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## A refinement protocol to determine structure, topology, and depth of insertion of membrane proteins using hybrid solution and solid-state NMR restraints

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### Abstract

To fully describe the fold space and ultimately the biological function of membrane proteins, it is necessary to determine the specific interactions of the protein with the membrane. This property of membrane proteins that we refer to as *structural topology* cannot be resolved using X-ray crystallography or solution NMR alone. In this article, we incorporate into XPLOR-NIH a hybrid objective function for membrane protein structure determination that utilizes solution and solid-state NMR restraints, simultaneously defining structure, topology, and depth of insertion. Distance and angular restraints obtained from solution NMR of membrane proteins solubilized in detergent micelles are combined with backbone orientational restraints (chemical shift anisotropy and dipolar couplings) derived from solid-state NMR in aligned lipid bilayers. In addition, a supplementary knowledge-based potential,  $E_z$  (insertion depth potential), is used to ensure the correct positioning of secondary structural elements with respect to a *virtual* membrane. The hybrid objective function is minimized using a simulated annealing protocol implemented into XPLOR-NIH software for general use.

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## Keywords

Hybrid method; Membrane protein; Molecular modeling; Structural topology; PISEMA; Solid-state NMR

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## Introduction

Nearly all high-resolution membrane protein structures deposited in the protein data bank (PDB) have been determined by X-ray crystallography, solution NMR, and cryoEM. While these techniques offer unparalleled atomic resolution that guides the interpretation of membrane protein biological function, the coordinates generated by the refinement procedures do not include a topological dimension. There are now several examples in the literature, including mechanosensitive channels (Vasquez et al. 2008; Wang et al. 2008), nucleobase-cation-symport-1 transporters (Weyand et al. 2008), small multidrug resistant proteins (Bay et al. 2008), and potassium ion channels, that show protein function to be dictated by changes in the relative orientation of secondary structural elements with respect to the membrane (structural topology) rather than changes in the protein's secondary structure. Furthermore, other classes of membrane proteins do not even adopt a compact tertiary structure; rather their fold space is determined by the intrinsic interactions with the lipid membrane. Since membrane protein topology plays a fundamental role in protein function, the failure to describe such interactions results in incomplete structural characterization (von Heijne 2006).

To mimic membrane protein environments, solution NMR spectroscopists utilize several different hydrophobic environments, including organic solvents, short chain lipid micelles, detergent micelles, or more recently, isotropic bicelles (Opella and Marassi 2004). The deleterious effects of organic solvents on membrane protein function are well-documented and researchers have now almost completely abandoned this avenue. On the other hand, detergent micelles present a viable alternative. There have been several examples of micellar reconstitutions that preserve membrane protein functional integrity, giving rise to high-quality NMR spectra for structural and interaction studies (Traaseth et al. 2008a, b; Zmoon et al. 2005). However, the intrinsic curvature of micelles, which may impact membrane protein structure and topology, is a significant concern (Chou et al. 2002). Small membrane proteins (30–200 residues), which account for most of the membrane proteins in several genomes, are especially vulnerable to changes induced by the membrane mimicking environment. While isotropic bicelles constitute an excellent membrane mimic and an alternative to micelles, the large size of the bicelle/protein complex results in substantially broader spectra than in detergent micelles (Poget and Girvin 2007; Poget et al. 2007; Sanders et al. 1994).

Solid-state NMR of membrane proteins in oriented lipid bilayers is one of the most accurate methodologies to directly investigate protein topology. Improvements in sample preparation and NMR methodology have provided higher resolution and sensitivity leading to topological and dynamic studies of multispan membrane proteins (De Angelis et al. 2006; Durr et al. 2007; Hallock et al. 2002; Opella and Marassi 2004; Park et al. 2006). However, these experiments rely on backbone measurements of chemical shift anisotropy (CSA) and dipolar coupling (DC) using separated local field (SLF) experiments [reviewed in (Ramamoorthy et al. 2004)] and heteronuclear correlation (HETCOR) spectroscopy (Ramamoorthy et al. 1999). The lack of side chain restraints has prevented complete structural characterization of membrane proteins using this technique.

A number of recent reports show that structure, topology, and oligomerization of membrane proteins are preserved both in lipid bilayers and in detergent micelles (Franzin et al. 2007; Mackenzie 2006; Moore et al. 2008; Stouffer et al. 2008). In some instances X-ray structures of membrane proteins are almost superimposable to NMR structures determined in micelles (Zhou et al. 2008). Therefore, a logical solution to the structure determination problem is to combine information from different experimental approaches into one unique structural refinement protocol that provides structure and topology simultaneously. Based on these considerations, we propose a method that combines the high-resolution information obtained from solution NMR on membrane protein samples reconstituted in detergent micelles with the structural data obtained by solid-state NMR on proteins reconstituted into lipid bilayers. In this work, we illustrate how to implement both distance and orientational restraints into a single target function. A depth of insertion potential from DeGrado and co-workers is also employed to embed helical segments into a virtual bilayer (Senes et al. 2007). The energy landscape is explored using the simulated annealing protocol implemented in XPLOR-NIH software (Schwieters et al. 2003) to determine the high-resolution structure and topology of membrane proteins simultaneously. We show the application of this method to monomeric phospholamban (PLN), a single pass membrane protein involved in cardiac muscle physiology (Traaseth et al. 2008a, b). To obtain monomeric PLN, we mutated the three cysteines in the transmembrane domain (C36A, C41F, and C46A). These mutations maintain the functional integrity of the protein (Zamoon et al. 2003). The solution NMR data was taken from the dodecylphosphocholine (DPC) studies from Zamoon et al. (2003), while the solid-state NMR restraints were derived from the work in mechanically aligned lipid bilayers by Traaseth et al. (2006, 2009) and Mascioni et al. (2002).

## Methods and results

### Energy terms in structural refinement

The hybrid solution and solid-state NMR target function ( $E_{\text{total}}$ ) is formulated as a linear combination of geometrical ( $E_{\text{chem}}$ ), solution NMR ( $E_{\text{sol-NMR}}$ ), and solid-state NMR ( $E_{\text{ssNMR}}$ ) terms:

$$E_{\text{total}} = E_{\text{chem}} + E_{\text{sol-NMR}} + E_{\text{ssNMR}} \quad (1)$$

The geometrical and solution NMR potentials were available in XPLOR-NIH force field version 2.18 (Schwieters et al. 2003).  $E_{\text{chem}}$  is the sum of bonding ( $E_{\text{bonds}}$ ,  $E_{\text{angles}}$ ,  $E_{\text{improper}}$ ) and non-bonding ( $E_{\text{vdw}}$ ) interactions, with adjustable weighting factors ( $w$ ),

$$E_{\text{chem}} = w_{\text{bonds}} E_{\text{bonds}} + w_{\text{angles}} E_{\text{angles}} + w_{\text{improper}} E_{\text{improper}} + w_{\text{vdw}} E_{\text{vdw}} \quad (2)$$

$E_{\text{sol-NMR}}$  is the sum of restraints from solution NMR experiments such as NOEs ( $E_{\text{NOE}}$ ), hydrogen bonds ( $E_{\text{HBON}}$ ), torsion angles ( $E_{\text{CDIH}}$ ), and an empirical torsion angle potential from a database ( $E_{\text{DB}}$ ; Kuszewski et al. 1996, 1997):

$$E_{\text{sol-NMR}} = w_{\text{NOE}} E_{\text{NOE}} + w_{\text{CDIH}} E_{\text{CDIH}} + w_{\text{DB}} E_{\text{DB}} + w_{\text{HBON}} E_{\text{HBON}} \quad (3)$$

To include both CSA and DC data obtained from SLF solid-state NMR experiments such as PISEMA (polarization inversion spin exchange at the magic angle) (Wu et al. 1994), we used the penalty function  $E_{\text{ssNMR}}$ :

$$E_{\text{ssNMR}} = w_{\text{CSA}} E_{\text{CSA}} + w_{\text{DC}} E_{\text{DC}} \quad (4)$$

$$= w_{\text{PISEMA}} (E_{\text{CSA}} + w_{\text{T}} E_{\text{DC}}).$$

where  $w_{\text{T}}$  is the relative weighting between CSA and DC and  $w_{\text{PISEMA}}$  is the weighting of PISEMA potential both for CSA and DC.

To implement both CSA ( $E_{\text{CSA}}$ ) and DC ( $E_{\text{DC}}$ ) potentials, we used flat-well penalty functions as reported by Bertram et al. (2000):

$$E_{\text{CSA}} = \begin{cases} \sum_i (|\sigma_{c,i} - \sigma_{o,i}| - \sigma_{\text{Error},i})^2 & \text{if } |\sigma_{c,i} - \sigma_{o,i}| > \sigma_{\text{Error},i} \\ 0 & \text{if } |\sigma_{c,i} - \sigma_{o,i}| \leq \sigma_{\text{Error},i} \end{cases} \quad (5)$$

$$E_{\text{DC}} = \begin{cases} \sum_i (|v_{c,i} - v_{o,i}| - v_{\text{Error},i})^2 & \text{if } |v_{c,i} - v_{o,i}| > v_{\text{Error},i} \\ 0 & \text{if } |v_{c,i} - v_{o,i}| \leq v_{\text{Error},i} \end{cases} \quad (6)$$

$\sigma_{o,i}$  ( $\sigma_{\text{Error},i}$ ) and  $v_{o,i}$  ( $v_{\text{Error},i}$ ) are the experimental data (error) for CSA and DC and  $\sigma_{c,i}$ ,  $v_{c,i}$  are calculated values associated with residue  $i$ . In our implementation, this potential energy function allows for structural refinement using different experimental errors and CSA tensor values for individual residues. Similar energy functions for CSA and DC have been used to refine the backbone structures of membrane proteins (Im and Brooks 2004; Ketchem et al. 1997, 1996; Lee et al. 2008; Nevzorov and Opella 2003).

## Calculation protocol

**Step 1: implementation of solution NMR restraints**—In step 1, we used solution NMR data with geometrical restraints to define the secondary structural elements of monomeric PLN (Fig. 1).

- A. The starting structure of monomeric PLN (the AFA-PLN mutant was used; C36A, C41F, C46A) was in an extended configuration, according to the simulated annealing protocol defined by Nilges et al. (1988).
- B. Simulated annealing was carried out using the  $E_{\text{chem}}$  and  $E_{\text{sol-NMR}}$  potentials from Eq. 1. The system was cooled from 6,000 to 0 K in 5 K increments using 200 steps of torsion angle dynamics at each temperature. Both temperature and number of steps were optimized to achieve the lowest number of violations in the conformers generated. Note: the dihedral angle restraints for these calculations were generated from N, HN,  $H\alpha$ ,  $C'$ ,  $C\alpha$ ,  $C\beta$  PLN chemical shifts using TALOS version 98.040.21.02 (Cornilescu et al. 1999).
- C. 200 steps of molecular dynamics were performed in torsion angle space followed by minimization in both torsion angle and Cartesian spaces.
- D. At the conclusion of step 1, we generated 200 structures of AFA-PLN, where the 100 lowest energy structures were used for conformational analysis.

Similar to the conclusions by Zamoon et al. (2003), step 1 generates an ensemble of conformers with good convergence for the secondary structure elements, but that lacks precision in the structural overlay (see companion paper in Traaseth et al. 2009).

**Step 2: implementation of solid-state NMR restraints**—In step 2, we refined the structures obtained in step 1 with the solid-state NMR data (DC and CSA) using tensor 1 defined in Table 1.

- A. Starting from the ensemble of 100 lowest energy structures, simulated annealing calculations were carried out using all potential energy terms in Eq. 1 [ $E_{\text{chem}}$ ,  $E_{\text{sol-NMR}}$ ,  $E_{\text{ssNMR}}$ ]. The initial annealing temperature was 3,000 K. Here, only torsion angle dynamics were applied to optimize the orientation of each individual peptide plane of AFA-PLN with respect to the  $Z$ -axis (corresponding to the lipid bilayer normal).
- B. Cross-validation of CSA and DC restraints. To avoid overfitting the experimental data, we first optimized the solid-state NMR weighting factors ( $w_{\text{PISEMA}}$  and  $w_r$ ) in Eq. (4) based on the cross-validation factor  $R$  described by Cross and co-workers (Kim et al. 2001).

$$R = \frac{1}{N} \sum_{i=1}^N \left( \frac{|\text{calculated}_i - \text{observed}_i|}{\text{error}_i} \right)^2 \quad (7)$$

$N$  is the number of data points used in the refinement protocol, while *calculated<sub>i</sub>*, *observed<sub>i</sub>* and *experimental error<sub>i</sub>* are the calculated value, experimental value, and experimental error in CSA and DC for residue  $i$ . We fixed the weighting factors for  $E_{\text{chem}}$  and  $E_{\text{sol-NMR}}$ , and randomly partitioned the experimental data into two datasets: working (80% of data) and free (20% of data). The working dataset is included in the simulations, while the free dataset is not used and is back-calculated from the model. The work  $R$  ( $R_{\text{work}}$ ) indicates the fitting quality of the data included in the refinement, while the free  $R$  ( $R_{\text{free}}$ ) indicates the agreement between experimental and calculated data for residues with no CSA or DC restraints in the refinement protocol. In Fig. S1, we varied two factors: (1)  $w_{\text{PISEMA}}$  from 0.1 to 1,000 for a total of ten independent calculations and (2)  $w_r$ , the relative weight ratio of the DC/CSA potential, from 1/1 to 9/1. The  $R_{\text{work}}$  and  $R_{\text{free}}$  for CSA and DC were analyzed based on statistics for the 20 lowest energy structures from each of the calculations, as shown in Fig. S1. As  $w_{\text{PISEMA}}$  becomes too large ( $>10$ ), the quality of the structures decreases ( $E_{\text{bonds}}$ ,  $E_{\text{angles}}$ ,  $E_{\text{improper}}$ , and  $E_{\text{vdw}}$  become high) and the improvement in the  $R$  values is minimal. Similarly, when  $w_r$  is  $>3$ , the geometrical and solution NMR penalties increase dramatically without improving the PISEMA  $R$  factor. When  $w_r = 1$ , there is a large DC  $R_{\text{free}}$  factor, indicating poor correlation with the experiments. Based on Fig. S1, the optimal weighting factor were found to be 5 for  $w_{\text{PISEMA}}$  and 3 for  $w_r$ , which is in agreement with the 10 ( $w_{\text{PISEMA}}$ ) determined by Cross and co-workers (Kim et al. 2001).

- C. Using optimized weighting factor, a total of 400 structures were generated with the 100 lowest energy conformers selected for further analysis.

To evaluate the effects of the solid-state NMR restraints on refinement of the solution NMR structural ensemble, we defined three angles shown in Fig. 2:  $\theta$  (tilt angle with respect to the  $Z$  axis or membrane normal),  $\rho$  (azimuthal rotation angle around helix axis), and  $\chi$  [interhelical angle between domain Ia (residues 1–16) and the helix comprised of domains Ib and II (residues 23–52)]. The distributions of  $\theta$  and  $\rho$  for the two helical domains in the 100 lowest energy conformers are shown in Fig. 3a. Noticeably, the orientation of the helix described by domains Ib and II is well-defined, with  $\theta_{\text{Ib,II}} = 23 \pm 3^\circ$  and  $\rho_{\text{Ib,II}} = 205 \pm 3^\circ$ . The tilt angle of domain Ia is somewhat less defined with  $\theta_{\text{Ia}} = 95 \pm 7^\circ$ . While the solution NMR ensemble (Fig. 3b) fails to define the  $\chi$  angle, the incorporation of CSA and DC

restraints drastically reduces the conformational space allowed, confining  $\chi$  between  $70^\circ$  and  $125^\circ$ . In contrast, the rotation angle ( $\rho_{1a}$ ) for domain Ia is not well-defined using PISEMA data alone. This originates from several different factors: (a) helices with tilt angles  $\sim 90^\circ$  have inherently clustered CSA (65–82 ppm) and DC (4–5 kHz) values, (b) helices with  $\theta \sim 90^\circ$  result in degenerate PISEMA spectra for  $\rho = 0^\circ$  or  $180^\circ$ , and (c) the dynamic nature of domain Ia results in broad resonances (Metcalf et al. 2005; Traaseth et al. 2006). While these spectroscopic difficulties were partially overcome using several selectively labeled samples (see Traaseth et al. 2009), in order to determine the depth of insertion of the protein and fully resolve the orientation of the  $\rho_{1a}$  angle, we needed to incorporate a knowledge-based  $E_z$  potential (step 3).

### Step 3: incorporation of the knowledge-based $E_z$ depth of insertion potential

—The knowledge-based  $E_z$  potential by Degradó and coworkers (Senes et al. 2007) was used to embed the helical domains of AFA-PLN into a virtual bilayer. The  $E_z$  pseudo-energy was derived from the propensity of each amino acid in helical membrane proteins to insert into a lipid bilayer. This knowledge-based potential was derived from the statistical analysis of 24 crystal structures of helical membrane proteins. In our calculation, the helical domains were independently optimized under the knowledge-based  $E_z$  potential using rigid-body minimization, where the only degree of freedom is the translation of an entire domain along the  $Z$ -axis. Note that the loop connecting domain Ia with domain Ib was allowed to move and adjust the depth of insertion for all regions of AFA-PLN. The incorporation of the  $E_z$  term into the total energy function for simulated annealing calculations resulted in erroneous insertion of the protein in the virtual membrane irrespective of the weighting factor used (data not shown). This results from the step-function shape of the knowledge-based  $E_z$  potential which acts on each residue only along the  $Z$  axis. While this potential applied to single residues is not sufficient to embed the protein in the virtual membrane, the sum of the  $E_z$  potentials for all of the residues in a helix results in a unique minimum. Below we summarize the steps used to incorporate this knowledge-based potential into our calculations.

- A. Definition of virtual bilayer. A virtual bilayer was built with the membrane normal parallel to the  $Z$ -axis and the origin in the center of the hydrocarbon core. The  $Z$  coordinate of the COM (center of mass) for each helical domain (using  $C_\beta$  atoms since  $E_z$  potential depends only on  $C_\beta$  coordinates) describes its insertion into the virtual bilayer.
- B. Starting with the 100 lowest energy structures obtained after step 2, we minimized the domain insertion along the  $Z$  coordinate by rigid-body minimization using the  $E_z$  potential (i.e., the helical regions of domain Ia (residues 3–15) and domains Ib and II (residues 23–47) were held rigid). The use of the  $E_z$  potential helped to resolve the rotation angle  $\rho_{1a}$  ambiguity. Only structures with  $\rho_{1a}$  within the range of  $60$ – $150^\circ$  gave reasonable minimized structures (Fig. S2). These angles are consistent with results from paramagnetic quenching experiments and NOEs connecting the hydrophobic face of domain Ia to the detergent micelle (Traaseth et al. 2008a, b; Zamoan et al. 2003). The introduction of the  $E_z$  potential confined the trans-membrane helix into a well-defined minimum, while the depth of insertion of the cytoplasmic domain was less defined with respect to the *virtual* membrane (Fig. 4). Regardless, the translational degeneracy was removed, resulting in one population of conformers with the hydrophilic residues pointing toward the bulk water, in agreement with the amphipathic nature of the cytoplasmic helix. Note that domain Ib in the N-terminal portion of the transmembrane domain protrudes outside the virtual membrane. This is in agreement with the hydrophilic nature of this domain and explains the hydrogen/deuterium exchange data (Zamoan et al.



2003). The flexible loop is also exposed to the solvent, which is in agreement with the dynamics probed by NMR spin relaxation measurements (Metcalf et al. 2004).

- C. In order to relax local geometries and improve the structure quality, we subjected the remaining structures to low temperature (300 K) simulated annealing using torsion angle and Cartesian molecular dynamics with full van der Waals interactions and all of the other restraints, but excluding the  $E_z$  potential. Although the knowledge-based  $E_z$  potential was not used during this annealing step, we verified that the structures did not deviate from the depth-of-insertion minima shown in Fig. 5b. From this ensemble, we selected 20 low energy conformers of PLN with no restraint violations, and deposited the structures into the protein data bank (*PDB 2KB7*, see Traaseth et al. 2009). The biological relevance of structures is discussed in our companion paper (Traaseth et al. 2009). Structural statistics are provided in Table S1.

### Effects of different tensor components

One debate in the field surrounds the  $^{15}\text{N}$  CSA and  $^1\text{H}$ - $^{15}\text{N}$  DC tensor components used to model solid-state NMR data obtained from PISEMA-type experiments. To address this issue, we carried out a systematic conformational search for AFA-PLN domain Ib and II utilizing different tensor components shown in Table 1. We performed Steps 1–2 with low temperature annealing as discussed in the “Methods and results” section for each of the three tensors given in Table 1. Following energy minimization, the 20 lowest energy structures are selected with good agreement between the experimental and calculated DC and CSA using the three different tensors (Fig. 6). While small changes in local structure may occur, differences in the tensor values only marginally affect the overall protein topology. As shown in Table 1, changing the tensor components (tensors 2 and 3) from those used for the majority of this work (tensor 1), caused the tilt ( $\theta$ ) and azimuthal ( $\rho$ ) angles to vary by  $2.3^\circ$  and  $1.1^\circ$ , respectively. Given the approximations necessary to calculate these angles for the transmembrane domain of PLN, these variations are negligible. In addition to the overall orientation, we also examined effects of different tensors on the local structures. Due to the nature of structural calculation, the bond angle, bond length and peptide plane planarity are all within acceptable errors of peptide geometry. We thus compared the Ramachandran angle ( $\phi$ ,  $\psi$ ) for the transmembrane domain shown in Fig. S3. Although structures remained helical from all sets of simulations, there were considerable variations in these angles. These are likely due to the inclusion of different CSA and DC tensors.

Note that for glycine residues, a different set of tensor components needs to be used (Straus et al. 2003).

### Sensitivity to PISEMA misassignments

Solid-state NMR methods to obtain orientational information from assigned PISEMA resonances are well-established. However, the assignment procedures for *real* membrane protein samples are still in their infancy. Although PISA wheels are a common way to assign resonances in well-dispersed spectra, most experimental PI-SEMA spectra contain inhomogeneous broadening of resonances (due to mosaic spread and protein dynamics) making the assignments challenging. In spite of several computational attempts, these effects have not been fully rationalized (Quine et al. 2006; Straus et al. 2003). Therefore, selective labeling techniques constitute an important resource for resonance assignment. Unfortunately, if residues are on the same face of a helix, degenerate frequencies will be expected, further hampering the assignment procedure. To assess the sensitivity of our method to this problem, we have carried out two different calculations with two equiprobable assignments for the [ $^{15}\text{N}$ -Ile] AFA-PLN PISEMA spectrum shown in Fig. 7b (see Traaseth et al. 2009). The assignment was carried out using a combinatorial search

routine as described by Buffy et al. (2006). Using our hybrid method, we found that the two resonance assignments result in two conformational ensembles with similar tilt angles ( $24 \pm 3^\circ$  and  $22 \pm 2^\circ$ ), but different azimuthal rotation angles ( $203 \pm 4^\circ$  and  $189 \pm 3^\circ$ ). The conformational ensemble with the incorrect assignment (black assignment in Fig. 7) displayed covalent geometry violations at Ile45 for more than half of the structures, while also giving consistently higher conformational energy. This demonstrates that our hybrid method is sensitive to misassignment, and that it can be used to assist the assignment process.

### Static helix approach

PISEMA data has been used to calculate the tilt and rotation angles of helices based on a static ideal helix (Marassi and Opella 2002). In Fig. 8, we show the agreement between the tilt and rotational angles with the PISEMA data using a static approach. The potential energy landscape represented in Fig. 8 was obtained by calculating the least square fit between the experimental PISEMA data and an ideal poly-alanine  $\alpha$ -helix. The lowest energy conformers generated with the hybrid method are located at the bottom of the minima defined by tilt and rotation angles for domains Ia and the helix comprised of domains Ib and II.

This indicates that (1) the lowest energy conformers generated with the hybrid method are in remarkable agreement with the experimental data and (2) there is negligible discrepancy between the static fitting of the PISEMA data and a more refined conformational search procedure such as simulated annealing. A corollary to these considerations is the ideal nature of transmembrane helices embedded in lipid bilayers, demonstrated in several studies by Cross and co-workers using numerous membrane proteins (Page et al. 2008). Nevertheless, compared to the fitting with a rigid, ideal helix, the hybrid approach has the advantage of obtaining a high-resolution structure for both the backbone and side chains.

### Discussion

The correlation between structure and function of complex biological systems requires the combination of several different techniques (biophysical and biochemical) to validate any mechanistic conclusions. The combination of methodologies, *hybrid methods*, is becoming quite common in structural biology (Cowieson et al. 2008). For instance, Clore and co-workers have used X-ray crystallographic data synergistically with solution NMR restraints to refine the structure of soluble proteins (Shaanan et al. 1992). More recently, solution NMR and small angle X-ray scattering (SAXS) data have been combined to refine the structures of large complexes and to resolve the intrinsic ambiguities of residual dipolar couplings (Gabel et al. 2006, 2008). Others have combined long-range distances from EPR with shorter range solution NMR distances for structure determination (Tamm et al. 2007). Finally, sparse solid-state NMR data have been used to complement X-ray structures to orient  $\beta$ -barrels within lipid bilayers (Mahalakshmi and Marassi 2008).

In this work, we combined structural restraints from solid-state and solution NMR into a global target function, which includes a depth of insertion potential (Senes et al. 2007). We applied this method to monomeric PLN, a small membrane protein that adopts an L-shaped conformation in lipids and whose tertiary fold is dictated by its strong interactions with the lipid membrane (see Traaseth et al. 2009). Solid-state and solution NMR have been previously used to study several membrane protein systems, demonstrating remarkable similarities of membrane protein structures in detergent micelles and lipid bilayers (Franzin et al. 2007; Gong et al. 2007; Hu et al. 2007; Opella and Marassi 2004; Page et al. 2007). Consistent with these findings, our calculations show that the structural restraints for both solution and solid-state NMR are compatible and can be used synergistically to define the



structure and topology of PLN in lipid membranes. In addition, the depth of insertion potential gives a clear picture of the membrane location with respect to the protein that can be used as a starting point to embed it into an explicit lipid bilayer for further refinement. In fact, this final refinement step of the structural ensemble has already been adopted for soluble proteins (Calhoun et al. 2008; Kordel et al. 1997; Linge et al. 2003; Xia et al. 2002), and will be especially important for membrane proteins to (1) better define the fold space and (2) image protein–lipid interactions at the atomic level. Table S1 reports the statistics on the conformers generated with the hybrid method and a comparison with the conformational ensemble derived from the solution NMR data alone. It is clear that the introduction of the solid-state NMR data does not significantly affect the covalent geometry. The changes in the ideal geometry are within the allowed errors and all of the conformers have passed the conformational filters of the PDB deposition site. Also, the resolution of each single helical domain does not change. However, the overall resolution of the backbone changes dramatically. While in the conformational ensemble generated from the solution NMR data alone the rmsd for all of the backbone atoms was  $\sim 4.4$  Å, the introduction of the solid-state NMR data drastically reduced the rmsd to  $\sim 2.3$  Å. It is apparent that the interactions of PLN with the lipid membrane exemplified by the solid-state NMR data and the knowledge-base  $E_z$  potential limit the conformational space defining the structural topology of PLN. Nonetheless, the resolution achieved underscores the dynamic nature of PLN, a property that allows this small membrane protein to interact with several different binding partners.

Several other restraints obtained from biophysical and biochemical techniques can be easily incorporated in this protocol along the lines of the HADDOCK approach (Dominguez et al. 2003). In particular, it is important to measure depth of insertion experimentally using either paramagnetic quenching (Buffy et al. 2003) or protein/lipid REDOR measurements (Hong 2006; Mani et al. 2006; Toke et al. 2004) to complete the structural and topological characterization of membrane proteins within lipid bilayers. Although this approach is demonstrated for a small membrane protein, we are currently applying this method to the PLN pentamer (30 kDa), which will demonstrate the applicability to other multispan membrane proteins and complexes.

## Conclusions

We developed a new computational protocol that is based on a total objective energy function ( $E_{\text{total}}$ ) that incorporates solution NMR data (distances, torsion angles, and hydrogen bonds), solid-state NMR data (orientational information: CSA and DC) and a depth of insertion knowledge-based potential ( $E_z$ ) to determine the high-resolution structure and topology of membrane proteins. This hybrid energy function has been implemented into an XPLOR-NIH protocol and made available for general use. While this method is demonstrated for a small single pass membrane protein, the rapid progress made in sample preparation and spectroscopy of medium size membrane proteins will allow application to larger systems. The XPLOR-NIH scripts for the hybrid method are available for download at the authors' website ([www.chem.umn.edu/groups/veglia](http://www.chem.umn.edu/groups/veglia)).

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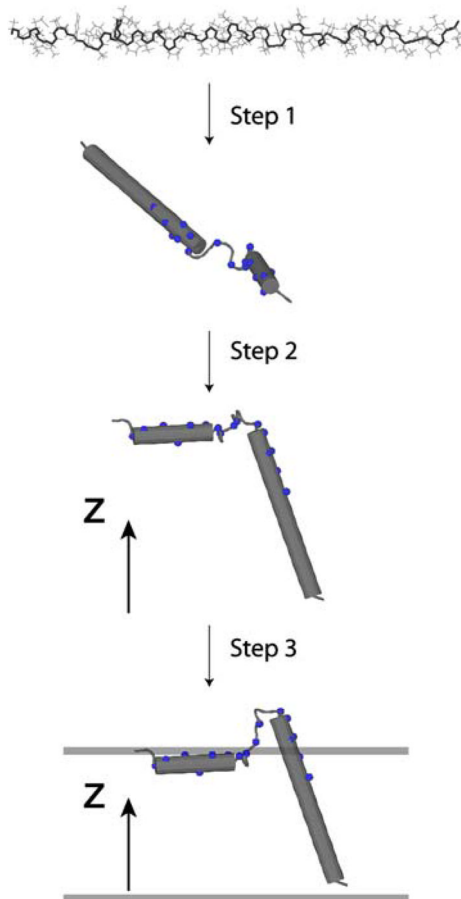
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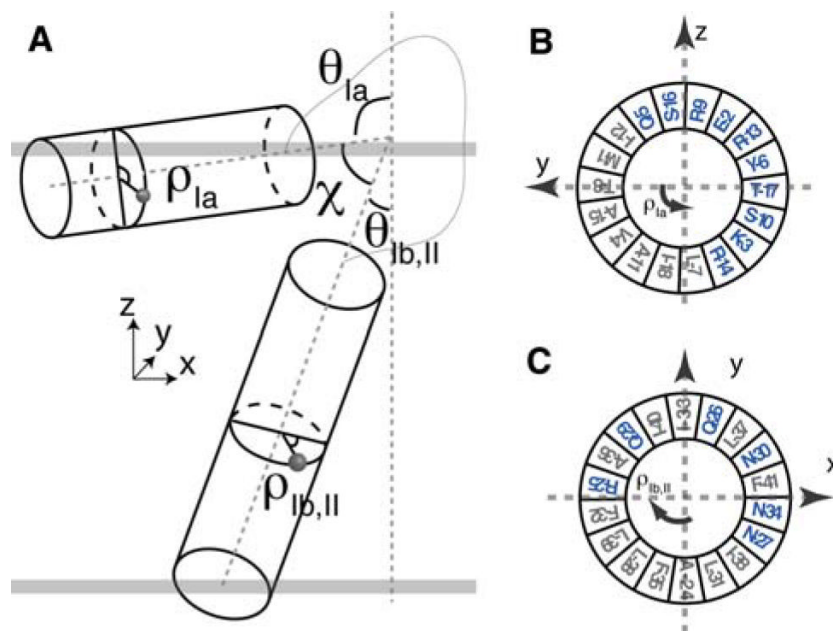
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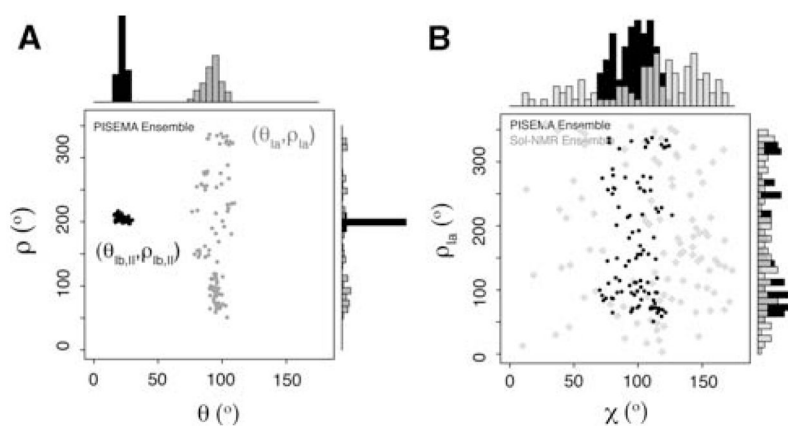


**Fig. 1.** Overview of the hybrid refinement protocol for the simultaneous determination of structure and topology of membrane proteins. *Step 1:* Starting from an extended structure, the simulated annealing protocol minimizes a target function containing only solution NMR data (NOEs, torsion angles, and hydrogen bonds). *Step 2:* The orientational constraints derived from solid-state NMR are included together with solution NMR restraints to obtain the correct orientation along the  $Z$  direction. *Step 3:* Depth of insertion is determined using rigid body minimization in the presence of the depth of insertion potential, keeping the helical orientation with respect to  $Z$  fixed. The resulting structural ensemble is refined using low temperature simulated annealing

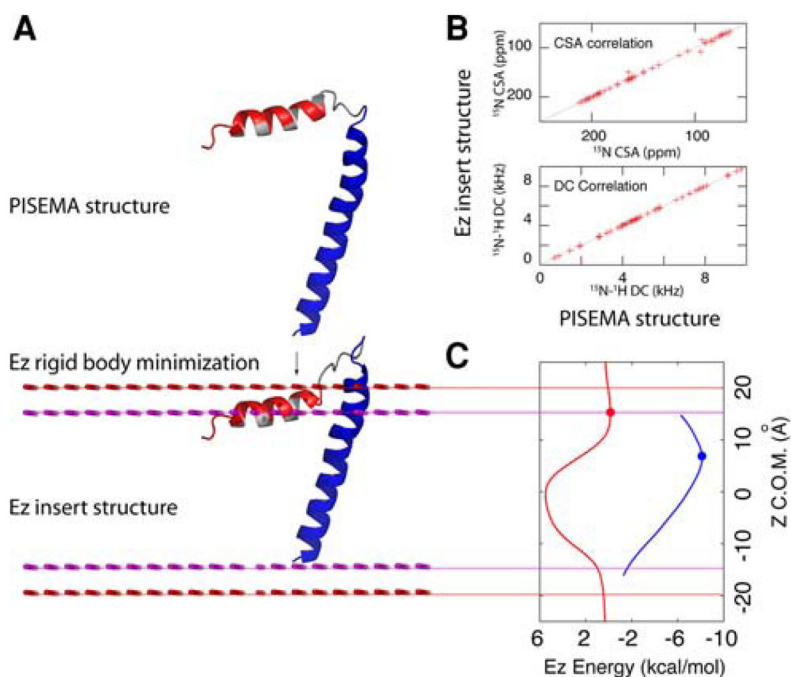




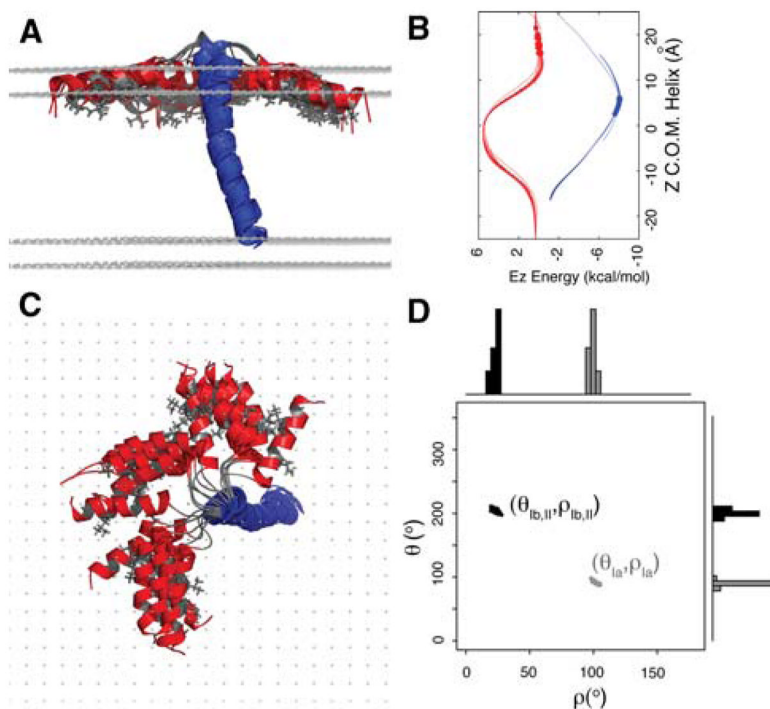
**Fig. 2.**  
**a** Definition of  $(\theta, \rho)$  describing the orientation of helix with respect to membrane normal  $Z$ .  $(\theta_{Ib, II}, \rho_{Ib, II})$  are the tilt and rotation angles for the domain Ib,II helix, while  $(\theta_{Ia}, \rho_{Ia})$  are the tilt and rotation angles for the domain Ia helix. The interhelical angle between the two domains is described by  $\chi$ . **b** Helical wheel representation of the reference orientation of domain Ia where  $\rho_{Ia}$  is defined to be zero. The N atom of T8 is aligned to  $+y$  axis.  $\rho_{Ia}$  rotates counterclockwise viewing from the top of  $y-z$  plane. At  $\sim 90$  degrees, the hydrophilic residues in *blue* point into bulk solution. **c** Reference orientation of transmembrane domain where  $\rho_{Ib, II}$  is defined to be zero



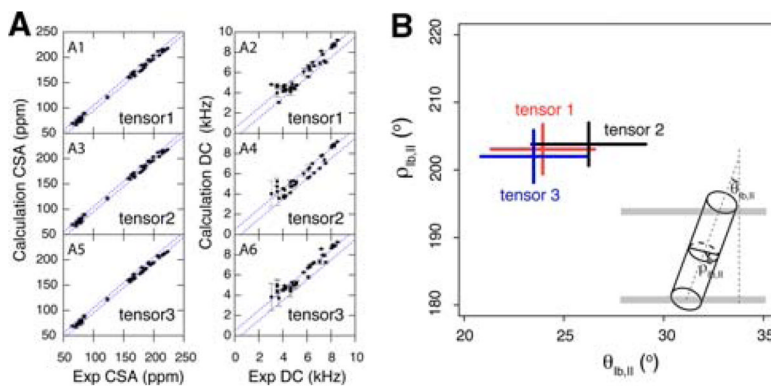
**Fig. 3.**  
**a** Distribution of  $\theta$  and  $\rho$  angles derived from domain Ia and domains Ib and II. **b** Distribution of interhelical angle  $\chi$  and  $\rho_{Ia}$  and the comparison with the solution NMR ensemble



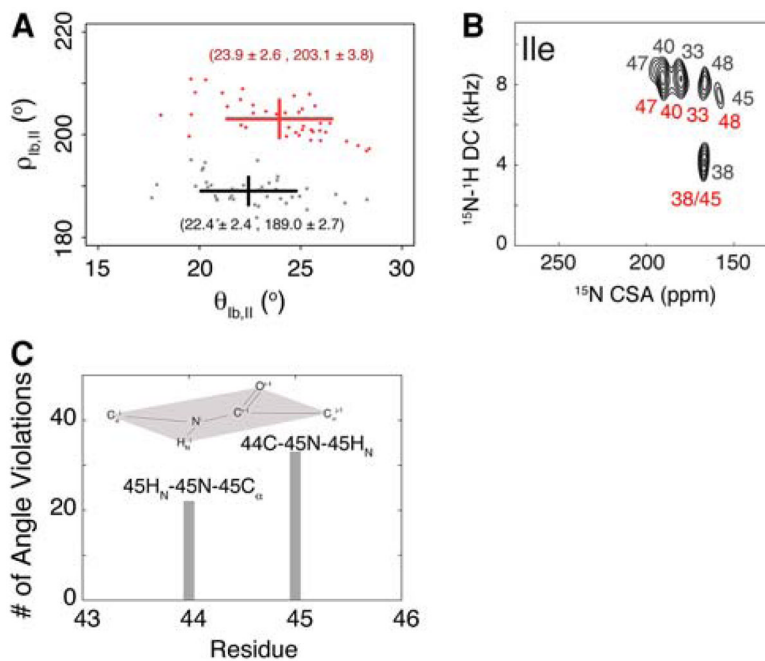
**Fig. 4.** Rigid-body minimization using the knowledge-based  $E_z$  potential. **a** Backbone cartoon representation of a selected PLN conformer before (*upper panel*) and after rigid-body minimization. **b** Comparison of the simulated CSA and DC before and after  $E_z$  minimization. The only discrepancies are due to residues located in the dynamic loop, which do not have CSA and DC restraints. **c** Representation of the  $E_z$  potential energy function for the domain Ib, II (*blue*) and domain Ia (*red*) helices. After minimization, both domains reside in the minima



**Fig. 5.**  
**a** Conformational ensemble (structures and topologies) representing the 20 lowest energy structures in the virtual bilayer. Domain Ia is colored in *red*, and domains Ib and II are in *blue*. Hydrophobic side chains of domain Ia are shown in *grey*. Structure overlay is performed by rotating along *Z* and translating along *X* and *Y*, resulting in no changes in  $E_z$  energy and PISEMA data. **b** Position of the cytoplasmic and transmembrane domains with respect to the depth of insertion  $E_z$  potential. **c** Top view of **a**. **d** Distribution of  $\theta$  and  $\rho$  angles in the final structural ensemble

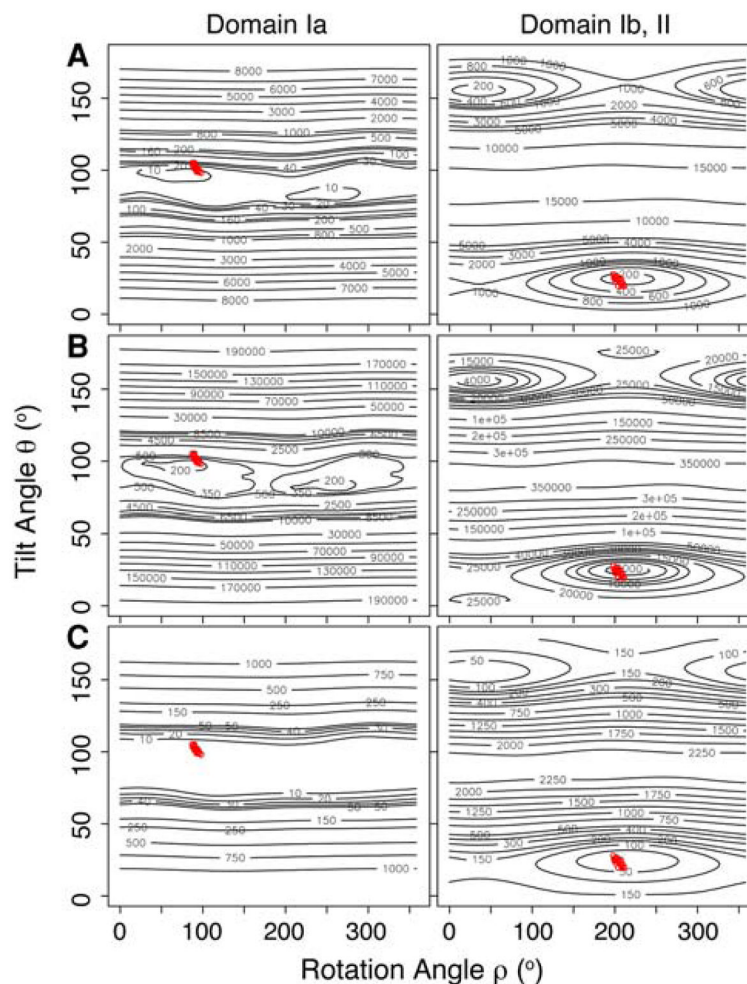


**Fig. 6.** **a** Comparison of the experimental and calculated CSA and DC for different values of the CSA tensor components shown in Table 1. The *diagonal errors* indicate ranges of  $\pm 5$  ppm and  $\pm 0.5$  kHz for CSA and DC, respectively. **b** Effects of the different CSA tensor components on the tilt and rotation angles for the helix defined by domains Ib and II of PLN. Tensors 1 (Wu et al. 1995), 2 (Page et al. 2008), and 3 are shown in *red*, *black*, and *blue*, respectively



**Fig. 7.** Effects of the resonance misassignment on the calculation of the structural topology of PLN. **a** Plots of the distribution of the rotation angles for the two ensembles obtained from simulated annealing calculations. The correct assignments give rise to an average rotation angle of  $203^\circ$ , while the incorrect assignment gives rise to an average rotation angle of  $189^\circ$ . **b** Experimental PISEMA spectra with the two equiprobable assignments for the isoleucine residues of PLN derived from the combinatorial assignment procedures (see text). **c** Histogram representing the number of violations obtained for residues 44 and 45 in the conformational ensemble generated with the incorrect assignments





**Fig. 8.**

Projection of  $(\theta, \rho)$  of the hybrid ensemble (20 lowest) for the two helical domains of PLN onto the PISEMA potential surfaces obtained using rigid helix fitting. Three different scoring functions are used to generate the potential surfaces:

**a**  $\text{score}_1 = \text{RMSD}_{\text{DC}} + \text{RMSD}_{\text{CSA}} \times \frac{\text{DC}_{\text{max}}}{\text{CSA}_{\text{max}}}$ , **b**  $\text{score}_2 = \text{RMSD}_{\text{DC}} \times \frac{\text{CSA}_{\text{max}}}{\text{DC}_{\text{max}}} + \text{RMSD}_{\text{CSA}}$  and

**c**  $\text{score}_3 = \frac{\text{RMSD}_{\text{DC}}}{\text{DC}_{\text{max}} - \text{DC}_{\text{min}}} + \frac{\text{RMSD}_{\text{CSA}}}{\text{CSA}_{\text{max}} - \text{CSA}_{\text{min}}}$ . These plots demonstrated that the topologies derived from the hybrid method correspond to the lowest energy minima identified in all of the PISEMA potential surfaces

**Table 1**

Different tensor values (tensor 1 (Wu et al. 1995) and 2 (Page et al. 2008)) and the resulting topological angles

Name	Tensor values		Topological angles	
	( $\sigma_{11}$ , $\sigma_{22}$ , $\sigma_{33}$ ) (ppm)	$\nu_{\text{NH}}$ (kHz)	$\theta_{\text{b,II}}$ ( $^{\circ}$ )	$\rho_{\text{b,II}}$ ( $^{\circ}$ )
Tensor 1	(64, 77, 217)	9.75	$24 \pm 2.6$	$203 \pm 3.7$
Tensor 2	(57.3, 81.2, 227.8)	10.735	$26 \pm 2.9$	$204 \pm 3.2$
Tensor 3	(64, 77, 217)	10.735	$24 \pm 2.6$	$202 \pm 3.9$

Note that tensor 3 is a combination of tensors 1 and 2