

Published in final edited form as:

*J Phys Chem A*. 2009 April 30; 113(17): 4930–4940. doi:10.1021/jp8093947.

## EPR Line Shifts and Line Shape Changes Due to Spin Exchange of Nitroxide Free Radicals in Liquids 6. Separating Line Broadening due to Spin Exchange and Dipolar Interactions

 Barney L Bales<sup>\*</sup>, Michelle Meyer, Steve Smith, and Miroslav Peric

Department of Physics and Astronomy and The Center for Supramolecular Studies, California State University at Northridge, Northridge, CA 91330

### Abstract

EPR spectra of perdeuterated 2,2,6,6-tetramethyl-4-oxopiperidine-1-oxyl (PDT) are studied as functions of molar concentration,  $c$ , and temperature,  $T$ , in water and 70 wt% glycerol in water. The increase of the intrinsic linewidth averaged over the three hyperfine lines,  $\langle B_{tot} \rangle$ , varies linearly with

$c$  with zero intercept in both solvents at all temperatures; therefore  $\frac{d}{dc} \langle B_{tot} \rangle$  is independent of  $c$ . The spin exchange induced dispersion, from which the spin exchange frequency,  $\omega_e$ , may be computed, increases linearly with  $\langle B_{tot} \rangle$ , passing through the origin in water and in 70 % glycerol at high temperatures; however, at low temperatures, where dipolar interactions broaden the spectra, linearity does not prevail until  $\langle B_{tot} \rangle > 1$  G due to a contribution of dipolar interactions to the dispersion. The

broadening constant due to spin exchange,  $\frac{d}{dc} \langle B_e \rangle$ , is found from the slope of the linear region,

permitting a computation of the dipolar constant,  $\frac{d}{dc} B_{dip} = \frac{d}{dc} \langle B_{tot} \rangle - \frac{d}{dc} \langle B_e \rangle$ . Thus, the separation of concentration broadening into spin exchange and dipolar contributions is effected without having to appeal to some supposed temperature dependence of the two interactions. The fractional broadening by spin exchange,  $\Omega(T)$ , is near unity at high temperatures in both solvents, decreasing to zero in 70 % glycerol at 273 K.  $\Omega(T)$  is a continuous function of the inverse rotational correlation time of PDT, but is discontinuous as a function of  $T/\eta$  where  $\eta$  is the shear viscosity.  $\Omega(T) = 0.5$ , where spin exchange and dipolar interactions contribute equally to the linewidth occurs at  $T/\eta = 20$

$\pm 1$  K/cP in 70 % glycerol. Hydrodynamic predictions of  $\frac{d}{dc} \langle B_e \rangle$  via the Stokes-Einstein (SE) equation are remarkably accurate in 70 % glycerol comparable with the results in a series of alkanes.

In water,  $\frac{d}{dc} \langle B_e \rangle$  is linear with  $T/\eta$  with zero intercept as required by the SE; however, with slope a

factor of 0.73 smaller.  $\frac{d}{dc} \langle B_{dip} \rangle$  is reasonably predicted by the SE only at very small values of  $\eta/T$  very quickly following an approximately logarithmic dependence rather than the linear prediction.

Values of  $\frac{d}{dc} \langle B_{dip} \rangle$  approach a plateau above  $\eta/T = 0.20$  cP/K that is about one-half the solid state limit. Line shifts due to spin exchange are not yet useful to deduce values of  $\Omega(T)$  due to a lack of knowledge of the time between re-encounters; however, they may be used to verify the values determined from line broadening and spin exchange induced dispersion. Some effects at low

\* Corresponding author. barney.bales@csun.edu. Webpage: <http://www.csun.edu/~vcphy00s/BBVita.html>.

Supplemental Information is available free of charge via the Internet at <http://pubs.acs.org>.

temperatures in 70% glycerol suggest that the effects of dipolar interaction are inadequately described by the widely accepted theory.

## Introduction

This is a continuation of series of papers (Parts 1 – 5)<sup>1-5</sup> whose ultimate goal is to measure the translational diffusion of nitroxide free radicals (nitroxides) in compartmentalized systems, such as micelles. We began the work 18 years ago<sup>6</sup> thinking that the theory was complete:<sup>7</sup>

spin exchange broadened and shifted the resonance lines when  $\frac{\omega_e}{\gamma} \ll A_0$  and narrowed them in the opposite extreme, where  $\omega_e$  is the spin exchange frequency,  $A_0$  the hyperfine coupling constant in the absence of spin exchange, and  $\gamma$  the gyromagnetic ratio of the electron. It was thought that  $\omega_e$  could be derived from any one of these three effects<sup>7</sup> and we thought that we could use a combination of the effects to deduce collision frequencies and statistical distributions of nitroxides in compartmentalized systems. However, the predicted line shifts turned out not to agree with experiment<sup>2,3,8-11</sup> and the narrowing is not practical to use in most interesting cases,<sup>7</sup> leaving line broadening as the only tool. It was known<sup>7</sup> that the line shape departed from Lorentzian as  $\omega_e$  increased, so we turned our attention to that aspect of the problem and learned that the line shape could be expressed as a sum of two components, an absorption and a “spin exchange induced dispersion.”<sup>1</sup> In Part 2<sup>2</sup> we showed that only two components were necessary to reproduce the EPR spectrum to very high values of  $\omega_e$  where  $\frac{\omega_e}{\gamma A_0}$  was near unity and the lines coalesced into a single line and in Part 3 we proposed that the line shift discrepancy was due to re-encounter collisions.<sup>3</sup> Part 4<sup>4</sup> provided a test of the line shift proposal and Part 5<sup>5</sup> studied re-encounters in a series of alkanes. In short, at high values of  $T/\eta$ , where  $T$  is the absolute temperature and  $\eta$  is the shear viscosity, we finally think that the theory is complete, although more work is needed to fully establish that the re-encounter mechanism is responsible for the line shift discrepancy.<sup>12</sup> We refer the reader to Parts 1 – 5 and the references therein for extensive discussion.<sup>1-5</sup> A major remaining problem arising as  $T/\eta$  decreases—separation of dipolar and spin exchange interactions—is the subject of the present work. This problem often occurs in the most interesting liquids and complex fluids because biologically important systems are often viscous enough to require such a separation.<sup>13</sup> Applications of spin-spin interactions between nitroxides in biological and model membranes are numerous.<sup>14</sup>

Both spin exchange and dipolar interactions broaden EPR lines linearly with concentration.<sup>13</sup> In the past, most of the efforts to separate the interactions had to rely on a supposed temperature dependence of the dipolar interaction often formulated from the Stokes-Einstein equation.<sup>13</sup> Unfortunately, a paper published in 1979 by Berner and Kivelson<sup>13</sup> showed that the line broadening method was not very promising because the linewidth did not show a hydrodynamic dependence at low values of  $T/\eta$ .

At first, our strategy to separate dipolar and spin exchange interactions was to measure spin exchange induced dispersion to deduce  $\omega_e$ , use this to compute the broadening due to spin exchange, and subtract from the total broadening to get the broadening due to dipolar interaction. Then we planned to use the line shifts to get an independent value of  $\omega_e$  to double check the procedure; however, this went awry as mentioned in the first paragraph. Therefore, we are left with just enough information to be able to separate the effects due to spin exchange and dipolar interactions at a single temperature not having to make any assumptions about the model of diffusion. Fortunately, we *are* able to utilize the line shifts to verify the results.

## Theory

### Heisenberg Spin Exchange

The theory has been presented in detail for EPR spectra due to  $^{14}\text{N}$  nitroxides undergoing spin exchange in low viscosity liquids.<sup>1-3,7</sup> We briefly summarize the results and refer the reader to the literature for more details and discussion.<sup>1-5</sup> The generally accepted expression describing the spectrum, given by Currin<sup>15</sup> in 1962, predicts that the lines broaden, shift toward the center of the spectrum, and eventually merge into a single line as  $\omega_e$  increases. We call Currin's<sup>15</sup> expression, derived by many workers,<sup>1</sup> the "rigorous theory." For a summary, see ref<sup>1</sup> and the Appendix of the same reference<sup>1</sup> for explicit expressions of the rigorous theory written for a first-derivative magnetic field-swept spectrum. The rigorous theory is in agreement with experiment not only for two- and three- but also for five-line spectra in every detail except one.<sup>2-4</sup> The line shifts are larger than predicted in every case<sup>2-4</sup> except one, aqueous solutions of Fremy's salt.<sup>1</sup> We showed that the line shapes in the first-derivative field-swept form predicted by the rigorous theory were reproduced to high precision by a simple analytical expression valid all the way up through line coalescence as follows<sup>2-4</sup>

$$Y(H, \omega_e) = \sum_{M_I} V_{pp}(M_I) S_{M_I}(H) + V_{\text{disp}}^{\text{ex}}(M_I) D_{M_I}(H) \quad (1)$$

where,  $H$  is the swept magnetic field,  $S_{M_I}(H)$  is the absorption, and  $D_{M_I}(H)$  is a spin-exchange induced "dispersion." The sum in eq 1 is over the nuclear quantum number  $M_I = +1, 0,$  and  $-1$ , labeling the low-, central-, and high-field lines, respectively. Equation 1 is identical to eq 1 of ref<sup>2</sup> except that the absorption of Lorentzian shape has been replaced by  $S_{M_I}(H)$  which is a Gaussian-Lorentzian sum function approximation to a Voigt shape given by eq 12 of ref<sup>16</sup>. The absorptions  $S_{M_I}(H)$  involve mixing parameters from which the separate Lorentzian and Gaussian contributions to the linewidths may be deduced and various parameters may be corrected for inhomogeneous broadening.<sup>16</sup>  $V_{pp}$  is the peak-to-peak amplitude of the absorption and  $M_I V_{\text{disp}}^{\text{ex}}(M_I)$  are the maximum values of  $D_{M_I}(H)$ . The expression for  $D_{M_I}(H)$  is given by eq 10 of ref<sup>2</sup>. The credit for the form of eq 1 ought to go to Molin et al.<sup>7</sup> who worked out the perturbation theory predicting the second term. Our contribution was to apply the term to nitroxides, develop non-linear least-squares fitting software allowing us to measure it, and permit the parameters in the original perturbation theory to vary with  $\omega_e$  in order to

preserve the simple form at high values of  $\frac{\omega_e}{\gamma A_0}$ .<sup>1,2</sup>

Tables 1 and 2 give the results of fitting spectra generated by the rigorous theory to eq 1 for  $^{14}\text{N}$  and  $^{15}\text{N}$  nitroxides, respectively. Inhomogeneous broadening by unresolved hyperfine structure has no effect on the results in Tables 1 and 2.<sup>2</sup>

In a conventional first-derivative presentation, such as that in Figure 1 below, where the integrated spectrum is positive,  $V_{\text{disp}}^{\text{ex}}(+)/V_{pp}(+) = -V_{\text{disp}}^{\text{ex}}(-)/V_{pp}(-) > 0$ , while for the center

line  $V_{\text{disp}}^{\text{ex}}(0) = 0$ . The ratios  $V_{\text{disp}}^{\text{ex}}(M_I)/V_{pp}(M_I)$  are monotonic functions of  $\frac{\omega_e}{\gamma A_0}$  that are well approximated by the following:

$$\frac{\omega_e}{\gamma A_0} = \pm \frac{3\sqrt{3}}{4} \left( \frac{V_{\text{disp}}^{\text{ex}}(\pm)}{V_{pp}(\pm)} \right) \left[ 1 - 0.707 \left( \frac{V_{\text{disp}}^{\text{ex}}(\pm)}{V_{pp}(\pm)} \right)^2 \right] \quad (2)$$

The leading term in eq 2 is predicted by perturbation theory<sup>1,7</sup> and the second, found by fitting the values in Tables 1 and 2, corrects<sup>3</sup>  $\frac{\omega_e}{\gamma A_0}$  to within 0.5% for  $\frac{\omega_e}{\gamma A_0} \leq 0.43$ . For more precision or to find higher values of  $\frac{\omega_e}{\gamma A_0}$ , interpolation of the values in Tables 1 and 2 may be employed.

In addition to the spin exchange induced dispersion component in eq 1, an “instrumental” dispersion component due to an imbalance in the microwave bridge often appears in the spectrum that may be corrected as follows:<sup>4</sup>

$$V_{\text{disp}}^{\text{ex}}(\pm)/V_{pp}(\pm) = V_{\text{disp}}(\pm)/V_{pp}(\pm) - V_{\text{disp}}(0)/V_{pp}(0) \quad (3)$$

There are two independent values of the right-hand side of eq 2, which we abbreviate as follows:

$$\left( \frac{V_{\text{disp}}^{\text{ex}}}{V_{pp}} \right)_{\pm}^{\#} \equiv \pm \left( \frac{V_{\text{disp}}^{\text{ex}}(\pm)}{V_{pp}(\pm)} \right) \left[ 1 - 0.707 \left( \frac{V_{\text{disp}}^{\text{ex}}(\pm)}{V_{pp}(\pm)} \right)^2 \right] \quad (4)$$

To save space, we shorten the terms “spin-exchange induced dispersion” to “exchange dispersion,” spin exchange to “exchange,” and dipolar interactions to “dipolar” when the meaning is clear.

## Dipolar

The effect of dipolar on spin relaxation in liquids at large values of  $T/\eta$  where motional narrowing theory is expected to hold, has been available for many years<sup>17</sup> and applied to electron spins at X-band.<sup>13</sup> The effect of dipolar in the static limit where  $T/\eta$  approaches zero, is also known.<sup>13</sup> In both of these limits, the lines are Lorentzian in shape and broaden linearly with the molar concentration,  $c$ . In recent years, Nevzorov and Freed have bridged the gap between the motional narrowing regime and the rigid limit by considering coupled systems of two spins<sup>18</sup> and many-body systems.<sup>19</sup> These later theories have not yet advanced to the point of being applicable to nitroxides; however, they confirm that the lines broaden in proportion to  $c$  and are approximately Lorentzian.

## Assumptions

Thus, summing up the older theories, we assume that dipolar

1. produces broadening that is proportional to  $c$  and is independent of  $M_I$ ;
2. produces Lorentzian lines;
3. contributes neither to  $V_{\text{disp}}$ ; nor to
4. line shifts; and
5. does not alter the relative intensities of the three lines.

The above assumptions follow from the widely accepted theory of the equations of motion of the transverse components of magnetizations that include only secular terms thereby preventing coherence transfer that is responsible for the interesting effects in spin exchange outlined above. We are indebted to one of the referees who called our attention to more recent work<sup>20</sup> where the authors argue that nonsecular terms are required in the dipolar interaction. Such terms would, in principle, also lead to dispersion, line shifts, etc.; however, with the important

difference that the signs of the dipolar nonsecular terms are opposite to those for spin exchange. The theory is not yet available for the three-line case pertinent to nitroxides; however, qualitatively, the lines are expected to shift toward one another and merge, but not narrow as the dipolar interaction increases.<sup>20</sup> Dispersion terms are expected of opposite sign to those under spin exchange. These two manifestations of dipolar interactions would negate assumptions 3 and 4. We do not yet know whether the second part of assumption 1 or assumption 5 are would be negated. The arguments in <sup>ref 20</sup> seem sound; however, we do not know of any experimental confirmation of the details of the theory. In this work, we are unable to confirm the theory but find intriguing evidence in favor of some of its aspects, which offers some interesting prospects for future work.

### Separation of Exchange and Dipolar

The broadening of the  $M_I$  line is computed from

$$B_{\text{tot}}(M_I) = \Delta H_{pp}^L(c)_{M_I} - \Delta H_{pp}^L(0)_{M_I} \quad (5)$$

where  $\Delta H_{pp}^L(c)_{M_I}$  and  $\Delta H_{pp}^L(0)_{M_I}$  are the peak-to-peak intrinsic Lorentzian linewidths in the presence and absence of concentration broadening due to exchange and dipolar, respectively. Exchange broadens the  $M_I = \pm 1$  lines slightly faster than the  $M_I = 0$  line (Table 1).<sup>3</sup> We define

$\langle B_e \rangle = \frac{1}{3} [B_e(+)+B_e(0)+B_e(-)]$  with  $\langle B_{\text{tot}} \rangle$  similarly defined. Line broadening due to the two interactions are additive,<sup>13</sup> thus

$$\langle B_{\text{tot}} \rangle = \langle B_e \rangle + B_{\text{dip}} \quad (6)$$

Exchange broadens the lines as follows:

$$\langle B_e \rangle = \frac{4}{3\sqrt{3}} \frac{\omega_e}{\gamma} = \pm A_0 \left( \frac{V_{\text{disp}}^{\text{ex}}}{V_{pp}} \right)_{\pm}^{\#} \quad (7)$$

Because both dipolar and exchange broaden the lines linearly with concentration, at a given temperature in a given solvent, the fractional broadening by exchange

$$\Omega(T) \equiv B_e/B_{\text{tot}} \quad (8)$$

is expected to be independent of the molar concentration,  $c$ . Near  $c = 0$ , it has been observed that linearity breaks down so following Berner and Kivelson,<sup>13</sup> we take the derivative of eq 6 with respect  $c$  and employ eq 7 to write

$$\frac{d}{dc} (\langle B_{\text{tot}} \rangle) = \pm \frac{d}{dc} A_0 \left( \frac{V_{\text{disp}}^{\text{ex}}}{V_{pp}} \right)_{\pm}^{\#} + \frac{d}{dc} (B_{\text{dip}}) \quad (9)$$

In terms of the broadening constants  $\frac{d}{dc}\langle B_{\text{tot}}\rangle$ ,  $\frac{d}{dc}\langle B_e\rangle = \pm \frac{d}{dc}A_0\left(\frac{V_{\text{disp}}^{\text{ex}}}{V_{pp}}\right)_{\pm}^{\#}$ , and  $\frac{d}{dc}B_{\text{dip}}$ , we write from eq 9

$$\frac{d}{dc}B_{\text{dip}} = \frac{d}{dc}\langle B_{\text{tot}}\rangle - \frac{d}{dc}\langle B_e\rangle. \quad (10)$$

The terms on the right-hand side of eq 10 are measured directly and independently allowing us to deduce  $\frac{d}{dc}B_{\text{dip}}$ . Note that the commonly quoted second-order spin exchange rate

constant<sup>7</sup> may be computed as follows:  $K_e = \frac{3\sqrt{3}\gamma}{4} \frac{d}{dc}\langle B_e\rangle$ .

In the region in which  $\Omega(T)$  is constant versus  $c$ , we may write

$$\Omega(T)_{\pm} = \frac{\frac{d}{dc}(\langle B_e\rangle/A_0)}{\frac{d}{dc}(\langle B_{\text{tot}}\rangle/A_0)} = \frac{\pm \frac{d}{dc}\left(\frac{V_{\text{disp}}^{\text{ex}}}{V_{pp}}\right)_{\pm}^{\#}}{\frac{d}{dc}(\langle B_{\text{tot}}\rangle/A_0)} \quad (11)$$

In eq 11, uncertainties in  $c$  and in the temperature yield uncertainties in  $\Omega(T)$  that can be minimized by observing that

$$\Omega(T)_{\pm} = \pm \frac{d}{d\langle B_{\text{tot}}\rangle} \left( \frac{V_{\text{disp}}^{\text{ex}}}{V_{pp}} \right)_{\pm}^{\#} \quad (12)$$

Therefore, independent measurements of the total broadening and the dispersion amplitude yield two independent values of  $\Omega(T)_{\pm}$ . Note that assumptions 1 and 3 are employed in the derivation of eq 12.

In principle, line shifts could also be used to provide another independent measure of  $\omega_e$  although of slightly lower precision.<sup>1</sup> There are two shifts: one that decreases the spacing between resonance fields,  $A_{\text{abs}}$ ; the other that decreases the spacing between the fields where the spectrum crosses the baseline,  $A_{\text{obs}}$ . See Figure 1, where  $2A_{\text{abs}}$  is depicted as the spacing between the maxima of the dispersion features; however, it is also the spacing between the absorption lines that are not shown.  $A_{\text{abs}}$  cannot be measured directly from the spectrum; it must be deduced from a fit to eq 1. The shifts leading to  $A_{\text{obs}}$  are about a factor of three larger than those leading to  $A_{\text{abs}}$ . See earlier papers for explicit expressions for these spacings as predicted by perturbation theory.<sup>1</sup> The expressions for  $A_{\text{abs}}$  are given in the footnotes to Tables 1 and 2. Note that the difference between the perturbation theory predictions of  $A_{\text{abs}}$  and the rigorous values are extremely small (compare columns 3 and 4 of Tables 1 and 2, respectively). The rigorous theory does not include the effect of exchange due to re-encounter collisions and is not in agreement with experiment except in one case, Fremy's salt.<sup>1</sup> The expression for  $A_{\text{abs}}$ , modified by inclusion of a parameter describing the effect of re-encounters on the line

shifts,  $\kappa = \frac{\sqrt{3}}{8}(1 + \sqrt{2})\sqrt{\gamma A_0 \tau_{RE}/2}$ , where  $\tau_{RE}$  is the mean time between re-encounter collisions is given in eq 8 of ref<sup>4</sup> while the same modification of  $A_{\text{obs}}$  is as follows:

$$\frac{A_{obs}}{A_0} = 1 - \left( \frac{9}{16} \Delta H_{pp}^L(0)_{\pm} - \kappa \right) \frac{B_e}{A_0} - \frac{27}{32} \left( \frac{B_e}{A_0} \right)^2 \quad (13)$$

where  $\Delta H_{pp}^L(0)_{\pm}$  is the average Lorentzian linewidth of the low- and high-field lines in the unbroadened spectrum. The additional term,  $-\kappa \frac{B_e}{A_0}$ , applies only in the case of strong exchange; i.e., when  $\omega_e$  is near the collision frequency.<sup>12</sup> We have focused almost entirely on  $A_{obs}$  in recent work because the effect due to re-encounter exchange is a larger fraction of the total shift, allowing us to study re-encounters.<sup>3-5</sup> Here, our motive is different: we wish to study

$A_{obs}$  because the term  $-\kappa \frac{B_e}{A_0}$  is quickly dominated by the final term in eq 13 masking the small temperature dependence of  $\kappa$ .<sup>4</sup> Therefore, a plot of  $\frac{A_{obs}}{A_0}$  vs.  $\frac{B_e}{A_0}$  will yield a universal curve for all temperatures and concentrations except near the origin. This will provide an independent

check on the values of  $B_e$  obtained from measurements of  $\pm \left( \frac{V_{disp}^{ex}}{V_{pp}} \right)_{\pm}^{\#}$ .

The translational diffusional coefficient of a Brownian particle diffusing in a bulk liquid, is given in the hydrodynamic limit by the Stokes-Einstein (SE) relation

$$D = kT / 6\pi a \eta \quad (14)$$

where  $a$  is the effective hydrodynamic radius and  $k$  the Boltzmann constant. Utilizing the SE relation, Berner and Kivelson<sup>13</sup> showed that

$$\frac{d}{dc} \langle B_e \rangle^{SE} = C_e T / \eta \quad (15)$$

and

$$\frac{d}{dc} \langle B_{dip} \rangle^{SE} = C_{dip} \eta / T \quad (16)$$

for low values of  $\eta/T$  and

$$\frac{d}{dc} \langle B_{dip} \rangle^{SE} = C_{dip}^{static} \quad (17)$$

in the solid state limit as  $\eta/T \rightarrow \infty$ . For nitroxides, evaluating the constants in Berner and Kivelson's<sup>13</sup> eqs 8, 10, and 11,  $C_e = 0.485cP \cdot G/M \cdot \text{deg}$ ,  $C_{dip} = 763 \text{deg} \cdot G/cP \cdot M$ , and  $C_{dip}^{static} = 49.03G/M$ . In evaluating  $C_e$ , we have assumed that the distance for exchange interaction is equal to the hydrodynamic diameter of the probe.

## Experimental

PDT (lot # A559P2, perdeuterated 2,2,6,6-tetramethyl-4-oxopiperidine-1-oxyl, 97 atom % d) was purchased from CDN Isotopes and used as received. ACS reagent glycerol (99.8%) was obtained from Sigma and used as received. Stock solutions of 19.9 mM in water and 93.6 mM in 70 wt% aqueous glycerol were prepared in MilliQ water by weight and diluted. Undegassed samples were sealed into 50- $\mu$ L pipettes and placed into sample tubes housed in Bruker's nitrogen gas-flow temperature control dewar. The accuracy of the concentrations is estimated to be  $\pm 3\%$  and the relative concentrations  $\pm 1\%$ . The temperature, accurate to  $\pm 1$  K, was measured with a thermocouple placed just above the active portion of the microwave cavity. During each spectrum, the temperature was stable to  $\pm 0.1$  K. Three EPR spectra were acquired, one after the other, with a Bruker 300 ESP X-band spectrometer interfaced with Bruker's computer employing a sweep time of 41 s; time constant, 10 ms; microwave power, 5 mW; sweep width, 50.2 G; and modulation amplitude, 0.2 G. The magnetic sweep width was measured with Bruker's NMR Gaussmeter operating in the 1-mG resolution mode and was averaged over a day's run. At low concentrations, the line shapes were accurately described by the Gaussian-Lorentzian sum functions,  $S_{ML}(H)$  from which the peak-to-peak Lorentzian,  $\Delta H_{pp}^L$  and Gaussian,  $\Delta H_{pp}^G$  line widths were determined and line height ratios were corrected for inhomogeneous broadening.<sup>16</sup> Least-squares fits of the spectra were carried out as detailed previously.<sup>1-3</sup> The Gaussian components, due to unresolved hyperfine couplings and the effect of field modulation yielded  $\Delta H_{pp}^G = 0.118 \pm 0.007$  G. After correcting for field modulation,<sup>21</sup> unresolved hyperfine structure contributed 0.088 G to  $\Delta H_{pp}^G$ .

## Results

Figures 1 – 3 show experimental spectra of 0.1 mM (first trace) and 93.6 mM (second trace) PDT in 70% glycerol at 343, 293, and 273 K, respectively. Fits of the second traces to eq 1 are indistinguishable from the spectra on the scales of Figures 1 – 3. The residuals, the spectra minus the fits, are shown as the noisy baselines. The third traces are the dispersion components of the fits. The spectrum and residual in Figure 1 may be compared with spectra of other nitroxides in various solvents at high values of  $T/\eta$  where exchange dominates in Figure 2 of ref<sup>1</sup> and Figures 2, 4, and 5 of ref<sup>2</sup>. Spectra simulated from the rigorous theory and their fits to eq 1 are given in Figure 1 of ref<sup>1</sup> and Figure 1 of ref<sup>2</sup>. In earlier publications, we have displayed all of the components of eq 1; here, we omit the absorption components to save space. In Figure 1 and even up to much higher values of  $\omega_e/\gamma$  leading to coalescence of the hyperfine lines into a single line,<sup>2</sup> only the absorption and dispersion components in eq 1 are needed to describe the spectra leading to undetectable differences in the spectra and the fits. Thus, the line shapes employed eq 1, Lorentzian-Gaussian sum functions at low values of  $\omega_e/\gamma$  and Lorentzian functions at high values, are in excellent agreement with experimental results.

In Figures 2 and 3, in contrast to the near-perfect fits in cases where exchange dominates, we observe that fits to spectra exhibiting significant dipolar broadening show definite residuals demonstrating that there are departures of the line shape from Lorentzian. The departures are not large; in the worst case, better than 97% of the doubly-integrated intensity of the spectra is accounted for by the Lorentzian shape. Nevertheless, an uncertainty that is difficult to estimate is introduced due to the residual. The fits to the spectra in Figures 1 - 3 lead to  $\Delta H_{pp}^G = 0$ ; thus, if the imperfect fits are due to inhomogeneous dipolar broadening, such broadening is not detected as being Gaussian. Note that the central lines in Figures 1 – 3 show dispersion components due to microwave bridge imbalances; these are used to correct the dispersion amplitudes using eq 3. Also, at 273 K, note the curious inversion of the dispersion components giving negative values at low field and positive at high, opposite of that expected from exchange. Note that the previous sentence was written before we received the input from



one of the referees suggesting that negative values could be expected from a more accurate theory of dipolar.<sup>20</sup> Fitting low-temperature spectra such as that in Figure 3 with only absorption components (setting  $V_{disp}(M_I) = 0$ ) yields residuals that are very similar to those in Figure 3; i.e., the line shape error is predominantly due to the absorption.

Figure 4 shows representative values of  $\langle B_{tot} \rangle$  at various temperatures as functions of  $c$  in water (4a) and in 70 % glycerol (4b); the straight lines are linear least-squares fits. In water, the coefficients of correlation,  $r$ , are better than 0.999 at all temperatures, while in 70 % glycerol the correlation is less satisfactory declining to reach a low of  $r = 0.986$  at 273 K.

The slopes of the total broadening,  $\frac{d\langle B_{tot} \rangle}{dc}$ , are given in Figure 5 as functions of  $T/\eta$ . In water, these slopes are linear with  $T/\eta$ ; however, they do not extrapolate to the origin. In 70 % glycerol,  $\frac{d\langle B_{tot} \rangle}{dc}$  shows an upturn as  $T/\eta$  decreases toward the origin as shown by the inset. In Figures 4–7 and 11, data from three spectra taken one after another are plotted separately to show the reproducibility. The appearance of a single datum means that the differences are too small to perceive on the scale of the figures.

Figure 6 displays values of  $-A_0 \left[ \frac{V_{disp}^{obs}}{V_{pp}} \right]_{-}^{\#}$ , corrected using eq 3, as functions of  $B_{tot}$  where the superscript *obs* denotes the observed value; i.e., that emerging from the fit to eq 1. Plots of  $+A_0 \left[ \frac{V_{disp}^{obs}}{V_{pp}} \right]_{+}^{\#}$  (not shown), are very similar to Figure 6; the error bars show one-half the difference in the two values. In water, Figure 6a, the plots are, as expected, straight lines of slope approximately equal to unity at higher temperatures, demonstrating that  $-A_0 \left[ \frac{V_{disp}^{obs}}{V_{pp}} \right]_{-}^{\#} = +A_0 \left[ \frac{V_{disp}^{ex}}{V_{pp}} \right]_{+}^{\#} = B_e$ . In contrast, in 70 % glycerol at lower temperatures, Figure 6b, plots of  $-A_0 \left[ \frac{V_{disp}^{obs}}{V_{pp}} \right]_{-}^{\#}$  do not become linear until  $B_{tot} > 1$  G. Additionally, at lower temperatures,  $-A_0 \left[ \frac{V_{disp}^{obs}}{V_{pp}} \right]_{-}^{\#}$  and  $+A_0 \left[ \frac{V_{disp}^{obs}}{V_{pp}} \right]_{+}^{\#}$  both become negative which cannot come from exchange. The straight lines in Figure 6a are linear least-squares fits constrained to the origin and in Figure 6b are fits to the highest 4 concentrations.

Clearly, the results in Figure 6b do not conform to assumption 3, that dipolar interactions do not contribute to the dispersion for values of  $B_{tot}$  near the origin. It is not clear from the fits

whether (1) dipolar yields negative values of  $\pm A_0 \left[ \frac{V_{disp}^{obs}}{V_{pp}} \right]_{\pm}^{\#}$  or (2) the imperfect fits evident in Figures 2 and 3 lead the search in parameter space to non-physical values. The input received from one of the referees suggests that (1) might be expected. Note that the correction term in eq 4 is negligible for the negative values so we drop the superscript #. If possibility (1) holds, then we conclude that the dipolar dispersion is not linear in  $c$  reaching a plateau at rather low

values of  $\langle B_{tot} \rangle$ . If possibility (2) holds, then the fact that  $-\left[ \frac{V_{disp}^{obs}}{V_{pp}} \right]_{+}^{\#} = \left[ \frac{V_{disp}^{obs}}{V_{pp}} \right]_{-}^{\#}$  even when these quantities are negative, is a curious coincidence. In either case, we may write

$$\left[ \frac{V_{\text{disp}}^{\text{obs}}}{V_{pp}} \right]_{\pm}^{\#} = \pm \left[ \frac{V_{\text{disp}}^{\text{ex}}}{V_{pp}} \right]_{\pm}^{\#} \pm \left[ \frac{V_{\text{disp}}^{\text{dip}}}{V_{pp}} \right]_{\pm} \quad (18)$$

where the superscript *dip* refers to the dipolar contribution be it due to possibility (1) or (2).

The data in Figure 6b and at all other temperatures show that values of  $\pm \left[ \frac{V_{\text{disp}}^{\text{dip}}}{V_{pp}} \right]_{\pm}$  achieve constant values away from the origin; thus, for these higher values of  $B_{\text{tot}}$ ,

$$\pm d \left[ \frac{V_{\text{disp}}^{\text{dip}}}{V_{pp}} \right]_{\pm} / dB_{\text{tot}} = 0, \text{ and}$$

$$\Omega(T)_{\pm} = \pm d \left[ \frac{V_{\text{disp}}^{\text{ex}}}{V_{pp}} \right]_{\pm}^{\#} / dB_{\text{tot}} = \pm d \left[ \frac{V_{\text{disp}}^{\text{obs}}}{V_{pp}} \right]_{\pm}^{\#} / dB_{\text{tot}} \quad (19)$$

allowing us to evaluate  $\Omega(T)_{\pm}$  from eq 12 and  $B_e$  from eq 8.

Therefore, from eq 18

$$\left[ \frac{V_{\text{disp}}^{\text{dip}}}{V_{pp}} \right]_{\pm} = \pm \left[ \frac{V_{\text{disp}}^{\text{obs}}}{V_{pp}} \right]_{\pm}^{\#} - \Omega(T)_{\pm} B_{\text{tot}} \quad (20)$$

Figure 7 shows values of  $\left[ \frac{V_{\text{disp}}^{\text{dip}}}{V_{pp}} \right]_{\pm}$ , which are average values of  $\pm \left[ \frac{V_{\text{disp}}^{\text{dip}}}{V_{pp}} \right]_{\pm}$ ; the error bars are the standard deviations of the two values. The dipolar contribution to the observed dispersion reaches a plateau above 40 mM. The solid lines in Figure 7 are non-linear least-squares fits to

$$\left[ \frac{V_{\text{disp}}^{\text{dip}}}{V_{pp}} \right]_{\pm} = \left[ \frac{V_{\text{disp}}^{\text{dip}}}{V_{pp}} (\infty) \right] \tanh(c/c^*) \quad (21)$$

where the constant  $\left[ \frac{V_{\text{disp}}^{\text{dip}}}{V_{pp}} (\infty) \right]$  is the plateau value at large  $c$  and  $c^*$  is a characteristic concentration at which the plateau is approached. There is no theoretical basis for eq 21; for now, we use it as a guide to the eye in Figure 7. For possible future reference, we plot the values

of the constants  $\left[ \frac{V_{\text{disp}}^{\text{dip}}}{V_{pp}} (\infty) \right]$  and  $c^*$  in the Supplemental Information.

Guided by the SE, eq 14, reasoning that at equal values of the diffusion coefficient,  $\Omega(T)$  ought to be the same, we plot values of  $\Omega(T)$  (the average of  $\Omega(T)_{\pm}$ ) versus  $\eta/T$  in Figure 8a for the

two solvents showing that the fractional broadenings by exchange are not accurately predicted by the SE in different solvents. Interestingly, there is enough information in the EPR spectrum to pursue the matter further. The rotational correlation time,  $\tau_R$ , was computed from line-height ratios using standard formulas<sup>22</sup> and corrected for inhomogeneous line broadening.<sup>16</sup> According to the Stokes-Einstein-Debye relation,  $\tau_R$  is proportional to  $\eta/T$ . In complicated solvents, neither  $D$  nor  $\tau_R^{-1}$  is proportional to  $T/\eta$ ; however, their ratio is constant.<sup>10,23</sup> In other words, the departures from the SE are similar for translation and rotation. This fact motivates the plot in Figure 8b in which values of  $\Omega(T)$  are plotted versus  $\tau_R^{-1}$  showing that the transition between the two solvents is continuous. Robinson et al.<sup>10</sup> write a modified SE employing a viscosity to a power  $< 1$  to unify the results. Another possible reason for this discrepancy might be the effect of steric hindrance on the spin exchange. Due to hydrogen bonding, the steric hindrance factor value<sup>7</sup> and thereby the efficiency of spin exchange may be affected by the solvation of the NO group in water more than in the water-glycerol mixtures.

Figure 9 shows values of  $\frac{d}{dc} \langle B_e \rangle$  as a function of  $T/\eta$  in the two solvents. The solid line is the SE prediction computed from eq 15, while the dashed line is a linear least-squares fit to the data in water yielding  $K_e = (0.352 \pm 0.001)T/\eta$  with  $r = 0.978$ . The SE is reasonably accurate in predicting values of  $\frac{d}{dc} \langle B_e \rangle$  in 70% glycerol as shown by the inset to Figure 9. The results in the inset are very similar to those from PDT in a series of alkanes: see Figure 5 of Part 5.<sup>5</sup> In water the ratio  $\frac{d}{dc} \langle B_e \rangle / \frac{d}{dc} \langle B_e \rangle^{SE} = 0.73$ . Berner and Kivelson<sup>13</sup> observed a similar result for di-tert-butyl nitroxide in water where  $\frac{d}{dc} \langle B_e \rangle / \frac{d}{dc} \langle B_e \rangle^{SE} = 0.65$  and pointed out that a difference in the hydrodynamic diameter of the probe and the exchange interaction distance in water could rationalize the result.

Values of  $\frac{d}{dc} \langle B_{dip} \rangle$  computed from eq 10 using the results in Figures 4 and 9 are plotted versus  $\eta/T$  in Figure 10. The straight line is the SE prediction, eq 16. The inset shows the results near the origin where the data in water are seen to be in reasonable agreement with the SE and the initial slope in 70% glycerol is also in reasonable agreement. The data in 70% glycerol rather quickly saturate, reaching only about one-half of the predicted static limit of 49.03G/M, eq 17. The solid line through the 70% glycerol data is a linear least squares fit of the data to a logarithm yielding  $\frac{d}{dc} \langle B_{dip} \rangle = 39.2 + 19.4 \log(\eta/T)$  in G/M with  $\eta$  in cP. The transition from low values of  $\eta/T$  to the static limit has been studied theoretically by Nevzorov and Freed<sup>18,19</sup> for a simplified system in which no hyperfine structure is included. Also, their treatment does not include various averaging effects, so their results can only be viewed qualitatively; nevertheless, Figure 6 of ref<sup>19</sup> shows that dipolar broadening decreases logarithmically with  $D$ . Near the static limit, the dipolar broadening becomes inhomogeneous.<sup>19</sup>

Values of  $A_{obs}/A_0$  for PDT in 70% glycerol are plotted versus  $\langle B_{tot} \rangle / A_0$  in Figure 11a and versus  $\langle B_e \rangle / A_0$  in Figure 11b where  $B_e$  is computed from eqs 19 and 8. The solid lines in Figure 11a are quadratic fits to guide the eye; in Figure 11b the dashed line is a fit of all of the data to eq 12 yielding  $\kappa = -0.045 \pm 0.001$ . The data at temperatures represented in Figure 11a are represented by the same symbols in 11b. All of the rest of the data are shown by small plus signs. Note in Figure 11a, at lower temperatures, where the dipolar interaction dominates,  $A_{obs}/A_0$  actually increases slightly with  $B_{tot}/A_0$  in contradiction to assumption 4 that dipolar interactions do not contribute to the line shifts. The formation of a common curve in Figure 11b supports the methods to evaluate  $\Omega(T)$  and  $B_e$ .

## Discussion

Our goal in this work was to demonstrate a simple means to separate exchange and dipolar based on the EPR line shape, not appealing to any supposed temperature dependence of the interactions. According to the assumptions, arising from older motional-narrowing theory, Figure 6b would have been a series of straight lines passing through the origin, and only two samples at high- and low-concentrations would be required; however, that simplicity has thus far eluded us. Obtaining the data in Figure 6b requires extensive sample preparation, data collection, and analysis. One can separate dipolar and exchange, but it is not yet simple. If, indeed the theory in <sup>ref 20</sup> proves to be an accurate representation of the experimental results, the separation will be complicated, however, with the added benefit of more information derived from the dipolar. We also carried out an extensive investigation employing <sup>15</sup>N PDT in 70 % glycerol at concentrations up to 85 mM which we are unable to report because that concentration is not sufficient to reach the linear portions of plots such as Figure 6b. Despite there being some advantages to using <sup>15</sup>N spin probes, they are inferior in these types of studies for several reasons: (1) for the same concentration, values of  $B_e$  are smaller by a factor of 3/4 because only collisions between different values of  $M_I$  produce broadening, (2) exchange dispersion which is a function of  $B_e / A_0$  is reduced by a factor of 1.4 because the hyperfine coupling constant is larger by that factor, and at least as important (3) one does not have the redundancy in the exchange dispersion at high- and low-fields as a measure of the uncertainty or the central line to correct instrumental dispersion. Nevertheless, we note that in the testing stages of the theory in <sup>ref 20</sup>, these probes could be quite valuable because they would not present complications inherent in the second part of assumption 1 or in assumption 5.

Note that it is simple to detect the *presence* of dipolar by (1) preparing plots comparable to those in Figure 6 and looking for slopes less than unity or (2) preparing plots such as those in Figure 11a and looking curves that do not fall on top of one another. The more demanding task is to quantitatively compute values of  $\Omega(T)$ .

The analysis requires non-linear least-squares fitting of the spectra to eq 1 which gives nearly perfect fits in the limit  $\Omega(T) \rightarrow 1$  but poorer fits as  $\Omega(T)$  decreases toward the dipolar limit  $\Omega(T) \rightarrow 0$ . These imperfect fits are due to our incomplete knowledge of the proper EPR shape of a line broadened by dipolar. There is a small, unknown component to the line shape in addition to the predominant Lorentzian shape. Interestingly, “Pake” doublets appear in the theoretical line shapes; see Figure 1 of <sup>ref 18</sup>. Further, the lines are broadened inhomogeneously.<sup>18</sup> We suspected that both of these would be detected by the fitting software as Gaussian contributions to the line shape; however, the best fits yield  $\Delta H_{pp}^G = 0$ . The discrepancy in Figures 2 and 3 is small; however, the experimental line is higher in the wings than a Lorentzian, reminiscent of the result in solids; see, e.g., p. 135 of <sup>ref 7</sup>.

Fortunately, the values of  $\left[ \frac{V_{disp}^{dip}}{V_{pp}} \right]$  induced either by fitting to an improper line shape or by a true negative dipolar contribution, reach constant values at about  $\langle B_{tot} \rangle \approx 1$  G allowing us to

evaluate  $\Omega(T)$  from the slopes of  $\left[ \frac{V_{disp}^{obs}}{V_{pp}} \right]_{\pm}^{\#}$  observed at higher values of  $\langle B_{tot} \rangle$  where the dependence is linear. That this approach is valid is supported by the fact that values of  $A_{obs} / A_0$  form a common curve when plotted against values of  $\langle B_e \rangle / A_0$  derived from those slopes (Figure 11).

Values of  $\frac{d}{dc} \langle B_e \rangle$  for PDT in 70% glycerol are in remarkable agreement with those predicted from hydrodynamic theory using eq 15. We were surprised by this result, because it was known that the rotation of probes in aqueous solutions of glycerol did not follow the SE prediction

very well.<sup>10</sup> These data together with those published recently show that  $\frac{d}{dc} \langle B_e \rangle$  is very similar in hexane, decane, hexadecane, and now 70% glycerol. So far, only water is different in that

$\frac{d}{dc} \langle B_e \rangle / \frac{d}{dc} \langle B_e \rangle^{SE} = 0.73$  although the variation with  $T/\eta$  is nearly linear and passes through the origin. It is known that in many instances the value of  $\frac{d}{dc} \langle B_e \rangle$  of a nitroxide is lower in water than aprotic solvents.<sup>7</sup> The steric factor, which can be determined by both the structure of the nitroxide and the solvation of the NO group, can reduce the spin exchange efficiency of nitroxide collisions. It is possible that the discrepancy we observe in water is due to the solvation of NO.

The results in Figure 8a show that values of  $\Omega(T)$  at a given value of  $T/\eta$  are solvent dependent in the case of water and 70% glycerol, so one cannot accurately predict its value from a knowledge of the viscosity of the system. Nevertheless, as a guide, using the results in 70% glycerol, dipolar and exchange contribute about equally ( $\Omega(T) = 0.5$ ) at  $T/\eta = 20 \pm 1$  K/cP which corresponds to  $\eta = 15$  cP at 300 K. Typical micelles present microviscosities, as estimated from rotational correlation times of spin probes, as follows: sodium dodecyl sulfate and dodecyl trimethylammonium bromide, 5 – 7 cP;<sup>24</sup> tetraalkylammonium dodecyl sulfate 6 – 8 cP;<sup>25</sup> and a sugar-based dodecyl sulfate 20 – 28 cP.<sup>26</sup> Therefore, exchange is expected to dominate in some micelles and dipolar in others.

Figure 10 shows that the SE, eq 16, overestimates the values of  $\frac{d}{dc} \langle B_{dip} \rangle$  in 70 % glycerol even at small values of  $\eta/T$  while the predicted static limit, eq 17, is about a factor of two larger than

the experimental result. The dependence of  $\frac{d}{dc} \langle B_{dip} \rangle$  on  $\eta/T$  becomes decidedly non-linear rather near the origin which explains why efforts to fit line broadening due to exchange and dipolar using the sum of eqs 15 and 16 have failed in liquids<sup>13</sup> and have led to discrepancies in model membranes.<sup>14</sup> The discrepancies in predictions of eqs 16 and 17 and the experimental results are not unanticipated because the theoretical predictions are upper value estimates as emphasized by Molin et al.,<sup>7</sup> because significant averaging of dipolar can occur by various mechanisms.

Dipolar begins to be significant even in water at lower temperatures. It has been impossible in the past, measuring only the linewidths to detect the onset of dipolar. Results like those in Figure 5 have misled workers through the years. For example, the marvelous linearity of

$\frac{d}{dc} \langle B_{tot} \rangle$  with  $T/\eta$  in water seemed to show that dipolar was negligible; however, the non-zero intercept is revealing. Such a non-zero intercept has been observed and various suggestions

have been proposed;<sup>27</sup> however, comparing Figure 9, where  $\frac{d}{dc} \langle B_e \rangle$  does pass through the

origin, with Figure 5 where  $\frac{d}{dc} \langle B_{tot} \rangle$  does not, shows that neglect of even minor dipolar contributions can explain the apparent anomaly. Even in 70% glycerol, we observe a reasonably

linear dependence of  $\frac{d}{dc} \langle B_{tot} \rangle$  on  $T/\eta$  for  $T/\eta > 10$  K/cP (inset, Figure 5), but again with a non-zero intercept, signaling a significant dipolar contribution.

Considering Figure 10, is the pessimism of Berner and Kivelson<sup>13</sup> that EPR is not very promising to study diffusion still valid? Above about  $\eta/T = 0.1$  cP/K,  $K_{dip}$  becomes rather insensitive to diffusion. A typical diffusion coefficient in a model membrane is  $1 \times 10^7$  cm<sup>2</sup>/s; see ref<sup>14</sup> and references therein. Using eq 14, employing a probe radius of  $3.5 \text{ \AA}^5$  for PDT, this corresponds to  $\eta/T = 0.21$  cP/K, which falls in a region where the dipolar line broadening is rather flat with respect to  $D$ , supporting their conclusion. Thus, in viscous liquids, it does not seem promising to study dipolar to estimate diffusion coefficients, nevertheless, one can still separate exchange and dipolar and rely on the well understood exchange. The precision obviously suffers as  $\Omega(T) \rightarrow 0$ .

### Secondary effects of exchange

There are two “secondary” effects of exchange in the case of <sup>14</sup>N nitroxides: As  $\omega_e$  increases, (1) the outer lines broaden slightly faster than the central line and (2) the doubly-integrated intensity of the center line increases at the expense of the outer lines (Table 1).<sup>1,2</sup> The differential line broadening effect is small but easily measurable as shown in Figure 7 of Part 2<sup>2</sup> where the outer and central linewidths are plotted separately and compared with the rigorous theory. The differential intensity is also easily measurable and can be rather substantial. See Figure 6 of ref<sup>2</sup> that shows a dramatic difference in the outer and central lines and Figure 7 that compares the differential intensities to the rigorous theory. In water and in 70% glycerol at higher temperatures, we find (not shown) these expected secondary effects. At low values of  $T/\eta$ , all three lines broaden equally and all three intensities are equal, within experimental uncertainty. This confirms assumptions 1 and 5. In principle,  $\omega_e$ , and therefore  $\Omega(T)$ , could be estimated by these secondary effects.

It would be very useful to find a theoretically sound additional line shape component when dipolar is significant. If the search for a better line shape continues, it may be worth keeping in mind the possibility of negative dipolar-induced dispersion, even though these results do not prove it. It will be straightforward to test proposed shapes by including them into an equation such as eq 1. Unfortunately, as we have learned from fitting exchange broadened lines, it is not straightforward to discover such a shape from residuals such as those in Figures 2 and 3.<sup>28</sup> The non-linear least-squares fit finds a minimum in parameter space such that the residual does not resemble the missing component.<sup>28</sup>

### Is the theory of ref<sup>20</sup> valid?

One of the referees believes that we have observed results predicted by ref<sup>20</sup> in contrast to assumptions 1 – 5. Reference 20 predicts that when dipolar dominates, negative values of

$\pm \left[ \frac{V_{disp}^{obs}}{V_{pp}} \right]_{\pm}^{\#}$  arise and increasing  $B_{dip}$  causes lines to shift toward one another. Figure 12 where values of  $A_{abs}/A_0$  for 46.7 mM PDT in 70% glycerol 273 – 352 K are plotted as a function of

the mean values of  $\pm \left[ \frac{V_{disp}^{obs}}{V_{pp}} \right]_{\pm}^{\#}$ , shows both of these effects. At low temperatures, negative values

of  $\pm \left[ \frac{V_{disp}^{obs}}{V_{pp}} \right]_{\pm}^{\#}$  are correlated with line shifts toward the center as predicted by the theory in ref<sup>20</sup>. Noting the scale of the ordinate in the inset to Figure 12 shows that we would not have much to work with using the present data even if we had the full theory in ref<sup>20</sup>. Nevertheless the results in Figure 12 encourage the design of experiments extending to lower values of  $T/\eta$  and the development of the theory in ref<sup>20</sup> for three lines.

## Conclusions

The most important result from this work is that dipolar and exchange may be separated, Figure 8, but not by studying the linewidth alone. Non-linear least-squares fits of the spectra to a model of absorption (Gaussian-Lorentzian sum functions) plus spin-spin interaction induced dispersion yield nearly perfect fits in the limit of  $\Omega(T) \rightarrow 1$  and good, but not perfect fits in the opposite extreme  $\Omega(T) \rightarrow 0$ . From these fits, values of  $\Omega(T)$  may be deduced from the slopes of the dispersion component with respect to  $\langle B_{tot} \rangle$  from which  $\langle B_e \rangle$  may be computed.

Complicating the procedure, perhaps due to our imperfect knowledge of the proper line shape, is the fact that values of the dispersion do not become linear functions of  $\langle B_{tot} \rangle$  until  $\langle B_{tot} \rangle > 1$  G. Nevertheless, by carrying out the measurements to sufficiently high values of  $c$ , one may eliminate the influence of dipolar on the dispersion. As expected, at high values of  $T/\eta$ , we

find that exchange dominates while in the opposite extreme dipolar dominates.  $\frac{d}{dc} \langle B_e \rangle$  varies approximately linearly with  $T/\eta$  passing through the origin. The ratio of the values of

$\frac{d}{dc} \langle B_e \rangle$  to the SE prediction,  $\frac{d}{dc} \langle B_e \rangle / \frac{d}{dc} \langle B_e \rangle^{SE}$  is near unity for 70% glycerol and 0.73 for

water. Perhaps unexpectedly, in 70% glycerol, variations of  $\frac{d}{dc} \langle B_{dip} \rangle$  with  $\eta/T$  depart from linearity at rather low values following an approximately logarithmic dependence at higher values. This fact renders dipolar less useful to study diffusion than exchange. Results of line shifts and secondary effects – differential line broadening and line intensities – are consistent

with the separation procedure. Negative values of  $\pm \left[ \frac{V_{disp}^{obs}}{V_{pp}} \right]_{\pm}^{\#}$  accompanied by line shifts toward the center at low temperatures in 70% glycerol suggest that the effects of dipolar interaction are inadequately described by the widely accepted theory<sup>17</sup> and could be explained by a more recent theory<sup>20</sup>.

## Supplementary Material

Refer to Web version on PubMed Central for supplementary material.

## Acknowledgments

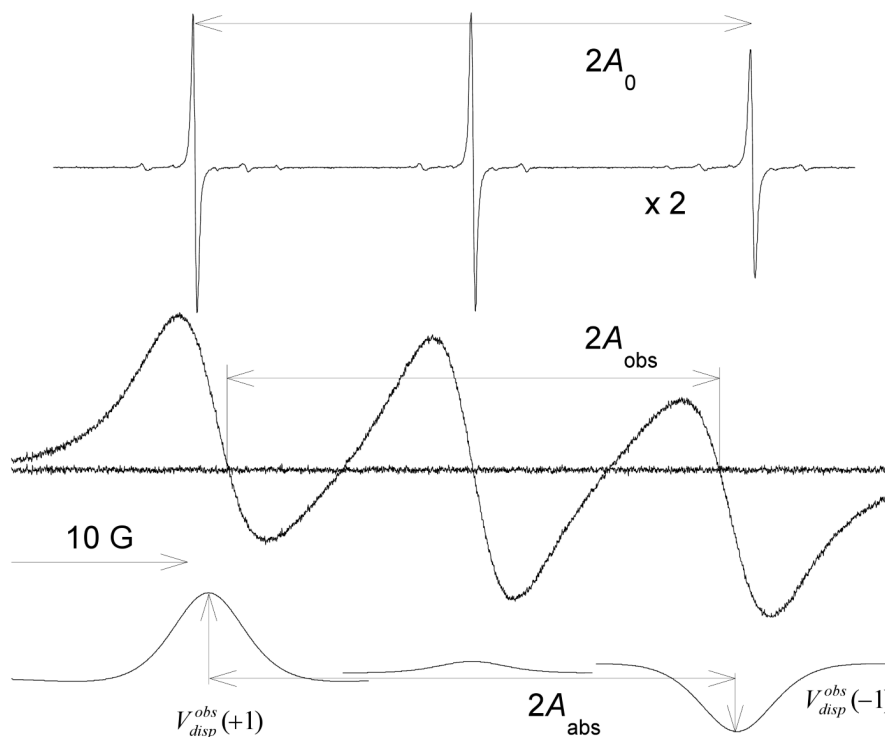
The authors gratefully acknowledge support from NIH grant 3S06 GM04680-10S1. We owe a debt of gratitude to two referees who obviously expended a great deal of effort to improve this paper and especially to one who alerted us to ref 20.

## References

1. Bales BL, Peric M. J Phys Chem B 1997;101:8707.
2. Bales BL, Peric M. J Phys Chem A 2002;106:4846.
3. Bales BL, Peric M, Dragutan I. J Phys Chem A 2003;107:9086.
4. Bales BL, Meyer M, Smith S, Peric M. J Phys Chem A 2008;112:2177. [PubMed: 18278887]
5. Kurban MR, Peric M, Bales BL. J Chem Phys 2008;129:064501. [PubMed: 18715079]
6. Bales BL, Stenland C. J Phys Chem 1993;97:3418.
7. Molin, YN.; Salikhov, KM.; Zamaraev, KI. Spin Exchange Principles and Applications in Chemistry and Biology. Vol. 8. Springer-Verlag; New York: 1980.
8. Halpern HJ, Peric M, Nguyen TD, Teicher BA, Lin YJ, Bowman MK. J Magn Reson 1990;90:40.
9. Nikonov AM, Nikonova SI. Sov J Chem Phys 1991;7:2408.
10. Robinson BH, Mailer C, Reese AW. J Magn Reson 1999;138:210. [PubMed: 10341124]
11. Robinson BH, Mailer C, Reese AW. J Magn Reson 1999;138:199. [PubMed: 10341123]

12. Salikhov KM. *J Magn Reson* 1985;63:271.
13. Berner B, Kivelson D. *J Phys Chem* 1979;83:1406.
14. Sachse JH, King MD, Marsh D. *J Magn Reson* 1987;71:385.
15. Currin JD. *Phys Rev* 1962;126:1995.
16. Bales, BL. Inhomogeneously Broadened Spin-Label Spectra. In: Berliner, LJ.; Reuben, J., editors. *Biological Magnetic Resonance*. Vol. 8. Plenum Publishing Corporation; New York: 1989. p. 77
17. Torrey HC. *Phys Rev* 1953;92:962.
18. Nevzorov AA, Freed JH. *J Chem Phys* 2000;112:1413.
19. Nevzorov AA, Freed JH. *J Chem Phys* 2000;112:1425.
20. Galeev RT, Salikhov KM. *Chem Phys Reports* 1996;15:359.
21. Bales BL, Peric M, Lamy-Freund MT. *J Magn Reson* 1998;132:279. [PubMed: 9632554]
22. Schreier S, Polnaszek CF, Smith ICP. *Biochim Biophys Acta* 1978;515:375.
23. Kovarskii AL, Wasserman AM, Buchachenko AL. *J Magn Reson* 1972;7:225.
24. Bales BL, Zana R. *J Phys Chem B* 2002;106:1926.
25. Bales BL, Benrraou M, Tiguida K, Zana R. *J Phys Chem B* 2005;109:7987. [PubMed: 16851934]
26. Bales BL, Ranaganathan R, Griffiths PC. *J Phys Chem B* 2001;105:7465.
27. Lang JC, Freed JH. *J Chem Phys* 1972;56:4103.
28. Peric M, Bales BL. *J Magn Reson* 2004;169:27. [PubMed: 15183353]



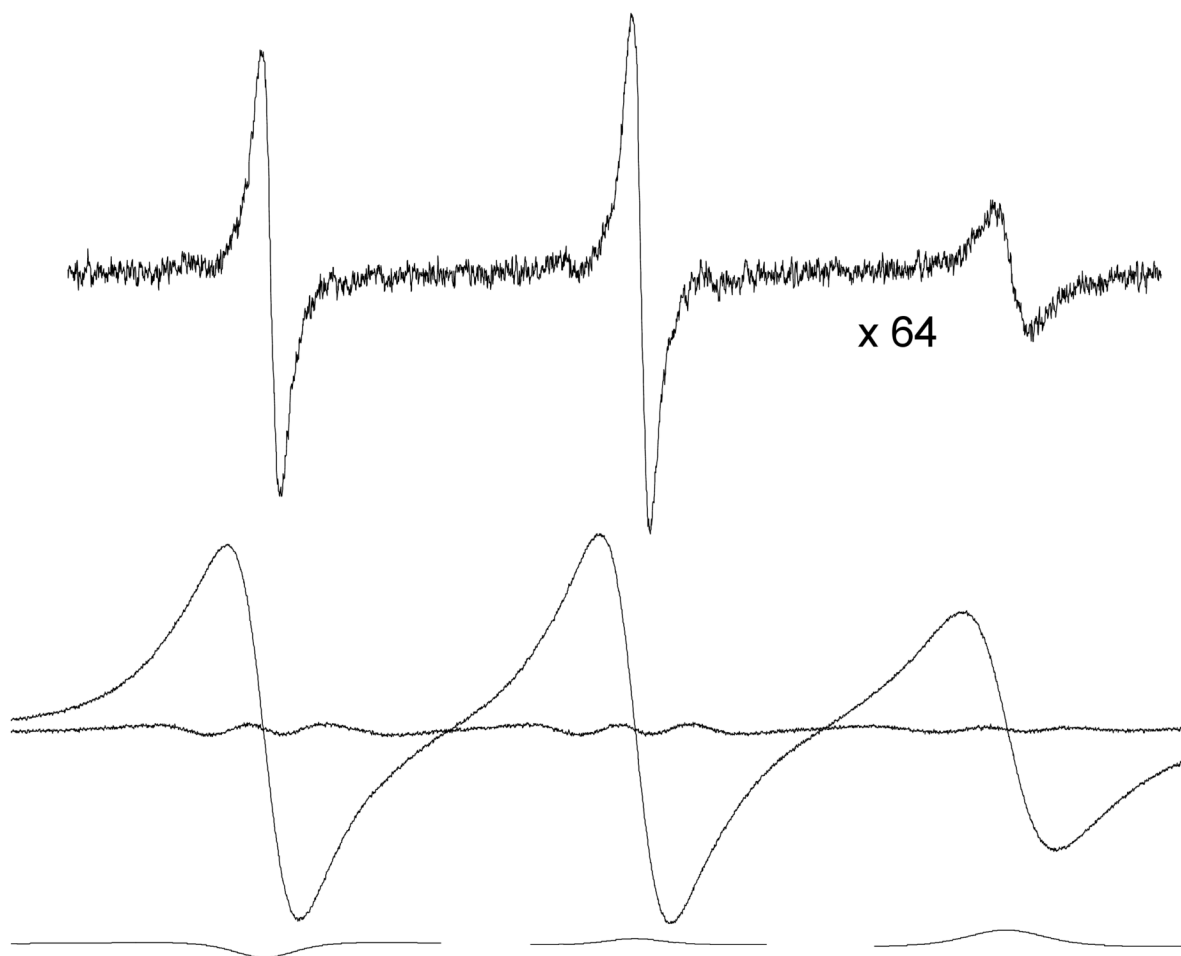


**Figure 1.**

EPR spectra due to 0.1 (top trace, at a relative gain of  $\times 2$ ) and 93.6 mM (middle trace) PDT in 70% glycerol at 343K. The fit of the middle trace to eq 1 is indistinguishable from the spectrum; the spectrum minus the fit gives the residual shown as the noisy baseline. The fit gives 3 absorption (not shown) and 3 dispersion lines that are given the lower trace. This is the “standard” presentation of the spectrum, yielding a positive absorption spectrum upon integration. The dispersion component on the central line is due to an imbalance in the microwave bridge; the component at low-field with a maximum value,  $V_{disp}^{obs}(+1) > 0$ , and the one at high-field with  $V_{disp}^{obs}(-1) < 0$  are expected from exchange. Exchange shifts the lines toward the center, resulting in line spacings  $2A_{obs}$  and  $2A_{abs}$  that are less than the spacing  $2A_0$  observed in the absence of exchange; the spectrum is dominated by exchange; the fractional broadening by spin exchange is  $\Omega(T) = 1.004 \pm 0.021$ .

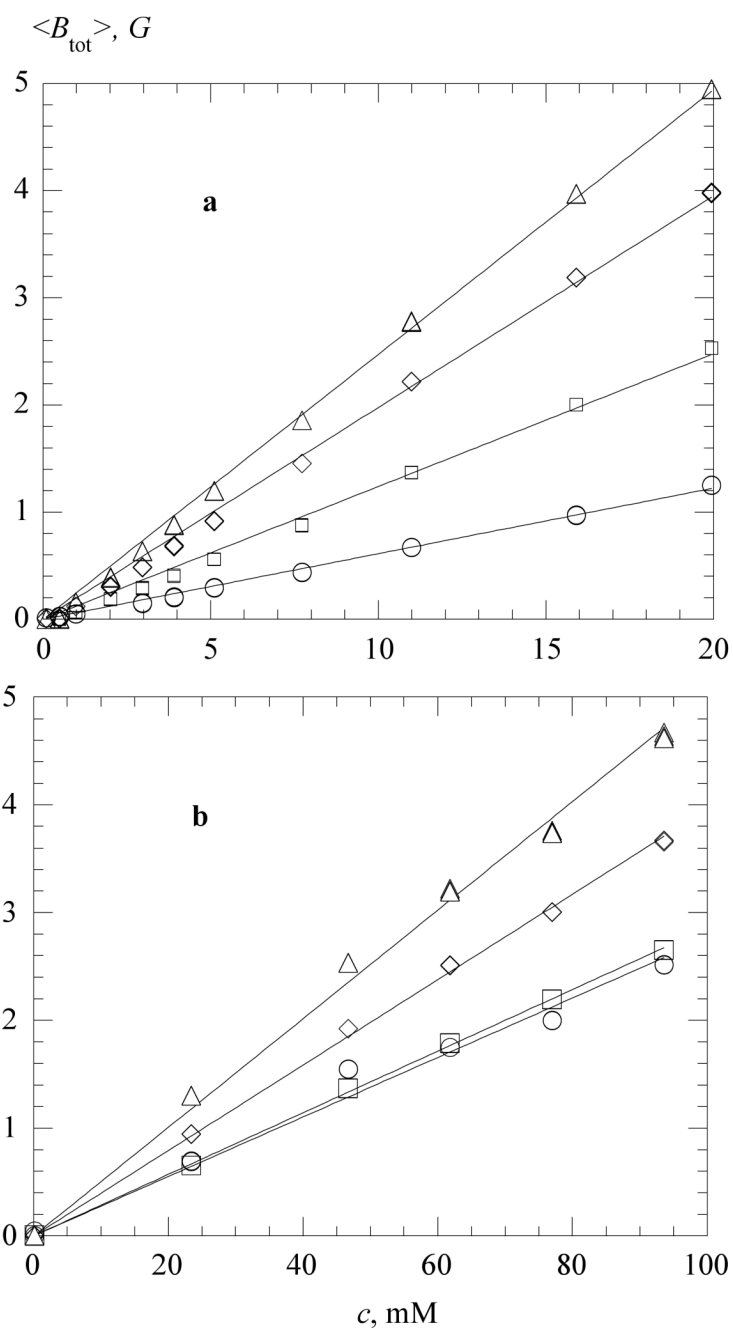


**Figure 2.** EPR spectra due to 0.1 (top trace, at a relative gain of  $\times 8$ ) and 93.6 mM (middle trace) PDT in 70% glycerol at 293 K. The fit and spectrum are indistinguishable; however, a definite residual is evident due to a significant dipolar interaction; the fractional broadening by spin exchange is  $\Omega(T) = 0.30 \pm 0.03$ .

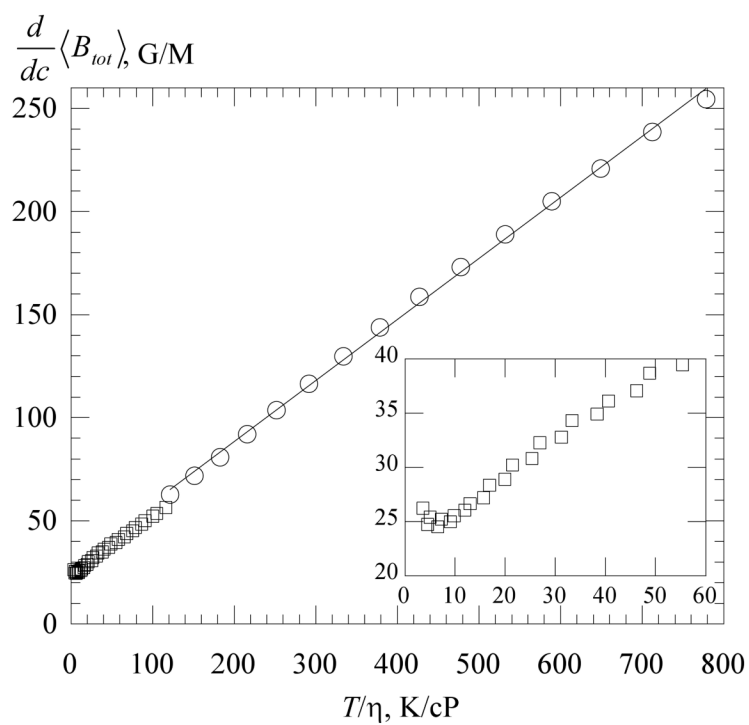


**Figure 3.**

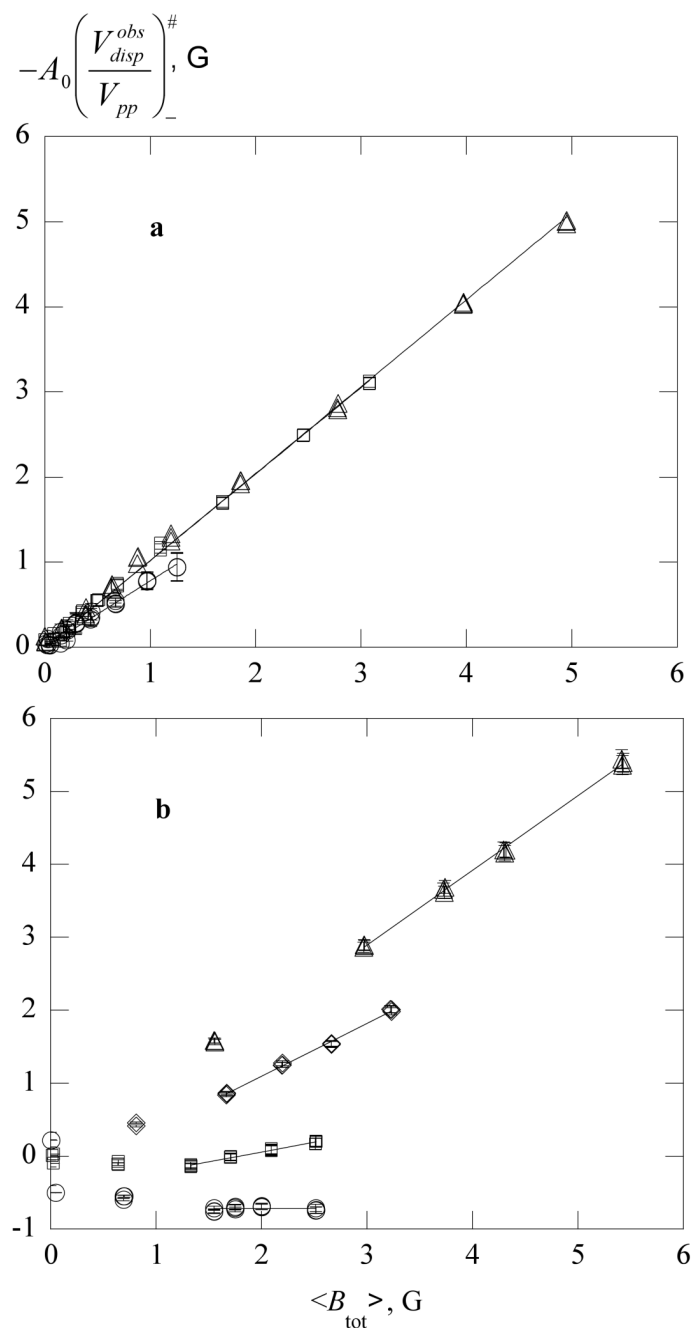
EPR spectra due to 0.1 (top trace, at a relative gain of  $\times 64$ ) and 93.6 mM (middle trace) PDT in 70% glycerol at 273 K. The fit and spectrum are indistinguishable; however, a definite residual is evident. Furthermore  $V_{\text{disp}}^{\text{obs}}(+1) < 0$  and  $V_{\text{disp}}^{\text{obs}}(+1) > 0$  the reverse of the signs expected from exchange. The contribution of exchange is negligible,  $\Omega(T) = 0.006 \pm 0.018$ .



**Figure 4.** Total linewidth versus PDT concentration (a) in water, at 268 K,  $\circ$ ; 298 K,  $\square$ ; 323K,  $\diamond$ ; and 338 K,  $\triangle$ ; and (b) in 70% glycerol, at 273 K,  $\circ$ ; 298 K,  $\square$ ; 323K,  $\diamond$ ; and 338 K,  $\triangle$ . The straight lines are linear least-square fits constrained to the origin. The slopes are plotted in Figure 5. Coefficients of correlation,  $r$ , are larger than 0.999 at all temperatures in water. In 70 % glycerol,  $r = 0.986$  (273 K),  $r > 0.999$  (298 and 323 K), and  $r = 0.986$  (338 K).



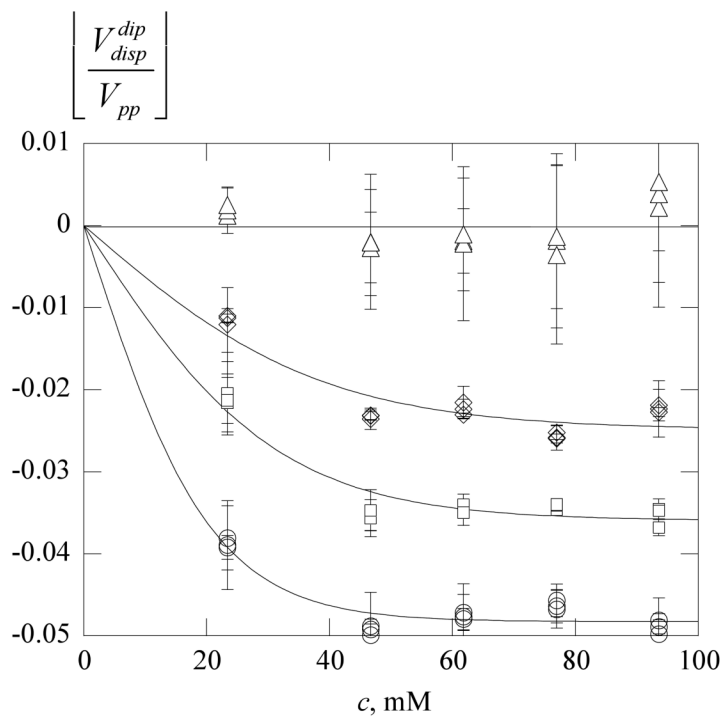
**Figure 5.** Total broadening constant versus for water,  $\circ$ ; and 70% glycerol,  $\square$ . The straight line through the water data is a linear least squares fit yielding  $\frac{d}{dc} \langle B_{tot} \rangle = (29.3 \pm 0.2) + (0.296 \pm 0.001) T / \eta$  in G/M with  $\eta$  in cP and  $r = 0.999$ . The data in 70 % glycerol show an upturn at low values of  $T / \eta$  (inset).



**Figure 6.**

Values of  $-A_0 \left( \frac{V_{disp}^{obs}}{V_{pp}} \right)^{\#}$  versus the total broadening in (a) water at 268 K,  $\circ$ ; 303 K,  $\square$ ; 338K,  $\triangle$ ; and in (b) 70% glycerol at 273 K,  $\circ$ ; 293 K,  $\square$ ; 313K,  $\diamond$ ; and 352 K,  $\triangle$ . The straight lines through the water data are linear least-squares fits constrained to the origin ( $r = 0.999$ , 303 and 338 K;  $r = 0.991$ , 268 K). Those through the 70% glycerol data are linear least-squares fits to the highest 4 concentrations ( $r = 0.994$ , 293 K;  $r = 0.998$ , 313 K; and  $r = 0.999$ , 352 K). In water at all temperatures and in 70 % glycerol above 323 K, linear fits constrained to the origin using all concentrations yield the same slopes as those using the highest four not constrained

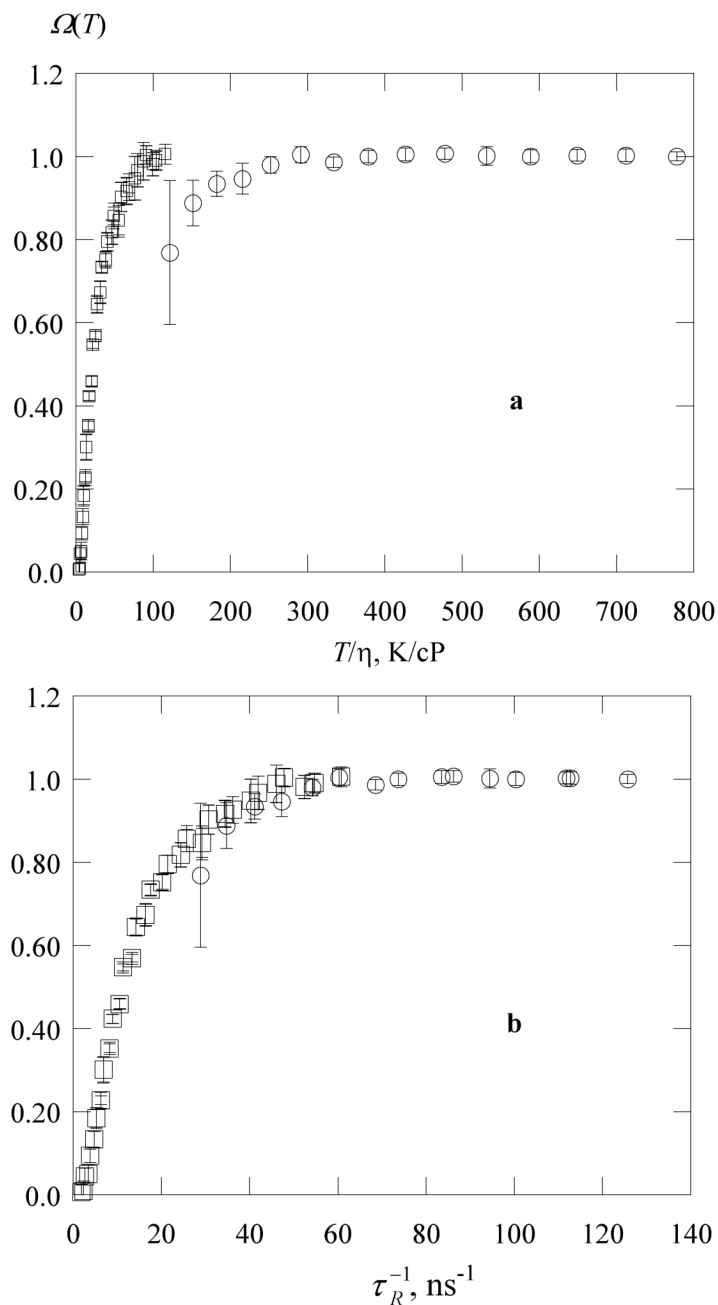
to the origin. Plots of  ${}^{+A_0} \left( \frac{V_{\text{disp}}^{\text{obs}}}{V_{pp}} \right)_+$  (not shown) are similar. The error bars are one-half the difference in  ${}^{\pm A_0} \left( \frac{V_{\text{disp}}^{\text{obs}}}{V_{pp}} \right)_\pm$ . Were it not for a dipolar contribution to the dispersion, the ordinate would be the broadening due to exchange,  $B_e$ , eq 7.



**Figure 7.** Dipole-induced dispersion in 70% glycerol versus concentration. Same symbols as Figure 6b. The average value from eq 20 is plotted; the error bars are one-half the difference. The solid lines are least-squares fits to eq 8 with  $c^* = 0, 21, 31,$  and  $39$  mM from high to low temperature,

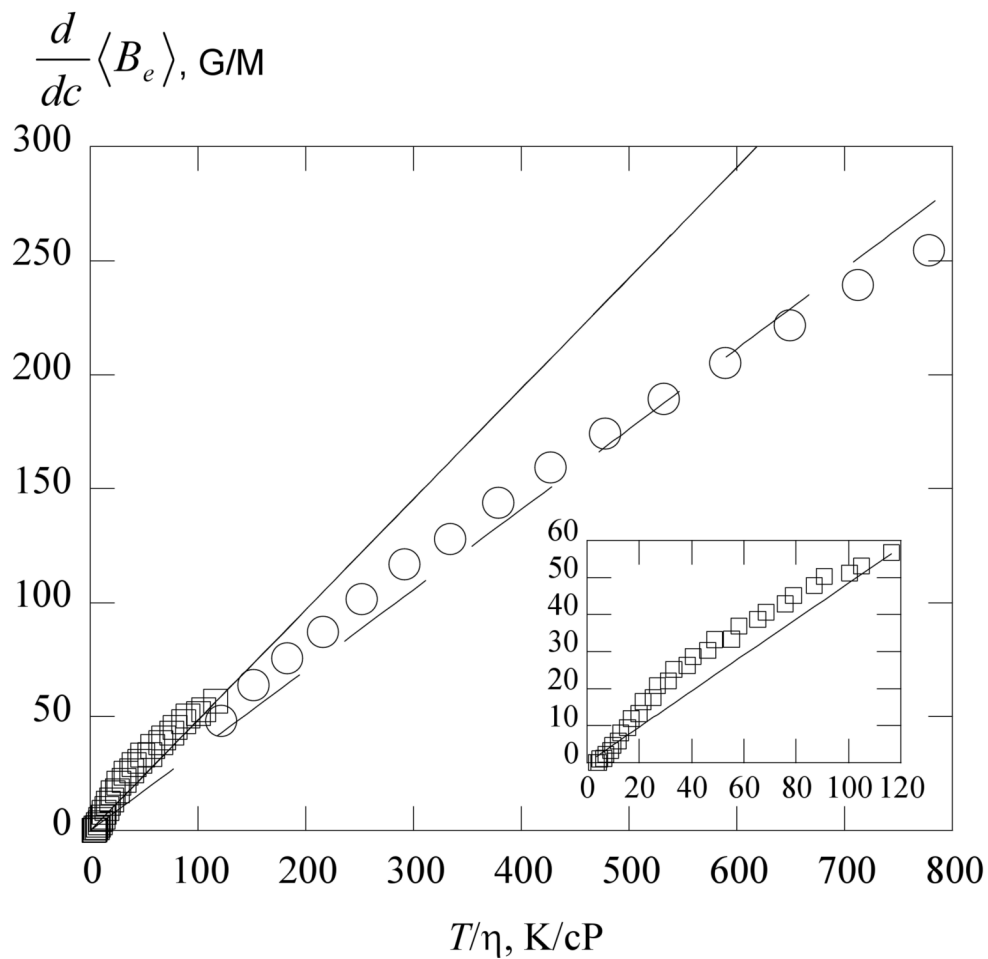
respectively. The plateau values  $\left[ \frac{V_{disp}^{dip}}{V_{pp}}(\infty) \right] = 0, -0.025, -0.036,$  and  $-0.048,$  respectively.



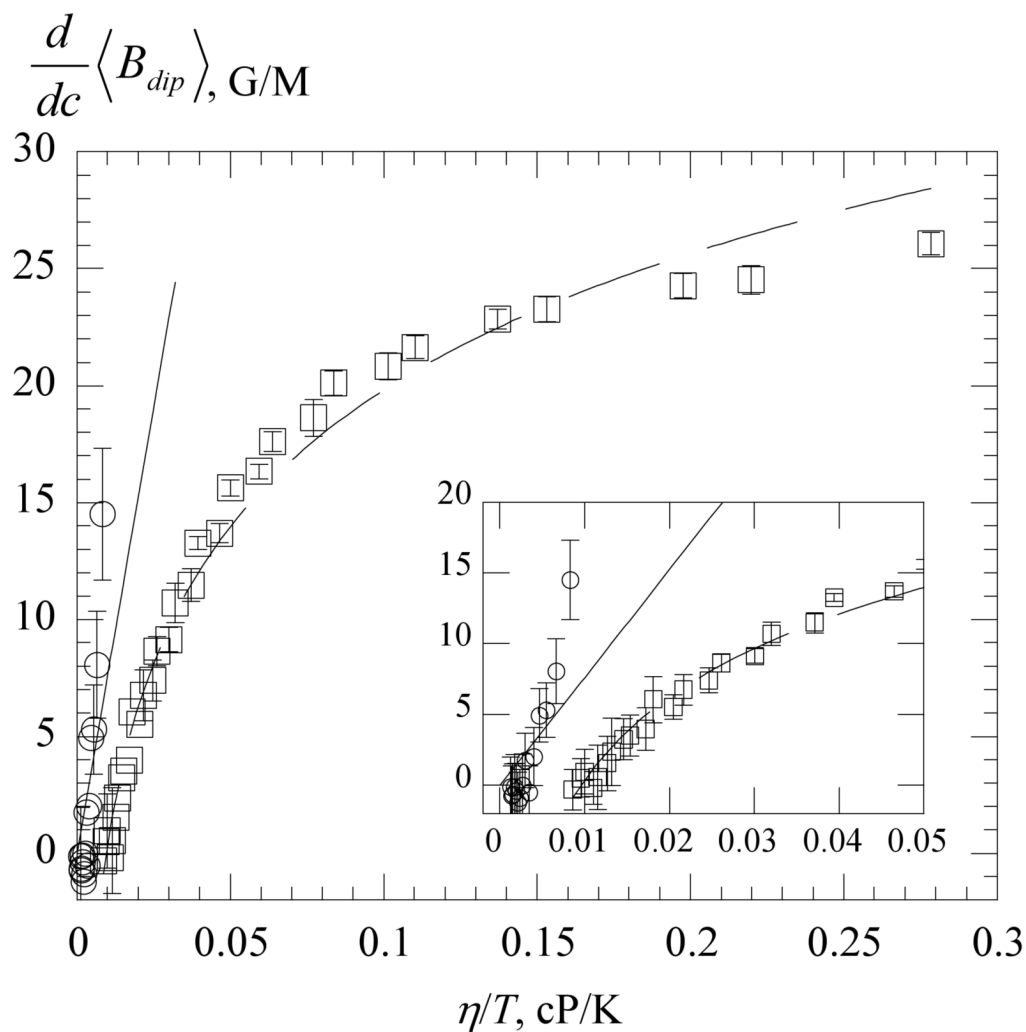


**Figure 8.**

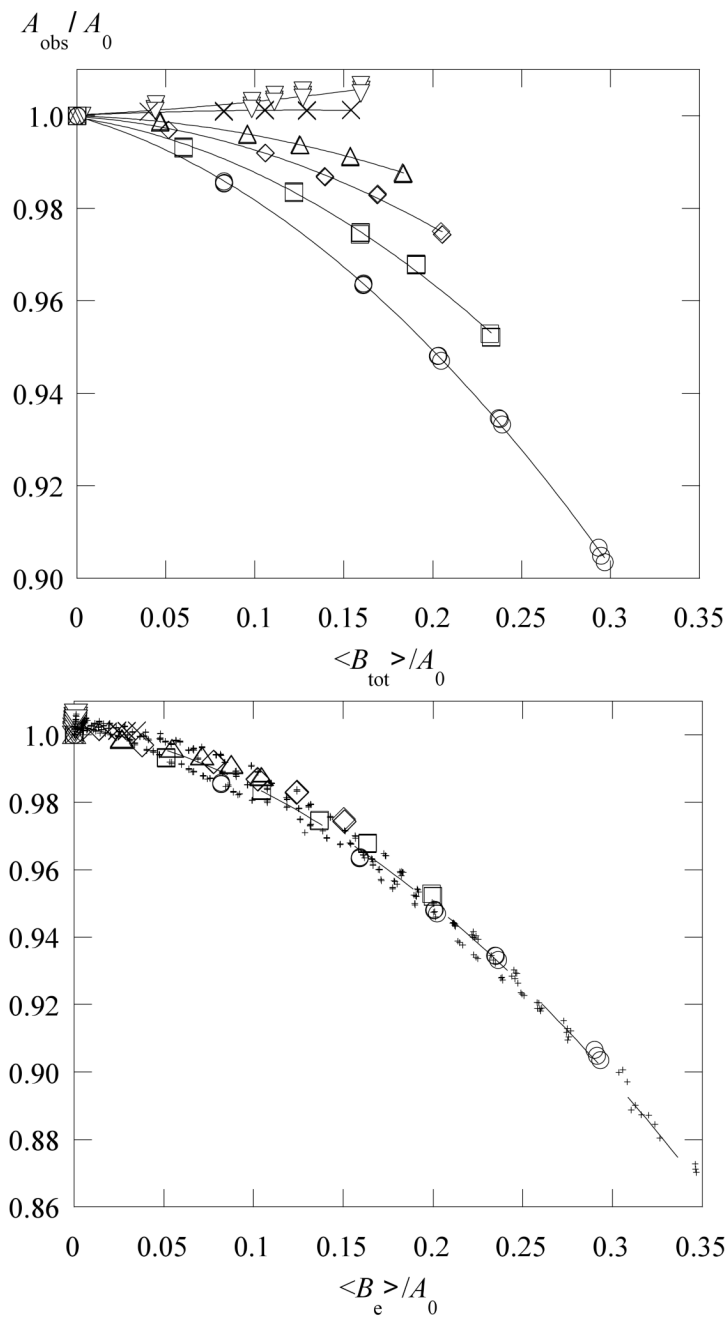
Fractional broadening by exchange of PDT in water, ○; and in 70% glycerol, □; versus (a)  $T/\eta$  and (b) the inverse rotational correlation time. The error bars are one-half the difference in the values of  $\Omega(T)_{\pm}$ . In 70% glycerol, exchange and dipolar broadening are equal at  $T/\eta = 20 \pm 1$  K/cP which corresponds to  $\eta = 15$  cP at 300 K.



**Figure 9.** Exchange broadening constant of PDT versus  $T/\eta$  in water, ○ and 70% glycerol, □. The solid line is the SE prediction, eq 15 and the dashed line is a linear least-squares fit to the water data.

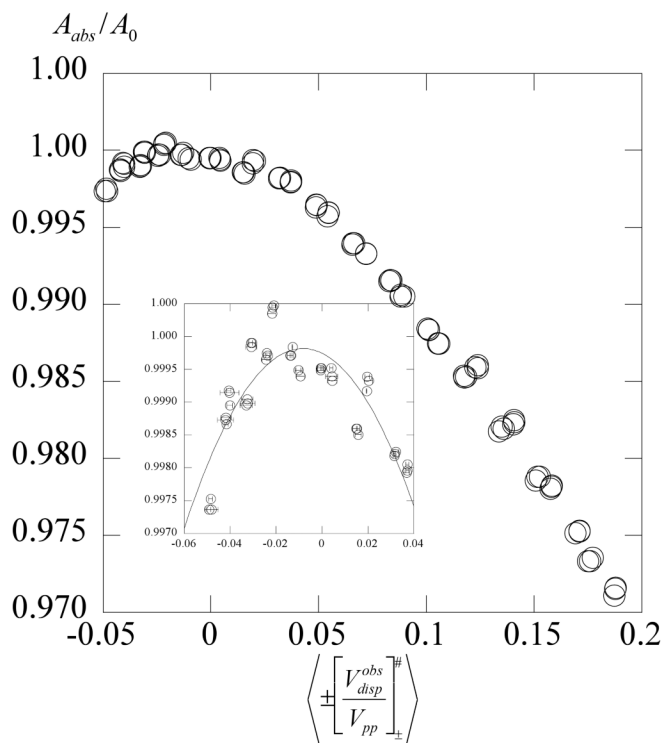


**Figure 10.** Dipolar broadening constant versus  $\eta/T$  of PDT water,  $\circ$  and 70% glycerol,  $\square$ . The solid line is the SE prediction, eq 16 and the dashed line is a linear least-squares fit of the 70% glycerol data to a logarithm. At 300 K, sodium dodecyl sulfate micelles fall at approximately 0.02 cP/K<sup>24</sup> and model membranes at about 0.21 cP/K.<sup>14</sup>



**Figure 11.**

The observed line spacing of PDT in 70% glycerol versus (a) the total line broadening at 273 K,  $\nabla$ ; 292 K,  $\times$ ; 307 K,  $\Delta$ ; 313 K,  $\diamond$ ; 323 K,  $\square$ ; and 342 K,  $\circ$ ; and versus (b) the exchange line broadening employing the same symbols. Data at all other temperatures are plotted in (b) as small plus signs. In (a), the solid lines are quadratic fits to guide the eye at each temperature and in (b) the dashed line is the fit of all data to eq 13. Equation 13 predicts a common curve versus the exchange broadening away from the origin as is observed.



**Figure 12.**

Line shifts versus the observed dispersion for 46.7 mM PDT in 70% glycerol 273 – 352 K. The abscissa is the mean of values measured from  $M_I = \pm 1$  corrected for instrumental dispersion; the error bars shown in the inset are the standard deviations. The solid line in the inset is a linear least-squares fit of the data to a quadratic. Spin exchange increases toward the right and dipolar toward the left. The lines shift toward one another both above and below the position of no shift near dispersion equal to zero. Negative values of the dispersion and line shifts toward the center of the spectrum at low temperatures are predictions of the theory in ref 20.

Table 1

Results of fitting spectra generated by the rigorous theory to eq 1. (a) <sup>14</sup>N nitroxides.

Dispersion and line shifts.					
$\frac{4}{3\sqrt{3}} \frac{\omega_e}{\gamma A_0}$	$\frac{\omega_e}{3\sqrt{3}} \frac{d}{\gamma A_0} \left( \frac{V_{ex}(\pm)}{disp(\pm)} \right) \frac{A_{abs}/A_0^b}{V_{pp}(\pm)}$	$\frac{A_{abs}/A_0^b}{b}$	$\frac{A_{obs}/A_0^c}{A_0}$	$I(\pm)^b$	$I(\pm)/I(0)_0$
0.0000	0.0000	1.0000	1.0000		
0.037890	0.037896	0.99960	0.99960	0.33334	1.0000
0.075781	0.076254	0.99838	0.99838	0.33307	0.99757
0.11367	0.11508	0.99636	0.99637	0.33225	0.99031
0.15156	0.15477	0.99352	0.99354	0.33088	0.97823
0.18945	0.19575	0.98985	0.98991	0.32893	0.96138
0.22734	0.23829	0.98536	0.98546	0.32636	0.93981
0.26523	0.28292	0.98002	0.98021	0.32314	0.91359
0.30312	0.32996	0.97382	0.97416	0.31920	0.88283
0.34101	0.38077	0.96676	0.96729	0.31447	0.84765
0.37890	0.43530	0.95881	0.95962	0.30887	0.80820
0.41679	0.49497	0.94995	0.95114	0.30227	0.76464
0.45468	0.56178	0.94018	0.94186	0.29456	0.71720
0.49257	0.63783	0.92946	0.93176	0.28555	0.66614
0.53046	0.72426	0.91779	0.92086	0.27505	0.61177
0.56835	0.82784	0.90515	0.90915	0.26280	0.55445
0.60624	0.95396	0.89153	0.89663	0.24853	0.49465
0.64413	1.1147	0.87693	0.88331	0.23187	0.43291
0.68202	1.3227	0.86137	0.86917	0.21242	0.36986
0.71992	1.6183	0.84490	0.85423		
0.75781	2.0847	0.82759	0.83849		
0.79570	2.9274	0.80958	0.82193		
0.83359	4.9254	0.79107	0.80457		
Broadening and intensities.					
$\frac{4}{3\sqrt{3}} \frac{\omega_e}{\gamma A_0}$	$\frac{\omega_e}{3\sqrt{3}} \frac{d}{\gamma A_0} \frac{B_e(0)}{A_0^b}$	$\frac{B_e(\pm)}{A_0^b}$	$\frac{B_e(\pm)/B_e(0)}{B_e(0)/A_0^b}$	$I(\pm)^b$	$I(\pm)/I(0)_0$
0.0000	0.0000	0.0000	1.0000	0.33334	1.0000
0.037890	0.037880	0.037896	0.037891	0.33307	0.99757
0.075781	0.075699	0.075822	0.075781	0.33225	0.99031
0.11367	0.11339	0.11381	0.11367	0.33088	0.97823
0.15156	0.15090	0.15189	0.15156	0.32893	0.96138
0.18945	0.18815	0.19010	0.18945	0.32636	0.93981
0.22734	0.22507	0.22848	0.22734	0.32314	0.91359
0.26523	0.26159	0.26706	0.26524	0.31920	0.88283
0.30312	0.29762	0.30588	0.30313	0.31447	0.84765
0.34101	0.33306	0.34500	0.34102	0.30887	0.80820
0.37890	0.36781	0.38446	0.37891	0.30227	0.76464
0.41679	0.40175	0.42433	0.41681	0.29456	0.71720
0.45468	0.43475	0.46468	0.45470	0.28555	0.66614
0.49257	0.46664	0.50558	0.49260	0.27505	0.61177
0.53046	0.49725	0.54712	0.53049	0.26280	0.55445
0.56835	0.52636	0.58941	0.56839	0.24853	0.49465
0.60624	0.55372	0.63258	0.60630	0.23187	0.43291
0.64413	0.57904	0.67678	0.64420	0.21242	0.36986

Dispersion and line shifts.			
$\frac{4}{3\sqrt{3}} \frac{\omega_e d}{\gamma A_0}$	$\left( \frac{V_{ex}(\pm)}{disp} \right) / \left( \frac{V_{pp}(\pm)}{b} \right)$	$A_{abs}/A_0^b$	$A_{obs}/A_0^c$
0.68202	0.60199	0.72217	1.1996
0.71992	0.62216	0.76895	1.2359
0.75781	0.63913	0.81735	1.2788
0.79570	0.65242	0.86759	1.3298
0.83359	0.66157	0.91992	1.3905
			0.18974
			0.16338
			0.15296
			0.098316
			0.059713
			0.61945
			0.67190
			0.73240
			0.80132
			0.12269
			0.068001

<sup>a</sup> Input values. Compare with column 7.

<sup>b</sup> Fitting simulated spectra to eq 1.

<sup>c</sup> Computed from the perturbation values  $A_{abs}/A_0 = 1 - \frac{9}{32} \left( \frac{B_e}{A_0} \right)^2$ .

Table 2

Results of fitting spectra generated by the rigorous theory to eq 1.  $^{15}\text{N}$  nitroxides.

$\frac{\omega_e}{\sqrt{3}\gamma A_0}$	$a$	$b$	$A_{abs}/A_0$	$A_{abs}/A_0$
0.0000	0.0000		1.0000	1.0000
0.011364	0.015155		0.99982	0.99981
0.022727	0.030326		0.99923	0.99923
0.034091	0.045530		0.99827	0.99826
0.045455	0.060802		0.99691	0.99690
0.056818	0.076124		0.99514	0.99516
0.068182	0.091555		0.99300	0.99303
0.079545	0.10706		0.99045	0.99051
0.090909	0.12276		0.98755	0.98760
0.10227	0.13857		0.98418	0.98431
0.11364	0.15456		0.98045	0.98063
0.12500	0.17074		0.97627	0.97656
0.13636	0.18710		0.97173	0.97211
0.14773	0.20368		0.96673	0.96727
0.15909	0.22069		0.96127	0.96204
0.17045	0.23795		0.95541	0.95642
0.18182	0.25547		0.94914	0.95041
0.19318	0.27332		0.94236	0.94402
0.20455	0.29156		0.93514	0.93724
0.21591	0.31033		0.92745	0.93007
0.22727	0.32973		0.91927	0.92252
0.23864	0.34928		0.91059	0.91458
0.25000	0.36999		0.90141	0.90625
0.26136	0.39069		0.89168	0.89753
0.27273	0.41272		0.88141	0.88843
0.28409	0.43512		0.87055	0.87894
0.30682	0.48331		0.84709	0.85879
0.31818	0.50850		0.83445	0.84814
0.32955	0.53506		0.82109	0.83710
0.34091	0.56311		0.80705	0.82567
0.35227	0.59276		0.79227	0.81386
0.36364	0.62421		0.77673	0.80165
0.37500	0.65721		0.76036	0.78906
0.38636	0.69319		0.74309	0.77608
0.40909	0.77295		0.70564	0.74897
0.43182	0.86739		0.66377	0.72030
0.45455	0.98294		0.61659	0.69008
0.47727	1.1309		0.56268	0.65832
0.50000	1.3333		0.50000	0.62500
0.52273	1.6416		0.42458	0.59013
0.54545	2.2188		0.32778	0.55372
0.56818	4.2669		0.17753	0.51575

<sup>a</sup> Input values. The broadening measured from the simulated spectra yields values of  $\left(\frac{B_e}{A_0}\right)$  that are within  $10^{-8}$  of  $\frac{\omega_e}{\sqrt{3}\gamma A_0}$ .

<sup>b</sup> Fitting simulated spectra to eq 1.

<sup>c</sup> Computed from the perturbation values  $A_{abs}/A_0 = 1 - \frac{3}{2}\left(\frac{B_e}{A_0}\right)^2$ .