

# Experimental Protocols Alter Phototransduction: The Implications for Retinal Processing at Visual Threshold

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Vision in dim light, when photons are scarce, requires reliable signaling of the arrival of single photons. Rod photoreceptors accomplish this task through the use of a G-protein-coupled transduction cascade that amplifies the activity of single active rhodopsin molecules. This process is one of the best understood signaling cascades in biology, yet quantitative measurements of the amplitude and kinetics of the rod's response in mice vary by a factor of  $\sim 2$  across studies. What accounts for these discrepancies? We used several experimental approaches to reconcile differences in published properties of rod responses. First, we used suction electrode recordings from single rods to compare measurements across a range of recording conditions. Second, we compared measurements of single-cell photocurrents to estimates of rod function from *in vitro* electroretinograms. Third, we assayed the health of the post-receptor retinal tissue in these different conditions. Several salient points emerge from these experiments: (1) recorded responses can be altered dramatically by how the retina is stored; (2) the kinetics of the recovery of responses to bright but not dim flashes are strongly sensitive to the extracellular concentration of magnesium; (3) experimental conditions that produce very different single-photon responses measured in single rods produce near identical derived rod responses from the electroretinogram. The dependence of rod responses on experimental conditions will be a key consideration in efforts to extract general principles of G-protein signaling from studies of phototransduction and to relate these signals to downstream mechanisms that facilitate visual sensitivity.

## Introduction

Rod photoreceptors provide the dominant input to the visual system under conditions in which photons are scarce. This requires that rods detect the absorption of single photons and that the resulting signals are effectively read out by the retinal circuitry (for review, see Sharpe and Stockman, 1999; Field et al., 2005). Our understanding of these processes relies on an accurate description of the rod's responses to single absorbed photons. Yet the amplitude and kinetics of the rod's responses vary by a factor of  $\sim 2$  across studies (Calvert et al., 2000; Mendez et al., 2000; Luo and Yau, 2005; Doan et al., 2006, 2009; Krispel et al., 2006; Fu et al., 2008; Wen et al., 2009; Gross and Burns, 2010; Okawa et al., 2010). These discrepancies hamper attempts to extract general principles of G-protein-coupled receptor (GPCR) function from the photoreceptors, and attempts to construct an understanding of rod vision from the molecular to the organismal level. The work described here aimed to understand the origin of these differences.

Studies of single-photon responses offer an unprecedented opportunity to understand how activation of single GPCRs leads

to the biochemical amplification that is a hallmark of G-protein signaling (Baylor et al., 1979b; Pugh and Lamb, 1993). Activation of rhodopsin leads, through the G-protein transducin and cGMP phosphodiesterase (PDE), to a decrease in cGMP concentration, channel closure, and a decrease in current. Recovery of the resulting response depends on inactivation of these components and resynthesis of cGMP by guanylate cyclase. This basic signaling scheme was worked out  $>20$  years ago; more recent work has focused on providing quantitative estimates of the relevant rate constants using a combination of biophysical and genetic approaches. Indeed, a large body of work has made rod phototransduction the best understood G-protein cascade in biology (Burns and Arshavsky, 2005; Burns and Pugh, 2010). However, differences in measured rod responses indicate that the resulting quantitative picture depends strongly on recording conditions (Doan et al., 2009; Gross and Burns, 2010).

Here we identify elements of standard recording protocols that alter the measured parameters of phototransduction. We show that conditions for storing the retina can strongly affect rod sensitivity and single-photon response kinetics, and can have dramatic consequences for the health of post-receptor retinal tissue. Differences in the extracellular magnesium concentration account for further quantitative differences in the kinetics of recovery from bright light flashes. We attempted to identify *in vitro* conditions replicating those found *in vivo* by recording electroretinograms (ERGs) from isolated tissue. Derived rod responses to both bright and half-saturating flashes were consistent with similar measures from *in vivo* ERG studies. The rod responses derived from the ERG did not, however, exhibit the clear dependence on recording conditions seen in single rod recordings. This

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**Table 1. Definitions of shorthand for protocols used in this study**

Designation	Color	Storage	Perfusion	Electrode
L-15/Locke's	Red	L-15 culture media on ice	Locke's, 36°C	HEPES Locke's, 36°C
Locke's/Locke's	Blue	Locke's, 32°C	Locke's, 36°C	HEPES Locke's, 36°C
Ames'/Ames'	Black	Bicarb. Ames, 32°C	Bicarb. Ames, 36°C	HEPES Ames, 36°C
Ames <sup>Ca,Mg</sup> /Ames <sup>Ca,Mg</sup>	Green	1.15 mM Ca <sup>2+</sup> , 1.2 mM Mg <sup>2+</sup> Bicarb. Ames, 32°C	1.15 mM Ca <sup>2+</sup> , 1.2 mM Mg <sup>2+</sup> Bicarb. Ames, 36°C	1.15 mM Ca <sup>2+</sup> , 1.2 mM Mg <sup>2+</sup> HEPES Ames, 36°C
Ames <sup>Mg</sup> /Ames <sup>Mg</sup>	Magenta	1.2 mM Ca <sup>2+</sup> , 2.4 mM Mg <sup>2+</sup> Bicarb. Ames, 32°C	1.2 mM Ca <sup>2+</sup> , 2.4 mM Mg <sup>2+</sup> Bicarb. Ames, 36°C	1.2 mM Ca <sup>2+</sup> , 2.4 mM Mg <sup>2+</sup> HEPES Ames, 36°C
		1.15 mM Ca <sup>2+</sup> , 2.4 mM Mg <sup>2+</sup>	1.15 mM Ca <sup>2+</sup> , 2.4 mM Mg <sup>2+</sup>	1.15 mM Ca <sup>2+</sup> , 2.4 mM Mg <sup>2+</sup>

"Designation" denotes storage and perfusion solutions in each protocol. The color scheme defined here is held constant throughout the figures.

**Table 2. Summary data of single rod measurements**

Protocol	peak (pA)	$t_p$ (ms)	$t_i$ (ms)	$\tau_r$ (ms)	$r_{max}$	$i_0$ (photons $\cdot \mu\text{m}^{-2}$ )	$t_p(i_0)$ (ms)	$\tau_D$ (ms)
L-15/Locke's	0.24 ± 0.03 (12)	127 ± 5	240 ± 11	174 ± 8	11.3 ± 0.5 (30)	57.4 ± 2.3 (15)	103 ± 4	210 ± 16 (21)
Warm L-15/Locke's	0.56 ± 0.08 (15)	159 ± 8	292 ± 25	217 ± 19	13.5 ± 0.9 (10)			
Locke's/Locke's	0.82 ± 0.09 (13)	161 ± 5	324 ± 22	235 ± 19	16.9 ± 0.8 (22)	39.0 ± 1.8 (15)	124 ± 1	159 ± 17 (7)
Ames'/Ames'	1.04 ± 0.06 (19)	181 ± 6	369 ± 18	265 ± 17	15.1 ± 0.6 (50)	32.6 ± 2.2 (19)	127 ± 4	376 ± 23 (23)
Ames <sup>Ca,Mg</sup> /Ames <sup>Ca,Mg</sup>	0.90 ± 0.10 (20)	187 ± 6	332 ± 24	230 ± 22	14.5 ± 0.7 (29)	34.4 ± 1.6 (23)	125 ± 3	173 ± 10 (15)
Ames <sup>Mg</sup> /Ames <sup>Mg</sup>	0.60 ± 0.05 (30)	190 ± 5	318 ± 13	217 ± 10	13.5 ± 0.5 (33)	33.8 ± 3.5 (23)	124 ± 7	209 ± 10 (16)

The first five columns summarize properties of responses to dim flashes: peak, peak amplitude of the estimated single-photon responses using a fixed collecting area of 0.5  $\mu\text{m}^2$ ;  $t_p$ , time-to-peak of single-photon response;  $t_i$ , integration time or area under the normalized single-photon response;  $\tau_r$ , exponential constant describing single-photon response recovery phase. Dark current,  $r_{max}$ , is calculated from saturated responses.  $i_0$  is the strength of the flash producing a half-maximal response, and  $t_p(i_0)$  is the time-to-peak of that response.  $\tau_D$  is the dominant time constant for decay of light-activated PDE activity (Fig. 5).

discrepancy limits the utility of the ERG in establishing conditions that mimic those found *in vivo*.

## Materials and Methods

**Mice.** Female C57BL/6-type mice, 4–6 weeks old, were purchased from either The Jackson Laboratory or Charles River Laboratories. *GCAP<sup>-/-</sup>* mice were obtained from Dr. Jeannie Chen, University of Southern California, Los Angeles, CA. Following 12–14 h dark adaptation, a mouse was killed and the eyes were removed under infrared illumination. The cornea, lens, and vitreous were removed from each eye, leaving the retina attached to the retinal pigmented epithelium and sclera. This eyecup was placed in a light-tight container in ~50 ml of storage solution for use throughout the day (7–10 h).

**Solutions.** The media and their temperatures used in these experiments are the primary variables in this study. Up to three solutions were used in a single experiment: one solution in which to store tissue, another to perfuse the tissue during recording, and a third to fill the recording pipette. Shorthand designations for particular combinations of solutions are defined in Table 1, and all figures maintain the defined color scheme.

Ames' medium (Sigma-Aldrich or Invitrogen) was developed for storage of retinal tissue of rabbits (Ames and Nesbett, 1981). Bicarbonate-buffered Ames' medium (Bicarb. Ames in Table 1) contained 22.6 mM sodium bicarbonate, 80,000 U/L penicillin, and 80 mg/L streptomycin, and was equilibrated with 95% O<sub>2</sub>/5% CO<sub>2</sub>. Electrodes were filled with HEPES Ames' medium in which sodium bicarbonate was replaced with 15 mM NaCl and 10 mM HEPES. The pH was adjusted to 7.4 with 1N NaOH, and the osmolality of each medium was adjusted to 280 ± 2 mOsm/kg. In some experiments an additional 0.05 mM CaCl<sub>2</sub> and/or 1.2 mM MgCl<sub>2</sub> was added to Ames' medium to match the divalent concentrations in Locke's medium.

Leibowitz L-15 culture medium (Invitrogen) was supplemented with 10 mM glucose and 0.1 mg/ml BSA (Sigma-Aldrich) and bubbled with 100% O<sub>2</sub>. Dissection began at least half an hour thereafter. The light-tight container was placed on ice for cold storage. We did not observe obvious differences in experiments in which the L-15 medium was either prechilled or placed on ice after the dissection. To separate the effects of L-15 and cold storage, we also performed a few experiments with retina stored in warm (32°C) L-15 medium (Table 2).

Locke's medium contained the following (in mM): 112.5 NaCl, 3.6 KCl, 2.4 MgCl<sub>2</sub>, 1.2 CaCl<sub>2</sub>, 20 NaHCO<sub>3</sub>, 10 HEPES, 0.02 EDTA, 3 Na<sub>2</sub>Succinate, 0.5 Na-L-glutamate, 10 glucose, 0.1 mg/ml BSA, 10 ml/L

Basal Medium Eagle (BME) 100× Vitamins, and 20 ml/L 50× BME Amino Acids (Schnapf et al., 1990). The solution was equilibrated with 95% O<sub>2</sub>/5% CO<sub>2</sub> and the pH adjusted to 7.4. Electrodes were filled with HEPES Locke's which lacked sodium bicarbonate, amino acids, and vitamins and contained 130 mM NaCl.

**Light stimuli.** Light-emitting diodes (LEDs) with peak emission wavelengths of either 465 or 575 nm provided light stimuli. The LED output was attenuated with neutral density filters and focused to a spot by either the microscope objective or the condenser. The LED output was converted to equivalent photons at 501 nm using the measured LED spectrum, calibrated power output, and measured rod spectral sensitivity (Baylor et al., 1984). Stimuli were generated by custom software.

**Single-cell recordings.** Suction electrodes were used to record the membrane currents of single rod outer segments (Yau et al., 1977; Doan et al., 2009). Glass micropipettes were pulled to tip inner diameters of ~2  $\mu\text{m}$ , then fire-polished to ~1.6  $\mu\text{m}$ , coated with butyltrichlorosilane to reduce cell-to-glass adhesion, and the tips were coated in wax to reduce capacitance. Electrodes used for recordings had resistances of 3–4 M $\Omega$ . Calomel electrodes and an active bath ground circuit facilitated stable recordings.

Between three and eight times over the course of the experiment day, ~1/12th of a single retina was transferred to a platform of Sylgard and finely chopped in a 350  $\mu\text{l}$  droplet of storage solution containing 10 U of DNase (Sigma-Aldrich). The resulting suspension was transferred to a recording chamber on the stage of an inverted microscope equipped with an infrared viewing system. Pieces of retina were allowed to adhere to the bottom of the chamber for 5 min, then perfused (2–3 ml/min, 37 ± 1°C) for the course of the experiment.

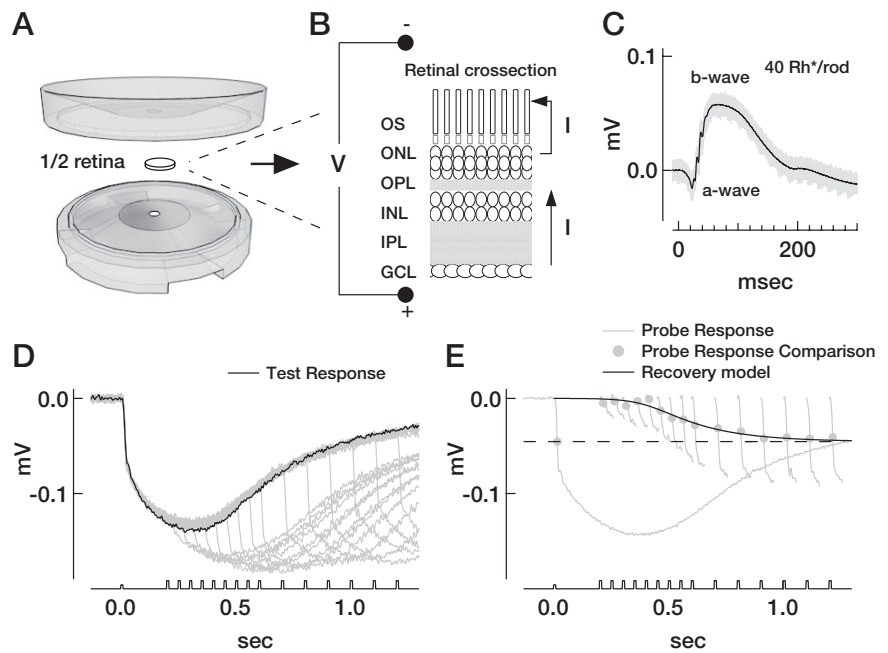
Cells were selected for extended recordings, lasting anywhere from 10 to 30 min, if their outer segment had not been obviously damaged by the suction procedure and if they had maximal light responses (i.e., dark currents) exceeding 8 pA. The dark current and kinetics of the light response were monitored and a recording was aborted if these metrics changed noticeably. Dark currents sometimes increased while the kinetics of the response remained unchanged, after which the recording was stable. We attribute this increase in dark current to an improved electrical seal at the mouth of the electrode. In these cases, responses before the change in dark current were not analyzed.

Single photon responses (response per Rh<sup>+</sup>) were estimated from the responses to at least 30 flashes at several flash strengths. Responses to a given flash strength were averaged and scaled by the flash strength. Time-to-peak was measured from the beginning of the 10 ms stimulus. Com-

paring the average responses across flash strengths checked linearity of the responses. We estimated a single photon response using two approaches. First, we scaled the photon density,  $I$  (units of photons  $\cdot \mu\text{m}^{-2}$ ), by an average rod collecting area of  $0.5 \mu\text{m}^2$  (Field and Rieke, 2002) (see Fig. 3G). Second, we extracted the scale factor that produced a best fit between the rising phases (10–90% of peak) of the mean squared response and the ensemble variance for a collection of responses to a single flash strength, this procedure is described below and illustrated in Figure 3G (Baylor and Hodgkin, 1973; Baylor et al., 1979b). The ensemble mean-squared response,  $\mu^2$ , to a particular flash strength is assumed to be given by  $\mu^2 = n^2 a^2$ , where  $n$  is the mean number of Rh\* and  $a$  is the average single photon response. If the time course of  $a$  varies negligibly, the response variance will be dominated by Poisson fluctuations in photon absorption, i.e.,  $\sigma^2 = na^2$ , where  $\sigma^2$  is the measured time-dependent ensemble variance. The ratio of the mean-squared response to the variance,  $\mu^2/\sigma^2$ , estimates the mean number of Rh\* per flash,  $n$ . The mean number of Rh\* divided by the photon density estimates the collecting area. We performed the scaling procedure for the response ensemble at each flash strength to obtain at least 3 values for  $n$  (see Fig. 3G). The collecting area was estimated from the slope of the best fit line when  $n$  was plotted against photon density.

**Electroretinograms.** ERGs measure the changes in the transretinal potential in response to light, reflecting the changes in flow of current across the retina (Fig. 1A–C). To measure electroretinograms of isolated tissue, approximately half of one retina was placed in a custom cylindrical plastic ERG chamber (Fig. 1A). The chamber provided a single electrical pathway across the retina, between a recording electrode at one side and a reference electrode at the other (Fig. 1B). To this end, the cylinder was composed of hollowed-out bottom and top pieces that joined at a ridge around the perimeter. With the top and bottom pieces adjoined, there remained a space between the pieces  $\sim 300 \mu\text{m}$  high and 12 mm in diameter to accommodate a flat-mounted piece of retina. Two discs of thin transparent plastic encased the retina, serving to separate solutions in the top and bottom of the ERG chamber. These distinct pools of solution were electrically isolated save through a pair of holes through the transparent discs, each hole 0.7 mm in diameter. For recordings, a circular filter paper was attached through dots of grease to the transparent plastic disc on the top piece of the ERG chamber and a hole was punched in the paper in line with the hole in the transparency. One half of one retina was placed ganglion-cell-side-down onto the filter paper, and the bottom and top pieces of the ERG chamber were sealed together by a ring of grease around the tissue. The ERG chamber was then placed into the microscope. Recordings with drifting baseline or unstable light responses, likely due to leak of current around the retina, were rejected.

Light stimuli were delivered as described above. The amplitude of the response to a constant stimulus varied across pieces, and depended on the quality of the electrical seal. It was therefore difficult to compare absolute amplitudes of the ERG responses across pieces of tissue, and normalized responses were typically used instead for comparison. Maintaining stable recordings at  $36^\circ\text{C}$  was greatly facilitated by perfusion of both sides of the retina. *In vitro* ERGs had a shape and sensitivity similar to *in vivo* measurements (see Fig. 9) (Saszik et al., 2002; Robson et al., 2004; Herrmann et al., 2010), including oscillations during the rising phase of the b-wave (Fig. 1C) elicited by a 0.2 ms flash (rather than the 10 ms flash used for most experiments).



**Figure 1.** ERG methods and description of paired-pulse derivation of rod responses. **A**, ERG recording apparatus; scale is approximate. The space between top and bottom where the retina was placed measured  $\sim 12$  mm in diameter and  $300 \mu\text{m}$  high. **B**, Schematic cross-section of retina and measurement of the ERG. OS, Photoreceptor outer segments; ONL, outer nuclear layer; OPL, outer plexiform layer; INL, inner nuclear layer; IPL, inner plexiform layer; GCL, ganglion cell layer. **C**, Twenty raw ERGs (gray) and average (black) following a 0.2 ms flash producing  $40 \text{ Rh}^*/\text{rod}$  in Ames/Ames' conditions. Reduction in current flowing along the photoreceptor outer segments in response to light results in the corneal-negative deflection of the ERG, termed the a-wave, while the b-wave reflects responses of second-order neurons. **D, E**, Paired-pulse analysis to derive rod responses in Ames/Ames', with 4 mM aspartate, which eliminates the b-wave (see Materials and Methods). **D**, Overlaid responses (gray traces) to paired test ( $840 \text{ Rh}^*/\text{rod}$ ) and probe flashes ( $1680 \text{ Rh}^*/\text{rod}$ ) presented at increasing intervals; the response to the test flash alone is shown in black. **E**, The probe response rising phases from **D** (short traces, baseline corrected) are compared with the response to the probe flash alone (complete trace). Comparison is made at the time point of fastest rate of change of the response to the probe alone (dot, black dashed line). The ratio of the rising phase amplitude, at this point, to the amplitude of the probe-alone response (dashed black line) maps out a derived rod response. The recovery points are then fit with a model recovery function ( $F'(t) = \{1 + \exp[a(t - t_{1/2})]\}^{-h}$ ,  $t_{1/2} = 405 \text{ ms}$ ,  $a = 14.4$ ,  $h = 0.3$ ). For illustration purposes in this panel, the black trace plots the function,  $p[1 - F'(t)]$ , where  $p = -45 \mu\text{V}$  is the amplitude of the response to the probe alone (dashed black line).

**Derivation of rod responses from ERGs.** To derive a rod response from the ERGs, we used a paired-pulse technique in which the response to a bright, saturating probe flash was measured as a function of a time separation between an initial test flash and the probe flash (Fig. 1D) (Lyubarsky and Pugh, 1996; Pepperberg et al., 1997; Hetling and Pepperberg, 1999). The derivation of a rod response involved comparing the response to a probe flash that followed a test flash with the response to the probe flash alone (Fig. 1D,E). Responses to the probe flash were isolated by subtracting the response to the test flash alone. The amplitude of the resulting responses was measured at the time of the a-wave's maximum rate of change in response to the probe alone, typically  $\sim 15$  ms (Fig. 1E). The probe flash strength and measurement during the rising edge of the resulting response were used primarily to avoid bleaching the rod photopigment with probe flashes sufficiently bright to produce an a-wave that peaked at  $< 15$  ms. Some recording conditions also produced dramatic alterations in the ERG b-wave, which could mask or reveal differing amounts of the a-wave; measuring the probe response during its leading edge minimized sensitivity to such alterations of the b-wave. Elimination or disruption of the b-wave occurred either pharmacologically, following addition of 4 mM aspartate to the perfusion (Fig. 1D,E) (Green and Kapousta-Bruneau, 1999), or as a result of the experiment media (see Fig. 8).

The two parameters of importance in the paired-pulse protocol are the strength of the test flash and the delay between the test and probe flashes. The probe flash strength is constant and is sufficiently bright to saturate the rod response and hence determine the dark current remaining at the time of the probe flash. The remaining dark current depends on the strength of the test flash and the delay between test and probe flashes. In



one extreme, the test flash saturates the rods and the presentation of the probe flash during this saturation produces no subsequent response. In the other extreme, the time between test and probe flashes is sufficient to allow the rods to recover fully, and the response to the probe flash is identical to that when the probe flash is presented alone. Cones are not expected to contribute substantially to the ERG a-wave (Lyubarsky and Pugh, 1996); consistent with this expectation, ERGs from rod transducin knock-out mice (*GNAT1*<sup>-/-</sup>) (Calvert et al., 2000) recorded in Ames' medium did not exhibit noticeable a-waves over the range of flash strengths used here (data not shown).

The ratio of the probe responses at each time delay,  $t_i$ , to the response to the probe alone estimates the fraction of remaining rod circulating current as a function of time,  $R(t_i)$ . We then define:

$$F(t_i) = 1 - R(t_i)$$

where  $F(t_i)$  gives the derived rod response. For bright test flashes that saturate the rod circulating current, we fit  $F(t_i)$  with the following model of the recovery phase:

$$F'(t) = \frac{1}{(1 + e^{a(t-t_{1/2})/h})^b}$$

where  $t > 0$  is time, and  $a$ ,  $t_{1/2}$ , and  $h$  are constants. We used this model to interpolate between the measured points and estimate the time required,  $T_{sat}$ , for the responses to recover to a criterion level of 0.5. We repeated this procedure for three flash strengths and estimated the dominant time constant for decay of light-activated PDE activity from the slope of the relation between  $T_{sat}$  and the logarithm of the flash strength.

We used different recovery criteria to determine dominant time constants for single-cell (criterion of 0.15) (Krispel et al., 2006; Burns and Pugh, 2009; Doan et al., 2009) and ERG recordings (criterion of 0.5) (Lyubarsky and Pugh, 1996; Hetling and Pepperberg, 1999) (see Fig. 6). The low criterion in the case of single rods minimizes sensitivity to the trajectory of the response recovery by "catching" the response soon after recovery begins. Because the ERG reflects activity of many rods, however, a similar low criterion would bias the dominant time constant toward those rods that recover most quickly. Instead, the criterion of 0.5 should better capture the average response. Consistent with this, the dominant time constant for the ERG-derived responses depended on the criterion that defined recovery, with lower criteria producing faster dominant time constants. Importantly, lower criteria produced faster dominant time constants in all recording conditions, and the similarity of the dominant time constant across recording conditions remained for criteria between 0.1 and 0.5 ( $p > 0.24$ , Mann–Whitney–Wilcoxon rank-sum tests (MWW test), comparison of  $\tau_D$  between Ames'/Ames' and Ames',Ca,Mg/Ames',Ca,Mg). Further, the differences between ERG-derived and single-cell recovery kinetics exceed those reported if similar criteria (0.15 or 0.5) are used to measure time in saturation for both types of data.

We also applied the paired-pulse technique to test flashes that produced approximately half-saturating responses ( $\sim 30$  Rh\*). We fit the normalized derived rod response with the model from Hetling and Pepperberg (1999):

$$F'(t) = \gamma[1 - e^{-a(t-t_d)^2}] \cdot e^{-t/\tau_w}, t > t_d,$$

where for  $t < t_d$ ,  $F'(t) = 0$ .

Derived responses with and without 4 mM aspartate added to the perfusion solution were similar, and hence the measurements were combined.

Experiments were designed such that the cumulative light exposure during a typical recording bleached <3% of the total rod photopigment; furthermore, the retina was allowed to recover for >25 s between test flashes. Data were rejected if responses to periodically delivered probe flashes exhibited systematic changes indicative of long term adaptation.

A concern about this approach is that bright test flashes could alter the kinetics of the probe response and hence bias the derived responses. Two observations indicated such effects were minimal. First, the impact of changes in the kinetics of the probe response would be expected to de-

pend on probe flash strength—with smaller effects for brighter probe flashes which more quickly saturate the rods. However, derived responses to probe flashes ranging from 840 Rh\*/rod to 9900 Rh\*/rod were similar ( $R^2 = 0.06$ ,  $p = 0.5$ ,  $n = 9$ ). Likewise, changes in the kinetics of the probe response could cause the derived measurements to depend on the rising phase time point where we compared probe responses. However, derived rod responses (see Figs. 6, 7) did not systematically depend on the time at which we measured the probe response ( $R^2 = 0.01$ ,  $p = 0.15$ ,  $n = 175$ , for all points  $\sim 12$ –20 ms in Ames'/Ames' control conditions,  $\sim 12$ –40 ms in Ames'/Ames' with 4 mM aspartate, Fig. 1E). As a further check, we performed identical experiments on individual rods, which should show similar sensitivity to changes in probe flash kinetics. Dominant time constants measured using the paired-flash approach and measured directly from the test flash response did not differ noticeably (see Fig. 6D,E).

**Data collection and analysis.** Recorded responses were filtered at 30 or 300 Hz (8-pole Bessel) and digitized at 1 kHz by an ITC16 or ITC18 A-to-D converter (HEKA). Digitized responses were collected and stored using custom recording software written in C and Igor. Single-cell data were digitally filtered at 15 Hz, while ERG data were low-pass filtered at 100 Hz and the 60 Hz sine wave that best fit baseline noise was subtracted, unless otherwise noted. Data were organized using custom software (Phyzion Consulting) and further analyzed using custom Matlab routines (MathWorks).

**Live/dead staining assay.** To test the effects of tissue storage solutions on the viability of the retina, we imaged retinal cell body layers following staining with a mammalian Live/Dead Cell assay kit (Invitrogen). The protocol assays the integrity of the cell membrane by staining with calcein AM and ethidium homodimer-1 (EthD-1). If the membrane is intact, the cell retains esterase function capable of hydrolyzing the AM group, and cells will retain the calcein dye in their cytosol. If the membrane is compromised, EthD-1 can access the nucleus and stain DNA. Before staining, 4  $\mu$ M EthD-1 and 2  $\mu$ M calcein AM dyes were added to 10 ml of Ames' medium. Approximately one third of a retina was placed in the staining solution for 30 min before adhering the tissue to filter paper (Anodisc, Whatman), securing the filter paper to the bottom of a recording chamber with vacuum grease, and placing the mounted tissue in the microscope. The mounted retina was then always perfused with Ames solution at room temperature ( $\sim 24^\circ\text{C}$ ). Confocal image z-stacks were taken through the layers of cell bodies with a Nikon laser scanning confocal microscope system using lasers with peak wavelengths of 488 nm to excite the calcein dye and 543 nm to excite the EthD-1 homodimer. Image analysis was performed in NIH ImageJ. The pixel intensity in each channel was adjusted *post hoc* to linearly span the entire range of pixel intensities. Membrane integrity, and cell health, was assessed by the appearance of EthD-1 staining in the nucleus. The cells visibly stained in the EthD-1 channel in ganglion cell layers were counted by hand using the NIH ImageJ Cell Counter plugin.

**Statistics.** All values reported in the text are given as mean  $\pm$  SEM. Significance was tested using the MWW test.

## Results

Differences in the reported properties of the single-photon responses of mouse rods are strongly correlated with differences in experimental protocols—specifically the solutions used to store and perfuse the retina. We start by determining how the responses of single rod photoreceptors depend on recording conditions. We then measure the effect of recording conditions on estimates of rod function derived from the electroretinogram, and compare those measures to the single-cell data. Finally we compare the physiological and morphological viability of the retina under these different conditions.

### Experimental conditions alter the kinetics and amplitude of responses of single rods

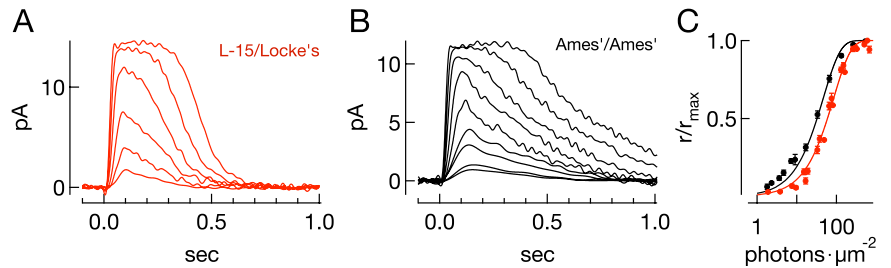
We recorded the light responses of single-rod photoreceptors in several experimental conditions (Table 1). Commonly used pro-

protocols include storing the retina in Leibowitz L-15 culture media and recording in Locke's solution (L-15/Locke's, Table 1) (Baylor et al., 1984; Xu et al., 1997; Makino et al., 2004; Krispel et al., 2006) or storing and recording in Ames' solution (Ames'/Ames', Table 1) (Field and Rieke, 2002; Doan et al., 2009; Okawa et al., 2010). Consistent with previous results, these protocols produced substantial differences in the kinetics and sensitivity of the rod responses (Table 2; Fig. 2A–C).

Differences in rod sensitivity could be due either to differences in the ability to capture incident photons and produce an electrical response, or to differences in the amplification of the signal produced by rhodopsin activation. To test for differences in effective photon capture, we used a frequency of seeing analysis (Baylor et al., 1979b) applied to responses of *GCAP<sup>-/-</sup>* mice (Mendez et al., 2001; Burns et al., 2002). These rods produce larger and more distinct single-photon responses than wild-type rods because of the lack of  $\text{Ca}^{2+}$ -feedback control of cGMP synthesis (Fig. 3A,D) This allowed accurate determination of the fraction of presentations of a fixed-strength flash resulting in a response (effective absorption of one or more photons) or failure (no effective photon absorptions) (Fig. 3B,E). The failure probability from Poisson statistics is  $p_0(I) = e^{-mI}$ , where  $I$  is the photon density ( $\text{photons} \cdot \mu\text{m}^{-2}$ ), and  $m$  is the collecting area of the rod in units of  $\mu\text{m}^2$ . Figure 3C plots the response probability ( $1 - p_0(I)$ ) against flash strength for *GCAP<sup>-/-</sup>* rods in L15/Locke's (red) and Ames'/Ames' (black); smooth fits assume Poisson statistics. Estimated collecting areas were not different (Fig. 3F,  $0.50 \pm 0.04 \mu\text{m}^2$ , mean  $\pm$  SEM,  $n = 5$ , vs  $0.47 \pm 0.02 \mu\text{m}^2$ ,  $n = 8$ ,  $p = 0.52$ , MWW test).

Collecting areas are often estimated by comparing the time-dependent variance of the responses to a repeated dim flash with the square of the mean response. Assuming that the time-dependent variance is dominated by Poisson fluctuations in photon absorption, the ratio of the square of the mean response to the variance provides an estimate of the mean number of effective rhodopsin activations,  $n$ , produced by the flash (see Materials and Methods). We used the initial part of the response to compute  $n$  to avoid variability in the falling phase of the response which could systematically bias the estimate. We applied this procedure to ensembles of responses to increasing flash strengths from wild-type rods in the L15/Locke's and Ames'/Ames' conditions (Fig. 3G, example cell in Ames'/Ames'). The thick traces in Figure 3G compare the scaled ensemble mean-squared response with the ensemble variance (thin traces) for three flash strengths. The collecting area,  $m$ , was determined from the slope of the  $n$  vs  $I$  data (Fig. 3G, main panel), where  $I$  is the flash strength in  $\text{photons} \cdot \mu\text{m}^{-2}$ . We then divided the mean response to each flash by  $n' = mI$ , the flash strength in terms of  $\text{Rh}^*$ , to arrive at an estimate of the single-photon response. Finally we averaged these estimates across flash strengths.

Figure 3, H and I, shows, for each cell, the area (charge) under the single-photon response derived from the variance analysis plotted against that derived using a fixed collecting area. For the Ames'/Ames' condition (Fig. 3H), single-photon responses estimated from variance analysis do not differ systematically from those estimated using a fixed collecting area, likely because variability of the initial  $\sim 200$  ms of the response is dominated by Poisson fluctuations (Doan et al., 2009). For the L15/Locke's



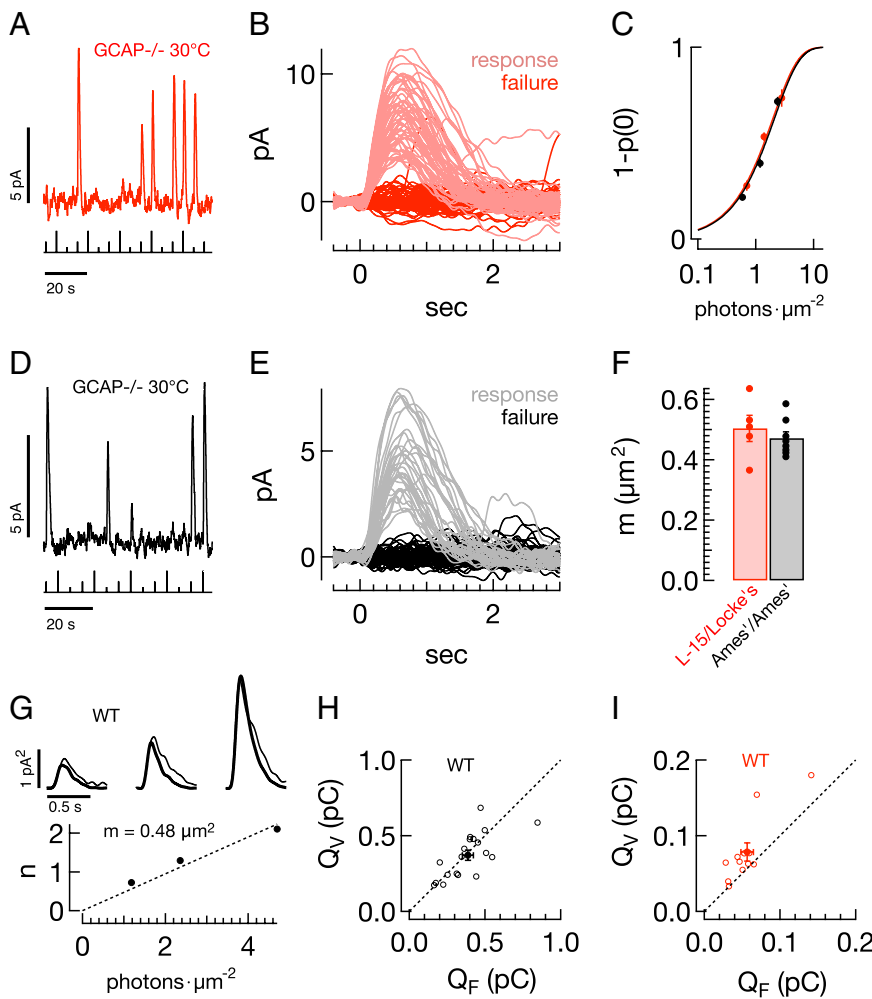
**Figure 2.** Rod sensitivity depends on the experimental protocol. **A, B**, Family of flash responses from a rod in L-15/Locke's (**A**) and in Ames'/Ames' (**B**) conditions. Flash strengths double over ranges of 17–524 (**A**) and 2–516 (**B**) photons  $\cdot \mu\text{m}^{-2}$ . **C**, The peak fractional current suppression ( $r/r_{\text{max}}$ ) as a function of the flash strength (mean  $\pm$  SEM,  $n = 19$  for Ames'/Ames' and 15 for L-15/Locke's). Smooth curves are the saturating exponentials that best fit the population data, with  $I_0 = 57$  photons  $\cdot \mu\text{m}^{-2}$  for L-15/Locke's (red), and 28 photons  $\cdot \mu\text{m}^{-2}$  for Ames'/Ames' (black).

condition, estimates of the single-photon response using variance analysis are consistently larger than those using the fixed collecting area (Fig. 3I). This larger single-photon response likely results because the assumed dominance of Poisson variation of rhodopsin activation does not hold, and variation in the response to a single  $\text{Rh}^*$  contributes to the time-dependent variance (see Materials and Methods). The results of Figure 3 raise a general note of caution about the use of variance analysis, i.e., that any manipulations of phototransduction that alter response variability could produce artifactual changes in the collecting area, resulting in errors in estimating the number of activated rhodopsin molecules and the single-photon response.

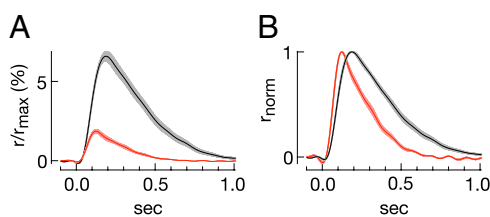
The similarity of the collecting areas determined by the frequency of seeing analysis in *GCAP<sup>-/-</sup>* mice suggests that the differences in sensitivity in L15/Locke's and Ames'/Ames' are due to differences in the single-photon response. We estimated the average responses of mouse rods to activation of single rhodopsin molecules under an array of conditions assuming a collecting area of  $0.5 \mu\text{m}^2$ . Single-photon responses using the L15/Locke's condition were smaller and shorter than those in Ames'/Ames' (Fig. 4A,B; Table 2). Both factors contribute to the reduced sensitivity (Baylor et al., 1979a).

Do the differences in kinetics and sensitivity of the L15/Locke's and Ames'/Ames' conditions reflect differences in retina storage or recording solutions? To answer this question we used warm Locke's medium for both storage and recording. The stimulus-response relation for the Locke's/Locke's condition fell between the red and black curves in Figure 2C (Table 1). The estimated single-photon response also showed intermediate amplitude and kinetics (Fig. 5A,B, blue curves), indicating that measured responses were sensitive to both storage and recording conditions. We separated the effects of temperature and L-15 storage by measuring dim flash responses in rods stored in warm L-15 medium (Table 2). In these rods, kinetics and dim flash response amplitude fell between those of rods from retina stored in cold L-15 and Locke's solution, suggesting that both the cold and the L-15 contribute to changes in kinetics.

The remaining differences between the Locke's/Locke's and Ames'/Ames' conditions appeared due to differences in divalent concentrations, as responses measured in Ames' and Locke's media with matched  $\text{Ca}^{2+}$  and  $\text{Mg}^{2+}$  concentrations were not distinguishable (Fig. 5A,B, green curve; Table 2).  $\text{Mg}^{2+}$  differs more substantially between Locke's and Ames', and Ames' medium with elevated  $\text{Mg}^{2+}$  alone also produced fractional responses and kinetics indistinguishable from those in Locke's/Locke's, though dark currents were reduced (Fig. 5A,B, magenta curve; Table 2).



**Figure 3.** The rod collecting area does not depend on the experimental protocol. **A–F**, Frequency of seeing analysis for *GCAP*<sup>-/-</sup> rods. **A**, Raw current trace from *GCAP*<sup>-/-</sup> rod in L-15/Locke's, showing responses to three repeated flash strengths (0.7, 1.4, or 2.8 photons · μm<sup>-2</sup>, stimulus trace below, 30°C). **B**, Superimposed responses to a flash with strength of 1.4 photons · μm<sup>-2</sup>, *p*<sub>0</sub> = 0.54. **C**, The average probability of observing a response (1 - *p*<sub>0</sub>), as a function of flash strength, for populations of *GCAP*<sup>-/-</sup> rods in both L-15/Locke's (red, *n* = 5) and Ames'/Ames' (black, *n* = 8). Smooth lines show the Poisson prediction, 1 - *p*<sub>0</sub>(*I*) = 1 - e<sup>-*mI*</sup>, with *m* = 0.50 μm<sup>2</sup> in L-15/Locke's and 0.47 μm<sup>2</sup> in Ames'/Ames'. **D**, Raw current trace for Ames'/Ames' condition (flash strengths: 0.6, 1.2, or 2.4 photons · μm<sup>-2</sup>). **E**, Superimposed responses to a flash with strength of 1.2 photons · μm<sup>-2</sup>, *p*<sub>0</sub> = 0.64. **F**, Collecting areas (CA) from individual *GCAP*<sup>-/-</sup> rods in L-15/Locke's (red, 0.50 ± 0.04 μm<sup>2</sup>), and Ames'/Ames' (black, 0.47 ± 0.02 μm<sup>2</sup>). **G–I**, Comparison of variance and frequency of seeing estimates. **G**, Top, Variance (thin trace) and scaled mean-squared response (thick trace) in a representative rod in Ames'/Ames' (see Materials and Methods) for flash strengths of 1.18, 2.36, and 4.72 photons · μm<sup>-2</sup>. Scale bars are 0.5 s and 1 pA<sup>2</sup>. The reciprocal of the scale factor applied to the mean-squared response gives the mean number of effective photon absorptions, *n*. Bottom, The collecting area, *m*, is the slope of the best fit line (dotted line) through the *n* vs *I* data, in this case, 0.48 μm<sup>2</sup>. **H**, Population data in Ames'/Ames' for the integral of the estimated single-photon response using the variance-to-mean-squared method, *Q*<sub>V</sub> (y-axis, pC) and using a fixed collecting area of 0.5 μm<sup>2</sup>, *Q*<sub>F</sub> (x-axis, pC). Filled circles give the average across the population ± SEM. **I**, Same as **H** for L-15/Locke's.



**Figure 4.** The single-photon response depends on the experimental protocol. **A**, Average single-photon responses as a percentage of suppressible current for rods in L-15/Locke's (red, *n* = 15) and Ames'/Ames' (black, *n* = 20). Contours are the SEM across a population of rods. **B**, Normalized responses from **A**.

**The dominant time constant of recovery depends on extracellular magnesium**

Differences in the time course of the single-photon responses in Ames'/Ames' and L15/Locke's (Fig. 4) could reflect differences in cGMP hydrolysis or synthesis or both. The kinetics of the decay of light-activated PDE activity can be monitored from the recovery from bright, saturating flashes (Pepperberg et al., 1992). Recovery of the circulating current following a bright flash requires that the rate of synthesis of cGMP exceeds the rate of hydrolysis by PDE. Under the assumptions, (1) that cyclase synthesis rate is constant while the rod remains saturated, and (2) that PDE deactivation follows an exponential decay, the time at which the rod current begins to recover, *T*<sub>sat</sub> should scale linearly with the logarithm of the flash strength (Fig. 5C,D). The slope of this initially (< 1000 Rh\*) linear relationship between *T*<sub>sat</sub> and the logarithm of the flash strength (Fig. 5D, best fit line) is the time constant of the assumed exponential decay of PDE hydrolysis and is often termed the dominant time constant of recovery, τ<sub>d</sub>.

Figure 5, E and F, collects measurements of the dominant time constant across conditions. The dominant time constants for rods recorded in L-15/Locke's are consistent with previously reported values (Fig. 5E,F; Table 2) (Chen et al., 2000; Calvert et al., 2001; Makino et al., 2004; Krispel et al., 2006; Gross and Burns, 2010). The dominant time constants for the Ames'/Ames' condition were significantly longer, again consistent with past work (Doan et al., 2009; Gross and Burns, 2010). Dominant time constants for the remaining conditions (Locke's/Locke's and Ames' with elevated divalents) were all significantly shorter than the Ames'/Ames' condition, and closer to the L-15/Locke's condition (Fig. 5F; Table 2). Thus, changes in Mg<sup>2+</sup> and Ca<sup>2+</sup> can significantly affect the dominant time constant and account for differences in dominant time

constant across recording conditions.

The dependence of the dominant time constant for the decay of light-activated PDE activity on divalent concentrations is striking given the relatively small changes in the single-photon responses. In the L-15/Locke's condition, the time constant for the decay of the dim flash response and the dominant time constant are similar. Such a similarity is expected if the rate of cGMP synthesis is sufficiently rapid that the cGMP concentration closely tracks the PDE activity, and hence the recovery of the dim flash response is limited by the rate of decay of PDE activity. Such rapid cGMP synthesis, however, appears not to always be the case; specifically, in Ames' with elevated divalents the cGMP con-



centration following a dim flash appears to recover to its dark value slowly compared with the decay of light-activated PDE activity.

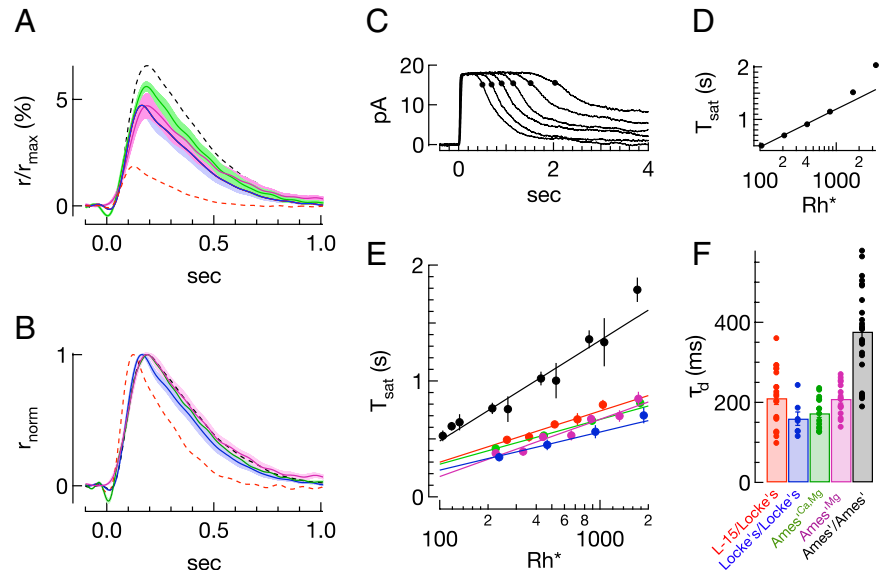
### Dominant time constants derived from the ERG are similar across experimental conditions

Rod responses *in vivo* are often assayed using the ERG. To determine whether we could use the ERG to identify *in vitro* conditions that closely replicate the *in vivo* rod responses, we derived rod responses from *in vitro* ERGs under a variety of conditions. We compared the derived rod responses to the single-cell data summarized above and to properties of *in vivo* ERGs.

We used a paired-flash approach to derive the time course of the rod response from the ERG (Lyubarsky and Pugh, 1996). In this approach, the response to a “test” flash is probed using a second, bright “probe” flash (see Materials and Methods). We then compare the a-wave, the leading edge of the ERG, in response to the probe flash either alone or following the test flash. We describe the other components of the ERG below; these differed substantially across recording conditions.

We performed the paired-flash analysis using two types of test flashes: bright flashes producing  $>200$  Rh\*/rod, and “dim” flashes producing  $\sim 30$  Rh\*/rod. Previous paired-flash studies using comparable flash strengths show that rod responses derived from *in vivo* ERGs recover quickly following bright flashes (dominant time constant of 140–200 ms (Lyubarsky and Pugh, 1996; Hetling and Pepperberg, 1999) and dim flashes (time-to-peak of 86 ms, decay time constant 163 ms (Hetling and Pepperberg, 1999)). The relatively fast kinetics of *in vivo* ERG-derived responses, and their resemblance to the kinetics of single cell currents in L-15/Locke’s (Table 2), support the conclusion that L-15/Locke’s responses are physiologically accurate (Burns and Pugh, 2010). This conclusion relies on the sensitivity of the rod responses derived from the ERG to alterations in the rod transduction currents.

We focused on comparing derived rod responses for conditions that produced pronounced differences in the responses of single rods. In the bright flash regime we compared the derived dominant time constant for recovery in Ames’ with that in Ames’ with elevated divalent concentrations; these conditions produced dominant time constants in single rods that differed by a factor of  $\sim 2$  (380 ms vs 170 ms, Table 2). Figure 6A superimposes raw traces from such an experiment. The response to the probe pulse (second flash,  $\sim 830$  Rh\*/rod) reflects the extent to which the rods have recovered following the test flash ( $\sim 830$  Rh\*/rod). Responses were aligned to the time of the probe flash and corrected by the amplitude at that point (Fig. 6B). Values of  $R(t_i)$  were evaluated at the point of maximum rate of change ( $\sim 16$  ms) of the response to the probe flash alone (thick black trace). Figure 6C plots response suppression ( $1 - R(t_i)$ ) as a function of the delay between test and probe flashes. The gray line is an empirical fit (see Materials and Methods) to the recovery used to interpolate between the measured points.

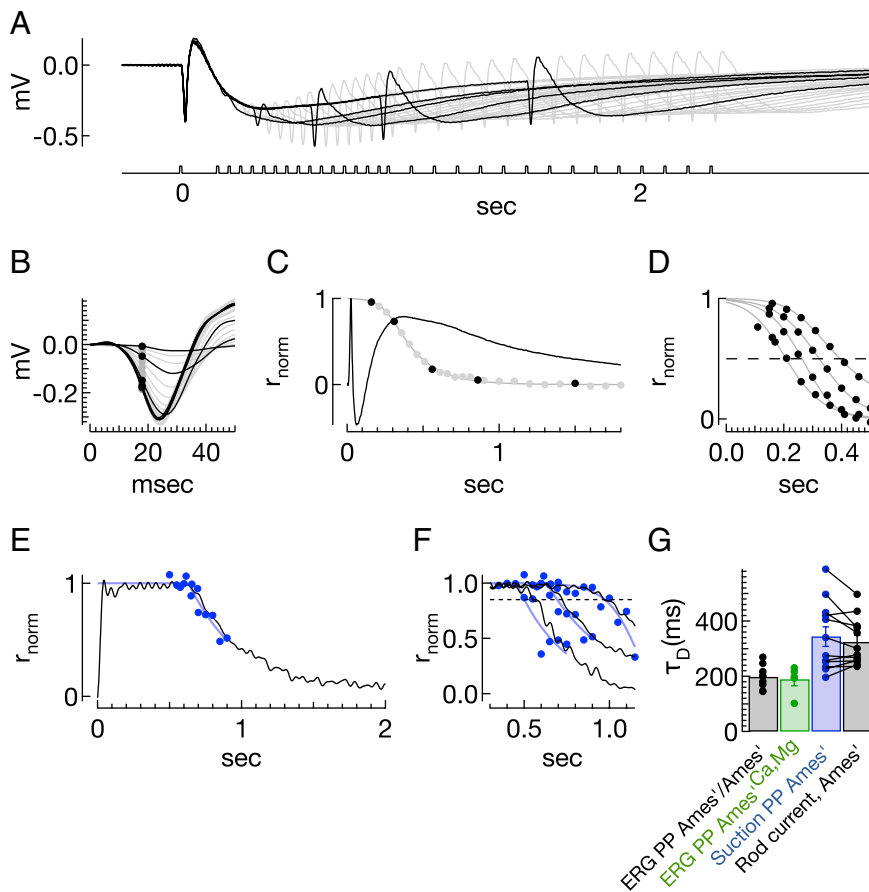


**Figure 5.** Dependence of single-photon response and dominant time constant of recovery on divalent concentration. **A**, Average single-photon responses in different conditions ( $\pm$  SEM), using a fixed collecting area: Locke’s/Locke’s (blue), Ames’<sup>Ca,Mg</sup>/Ames’<sup>Ca,Mg</sup> (green), and Ames’<sup>Mg</sup>/Ames’<sup>Mg</sup> (magenta). The Ames’ conditions match either Mg<sup>2+</sup> alone or both Ca<sup>2+</sup> and Mg<sup>2+</sup> concentrations in Locke’s. **B**, Normalized single-photon responses in the same conditions and colors. Dashed lines in **A** and **B** show single-photon responses from Figure 4: L-15/Locke’s (red), Ames’/Ames’ (black). **C–F**, Illustration of the technique. For a family of responses to bright flashes (Ames’/Ames’; 105, 208, 415, 829, 1658, 3315 Rh\*), the time at which a response recovers 15% of the dark current (**C**) is plotted against the logarithm of the flash strength (**D**). The slope of this linear relationship estimates the dominant time constant of recovery, in this case, 312 ms. **E**, Average ( $\log(i)$ ,  $T_{sat}$ ) pairs for all cells in a given condition; only points below 1000 Rh\* are fit. Slopes of best fit lines are as follows: L-15/Locke’s, 192 ms; Locke’s/Locke’s, 143 ms; Ames’<sup>Ca,Mg</sup>/Ames’<sup>Ca,Mg</sup>, 168 ms; Ames’<sup>Mg</sup>/Ames’<sup>Mg</sup>, 214 ms; Ames’/Ames’, 371 ms. **F**, Summary of dominant time constants from fits to data from individual rods as in **D** (see also Table 1).

Two control experiments indicated that the derived responses were dominated by the rods. First, adding 4 mM aspartate to the perfusion solution to eliminate responses produced downstream of the rods did not significantly change the results. Second, in retinas from *GNAT<sup>-/-</sup>* mice (Calvert et al., 2000), bright flashes ( $>1000$  Rh\*/rod) produced ERGs  $\sim 100$  times smaller than those of wild-type mice (data not shown), indicating that mouse cones make a negligible contribution to the a-wave at these flash intensities.

We repeated the procedure of Figure 6A–C for at least three test flash strengths and used the resulting measures of  $T_{sat}$  (Fig. 6D, inset) to determine the dominant time constant  $\tau_d$ . In the standard Ames’ condition (Ames’/Ames’),  $\tau_d$  of the ERG-derived rod response was  $200 \pm 15$  ms (mean  $\pm$  SEM,  $n = 9$ , Fig. 6G); this is within the range of *in vivo* ERG measurements, but is substantially different ( $p < 0.001$ , MWW test) than the average dominant time constant measured from single rods ( $376 \pm 23$  ms, Fig. 5F). Increasing the divalent concentration in Ames did not change  $\tau_d$  (green column in Fig. 6G,  $n = 5$ ,  $p = 1$  MWW test). Thus the  $\sim 2$ -fold change in  $\tau_d$  of the single-rod responses is not reflected in the ERG-derived rod responses.

Could the differences in  $\tau_d$  measured using paired-flash ERGs and single-cell recordings reflect differences in experimental protocol? To check this, we measured  $\tau_d$  from single rod outer segments using both saturation time following single bright flashes (as in Fig. 5) and paired-pulse experiments in the same cell (Fig. 6E–G, Ames’/Ames’,  $n = 12$ ). The two approaches provided indistinguishable ( $p = 0.9$ , MWW test) and highly correlated measures of  $\tau_d$  (connected points in Fig. 6G;  $r = 0.94$ ). In both cases the average  $\tau_d$  exceeded that derived from the ERG ( $p < 0.005$  for both methods, MWW test). The average  $\tau_d$  of single rods using



**Figure 6.** Dominant time constants of ERG-derived rod responses show little dependence on divalent concentration. **A**, ERGs in response to increasing delays between a test (834 Rh\*) and probe (834 Rh\*) flash (stimulus trace below). Black responses are followed through the subsequent panels. **B**, Rising phases of the responses to the probe flash, aligned to the time of the flash, when presented alone (thick black trace), and when following a test flash. Dots represent points of comparison to the thick curve, at the point of fastest change (18 ms). This comparison measures the fraction of the remaining rod response,  $R(t)$ . **C**, Fractional response suppression,  $F(t) = 1 - R(t)$ , as a function of the delay between test and probe flashes; the smooth gray curve is the best fit recovery function (see Materials and Methods) ( $t_{1/2} = 210$  ms,  $a = 30.6$ ,  $h = 0.1$ ). The thin black trace is the ERG in response to the test flash alone. **D**, Four derived rod responses (298, 417, 596, and 834 Rh\*). The dashed line plots the 50% criterion level used to determine the dominant time constant. **E**, Paired-pulse analysis in single rod outer segments (664 Rh\*) ( $F$ , 332, 664, 1327 Rh\*). **G**, Comparison of dominant time constants for ERG-derived responses in Ames'/Ames' ( $n = 9$ ) and Ames'<sup>Ca,Mg</sup>/Ames'<sup>Ca,Mg</sup> ( $n = 5$ ) conditions and single-rod recordings using paired-pulse analysis or directly from the current response (as in Fig. 5). Lines for single-cell recordings connect measures from the same cell.

the paired-pulse approach does not differ significantly from the population of rods in Figure 5 ( $p = 0.6$ , MWW test), indicating that the paired-pulse approach probes the rod response as expected.

#### Kinetics of ERG-derived half-saturating responses in Ames' are consistent with *in vivo* measurements

We also used the paired-pulse approach to derive the time course of the rod response to a subsaturating flash (Fig. 7). The test flash (~30 Rh\*) used in these experiments suppressed, at peak, approximately half of the response to the probe flash, reflecting a compromise between delivering few photons and eliciting a response large enough to produce measurable current suppression. This test flash is approximately equal to that used in previous ERG measurements (Hetling and Pepperberg, 1999).

In comparing rod outer segment current responses to ERG-derived estimates, it is important to note that in single rods, half-saturating current responses differ in kinetics substantially from single-photon responses (Fig. 7A,B; Table 2) (Sampath et al.,

2005; Burns and Pugh, 2010). Furthermore, differences across solutions are not as large when comparing half-saturating responses: i.e., the range of values for the time-to-peak across recording conditions is larger for single-photon responses (Fig. 7A, x-axis), than for half-saturating responses ( $y$ -axis). Nonetheless, the half-saturating response measured in L-15/Locke's still reaches peak earlier than in Ames' (Table 2; Fig. 7B). Consequently, we compared ERG-derived half-saturating responses in L-15/Locke's and Ames'/Ames' conditions. As in the bright flash regime, we measured the suppression of the probe response for a range of delays between test and probe flashes (Fig. 7C,D, bottom). We averaged the responses at each time point across retinæ for each condition and fit the population average with a model for *in vivo* ERG-derived responses (Hetling and Pepperberg, 1999) (Fig. 7E, thin traces). The addition of 4 mM aspartate to the perfusion media did not systematically affect the parameters of the fit, so measurements obtained in the absence and presence of aspartate are combined. The derived rod response in L-15/Locke's closely tracks that in Ames'/Ames', with peaks at 82 and 76 ms, respectively, and with recovery time constants of 191 ms in L-15/Locke's and 139 ms in Ames'/Ames'. For comparison, the gray stippled line shows the best fit model from the Hetling and Pepperberg *in vivo* ERG recordings (Hetling and Pepperberg, 1999) with a time-to-peak of 87 ms and recovery time constant of 163 ms. Thus, derived rod responses in Ames'/Ames' and L15/Locke's are not readily distinguished, and both are consistent with *in vivo* measurements.

#### Experimental conditions alter retinal morphology, health, and physiology

The results of the paired-pulse analysis indicate that storage and perfusion solutions do not significantly affect ERG-derived rod responses. We were also interested in how well different experimental conditions maintained post-rod retinal processing, permitting rod responses to be tracked through retinal circuits. We found profound differences in the responses and health of the post-rod circuitry under different experimental conditions.

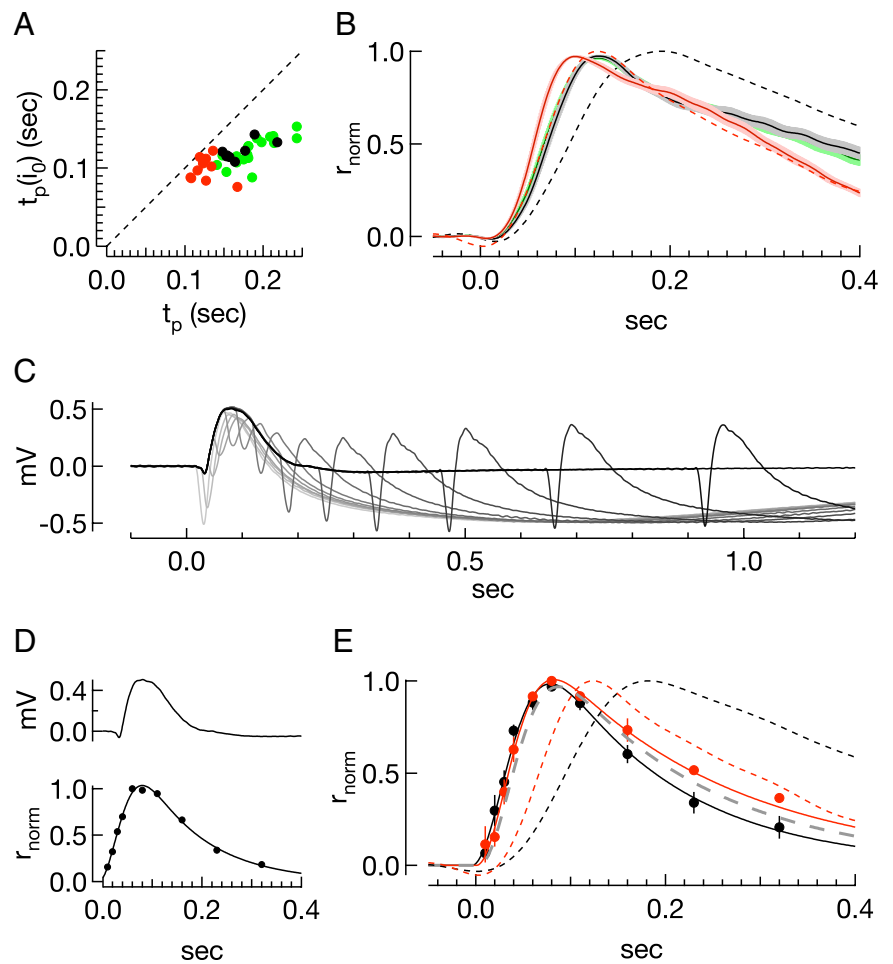
We recorded electroretinograms of isolated retina to judge overall retinal physiology under different conditions. We subjected each eye to a different protocol and always recorded last from tissue from the Ames'/Ames' condition to ensure that abnormal responses were not due to prolonged storage. Normalized average responses to flash families recorded using L-15/Locke's (Fig. 8A) were highly variable, whether compared with tissue from the same mouse stored in Ames' (Fig. 8C, matched panels), or to other tissue samples treated with L-15/Locke's. The traces in



Figure 8A are the averages of 4–6 responses; each individual response is of a stereotyped form (Fig. 8A, bottom), indicating that the abnormal form of the ERG does not change during recording. The responses in the middle column (Fig. 8B) come from retinas treated with the various “intermediate” conditions and largely resemble responses in the Ames/Ames’ condition. Comparisons of responses from the L-15/Locke’s and Locke’s/Locke’s conditions (Fig. 8B, blue traces) indicate that L-15 cold storage in particular is responsible for the response variability.

We further assessed the state of retinal tissue by observing tissue morphology in the ganglion cell layer in flat-mounted preparations. Most (>95%) of the ganglion cells in flat-mounted pieces of retina stored in Ames’ appeared healthy (Fig. 8F) and produced robust changes in the rate of action potential generation in response to light (data not shown). After storage in cold L-15 (Fig. 8D) the morphology of the retinal ganglion cells was abnormal: nuclei appeared bloated and distinct (white arrow), and the retina lacked a smooth quality that accompanies typically healthy and responsive tissue (circles). In addition, attempts to record light responses failed because the cell membranes were too fragile to seal onto with a patch electrode. Ganglion cells in tissue stored in warm Locke’s solution (Fig. 8E) lacked the dramatically abnormal appearance of tissue stored in cold L-15, and approximately half of the cells were intact and suitable for recording. The distinct differences in ganglion cell health caused us not to pursue experiments to compare the impact of altered solutions on retinal output.

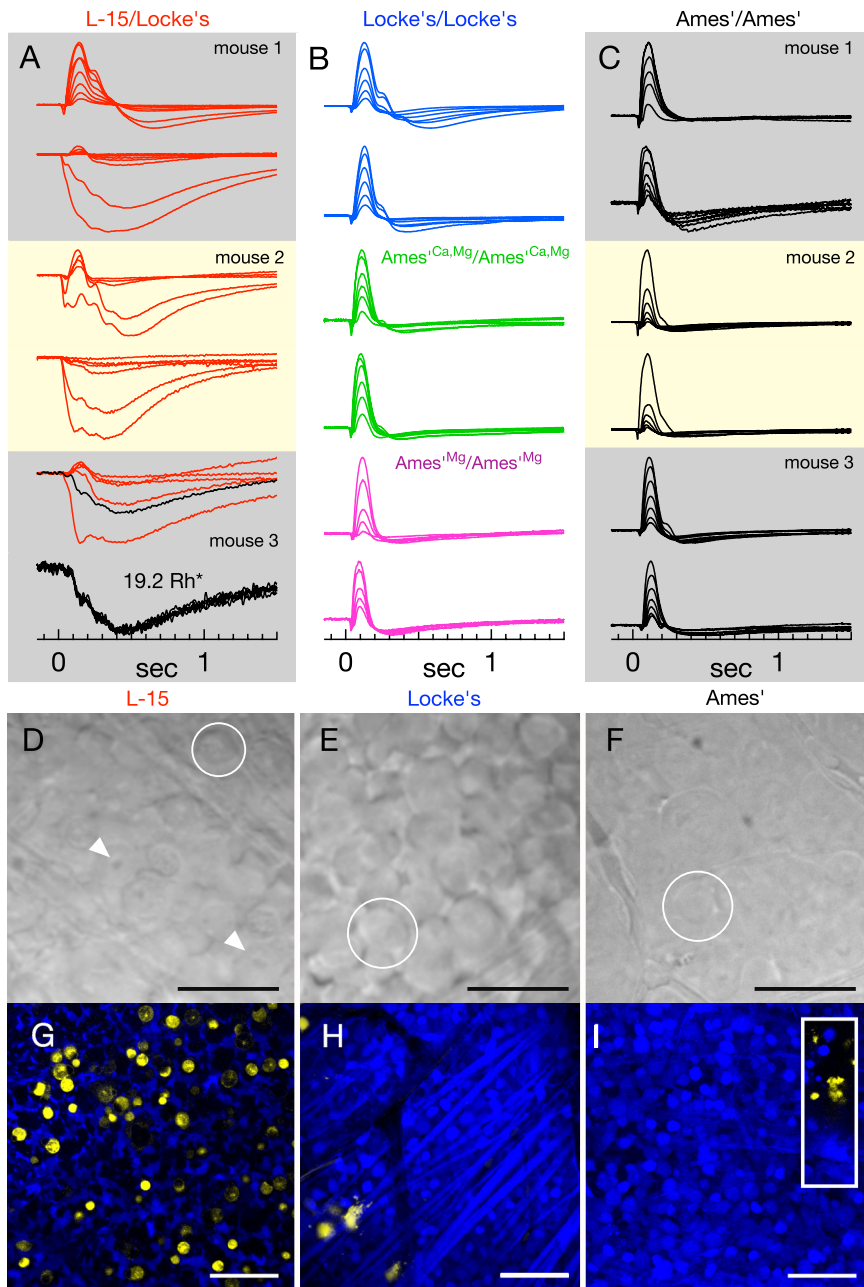
To better assess the state of the tissue in each condition, we used a cell staining assay that discriminates between living and dead cells based on the integrity of the cell membrane (see Materials and Methods). Figure 4G–I shows confocal images of the ganglion cell layer, where live cells are labeled with calcein AM (blue) and dead cells with EthD-1 (yellow). Retina from the Ames/Ames’ condition was imaged after retina stored in L-15, and hence served as a control for the effects of storage duration. Cold storage in L-15 results in more yellow-stained cells than any of the other storage conditions: L-15,  $119 \pm 32$  ( $n = 8$ ); Locke’s,  $5 \pm 1.7$  ( $n = 4$ ); Ames<sup>Mg</sup>,  $2$  ( $n = 2$ ); Ames’,  $2.6 \pm 0.9$  ( $n = 5$ ). L-15 storage produced noticeable, but less dramatic, differences in the bipolar cells in both differential-interference contrast and live/dead cell images (data not shown). Rod bipolar light responses from retina stored in L-15 were highly variable, and thus we did not attempt to compare them with those from retina stored in Ames’ medium. Retinas imaged following warm L-15 storage also had extensive yellow staining ( $130 \pm 40$  cells,  $n = 3$ ). Thus L-15



**Figure 7.** Single-rod and ERG-derived responses to half-saturating flashes exhibit similar kinetics across conditions. **A**, The time-to-peak of the half-saturating response ( $y$ -axis) for single-rod recordings plotted against time-to-peak of the single-photon response ( $x$ -axis). **B**, Population mean normalized half-saturating outer segment current responses in L-15/Locke’s, Ames<sup>Ca,Mg</sup>/Ames<sup>Ca,Mg</sup>, and Ames/Ames’, and the normalized single-photon response in L-15/Locke’s and Ames/Ames’ (red and black dashed lines; see Fig. 4). **C**, ERGs in response to increasing delays between test (27 Rh\*) and probe (2400 Rh\*) flashes. **D**, Isolated ERG response to a flash producing 27 Rh\* in Ames’ and the normalized derived rod response using the paired-pulse technique to the same test flash strength. Smooth curve fit to the data is  $r(t) = \gamma[(1 - \exp(-\alpha(t - t_d)^2))\exp(-t/\tau)]$ ,  $t_d = -7.8$  ms,  $\gamma = 2.4$ ,  $\alpha = 2.3 \cdot 10^{-4}$  ms<sup>-2</sup>,  $\tau_\omega = 121$  ms. **E**, Average derived rod responses in Ames/Ames’ (black,  $t_d = -4.1$  ms,  $\gamma = 1.7$ ,  $\alpha = 3.0 \cdot 10^{-4}$  ms<sup>-2</sup>,  $\tau_\omega = 134$  ms) and L-15/Locke’s (red,  $t_d = 0.56$  ms,  $\gamma = 1.7$ ,  $\alpha = 3.8 \cdot 10^{-4}$  ms<sup>-2</sup>,  $\tau_\omega = 191$  ms). The gray dashed line is the same model using the *in vivo* ERG-derived parameters (Hettinger and Pepperberg, 1999) ( $t_d = 3.1$  ms,  $\gamma = 1.84$ ,  $\alpha = 3.23 \cdot 10^{-4}$  ms<sup>-2</sup>,  $\tau_\omega = 163$  ms). The dashed lines plot the normalized single-photon response in L-15/Locke’s and Ames/Ames’ (red and black dashed lines; see Fig. 4).

storage compromises retinal health and is not ideal for studies of responses downstream of the photoreceptors.

As a test for more subtle variations in retinal processing, we compared complete ERG waveforms across conditions with properties of *in vivo* ERGs. Average flash responses exhibit a form resembling the *in vivo* ERG response over the scotopic range (Saszik et al., 2002), showing what appeared to be a scotopic threshold response and developing both an a-wave and a b-wave as the flash strength increases (Fig. 9A). The normalized b-wave amplitudes (from baseline) show similar sensitivity across conditions, reaching a half-maximal response near 3 Rh\*/rod (Fig. 9B); this is close to the half-maximal flash strength measured directly from rod bipolar cells (Field and Rieke, 2002). Accounting for the optics of the eye (Gross and Burns, 2010; Herrmann et al., 2010), 3 Rh\*/rod corresponds to  $\sim 0.006$  cd s/m<sup>2</sup>, similar to the half-maximal flash strength of 0.003 cd s/m<sup>2</sup> measured *in vivo* (Herrmann et al., 2010). The time-to-peak of the ERG b-wave in Ames’



**Figure 8.** The viability of retinal tissue under different conditions: *in vitro* ERGs, morphology, and Live/Dead staining. **A–C**, Normalized ERGs in response to flash families from individual isolated retinas ( $\sim 0.5$ – $40$   $Rh^*/rod$ ). ERGs measured from half of one retina placed in L-15/Locke's conditions (**A**) (both halves in one box) were followed by ERG recordings from the other retina in Ames'/Ames' (**C**) (labels identify retina from the same mouse). Each trace is the average of  $\sim 5$  responses. The bottom panel depicts individual responses comprising the trace in black in the above flash family ( $19.2$   $Rh^*$ ). **B**, Representative ERG families from tissue in Locke's/Locke's (blue), Ames'<sup>Ca,Mg</sup>/Ames'<sup>Ca,Mg</sup> (green), and Ames'<sup>Mg</sup>/Ames'<sup>Mg</sup> (magenta). ERGs of retina in L-15/Locke's exhibit greater variability than other conditions. **D–F**, Differential interference contrast images of flat-mounted retinal tissue following storage in cold L-15 (**D**), warm ( $32^\circ C$ ) Locke's medium (**E**), and warm ( $32^\circ C$ ) Ames medium (**F**). Scale bars,  $25$   $\mu m$ . Circles indicate ganglion cell somas and arrows in **D** illustrate prominent, pockmarked nuclei. **G–I**, Confocal images through the ganglion cell layer of live (blue) and dead (yellow) tissue following storage in cold L-15 (**G**), warm Locke's (**H**), or Ames' (**I**) (inset, yellow stained cells at the edge of the retina piece). Scale bars,  $50$   $\mu m$ . Counts of yellow nuclei (mean  $\pm$  SEM) in images of retina stored in cold L-15,  $119 \pm 32$  ( $n = 8$ ); warm Locke's,  $5 \pm 1.7$  ( $n = 4$ ); Ames' with elevated  $Mg^{2+}$ ,  $2$  ( $n = 2$ ); Ames',  $2.6 \pm 0.9$  ( $n = 5$ ). An image of healthy retina on this scale contained  $286 \pm 16$  ganglion cells.

medium (Fig. 9C, black curve) ranged from 102 to 114 ms over the light intensities in Figure 8, comparable to *in vivo* measurements of 110 ms over the scotopic range (Saszik et al., 2002); the time-to-peak was insensitive to changes in divalent concentration. The b-wave measured in Locke's medium was significantly

slower across flash intensities ( $p < 10^{-11}$ , MWW test, all b-wave time-to-peak values in Ames'/Ames' vs Locke's/Locke's). This may be a result of altered synaptic processing due to the presence of HEPES in Locke's solution (Hirasawa and Kaneko, 2003; Davenport et al., 2008; Fahrenfort et al., 2009).

## Discussion

The substantial differences between measured properties of single-photon responses across studies hinder progress in two broad areas. First, rod phototransduction is the best-studied example of how G-protein signaling cascades amplify activity of single active receptors (for review, see Pugh and Lamb, 1993; Rieke and Baylor, 1998; Burns and Arshavsky, 2005); substantial differences in single-photon responses indicate a previously unappreciated and surprising sensitivity of this process to experimental conditions. Second, comparison of signal and noise properties of rod responses with those inferred from behavioral measurements has shaped investigation of the retinal circuits dedicated to low light level vision; this work has revealed several signal processing mechanisms that appear near-ideally matched to properties of the single-photon response (Bialek and Owen, 1990; Field and Rieke, 2002; Armstrong-Gold and Rieke, 2003; Berntson et al., 2004; Sampath and Rieke, 2004). These mechanisms, as well as other aspects of our understanding of retinal processing, will need to be reevaluated if they are based on nonphysiological photoreceptor responses. Below we discuss our results in the context of these issues.

### Cold L-15 medium accelerates the single-rod responses and adversely affects the health of retinal tissue

Storing the retina in L-15, either warm or cold, had dramatic effects on both the responses of single-rod photoreceptors and the general health of the retina. Responses of rods recorded in Locke's medium retained a dependence on how the retina was stored: cold L-15 storage resulted in smaller and faster single-photon responses than other storage conditions, even when the same conditions were used during actual recordings. Assuming the transduction cascade operates linearly at low light levels, its kinetics are controlled by two factors—the rate of decay of light-activated PDE activity and the dark PDE activity level (Rieke and Baylor, 1996). The dominant time constant for bright flashes, however, had at most a weak dependence on how the retina was stored. Assuming the same process dominates the rate of decay of

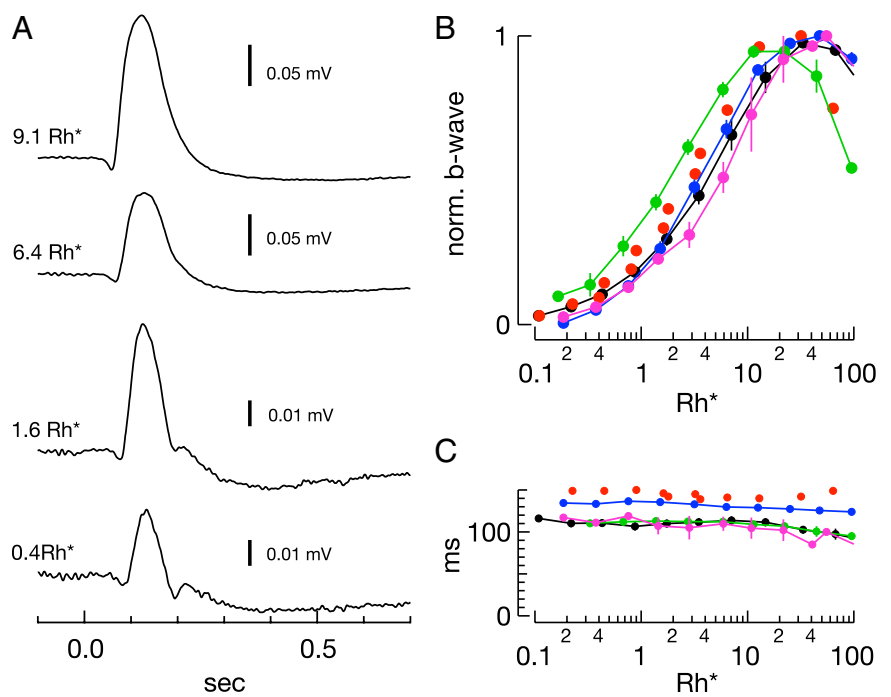
PDE activity for bright and dim flashes, the weak dependence of the dominant time constant on storage conditions makes changes in the decay of PDE activity unlikely to explain the effect of L-15 on the single-photon response. Instead our results are most consistent with an increased rate of dark PDE activity and corresponding speeding of the response of the transduction cascade to perturbations in cGMP concentration; simple transduction cascade models would require a  $\sim 3$ -fold increase in dark PDE activity to account for the difference in kinetics between the responses from retina stored in Locke's vs cold L-15 (Nikonov et al., 2000).

While the rods appeared healthy and generated robust light responses in retina stored in L-15, the remainder of the retina showed clear morphological and physiological abnormalities. In particular, L-15 storage precluded comparing rod responses with those of downstream neurons. This effect of L-15 storage is consistent with observations by Ames and Nesbett (Ames and Nesbett, 1981) during their original formulation of Ames' medium: they noted that typical culture solutions did not preserve the morphology, metabolism, or light responsiveness of rabbit retina. Winkler et al. (1977) similarly found that the ERG b-wave disappeared following storage in solutions with phosphate or Tris pH buffering. Together, these studies suggest that bicarbonate, which is absent in L-15, is critical for maintaining normal retinal morphology and function.

### The ERG a-wave is insensitive to alterations in rod phototransduction

The electroretinogram is often used to infer *in vivo* properties of the rod responses (Lyubarsky and Pugh, 1996; Pepperberg et al., 1997; Hetling and Pepperberg, 1999; Friedburg et al., 2001). Our results raise two cautions about the generality of such an approach. First, manipulations that produced clear changes in measured rod responses had surprisingly little effect on the responses derived from the ERG. In particular, dominant time constants from ERG-derived responses were insensitive to alterations in the  $Mg^{2+}$  concentration, unlike the dominant time constants measured from single rods. Previous *in vivo* and *in vitro* ERG measurements also suggest that the ERG a-wave may not be directly proportional to the rod photocurrent; specifically, voltage changes measured along the longitudinal axis of the outer segments in response to bright flashes show a transient component that is absent in the photocurrent and likely originates in the inner retina (Green and Kapousta-Bruneau, 1999). A similar effect may reduce the sensitivity of the ERG to alterations in the rod photocurrents. Alternatively, the suction procedure itself may alter rod photocurrents in such a way as to eliminate important components comprising the ERG, though we did not see a corresponding transient component in voltage-clamp recordings from the rod soma (data not shown; see also Baylor and Nunn, 1986).

Second, paired-flash protocols used to derive rod responses to subsaturating lights rely on the first flash producing a substantial



**Figure 9.** Properties of the *in vitro* scotopic ERG. **A**, Example Ames'/Ames' ERG traces at various flash strengths. **B**, Sensitivity of the b-wave amplitude across conditions for samples shown in Figure 8. Average half-saturating strengths in  $Rh^*/rod$ : Locke's/Locke's (blue), 3.2  $Rh^*/rod$ ; Ames'/Ames' (black), 4.0  $Rh^*/rod$ ; Ames'<sup>Ca,Mg</sup>/Ames'<sup>Ca,Mg</sup> (green), 1.8  $Rh^*/rod$ ; Ames'<sup>Mg</sup>/Ames'<sup>Mg</sup> (magenta), 5.7  $Rh^*/rod$ . Red dots in **B** and **C** represent the b-wave amplitudes of the one sample recorded in L-15/Locke's that exhibited a b-wave (Fig. 8*H*, top family). Half-saturating flash strength for this sample was 2.8  $Rh^*/rod$ . **C**, Average time-to-peak of the b-wave across conditions and flash strength for samples shown in Figure 8.

suppression of current. Typically such measurements are made for initial flashes producing close to a half-maximal response (Pepperberg et al., 1997; Hetling and Pepperberg, 1999; Friedburg et al., 2001). When measured in single cells, however, the kinetics of half-maximal responses differ considerably from single-photon responses (Baylor et al., 1979a). Consequently, derived subsaturating flash responses from the ERG are not comparable to the rod's single-photon response. Furthermore, half-maximal responses of single rods showed little dependence on recording conditions; the same was true for derived subsaturating responses from the ERG. The striking difference in single-photon responses and minimal difference in half-saturating responses across conditions could be at least partially explained if calcium feedback to the rate of cGMP synthesis operates strongly and nonlinearly for brighter flashes (Burns et al., 2002), leading to a shift in the time constants controlling changes in cGMP.

We hoped to use the rod responses derived from the ERG to identify *in vitro* conditions producing rod responses that replicated those *in vivo*. This did not prove possible because of the insensitivity of the *in vitro* ERGs to manipulations that produced substantial changes in single rod responses. In particular, the similarity of ERG-derived rod responses prevented us from identifying  $Mg^{2+}$  and  $Ca^{2+}$  concentrations required to replicate *in vivo* conditions.

### Increases in divalent ions accelerate recovery kinetics in responses to bright flashes

Under some conditions, the dim flash response recovers with a time course similar to the dominant time constant for the decay of light-activated PDE activity. However, several manipulations cause the time constant of the recovery of the dim flash response to differ from the dominant time constant. For example, under-



expression of rhodopsin kinase 1 in recording solutions with high divalent concentrations produces a clear slowing of the dim flash response with little change in the dominant time constant (Chen et al., 2010). Similarly, lack of recoverin speeds the dim flash response with a less dramatic change in dominant time constant (Makino et al., 2004). Differences in responses at low and high  $Mg^{2+}$  concentrations provide a dramatic example: dim flash responses recovered  $\sim 10$ – $15\%$  more quickly at high  $Mg^{2+}$  while the dominant time constant was nearly twofold smaller (Fig. 5).

Changing the  $Mg^{2+}$  concentration also had unexpected effects on the response area (i.e., integrated response). If the measured transduction currents are linearly proportional to the light-activated PDE activity (Burns et al., 2002), then a twofold change in the dominant time constant should produce a twofold change in response area. Whether this change in area is produced by a change in response amplitude, kinetics or both will depend on the cGMP turnover rate and how closely the cGMP concentration tracks the PDE activity (Nikonov et al., 2000). Single-photon response areas, however, showed a much weaker dependence on  $Mg^{2+}$  concentration than the dominant time constant. This indicates either that (1) the gain and/or time course of the PDE activity estimated from bright flashes fails to apply to dim flashes, or (2) changing  $Mg^{2+}$  has multiple effects, not all of which are captured by the dominant time constant. The rate of cGMP synthesis is known to depend on magnesium (Peshenko et al., 2008), but a simple scaling of the synthesis rate would equally affect the time in saturation at all flash strengths (Burns and Pugh, 2009), leaving the estimate of the dominant time constant unchanged. Instead, the sensitivity of  $\tau_d$  to  $Mg^{2+}$  is unexpected, and it will be interesting to understand this effect in more mechanistic detail.

### Importance for visual sensitivity

The behavioral performance of the visual system of many vertebrates approaches limits set by the division of light into discrete photons and noise in the responses of the rod photoreceptors (Hecht et al., 1942; Sakitt, 1972; Aho et al., 1988; Naarendorp et al., 2010). This observation has motivated considerable theoretical and experimental work to understand retinal processing at visual threshold. For example, absolute visual sensitivity places at least two requirements on the transmission of signals from rods to rod bipolar cells. First, the reduction in transmitter release produced during the hyperpolarizing single-photon response must be distinguishable from synaptic noise (Rao-Mirotnik et al., 1998; Schein and Ahmad, 2005, 2006). Fluctuations in vesicle release in the dark will produce random pauses or slowing in release that could be confounded with true single-photon responses. The effect of such synaptic noise on visual sensitivity will depend strongly on how much and for how long the single-photon response suppresses the release rate, which in turn depends on the amplitude and kinetics of the single-photon response. Second, transmission from rods to rod bipolar cells must separate single-photon responses in a minority of the converging rods inputs from continuous noise generated in the majority (Baylor et al., 1984; van Rossum and Smith, 1998). A thresholding nonlinearity at the rod-to-rod bipolar synapse appears to serve such a role (Field and Rieke, 2002; Sampath and Rieke, 2004). The effectiveness of such a mechanism has a strong dependence on how well it is matched to the signal and noise properties of the rod responses (Okawa et al., 2010). Thus the effect of experimental conditions on the single-photon response amplitude and kinetics is important for our understanding of the operation both of phototransduction and of downstream processes critical in preserving the fidelity of rod vision.

### Notes

Additional material for this article is available at <http://rieke-server.physiol.washington.edu/Publications/publications.htm>. This material has not been peer reviewed.

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