

Evaporation kinetics and phase of laboratory and ambient secondary organic aerosol

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Field measurements of secondary organic aerosol (SOA) find significantly higher mass loads than predicted by models, sparking intense effort focused on finding additional SOA sources but leaving the fundamental assumptions used by models unchallenged. Current air-quality models use absorptive partitioning theory assuming SOA particles are liquid droplets, forming instantaneous reversible equilibrium with gas phase. Further, they ignore the effects of adsorption of spectator organic species during SOA formation on SOA properties and fate. Using accurate and highly sensitive experimental approach for studying evaporation kinetics of size-selected single SOA particles, we characterized room-temperature evaporation kinetics of laboratory-generated α -pinene SOA and ambient atmospheric SOA. We found that even when gas phase organics are removed, it takes ~24 h for pure α -pinene SOA particles to evaporate 75% of their mass, which is in sharp contrast to the ~10 min time scale predicted by current kinetic models. Adsorption of "spectator" organic vapors during SOA formation, and aging of these coated SOA particles, dramatically reduced the evaporation rate, and in some cases nearly stopped it. Ambient SOA was found to exhibit evaporation behavior very similar to that of laboratory-generated coated and aged SOA. For all cases studied in this work, SOA evaporation behavior is nearly size-independent and does not follow the evaporation kinetics of liquid droplets, in sharp contrast with model assumptions. The findings about SOA phase, evaporation rates, and the importance of spectator gases and aging all indicate that there is need to reformulate the way SOA formation and evaporation are treated by models.

single-particle mass spectrometry | morphology

Atmospheric particles have a strong, yet poorly characterized effect on climate (1). Organic aerosols (OA) comprise 20–90% of atmospheric dry particles mass (2), the majority and least understood of which is secondary organic aerosol (SOA), formed from oxidation of gas phase organic vapors in the atmosphere (3–6). Despite an ongoing intense research effort aimed at understanding the formation and atmospheric evolution of OA, current models severely underestimate the formation of SOA in the atmosphere (5, 7). The effort to resolve the persistent discrepancy between field measurements and the amount of SOA predicted by atmospheric chemistry models has mostly focused on improving the understanding of SOA formation yields and finding new sources (8–11). In contrast, present models maintain the following fundamental assumptions: (i) Gas-particle partitioning is modeled assuming that all organics form a pseudoideal solution in the condensed particle phase, (ii) SOA particles remain liquid-like throughout their lifetime in the atmosphere, (iii) reversible thermodynamic equilibrium exists between gas and particle phases, and (iv) adsorption of other organic species and their effects on SOA properties and evaporation are ignored.

The assumption that particles are liquid is central to all modeling studies and affects calculations of condensation and evaporation of particles in the atmosphere. Properties and behavior of multicomponent liquids are clearly very different from multicomponent solids. Seven orders of magnitude lower diffusion rates within a solid phase limit most gas-particle interactions to surface layer for solid particles (12). As mixed solid particle

evaporation proceeds, the solid surfaces change resulting in dynamic changes in evaporation rates.

As we show below, the evaporation kinetics of laboratory-generated and ambient SOA indicate that these particles are not liquid and do not behave as expected by the models. These findings are in accord with a recent report, in which SOA particle bouncing behavior was used to conclude that SOA particles are amorphous solids (13).

In their box modeling study of the Mexico City plume, Dzepina et al. (6) found that adding precursors and SOA formation mechanisms (9) increases modeled SOA mass, bringing model predictions closer to measurements. However, when evaporation is considered, their models predict that more than 3/4 of the SOA mass evaporates upon dilution by a factor of 10, in sharp contrast with measurements. The same study attempted to model measured SOA evaporation in thermofluidizer (TF) and concluded that to get better agreement with field measurements requires SOA evaporation coefficients that are 1–3 orders of magnitude below unity (6), which leads to inconsistency with smog chamber SOA formation measurements.

Nearly all SOA evaporation studies rely on TD, which involve measurements on particles exposed for a short time to elevated temperatures. Using TD, Huffman et al. (14) observed that laboratory-generated SOA is more volatile than ambient SOA. Most recent analysis of field TD datasets concludes that a significant fraction of the atmospheric OA consist of nonvolatile components (15). Although TD produces valuable data, their interpretation is challenging. There are issues related to non-linear coupling effects between various parameters such as initial aerosol compositions, enthalpies of vaporization, volatility distributions, and evaporation coefficients governing particle evaporation rates (16). Most importantly, if the SOA particles are not liquid droplets, increasing particle temperature may result in phase change and with it a change in evaporation behavior.

Grieshop et al. (17) studied the kinetics of dilution-induced evaporation of polydisperse SOA in a smog chamber at room temperature and concluded that the system returns to equilibrium by evaporation within ~2.5 h. They attributed this "surprisingly slow" evaporation rate to the formation of low-volatility oligomers (17). However, comparison between the evaporation kinetics observed in Grieshop et al. (17) and in our laboratory indicates that 2.5 h is most likely the point at which the evaporation rate slows down and not where it ends. It has also been suggested that this dilution-induced evaporation type of experiment could be strongly affected by material deposited on smog chamber walls (18).

In this paper, we present results from a comprehensive investigation of SOA evaporation kinetics using a newly developed,

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highly sensitive experimental setup that allowed us to conduct a number of measurements that were not made before. We (i) characterized the evaporation of size-selected SOA particles at room temperature, (ii) monitored SOA evaporation for over 24 h, (iii) investigated the effect of hydrophobic “spectator” organic vapors on SOA evaporation, (iv) quantified the effect of aging on the evaporation kinetics of pure and coated SOA particles, and (v) studied the evaporation kinetics of size and chemically resolved ambient SOA particles at room temperature.

Results and Discussion

The schematic of the experimental setup is shown in Fig. 1. Laboratory SOA particles were generated in a Teflon bag by the reaction of α -pinene with ozone. Once reaction was complete, they were classified using a differential mobility analyzer (DMA), and the resulting monodisperse aerosol was transported through two room-temperature activated charcoal denuders and loaded into an evaporation chamber containing activated charcoal. The fact that particles were generated in a separate chamber and that gas phase organics in the evaporation chamber were continuously removed eliminates the issue of gas-wall repartitioning. Low particle number-concentrations in the chamber (~ 10 – 200 particles cm^{-3}) provide the means to follow kinetics of single-particle evaporation. SOA evaporation kinetics were monitored for over a ~ 24 -h period by periodically measuring changes in particle vacuum aerodynamic diameter, d_{va} , distributions, and mass spectra using our single-particle mass spectrometer, SPLAT II (19). In experiments involving spectator gases, small amounts of dioctyl phthalate (DOP), dioctyl sebacate (DOS), and pyrene (PY), or a mixture of organics were placed at the bottom of the Teflon bag and allowed to equilibrate before reaction started. As evident by the particle mass spectra, SOA formed under these conditions acquire a coating of the spectator organic compound (20). Ambient particles were dried and then treated in the same manner as the laboratory particles.

Evaporation Kinetics of Single-Component Test Particles. We first present, in Fig. 2, results of evaporation studies for single-component particles composed of DOP. Fig. 2*A* shows how the d_{va} of DOP particles with initial diameter $d_0 = 500$ nm decreases as a function of time as the particles evaporate. Fig. 2*B* shows how particle diameters, d , decrease during evaporation in a plot of d^2/d_0^2 vs. time. The data agree with traditional models of single-component liquid droplet evaporation, which predict that the square of the particle diameter decreases nearly linearly with time (21) and that the slope of d^2/d_0^2 vs. time is size-dependent; that is, smaller particles evaporate faster. This size-dependent evaporation behavior also quantitatively explains the observed increase in the width of the d_{va} distributions as the particles evaporate. DOP particles on the left edge of the size distribution (slightly smaller) evaporate faster than the larger particles on the right edge of the distribution. Applying the analysis of Zhang et al. (22) to these data, we obtain, from the slope of d^2/d_0^2 vs. time, a vapor pressure of 10^{-7} Torr, in excellent agreement with previously reported values (23).

These measurements validate the experimental methodology and introduce the time scales expected for particle evaporation: 300-nm DOP particles, whose vapor pressure is 10^{-7} Torr, lose more than 70% of their volume by evaporation in 1 h, whereas 700-nm DOP particles lose less than 20% of their volume over the same time.

Evaporation Kinetics of Pure Laboratory SOA. We proceed to investigate the evaporation kinetics of α -pinene SOA particles. The results of these measurements are summarized in Fig. 3. Fig. 3*A* shows how the d_{va} of SOA particles with $d_0 = 160$ nm, i.e., $d_{va0} = 187$ nm, decreases during evaporation. After comparison with Fig. 2, two observations are notable. First, the diameters of SOA particles decrease with time significantly slower, and second, the d_{va} distributions remain narrow during SOA particles evaporation. Fig. 3*B* shows d^2/d_0^2 as a function of evaporation time for SOA particles with different initial diameters. Evidently, SOA particles undergo evaporation in two distinct stages. In the fast evaporation stage particles lose $\sim 50\%$ of their volume in less than 100 min, whereas in the slow stage, it takes an additional 1,400 min to lose another $\sim 25\%$ of their initial volume. Moreover, Fig. 3*B* shows that the evaporation kinetics of SOA particles is nearly size-independent.

Fig. 3*C* presents a comparison between modeled and observed SOA evaporation behavior, plotted here and in the rest of the figures as d_{va}^3/d_{va0}^3 , i.e., the remaining volume fraction, as a function of time. Calculations for three particle sizes (100, 160, and 251 nm) are performed using a detailed kinetic model that employs the seven-product volatility basis sets developed to fit SOA formation data during α -pinene ozonolysis (24), and using two different mass accommodation coefficients (α) of 1.0 and 0.05 (16). Note that although $\alpha = 1$ is consistent with observed SOA formation rates, $\alpha = 0.05$ leads to SOA formation rates that are inconsistent with smog chamber measurements. Nevertheless, we use both values to illustrate that even unreasonably low α do not yield agreement with experiment. There are two important differences between modeled and observed evaporation kinetics. First, calculated evaporation rates are much faster than observed. For example, according to the model, 151-nm particles will lose 75% of their volume in 1 to 20 min of evaporation, depending on the mass accommodation coefficient, whereas the data show that in reality it takes over 24 h, independent of particle size. Second, calculated SOA evaporation behavior shows the expected size dependence, with larger particles evaporating significantly slower. In contrast, the observed data are nearly size-independent. It should be noted that all models based on fits to smog chamber data yield similar fast evaporation rates and size-dependent behavior. The simple fact that SOA evaporation does not follow the size-dependent evaporation behavior expected of liquid droplets indicates that these particles are not liquid-like.

Evaporation Kinetics of Coated Laboratory SOA. In the real atmosphere SOA particles form in the presence of a mixture of organic vapors, some of which are hydrophobic. We have previously

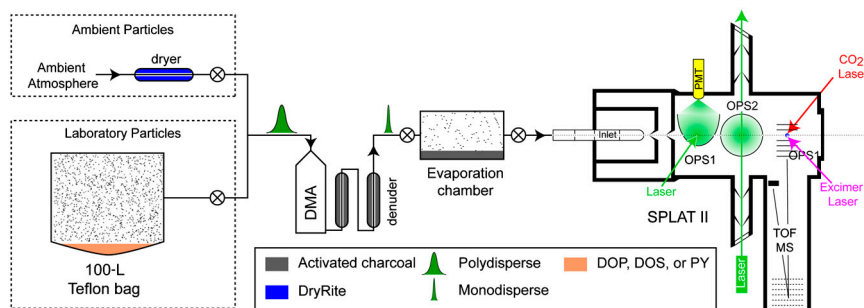


Fig. 1. Schematic of the experimental setup.

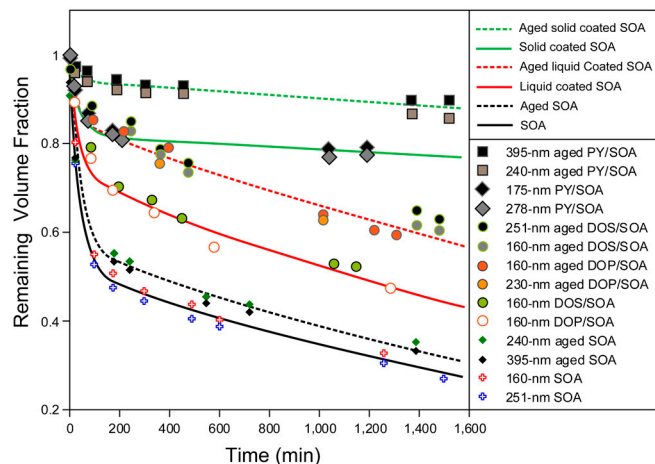


Fig. 4. Evaporation kinetics (reported as remaining volume fraction vs. time) for a number of the particle types indicated in the legend. All lines represent biexponential fits to the data. For each particle type, a single line is used to fit different sizes, and a single line is used to fit the DOP- and DOS-coated particles. Dashed lines represent evaporation of aged particles.

changes for m/z smaller than 200 amu, disappearance of the peak at $m/z = 201$ after 2 min of evaporation, and, at later times, slow increase of intensity of peaks with m/z higher than 202, which could indicate oligomer formation (29). However, we note that aging alone produces no observable changes in particle mass spectra. The changes in mass-spectral intensities with evaporation time are illustrated in Fig. S1. The intensity of the DOP-characteristic peak ($C_6H_4(CO)_2OH^+$, $m/z = 149$) in DOP-coated SOA particles, rapidly drops with evaporation time, and then remains constant. The PY parent ion peak ($m/z = 202$), in PY-coated SOA, decreases slowly, such that $\sim 50\%$ of it remains even after 25 h of evaporation. The comparison between PY and DOP suggests that there are differences in morphological distribution of solid PY and liquid DOP in the coated SOA.

Evaporation Kinetics of Atmospheric SOA. To address the question of how closely the laboratory observations described above reflect reality in the atmosphere, we characterized the evaporation kinetics of atmospheric SOA particles sampled in situ during the recent Carbonaceous Aerosols and Radiative Effects Study (CARES) field campaign in Sacramento, CA, in June 2010. Because in the real atmosphere aerosol compositions change throughout the day, the first task was to identify the point at which SOA dominates the particle compositions. This was accomplished by following the diurnal evolution of particle number concentrations, size distributions, compositions, and densities. Early

in the morning aerosol loadings and number concentrations were low, and most particles were composed of organics mixed with a significant fraction of sulfate, with larger particles containing a larger fraction of sulfate (30). By $\sim 9:00$ AM, the number concentrations of very small particles (<14 nm) started to increase, indicating the beginning of SOA nucleation events. As the day progressed and emissions of the volatile organic compounds were processed, SOA particles increased in size, making it possible to characterize their size, composition, and density with SPLAT II. By early afternoon aerosol composition was dominated by oxygenated organics mixed with a small amount of organic amines and sulfate ($\sim 12\%$ volume fraction), and the density of both 100-nm and 151-nm particles was measured to be 1.32 ± 0.02 g cm^{-3} . At this point size-selected 100- and 151-nm particles were loaded into the evaporation chamber and their evaporation behavior was studied.

Fig. 5 shows the changes in the measured d_{va} distributions of 100-nm ambient particles measured on June 26 over 208 min of evaporation. At $t = 0$ the particles average d_{va} is 133 nm, which yields a density of 1.33 g cm^{-3} (see *SI Text*). With time in the chamber, the particles evaporate and after 208 min they shrink to 125 nm, i.e., lose 17% of their volume. Note that as the organic fraction of the particle evaporates, its sulfate fraction, and hence density, slightly increases, which when taken into account means that the actual evaporated volume of organics is 20%. Fig. 5B shows the remaining volume fraction for 100-nm and 151-nm particles as a function of time. Notably, the evaporation behavior of the atmospheric SOA is similar to that of laboratory-generated coated SOA. It is slow, size-independent, and can be fit with a biexponential curve that represents the fast and slow evaporation phases (see *SI Text*). Here, only a small fraction of the volume (0.135) is lost during the fast evaporation phase.

Although it is true that the vast majority of atmospheric particles were SOA, it is worth keeping in mind that the data in Fig. 5 are an average of all sampled particle types. It is possible to refine the analysis by classifying the data according to particle compositions and examine the evaporation of the individual classes. The mass spectra show that in addition to SOA-dominated particles, a small fraction of biomass burning and soot particles were also present, but their small numbers prevent us from analyzing their behavior in the evaporation chamber.

Two distinct types of SOA particles mixed with sulfate were observed on June 25, 26, and 27. Both types, labeled Type 44 and Type 43 to denote which of these two mass-spectral peaks has higher intensity prior to particle evaporation (Fig. S2), are dominated by oxygenated organics, have identical density, and are mixed with the same amount of sulfate ($\sim 12\%$). Type 43 particles contain in addition to oxygenated organics and sulfate a small amount of organic amines.

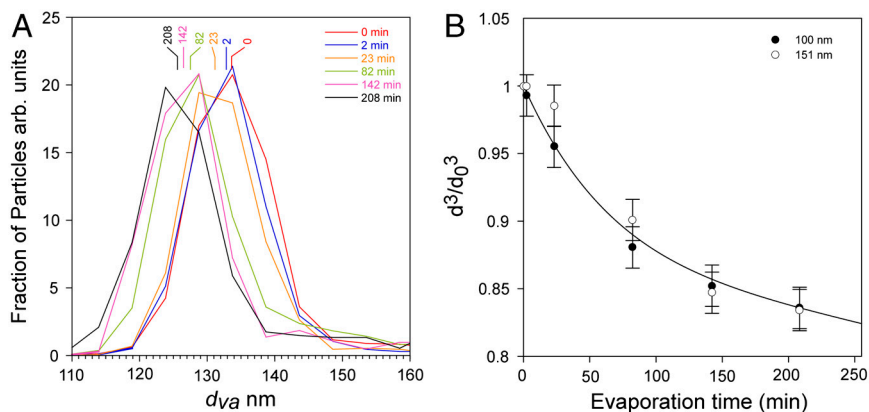


Fig. 5. Evaporation kinetics of ambient particles. (A) d_{va} distributions of 100-nm ambient particles at the marked evaporation times (in minutes). (B) Evaporation kinetics (reported as remaining volume fraction) for ambient 100-nm and 151-nm particles. The solid line is a biexponential fit to the data.

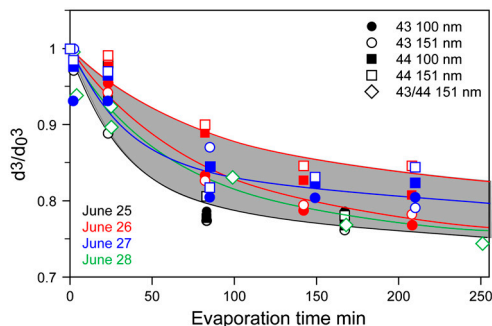


Fig. 6. Evaporation kinetics of 100-nm and 151-nm ambient SOA particles classified into two particle types (43 and 44). The solid lines are biexponential fits to the data and the gray shaded area marks the upper and lower limits of all observed data.

Fig. 6 shows the evaporation data acquired during the four-day study. Comparison between the evaporation data of 100-nm and 151-nm particles of both types shows very small differences. We note again that like laboratory SOA, evaporation of ambient SOA is nearly size-independent and conclude that, even when mixed with sulfate, ambient SOA particles do not behave like liquid droplets.

Aerosol loadings for June 28 were a factor of two higher and present a slightly different scenario. Because on this day the SOA particles mass spectra show continuous range of compositions spanning between the two particle types, they were treated as one class. Because the number of 100-nm particles in the evaporation chamber was insufficient to yield usable data, we present for this day data for 151-nm particles only.

In addition to the data points, Fig. 6 displays a gray-shaded region that marks the area between the fastest and slowest evaporation rates observed in this study, and biexponential fits to the data points given in *SI Text*. Note that there is no reason to expect the evaporation curves on different days to be identical. Variation in temperature, solar flux, cloud cover, wind speed, and wind direction are expected to produce particles with somewhat different compositions and hence evaporation rates. Nevertheless, a comparison of the data from these four days points to a remarkably similar evaporation behavior.

Examination of changes in mass-spectral intensities of ambient SOA particles with evaporation time show that changes are relatively small, and that unlike in the case of laboratory SOA, no increase in the mass-spectral intensity at high m/z is observed. In addition, the intensity of peaks 44 and 73, for both particle types, decrease rapidly with evaporation and then remain constant, which could indicate that some of their intensity belongs to surface coatings (shown in Fig. S3).

Fig. 7 shows a comparison between the results for laboratory-generated SOA shown in Fig. 4, the measured evaporation kinetics of ambient atmospheric SOA particles as presented by the gray-shaded region, and the calculated evaporation kinetics presented in Fig. 3C. Based on Fig. 7 we conclude that the evaporation of atmospheric SOA is very similar to that of laboratory-generated α -pinene SOA coated either with PY, organic mixture, or aged SOA coated with DOP or DOS. This underscores the importance of the laboratory investigations of the coated SOA particles and suggests that the coated SOA is in a sense similar to real-world atmospheric SOA. The differences between the observed evaporation behaviors and that predicted by models are rather extreme. Even when using unreasonable mass accommodation of 0.05, models predict very fast evaporation, whereas the data indicate that evaporation is slow. Assuming liquid-like SOA behavior, models predict that evaporation of small particles is faster than that of larger ones, but the data clearly show that evaporation is nearly size-independent indicating that these particles cannot be liquids.

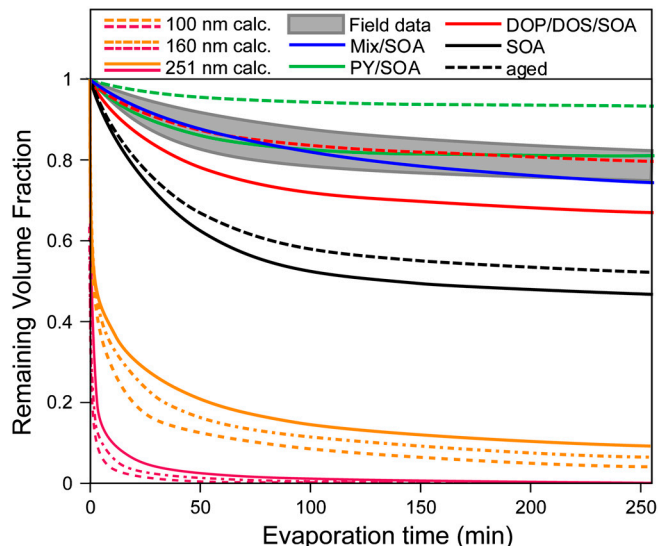


Fig. 7. Comparison between the observed evaporation kinetics of ambient organic particles (gray area), the laboratory-generated α -pinene SOA particles with and without various coatings and aging, and evaporation kinetics calculated based on current models.

Conclusions

We summarize our findings as follows: (i) Overall, SOA evaporation is more than ~ 100 times slower than expected by current models. It proceeds in two stages: a fast stage that takes ~ 100 min, followed by very slow evaporation that lasts more than a day. (ii) SOA evaporation is nearly size-independent and does not follow the evaporative behavior of liquid droplets, indicating that SOA particles are not liquids. This conclusion is consistent with the results of our previous study that showed that particles with two reverse morphologies—DOP-coated SOA particles and SOA-coated DOP particles—remain stable for many hours, which would be incompatible with the rapid diffusion rates expected if both DOP and SOA are liquids (20). (iii) SOA particles formed in the presence of vapors of hydrophobic organic compounds exhibit lower evaporation rates. (iv) Aging decreases evaporation rates in all cases; however, its effect on coated particles is significantly more pronounced. (v) The evaporation of ambient SOA is similar to that of coated and aged laboratory-generated α -pinene SOA—it is slow and size-independent. (vi) Changes in mass-spectral intensity with evaporation suggest that ambient SOA acquires coatings, much like laboratory SOA.

Because in our experiments evaporation is induced by sudden and complete removal of the vapor phase, it is expected to be significantly faster than in the atmosphere, where the vapor phase concentrations decrease in response to dilution and chemical processing, both of which occur on a time scale of a few hours. Moreover, in the atmosphere, where a large number of organic vapors are present, as some coatings evaporate, others are adsorbed to form new coatings.

These observations have important implications to the way we view and model the formation and fate of SOA particles. The data show that the assumptions of liquid-like SOA behavior and instantaneous reversible equilibrium gas-particle partitioning are seriously flawed. They show that adsorption of organic vapors, present at pressures equal to or lower than their vapor pressure, increases organic mass and inhibits evaporation, an effect that is enhanced with particles aging. In contrast with current air-quality models, in which SOA particles rapidly evaporate due to dilution, our findings indicate that evaporation due to transport to cleaner environments is almost negligible. This has the effect of increasing the contribution of SOA to organic aerosol in both urban and remote locations. Moreover, the observed

size-independent evaporation behavior implies that, in contrast with models, small (~100 nm) particles do not evaporate faster and thus retain their potential to serve as cloud condensation nuclei. Because SOA formation is modeled on the assumptions that these particles are liquids and that spectator gases play no role, the findings presented here have implications for this aspect of SOA modeling as well. This work could provide an essential missing link to process representation of organics aerosols in models.

Materials and Methods

Particle Generation. DOP and DOS particles were generated by atomization of the neat liquids. Homogeneous nucleation of SOA was initiated by injecting 200 ppb of α -pinene, ~500 ppb of O_3 , and ~250 ppm of cyclohexane, used as an OH scavenger, into a clean 100-L Teflon bag filled with zero air. Coated particles were generated by carrying out the SOA generation procedure in a Teflon bag that contained a small amount of bulk coating material. Ambient particles were sampled in situ, dried, and characterized using SPLAT II and a DMA, as described elsewhere (31).

Particle Analysis. Particles were size-selected with a DMA, passed through two charcoal denuders, and loaded in the stainless steel evaporation chamber, partially filled with activated charcoal.

A single-particle mass spectrometer, SPLAT II, was used to measure individual particle composition, vacuum aerodynamic diameter, d_{va} , density, and shape, as described elsewhere (19, 32). Particle evaporation kinetics was quantified by measuring changes in particle d_{va} with 0.5% precision. A more detailed description of the experimental setup is provided in *SI Text*.

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