rm SERCA}c^{(i)}, \qquad (13)$$

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where α converts current to flux, p_{leak} is the permeability of the ER membrane, and k_{PMCA} and k_{SERCA} are the plasma membrane and ER Ca²⁺ ATPase pump rates, respectively. Since *c* and c_{er} represent the free Ca²⁺ concentration in the cytosol and the ER, respectively, we multiply the corresponding fluxes in Eqs. 3 and 4 by the fraction of free to total cytosolic Ca²⁺, f_{cyt} , and ER Ca²⁺, f_{er} . For the ER Ca²⁺ concentration, the flux $J^{(i)}_{\text{er}}$ has also to be scaled by the ratio of the volumes of the cytosolic compartment, V_{cyt} , and the ER compartment, V_{er} .

In Eq. 8, the K_{ATP} conductance is given by

$$g_{\mathrm{K}_{\mathrm{ATP}}} = \overline{g}_{\mathrm{K}_{\mathrm{ATP}}} a^{(\mathrm{i})}, \qquad (14)$$

where $\overline{g}_{K_{ATP}}$ is the maximum K_{ATP} channel conductance and *a* is the fraction of open K_{ATP} channels. The open fraction depends on the nucleotide concentrations and can be interpreted roughly as the ratio of ADP to ATP in the cell, which is assumed in the model to be regulated dynamically by the Ca²⁺ concentration. Thus, one more equation is added to Eqs. 1–4:

$$\frac{da^{(i)}}{dt} = \frac{a^{(i)}_{\infty}(c^{(i)}) - a^{(i)}}{\tau_{a}}.$$
 (15)

The time constant τ_a is large, so *a* slowly follows the steady-state function $a_a^{(i)}$, which has an increasing sigmoidal dependence on cytosolic Ca²⁺ concentration:

$$a_{\infty}^{(i)}(c^{(i)}) = \left(1 + \exp\left(\frac{r - c^{(i)}}{s_{a}}\right)\right)^{-1}$$
. (16)

This incorporates the effect of increasing glucose concentration through the parameter *r* and admits a phenomenological interpretation of the hypothesized negative feedback of Ca²⁺ on the ATP concentration as resulting from either stimulated ATP hydrolysis or inhibited ATP production (18,25–29). In either case, K_{ATP} conductance provides slow negative feedback to increases in *c*, which complements the rapid negative feedback provided by $g_{K_{Ca}}$.

The model considered has been incorporated as a subset into a more recent model that includes, in addition, oscillations in glycolysis to produce slow bursting (period up to five minutes) and combined fast and slow oscillations (30). We use the submodel because we only have data on fast oscillations in AAA islets to which to compare.

The homogeneous islet model (representative single cell) is appropriate for studying the effects of K_{ATP} mutations on cell dynamics and bifurcation structure, but we are also interested in how much gap junctional coupling is needed to synchronize β -cells with different expression levels of functional K_{ATP} channels. To assess this, we incorporate both electrical and calcium coupling by adding the following gap-junctional conductance terms to the equations for V and c, respectively, Eqs. 1 and 3,

$$I_{c}^{(i)}(V^{(i)}) = \sum_{j \in \Gamma_{i}} g_{c,V}(V^{(i)} - V^{(j)}), \qquad (17)$$

$$J_{\rm diff}^{(i)}(c^{(i)}) = \sum_{j \in \Gamma_i} g_{\rm c,Ca}(c^{(i)} - c^{(j)}), \qquad (18)$$

where $g_{c,V}$ and $g_{c,Ca}$ are the coupling conductance and the gap-junctional permeability for Ca²⁺, respectively, and the sum is taken over the set Γ_i of nearest neighbor cells to which cell (*i*) is coupled. In addition to the dependent variables and other expressions noted above, $\overline{g}_{K_{ATP}}$ is indexed by the cell number (*i*) to represent mosaicism.

The model Eqs. 1–4 and 15, for the homogeneous islet case, were solved numerically using the software package XPPAUT (31). The bifurcation analysis was performed with AUTO 2000 (32). For the simulations of many coupled cells, we used a fourth-order Runge-Kutta method implemented in Fortran 95. The computer code used in the simulations can be downloaded from http://lbm.niddk.nih.gov/sherman/.

RESULTS

Compensation for reduced $K_{\mbox{\scriptsize ATP}}$ conductance in AAA islets

In Fig. 1 we simulate the behavior of WT and AAA (1,2) islets, assuming that they are sufficiently tightly coupled that each β -cell behaves more or less as the rest of the cells, and the islet can be approximated by a single representative β -cell (33). This simplification allows us to use the system of Eqs. 1–4 and 15, with (*i*) = 1, i.e., in the case of a single cell, to simulate islet behavior.

The parameter $\overline{g}_{K_{ATP}}$ represents the maximum nucleotidesensitive K^+ conductance or, in other words, the K_{ATP} conductance when all $K_{\rm ATP}$ channels are open. Since in AAA islets the K_{ATP} channels in ~70% of the β -cells are nonfunctional whereas the remaining cells have normal K_{ATP} function (1,2), we simulate the perfectly coupled AAA islet by reducing the total K_{ATP} conductance by twothirds from $\overline{g}_{K_{ATP}} = 30,000 \text{ pS}$ (WT) to $\overline{g}_{K_{ATP}} = 10,000 \text{ pS}$ (AAA). The simulations in Fig. 1 *a* indicate that the threefold reduction of the KATP conductance is compensated by an increase in the fraction of open KATP channels. With glucose concentration, represented indirectly by the parameter r, kept fixed here at 0.2 μ M, the open fraction is increased a little less than threefold, resulting in a slightly smaller active KATP conductance and bursting with slightly increased plateau fraction (Fig. 1 b). If the value of r is reduced slightly, the open fraction and plateau fraction would be equalized, which can be interpreted as a shift in the glucose dose-response curve, as discussed in the next section. The increase in KATP channel open-fraction is mediated by a small rise in cytosolic Ca^{2+} (Fig. 1 c), which is sufficient to raise the ADP/ATP ratio via the mechanism represented in Eqs. 15 and 16. The increase in Ca^{2+} results as well in an increase in $c_{\rm er}$ (Fig. 1 d).

The behavior of the model when the K_{ATP} conductance decreases is consistent with the spare-channel hypothesis (11). Since most of the K_{ATP} channels are closed under normal physiological stimulation (8–14), there are still plenty of spare channels that can be opened to provide K_{ATP} current sufficient to produce bursting electrical activity even in a situation like the AAA islets, in which approximately two-thirds of the K_{ATP} channels are not functional (1,2).

Glucose dose-response curve

Having shown that the model can explain the compensation for reduced K_{ATP} conductance, we continue by investigating whether the model can reproduce the reported modest shift in the glucose dose response of AAA islets compared to WT (1,2,21). To produce dose-response curves with the model, we need to map glucose concentration to values of *r*. We do this by mapping burst plateau fraction to *r*, as done in Himmel and Chay (23), shown in Fig. 2 *a*, and then mapping



FIGURE 1 Model simulations of WT and AAA islets. (*a*) Fraction of open K_{ATP} channels, *a*; (*b*) membrane potential, *V*; (*c*) cytosolic Ca²⁺ concentration, $[Ca^{2+}]_{i;}$; (*d*) ER Ca²⁺ concentration, $[Ca^{2+}]_{ER}$; and r = 0.2 for both islets.



FIGURE 2 Glucose dose-response curves for simulated of WT and AAA islets. (*a*) Mapping of parameter r to plateau fraction. (*b*) Mapping of glucose to plateau fraction via r. (*c*) Mapping of glucose to r implied by panels a and b.

glucose concentration to burst plateau fraction (22), shown in Fig. 2 *b*. The mapping from glucose to *r*, which is assumed to be the same for both WT and AAA, is shown in Fig. 2 *c*. Although secretion depends on a metabolic coupling factor or factors and can increase in constant Ca^{2+} (34–36), plateau fraction is an appropriate surrogate for secretion averaged over the period of 1 h in static incubation. In agreement with experimental results (see Fig. 4 *A* in (1), Fig. 5 in (21), and Fig. 1 *E* in (2)), we find only a modest shift of 2–3 mM glucose in the dose-response curve for the AAA islet compared to the WT islet.



FIGURE 3 Fast-slow analysis of the model. (*a*) Bifurcation diagrams for WT and AAA islets, showing the maximum and the minimum of the periodic orbits for $[Ca^{2+}]_i$ as a function of *a*. Superimposed on the bifurcation diagrams are the corresponding trajectories and the *a* nullclines for $r = 0.2 \ \mu$ M and $r = 0.1885 \ \mu$ M. (*b*) Bifurcation diagrams for WT, AAA islets, and an islet where 95% of the β -cells have nonfunctional K_{ATP} channels. The broken lines in the *z* curves denote instability.

Whereas Fig. 1 shows that at a given value of r, there is only a small change in plateau fraction as a result of the reduction in $\overline{g}_{K_{ATP}}$, Fig. 2 shows that compensation for the reduction in $\overline{g}_{K_{ATP}}$ requires only a small change in r to achieve the same plateau fraction. This small change represents the extent to which the rise in Ca²⁺ falls short of full compensation and which must be made up by a reduction in glucose concentration. In the next section, we examine the model from a geometrical viewpoint to see how these features come about.

Bifurcation analysis of the single-cell model

To gain a deeper understanding of which properties of the model give it the ability to adapt to a significantly reduced number of functional K_{ATP} channels with little change in the calcium dynamics and the glucose dose response curve, we carry out a fast-slow bifurcation analysis of Eqs. 1–4 and 15.

Since this system has more than one slow variable, there are several possible ways to divide it into fast and slow subsystems. We treat the fraction of open KATP channels, a, and the ER calcium concentration, $c_{\rm er}$, as bifurcation parameters, and, in contrast to the approach in Bertram and Sherman (3), we treat the cytosolic Ca^{2+} concentration, *c*, as a fast variable; the fast subsystem then consists of Eqs. 1-3. Fig. 3 a shows the bifurcation diagram with the fraction of open KATP channels, a, as the main bifurcation parameter and c as the output variable for both wild-type ($\overline{g}_{K_{ATP}} = 30,000 \text{ pS}$) and mutant $(\overline{g}_{K_{ATP}} = 10,000 \text{ pS})$ islets. The slow variation of c_{er} sweeps the z-shaped slow manifold to the right as $c_{\rm er}$ increases during the active phase and back to the left as $c_{\rm er}$ decreases during the silent phase. To indicate this, the z curves are plotted for three different values of $c_{\rm er}$, corresponding to the minimum, average, and maximum attained during the oscillations. To keep the figure simple, the periodic branches are shown only for the average values of $c_{\rm er}$. Finally, the Ca²⁺ versus a trajectories are superimposed and are seen to follow the paths predicted by the locations of the left knee and the terminations of the periodic branches.

The most prominent feature of the diagram is that decreasing $\overline{g}_{K_{ATP}}$ shifts the steady-state curves to the right

and slightly upward, corresponding to the threefold increase in the fraction of open K_{ATP} channels and the small increase in the level of Ca^{2+} that produces it. The behavior of the system is also determined by the nullcline for *a*, particularly its intersection with the z curve, which determines the unstable steady state of the full system, Eqs. 1-4 and 15. Bursting oscillations are obtained approximately when the nullcline crosses the z curve along the middle branch of saddle points, the portion that slopes up and to the right. When the nullcline intersects the bottom branch, the solution is a steady state, and, roughly, when it intersects the periodic branch, the solution takes the form of continuous spiking. The *a* nullcline translates upward as the glucose-sensing parameter r increases, and the steeply rising portion of the plateau fraction versus r curve in Fig. 2 a thus corresponds to the interval of r values for which the intersection of the a nullcline and the z curve traverses the region between the left knee and the periodic branch termination. The plateau fraction increases monotonically as the intersection moves up to the right because the flow becomes slower in the active phase and faster in the silent phase.

Examination of the bifurcation diagram with this in mind suggests that the modest shift in the glucose dose response curve shown in Fig. 2 b is a consequence of the nearly horizontal slope of the *a* nullcline. That is, for fixed *r*, the system with reduced $\overline{g}_{K_{ATP}}$ has an intersection that lies closer to the periodic branch, and hence exhibits larger plateau fraction, but only a small decrease in r is required to move the intersection back down to a level that corresponds to the same plateau fraction. Fig. 3 b shows that shifting r from 0.2 to 0.1885 μ M shifts the *a* nullcline down so that the it intersects the *z* curve for $\overline{g}_{K_{ATP}} = 10,000 \text{ pS}$ at the same Ca²⁺ value as the unshifted nullcline intersects the z curve for $\overline{g}_{K_{ATP}} =$ 30,000 pS. This results (not shown) in approximately the same average level of active $\overline{g}_{K_{ATP}}$ conductance ($\approx 245 \text{ pS}$) and approximately the same plateau fraction (≈ 0.25). The small shift in r required to equalize the plateau fraction reflects the small shift in the glucose dose-response curve produced by the model.

Fig. 3 b suggests that at $r = 0.2 \mu M$, cytosolic Ca²⁺ oscillations will still be obtained even when the maximum KATP conductance is reduced by 98% ($\overline{g}_{K_{ATP}} = 600$ pS). In this case, however, the levels of cytosolic Ca^{2+} and the plateau fraction would be elevated to a much greater extent. At $\overline{g}_{K_{ATP}} = 500 \text{ pS}$, bursting is no longer possible for this value of r, because the a nullcline intersects too deeply into the periodic branch, and if $\overline{g}_{K_{ATP}}$ is reduced further to 200 pS, only continuous spiking can be obtained, regardless of the value of r. Note, however, that even at $\overline{g}_{K_{ATP}} = 600 \text{ pS}$, bursting may be prevented by a further, biophysical restriction. This would require a to have an average value of ~0.38, whereas the experiments in the literature (8,13,25,37) indicate that even in 0 glucose, only ~20% of the total $\overline{g}_{K_{ATP}}$ conductance in the cell would be opened. (Washing out all the cell's ATP can open up the rest.) If



FIGURE 4 Glucose dose-response curves for simulated WT and AAA islets with s_a increased from 0.01 to 0.1 μ M (compare Fig. 2). (*a*) Mapping of *r* to plateau fraction. (*b*) Mapping of glucose to plateau fraction via *r*.

this bound is to be respected, then the model predicts bursting only for $\overline{g}_{K_{ATP}} > \sim 1200 \text{ pS}$.

If the *a* nullcline is made less flat by increasing s_a from 0.01 to 0.1, similar results are obtained (Fig. 4). Oscillations with little change in Ca²⁺ levels and plateau fraction can again be obtained when $\overline{g}_{K_{ATP}}$ is reduced by two-thirds (not shown), but the shift in the dose response curve is approximately doubled to ~4 mM.

To investigate further the robustness of the small dose response shift, we also considered the possibility that the wild-type value of $\overline{g}_{K_{ATP}}$ is four-times smaller (7500 pS) than in Fig. 1, with glucose remapped to *r* such that the mean K_{ATP} open fraction is fourfold higher than in Fig. 1 for the same glucose level. The shift in the dose response curve (not shown) is only slightly larger than in Fig. 2; it shifts ~2–3 mM in glucose for the AAA case, which now corresponds to $\overline{g}_{K_{ATP}} = 2500$ pS.

Coupled mosaic islets

The preceding sections have shown that a mosaic islet in which one-third of the cells have functional K_{ATP} channels



FIGURE 5 Simulations of a mosaic AAA islet of 216 coupled cells with two-thirds of β -cells having nonfunctional KATP channels, one-third with oscillating KATP conductance. The coupling strength used in the simulations is indicated in the top panels for *a* and *b*. The gray curves are for a representative mutant cell, and the black curves are for a representative wild-type cell. In the bottom panels, the nucleotide ratio *a* is plotted; for the wild-type cell only, this also represents the open probability of the K_{ATP} channels.

can exhibit nearly normal Ca²⁺ oscillations provided it is very tightly coupled, such that it behaves like a single cell with one-third the normal conductance. It remains to be shown that plausible coupling strength can achieve sufficiently tight coupling. To investigate this, we simulate AAA islets with 216 coupled β -cell that are arranged in a $6 \times 6 \times 6$ cube. Each cell is coupled to its nearest neighbor and has dynamics governed by Eqs. 1-4 and 15. In this cubic islet, the interior cells have six neighbors, whereas the peripheral cells have three, four, or five neighbors depending on their location on the islet surface. To achieve the mosaic expression of β -cells with functional K_{ATP} channels as reported in the transgenic $K_{ir}6.2[AAA]$ islet (1,2,21), we assume that two-thirds (a quantity of 144) of the cells in the islet, which are randomly (uniformly) distributed, are lacking functional K_{ATP} channels (have $\overline{g}_{K_{ATP}} = 0$ pS). The rest of the β -cells (a quantity of 72) are considered to be WT (have $\overline{g}_{K_{ATP}} = 30,000 \text{ pS}$). The cells are assumed to be coupled both electrically (38-41) and through diffusion of Ca^{2+} via gap-junctions (42–44).

In Fig. 5 we present model simulations of such an islet where the traces are taken from a WT β -cell and a β -cell in which there are no functional K_{ATP} channels. Two different coupling strengths are represented in Fig. 5. The coupling strength for membrane potential in Fig. 5 *b* is comparable to the experimentally measured values (45,46), whereas that in Fig. 5 *b* is increased greater than fourfold. The permeability of the gap junctions for Ca^{2+} is not known, and has been chosen sufficiently large to synchronize as described below.

Consistent with the experiments in Rocheleau et al. (2), Fig. 5 b demonstrates that strong electrical coupling in the physiological range is sufficient to synchronize the responses in the mosaic islet. Furthermore, the top panel of Fig. 5 bshows that the model can reproduce the experimental data showing that Ca²⁺ levels are equalized between mutant and wild-type cells within the islet (see Fig. 1 C in (2)). This result is not obtained in the absence of diffusion of Ca²⁺ through the gap junctions, even with exaggerated coupling through membrane potential (not shown). The bottom panels in Fig. 5 show that coupling through Ca^{2+} is also able to synchronize the nucleotide ratio a and that the stronger value of diffusion of calcium through gap-junctions is able to equalize the nucleotide ratio between the WT and mutant cells in the islet; this is a consequence of the feedback of Ca^{2+} onto ATP production (Eq. 15), which is assumed to operate whether or not KATP channels are present. In both Fig. 5, a and b, the open fraction of the WT cells agrees with the values found in the perfectly coupled islet (Fig. 1 d).

The most striking finding of our analysis is that although coupling strengths in the physiological range for normal islets (45,46) are sufficient to synchronize the responses in an AAA islet (Fig. 5 b), they are not sufficient to produce normal bursts in the membrane potential (Fig. 5 b, middle panel). In the simulations, the fast spikes within a burst of the membrane potential disappear by oscillator death because of the extreme heterogeneity of the mosaic islets (47). However, if the electrical and Ca^{2+} coupling are increased further (Fig. 5 a), the spikes in the membrane potential during the active phases of the bursts are restored even when the cells in the mutant islet are so different. In contrast, if $\overline{g}_{K_{ATP}}$ is reduced uniformly by two-thirds in all the cells, synchrony, equalized Ca^{2+} , and normal spike amplitude can be achieved with coupling only on membrane potential equal to the weaker value in Fig. 5 (not shown). This case is also of physiological interest as it corresponds to islets from mice that are heterozygous for the K_{ir}6.2 or SUR1 genes (21).

DISCUSSION

In this study we have sought to account for the nearly normal glucose sensitivity observed experimentally (1,2,21) in transgenic Kir6.2[AAA] islets that are characterized by mosaic structure in which ~70% of their β -cells lack functional K_{ATP} channels. We adapted a previously published mathematical model (3) of β -cell electrical activity that incorporates oscillations in the fraction of open KATP channels, which participate in the pacemaking of β -cells electrical activity together with Ca²⁺-sensitive K⁺ channels. The oscillations in KATP conductance reflect oscillations in the nucleotide ratio ADP/ATP that are driven by the rise and fall of cytosolic Ca²⁺. The model is agnostic as to whether this effect is caused by increased ATP consumption (48-50) or decreased ATP production (51) or both, so our results do not depend on the answer to this question. Under this generic assumption, a threefold reduction in KATP conductance in AAA islets leads to a nearly threefold increase in the fraction of open KATP channels, which restores nearly normal oscillations (Fig. 1) as observed experimentally. This result is achieved in the model because elevated Ca²⁺ persists until the total conductance of open KATP channels rises to the range (≈ 250 pS) needed to turn the activity off. This range is in agreement with that reported in wild-type (WT) islets (12, 18, 52, 53).

To achieve an exact threefold rise in active K_{ATP} conductance, glucose must be lowered in the AAA case compared to WT, that is, the glucose dose response curve shifts to the left. To estimate the size of the shift, we mapped the $\overline{g}_{K_{ATP}}$ open fraction obtained in the model as a function of the glucosesensing parameter, *r* in Eq. 15. We found (Fig. 2) that the shift was ~2 mM. We then showed geometrically that this small shift was possible because of the assumed relatively steep increase in steady-state $\overline{g}_{K_{ATP}}$ open fraction with Ca²⁺ (Fig. 3). The result is not excessively sensitive to the slope of this increase, as a 10-fold reduction in the slope results in only a twofold increase in the shift of the glucose dose response curve (Fig. 4). We also found that the modest shift in the dose response curve is not very sensitive to the assumed level of $\overline{g}_{K_{ATP}}$ (not shown).

The model predicts that islets can compensate for much larger reductions in $\overline{g}_{K_{ATP}}$. The most fundamental limitation is that the K_{ATP} open fraction cannot exceed 1; this would predict a minimum $\overline{g}_{K_{ATP}}$ value at which oscillations could ever be obtained of only 100 pS, given the other parameters of the model's fast subsystem. However, long before this limit was reached, bursting would be prevented because the *a* nullcline would intersect the branch of periodic solutions (Fig. 3); this is predicted to occur at $\sim \overline{g}_{K_{ATP}} = 500$ pS. Moreover, there is experimental evidence (8,13,25,37) that the fraction of open K_{ATP} channels in 0 glucose is ≈ 0.2 . This bound would limit the residual K_{ATP} conductance for a mutant to ~ 1200 pS. Even so, that represents impressive tolerance to a decrease of 96% from 30,000 pS, or 84% from 7500 pS.

We note that Ca^{2+} oscillations have been observed in islets from mice in which all K_{ATP} channels are lost (54,55). The mechanism described in this article cannot account for such behavior, and instead the existence of an alternative metabolically sensitive K⁺ channel has been suggested to compensate partially for the loss of K_{ATP} channels (54). It is possible that the compensation in AAA mutants is also due to an alternative channel. No such channel has yet been identified, however, and our model shows that it may not be required to account for the partial loss of K_{ATP} .

A model with constant fraction of open KATP channels would also be able to compensate for reduction in total K_{ATP} conductance. In this case, β -cell electrical activity would be driven by other currents, such as K_{Ca} current, and the role of KATP current would be only to trigger the bursts in the membrane potential rather than contribute to pacemaking. In this case, there would be no increase in the fraction of open K_{ATP} channels as a result of Ca^{2+} feedback and all of the burden of compensation would be borne by a reduction in glucose to increase the ratio of ADP/ATP. Thus, other things being equal, one might expect that the shift in the dose response curve when $\overline{g}_{K_{ATP}}$ is reduced would be larger than in a model with oscillating $\overline{g}_{K_{ATP}}$. However, our model is very sensitive to the open fraction of K_{ATP} channels when that fraction is fixed independent of Ca^{2+} . and we have not found any other satisfactory model with which to test this conjecture quantitatively. We can nonetheless conclude from the findings presented here that the ability of islets to compensate for substantial reductions in KATP conductance is not per se a reason to prefer models with nonoscillating K_{ATP} to models with oscillating K_{ATP} .

We also examined whether gap-junctional coupling of plausible strength can produce the synchronized response in mosaic islets assumed in the above calculations. We simulated islets in which two-thirds of the cells had no K_{ATP} channels while one-third had the normal complement. Although in our simulations we used only 216 cells

(arranged in a $6 \times 6 \times 6$ cube), we have obtained qualitatively similar results (not shown) using 1728 cells (arranged in a $12 \times 12 \times 12$ cube). As has been previously shown (56), the gap-junctional conductance required to synchronize an islet scales like the diameter of the islet. Thus, increasing the islet size from $6 \times 6 \times 6$ to $12 \times 12 \times 12$ would increase the conductance needed to synchronize its activity by a factor of approximately two. The simulations (Fig. 5) showed that synchrony and equalized Ca²⁺ between WT and mutant cells within the islet are possible only if coupling via diffusion of Ca^{2+} is assumed in addition to the standard coupling via membrane potential. This confirmed the suggestion that gap junctional coupling could account for the nearly normal glucose sensitivity observed experimentally in AAA islets (2). However, to obtain voltage bursts with spikes, rather than plateaus, the coupling for membrane potential had to be at least fourfold greater than that measured physiologically. At this time, membrane potential oscillations in AAA islets have not been studied, so we do not know if the coupling strength is upregulated or if in fact the spikes in the burst are eliminated. On the other hand, it is possible to obtain synchronized voltage bursts without spikes in the active phase as in Fig. 5 b if the coupling conductance is reduced to 0.2 nS because it is much easier to synchronize the slow waves than the spikes. In islets in which coupling is reduced uniformly in every cell, as in $K_{ir}6.2$ +/- or SUR1 + /- mice (21) or islets with partial but uniform pharmacological blockade of KATP channels, the model indicates that synchrony is not a problem.

We have focused on transgenic K_{ir}6.2[AAA] (1,2,21) islets and showed that reducing $\overline{g}_{K_{ATP}}$ results in a left shift of the glucose dose-response curve, which is partly compensated by the rise in Ca²⁺. Diabetes has been associated in genome scans with reduced K_{ATP} channel affinity for ATP. In such a case, one can expect a right shift in the glucose dose-response curve, as has been observed in transgenic mice (57). The extent of the shift as a function of the reduction in ATP sensitivity and whether the shift may be partly compensated by a fall in Ca²⁺ are interesting questions for future modeling studies.

CONCLUSIONS

The results of the study can be summarized in the following predictions about the fraction of open K_{ATP} channels in β -cells from WT and AAA islets, and about the effect of electrical gap-junctional coupling and calcium diffusion through gap-junctions on the collective behavior of coupled heterogeneous β -cells. We propose independent of the details of the model for electrical activity:

Prediction 1. The fraction of open K^+ channels in transgenic $K_{ir}6.2[AAA]$ islets (1,2,21) is approximately threefold larger than in wild-type islets at equivalent activity (Fig. 1 *a*).

- Prediction 2. Calcium diffusion through gap-junctions among the β -cells is needed to equalize $[Ca^{2+}]_i$ levels in the transgenic K_{ir}6.2[AAA] islets (1,2,21) due to widely discrepant K⁺ conductances (Fig. 5 *b*).
- Prediction 3. In K_{ir}6.2[AAA] islets, either coupling is stronger than in wild-type islets (Fig. 5 *a*) or spiking is lost (Fig. 5 *b*). K_{ir}6.2 +/- islets would not require such strong coupling because of lesser heterogeneity.

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