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Bone Mineral Imaged *In Vivo* by ^{31}P Solid State MRI of Human Wrists

Yaotang Wu, PhD^{1,2,3,*}, Timothy G. Reese, PhD^{2,3}, Haihui Cao, PhD^{1,2,3}, Mirko I. Hrovat, PhD⁴, Steven P. Toddes, MS⁵, Rostislav A. Lemdiasov, PhD⁵, and Jerome L. Ackerman, PhD^{1,2,3}

¹ Department of Orthopedic Surgery, Children's Hospital, Boston, MA 02115

² Athinoula A. Martinos Center for Biomedical Imaging Department of Radiology Massachusetts General Hospital, Charlestown, MA 02129

³ Harvard Medical School, Boston, MA 02115

⁴ Mirtech, Inc., Brockton, MA 02301

⁵ InsightMRI, Leominster, MA 01453

Abstract

Purpose—To implement solid state ^{31}P MRI (^{31}P SMRI) in a clinical scanner to visualize bone mineral.

Materials and Methods—Wrists of seven healthy volunteers were scanned. A quadrature wrist ^{31}P transmit/receive coil provided strong B_1 and good signal-to-noise ratio (SNR). A ^1H - ^{31}P frequency converter was constructed to enable detection of the ^{31}P signal via the ^1H channel. Data points lost in the receiver dead time were recovered by a second acquisition with longer dwell time and lower gradient strength.

Results—Three dimensional ^{31}P images, showing only bone mineral of the wrist, were obtained with a clinical 3T scanner. In the best overall case an image with isotropic resolution of ~ 5.1 mm and SNR of 30 was obtained in 37 min. ^{31}P NMR properties (resonance line width 2 kHz and T_1 17–19 s) of *in vivo* human bone mineral were measured.

Conclusion—*In vivo* ^{31}P SMRI visualization of human wrist bone mineral with a clinical MR scanner is feasible with suitable modifications to circumvent the scanners' limitations in reception of short- T_2 signals. Frequency conversion methodology is useful for implementing ^{31}P SMRI measurements on scanners which do not have multinuclear capability or for which the multinuclear receiver dead time is excessive.

Keywords

bone; solid state MRI; ^{31}P SMRI; UTE; osteoporosis; WASPI

INTRODUCTION

Osteoporosis affects roughly 10 million Americans over the age of 50 (1). In osteoporosis, resorbed bone is replaced by a lesser volume of normal bone, causing the loss of trabecular connectivity and an increase in cortical porosity. This results in serious long-term disability

Address of Corresponding Author: Yaotang Wu, Department of Orthopedic Surgery, Children's Hospital, 300 Longwood Avenue, Boston, MA 02115, Phone: 617-919-2060, Fax: 617-730-0122, yaotang.wu@childrens.harvard.edu.

and death for a significant number of patients. Measurement of areal bone mineral density (BMD) and the associated T score by Dual Energy X-ray Absorptiometry (DXA) is presently the standard method for diagnosing osteoporosis. However, the areal BMD measured by DXA is dependent on bone size and does not distinguish between cortical and trabecular bone. Even though Quantitative Computed Tomography (QCT) and peripheral QCT (pQCT) enable the volumetric scans of skeletal sites and provide separate information on cortical and trabecular bone, the significantly larger dose of ionizing radiation associated with QCT and its higher cost compared to DXA restricts its use in routine clinical screening (2, 3); pQCT is relatively new and clinical experience with it is limited.

Magnetic resonance imaging can provide true three-dimensional volumetric measurement while completely avoiding exposure of the subject to ionizing radiation. However, the very short T_2 and broad resonance of MR signals from ^{31}P nuclei in bone mineral and protons in solid bone matrix make the MRI measurement difficult. Recently, three MRI methods that address these difficulties have been developed, namely, Ultra short TE (UTE) imaging (4–6); Solid State MRI (^{31}P SMRI for its ^{31}P version (7,8) and water and fat suppressed projection imaging, WASPI, for its ^1H version (8–10)); and Sweep Imaging with Fourier Transformation (SWIFT) (11, 12). The 3D versions of these methods employ radial acquisition of free induction decays (FIDs). Two dimensional UTE uses slice selective half RF pulses and 2D radial acquisition of FIDs. Both 3D UTE and SMRI sequences employ non-selective RF excitation. The major difference between 3D UTE and SMRI is the timing of the RF excitation pulse relative to the projection gradient pulses and the sampling of the FID under either a rising gradient (UTE) or a constant gradient (SMRI). In UTE, the RF excitation pulse is applied before the projection gradients, while in SMRI, the projection gradient is applied and stabilized before an RF excitation pulse. These differences make SMRI more suitable for imaging short T_2 signals (13). In the SWIFT sequence, the excitation pulse is divided into N sub-pulses with amplitude and frequency modulations equivalent to an adiabatic sweep through resonance, making possible the excitation of broad resonances with limited RF power. The data acquisition is performed immediately after each sub-pulse under a constant gradient, which is turned on before the first sub-pulse and kept constant until the end of the last data acquisition.

It has previously been shown in ex vivo bone specimens that solid state ^{31}P MRI is as good as or better than DXA in measuring bone mineral density by validation against chemical analysis (7). That study utilized a Spectroscopy Imaging Systems (Varian/SISCO, Palo Alto, CA, USA) spectrometer, which includes solid state MR imaging and spectroscopy among its repertoire of measurements. A subsequent study demonstrated that it was possible, using a Bruker (Karlsruhe, Germany) electronics console interfaced to a clinical magnet, to obtain solid state ^{31}P MR images of human volunteers (14). As a first step toward translating solid state MRI into a practical clinical tool, in the present study we implemented ^{31}P SMRI in a Siemens 3T Trio scanner. The special considerations involved in acquiring very short T_2 signals on a clinical MR scanner are discussed.

MATERIALS AND METHODS

Hydroxyapatite Phantoms, Bone Samples, and Study Subjects

All animal and human images were obtained in full compliance with federal regulations and the guidelines of the respective clinical research committees of the relevant institutions.

Hydroxyapatite powder (Aldrich Chemical Company, Inc, Milwaukee, WI, USA) was densely packed into plastic tubes (inside diameter 14 mm, length 40 mm) and used in phantoms for assessing SNR and spatial resolution. To analyze the image resolution, two such tubes were bound together with a plastic spacer between them to create a resolution

phantom. The space separating the two HA columns (taking into account the tube wall thickness and the spacer) was 2.9 mm.

Whole pig legs, disarticulated at the knee, were obtained from ~3 month old Yorkshire pigs (Animal Research, Children's Hospital, Boston, MA, USA) following euthanasia of the animal. The legs were intact and kept frozen until scanned by MRI to test imaging protocols.

Seven healthy volunteers (5 males and 2 females), ages 28–60 years, participated in this study. During MRI scanning, the volunteer lay still in the scanner, with the right or left wrist (their choice) positioned at about 15° to B_0 in the horizontal plane (for comfort) inside a ^{31}P MRI transmit/receive quadrature birdcage coil from InsightMRI (Leominster, MA, USA). The coil package includes a fixturing system for comfortably constraining the wrist (Fig. 1) and an enclosure containing a quadrature hybrid transmit/receive switch and a low noise narrowband preamplifier. The coil has an elliptical shape to better conform to the shape of the human wrist. The shape and relatively tight fit permits a high filling factor to maximize the signal-to-noise ratio. The coil was designed to be positioned near the isocenter of the magnet (to simplify the radial reconstruction), without compromising patient comfort (and causing involuntary movements due to pressure points). To securely hold the patient's arm, a unique fixturing system was developed to secure the hand and forearm, thereby indirectly, though securely, fixing the position of the wrist. ^{31}P spectra, T_1 measurements, and images of bone mineral were acquired with three pulse sequences: single pulse, progressive saturation, and SMRI (acquisition of FIDs and 3D radial coverage of k-space), respectively.

The detailed SMRI method is described in (7). Briefly, this method utilizes a projection MRI pulse sequence consisting of a single, fixed-amplitude magnetic field gradient pulse, during which an excitation RF pulse is issued after a suitable delay following the start of the gradient pulse to allow eddy currents to decay. The free induction decay (FID) is acquired as soon as possible after the RF pulse. This sequence is repeated for a set of gradient directions distributed in a uniform pattern about the unit sphere.

In the following sections, the implementation of the ^{31}P SMRI in a Siemens 3T Trio clinical scanner is described.

MR Signal Excitation

A strong B_1 field is required to excite the broad resonance of the solid MR signals, which the body coil of clinical scanners cannot provide. A ^{31}P quadrature low-pass birdcage coil was constructed for human wrist imaging, which provides hard pulses as short as $10\ \mu\text{s}$ for a 10° flip angle and high filling factor for better signal to noise ratio (Fig. 1).

MR Signal Acquisition

In quantitative solid state MR imaging or spectroscopy measurements based on FIDs, it is critical to reduce the receiver dead time to as short a time as possible to minimize distortion or substantial loss of the initial FID samples. The SMRI pulse sequence therefore uses a single very short, intense, rectangular RF excitation pulse with no gradient switching of any kind between the RF pulse and signal acquisition. Sampling of the signal is initiated as soon as possible following the end of the RF pulse. To minimize the delay between the end of the RF pulse and the start of data sampling, the scanner's lower bound for this delay was reduced from the default value of $20\ \mu\text{s}$ to $10\ \mu\text{s}$ with the assistance of the scanner manufacturer's engineers. Because no slice selection is possible, a 3D data set must be acquired. To minimize the electronic delay for the switching of the transmit/receive (T/R) switch, the standard PIN diode T/R circuit was eliminated. Transmitter-receiver signal isolation was achieved with the use of a quadrature hybrid connecting the transmitter, receiver and the two ports of the quadrature wrist coil. Additional transmitter isolation was

achieved by inserting several stacks of back-to-back 1N914B switching diodes in series with the transmitter output. Just enough diodes were used to completely eliminate any receiver noise. Back-to-back diodes to ground were placed in front of the preamplifier to complete the T/R switch. These diodes are switched passively by the transmit pulse and need no external drive signal; the switching speed is on the order of nanoseconds.

Frequency Converter

Even with the reduction in the software receiver dead time and elimination of the PIN diode T/R switch, it was initially found that the intrinsic receiver recovery time of the ^{31}P channel of the scanner is over 200 μs , far too long for SMRI (Fig. 2a). This is intrinsic to the design of the multinuclear receiver channel and/or the digital signal processing (it is similar for other heteronuclei as well). A frequency converter was therefore developed to excite and observe the ^{31}P signal through the ^1H channel, taking advantage of the very short receiver recovery time of the ^1H channel. Details of this design are presented in Fig. 3. The ^1H transmit pulse from the scanner is directed into a Bird Electronic (Solon, OH, USA) 50 ohm high power termination, sampled with a Bird Electronic RF directional coupler (with an RF rather than DC output). The sampled ^1H transmitter pulse is mixed with a local oscillator (LO) frequency to translate the pulse to the ^{31}P frequency with all amplitude and phase characteristics preserved. The ^{31}P pulse is amplified by an Electronic Navigation Industries (Rochester, NY, USA) LPI-10 RF power amplifier (1 kW maximum output) and sent to the coil for ^{31}P MR excitation. The LO frequency is provided by a Programmed Test Sources (Littleton, MA, USA) PTS-160 frequency synthesizer phased locked to the 10 MHz reference clock of the scanner. The external RF power amplifier is gated (blanked) by the pulse sequence to suppress the noise from the RF power amplifier final stage. The unblanking pulse is timed to allow enough time prior to the RF pulse for the RF power amplifier to reach full gain, and following the RF pulse to squelch amplifier noise in time for the switching transients to die away before data is acquired. The timings must take into account the small but significant differential delays between the RF and digital logic channels of the scanner and frequency converter circuitry. The ^{31}P MR signal is amplified, mixed with the local oscillator frequency to translate it to the ^1H frequency, again with all amplitude and phase characteristics preserved. The signal is then applied to the scanner ^1H receiver. Although this bypasses the standard SAR monitoring of the scanner, careful testing of the RF coil prior to IRB approval of the study demonstrated that the SAR is very small (the ^{31}P frequency is well below the ^1H frequency and the pulse sequence has a low mean RF power level).

After the development of the frequency converter, and with additional consultation with the scanner manufacturer's engineers, it was discovered that the excessive receiver delay is caused by electronic switching between RF sidebands in the internal frequency converter of the scanner, and that following the initial multinuclear acquisitions during the execution of the pulse sequence, subsequent acquisitions occur without excessive delay. The problem only occurs for certain nuclei, including ^{31}P . This suggests that this problem can be circumvented in the future with modifications in the pulse sequence and processing algorithms to discard initial acquisitions and retain only those which are acquired with short receiver delay. However, we describe the frequency conversion methodology because it enables the measurements for scanners without multinuclear capability or for which the RF hardware limitations cannot otherwise be circumvented.

The ^{31}P preamplifier is an Advanced Receiver Research (Burlington, CT, USA) low noise (0.5 dB noise figure) custom-tuned low-magnetic communications preamplifier modified with 1N914B crossed (back-to-back) passively switched (by RF level) silicon diodes to ground on the input for preamp protection, and preamp bias power provided from the preamp RF output connection as is standard for clinical scanner preamps. A stack of

multiple crossed diodes in series was connected at the RF power amplifier output to eliminate any residual amplifier noise detected in the receiver signal. To reduce the number of cables required, 12 volts DC is imposed on the center conductor of the cable connecting the preamplifier output to the frequency converter to provide preamplifier bias. The various amplifiers, filters, attenuators and signal routing in the frequency converter ensure that maximal dynamic range and sufficient unwanted sideband suppression are achieved, while keeping low level RF signals within safe limits for the scanner and RF power amplifier.

The actively switched PIN diodes normally employed in the transmit/receive (T/R) switch in clinical scanners are relatively slow in switching speed, due to the inherent properties of the PIN junction as well as to the switching speed of the PIN diode drive signal. The above scheme eliminates all PIN diode switching along with the increased dead time contributed by this type of T/R switch. Typical conventional switching diodes like the 1N914B switch in ns. Transmit/receive isolation is determined in the modified design by the isolation of the quadrature hybrid (roughly 20 dB) and the passive switching diodes in the transmitter and receiver circuits.

The success of this design was demonstrated by the following experiment. A CW (continuous wave) sine wave signal from a frequency synthesizer at a slight offset to the nominal ^{31}P frequency was applied through an attenuator (ranging from 60–100 dB) to the quadrature hybrid where the RF coil would be connected. The scanner (^1H) RF pulse amplitude was set to zero. The output signal from the converter was sent to the ^1H receiver of the scanner and recorded by the scanner. There was no observable distortion of the received synthetic signal at the beginning of the recorded trace (Fig. 2b), and a ^{31}P signal linear dynamic range (thermal noise floor to 1 dB compression) of greater than 10^6 was achieved.

The RF ring-down time of the coil also affects the receiver dead time. When the T/R coil is well tuned and matched, the ring-down time is longer, which increases the total time the receiver is blocked (Fig. 4). This problem is more significant in ^{31}P SMRI than in proton WASPI, since the ^{31}P frequency is 2.5 times lower. One interpretation of coil quality factor Q is that it is numerically equal to the number of RF cycles in the exponential time constant L/R for the ring down ($Q = 2\pi f L/R$). For similar ^1H and ^{31}P coil Q s, the ^{31}P coil will take 2.5 times longer to ring down. At typical values of Q , $L/R \sim 1/6 \mu\text{s}$. In this study, we used $20 \mu\text{s}$ for the total ^{31}P receiver dead time, including the RF ring-down time which occurs roughly simultaneously with the receiver recovery. The software delay between the end of the RF pulse and the start of data acquisition was therefore set to $20 \mu\text{s}$.

Image Reconstruction with Missing Data Points recovered by a Second Acquisition

To observe the short T_2 signal, fast acquisition (dwell time as short as $5 \mu\text{s}$) is required to provide a broad acquisition spectral width. Even with the shortest achievable receiver dead time (e.g., $10 \mu\text{s}$ for ^1H and $20 \mu\text{s}$ for ^{31}P), the first few FID samples corresponding to the positions in the k-space close to the origin are still lost. Reconstructing an image without any correction for this loss means that the time domain data is shifted in k-space, in a manner not equivalent to simply translating the time origin (which would be correctable in post-processing), but rather as data points that are moved radially toward the origin, resulting in the Fourier transform being multiplied by position-dependent phase factors. The intensity of the resulting image, which is a magnitude mapping of the Fourier transform of the time domain data, would hence be seriously distorted by the various phase shift factors at different space positions (15).

The approach to correcting this problem by placing sampled data in the correct k-space positions and assigning zeros to the k-space positions of missing data is equivalent to taking

a convolution of the true image with a radial sinc function, which would depress the baseline of the image, resulting in a dark halo surrounding the image, and would distort the signal intensity.

Instead, our approach is to recover the missing data directly by means of a single additional acquired data set, in which the dwell time t_{dw} is increased by a factor C , so that the new dwell time Ct_{dw} is equal to the receiver dead time, making all data detectable except for the very first one (k-space origin). The sample exactly at the origin of k-space does not contribute to the reconstruction, so it does not matter if it is lost. The gradient strength is reduced in the second acquisition by that same factor C , making the second set of data points correspond precisely to the missing k-space points near the origin. The second acquisition requires many fewer projections than the first (regular) set because the previously missing data does not extend far into k-space when the minimal receiver dead time is applied. In the reconstruction, the second data set is re-sampled onto the same Cartesian coordinate system of the first data set. The second acquisition adds a trivial amount of time to the overall scan duration. Because the additional k-space points are measured in the second acquisition farther down the material's intrinsic T_2 decay curve than they would have been in the main acquisition, these points carry a somewhat greater T_2 weighting; however this does not seriously impact the overall measurement accuracy or introduce observable artifacts (15).

Imaging Resolution

The resolution-limiting factors of SMRI are the projection pixel size Δx_p and intrinsic pixel size Δx_i .

Δx_p is determined by the imaging parameters (15):

$$\Delta x_p = \text{FOV} \sqrt{\frac{\pi}{P}} \quad [1]$$

where FOV is the field of view and P the number of projections. FOV is determined by the acquisition spectral width W_s , the projection gradient strength G , the reconstruction matrix size M_{mat} and the actual number of data points used in the reconstruction N_{ps} :

$$\text{FOV} = (M_{\text{mat}}/N_{\text{ps}})W_s / (2\gamma G) \quad [2]$$

Δx_i is determined by the inherent MR properties of the imaged nuclei, mainly the resonance line width W_L , and the applied projection gradient strength G :

$$\Delta x_i = (2\pi W_L) / \gamma G \quad [3]$$

The overall resolution can be expressed as Δx_{all} :

$$\Delta x_{\text{all}} = \sqrt{\Delta x_p^2 + \Delta x_i^2} \quad [4]$$

The point spread function (PSF) is a useful concept to analyze resolution limitations (16):

$$\widehat{I}(x) = I(x) * h(x) \quad [5]$$

where $I(x)$ is the object, $\widehat{I}(x)$ is its image and $h(x)$ the PSF. The symbol denotes * convolution. In this study, a one dimensional profile was generated from the transverse

image of the two-tube phantom acquired by ^{31}P SMRI and deconvolved to obtain the corresponding $h(x)$ using the blind deconvolution method with MATLAB (Mathworks Inc, Natick, MA, USA). The width of $h(x)$ was then compared with Δx_{all} .

T1 Measurements

The T1 of bone mineral has been known to be long (14). To make the T1 measurement in a single scanning session duration tolerable to the subject, the measurement was carried out either by spectroscopy or SMRI with sharply reduced number of projections. Since inversion-recovery is not practical in measuring T1 of MR signals with very short T2, progressive saturation was used. TRs were 0.5, 1, 2, 3, 5 s in both measurements. For spectroscopy measurement, one hundred dummy scans were applied before each data acquisition to ensure MR steady states. Twenty acquisitions at each TR were averaged. For SMRI measurements, the number of projections was 184 with $G = 28$ mT/m, and dwell time = 7 μs , without acquisition average.

The ^{31}P SMRI experiments were carried out with a Siemens 3T Trio whole body scanner, with a maximum available gradient strength of 30 mT/m.

RESULTS

The ^{31}P resonance line width of hydroxyapatite was 2.0 kHz (Fig. 5a), and the T1 was 1.5 s. Based on these results, images of ^{31}P SMRI of the two-tube phantom (Figs. 5b and c) were obtained with the following protocol: FOV 150 mm; excitation pulse 10 μs for 10° flip angle; dwell time 7 μs ; gradient strength 28 mT/m; projection number 8148; receiver dead time 20 μs ; TR 50 ms. The second data acquisition to recover k-space points missed in the receiver dead time of the regular acquisition was carried out with lower gradient strength (7 mT/m), longer dwell time (20 μs), and fewer projections (20) while the other parameters were the same as those in the regular acquisition. The total scan time was about 7 min for both acquisitions combined.

Images of the two-tube phantom were reconstructed in two different ways: a) without any correction of the time domain data; b) with correction of the time domain data position in k-space and with the second data set to compensate for the missing data points. The results in Fig. 6 demonstrate good correction of the artifact due to missing data points.

The line profile of the image of the two-tube phantom in Fig 6b was generated across the center of the transverse view (Fig. 7a), which showed that the intensity of the gap between the two tubes dipped, but did not reach the background level. A deconvolution operation of the line profile was performed (Fig. 7b) and the PSF was obtained and plotted in Fig 7c. The full width at half maximum of the PSF was measured as 5.0 mm. According to Eqs. 1 and 3, for FOV = 150 mm and $P = 8148$, projection pixel size Δx_p was 3.0 mm; and for $W_L = 2$ kHz and $G = 28$ mT/m, intrinsic pixel size Δx_i was 4.1 mm. The Δx_{all} was 5.1 mm according to Eq. 5, in very close agreement to the measured 5.0 mm value.

In vivo ^{31}P MR study of seven volunteers is summarized in Table 1.

The ^{31}P resonance line width and T1 of bone mineral of two subjects was measured *in vivo* as ~2.0 kHz and ~17.2 s (by spectroscopy) – 19.2 s (by SMRI) at 3T, respectively (Fig. 8, Tab. 1). The bone mineral line width is similar to that of HA, but the T1 is much longer. Three dimensional ^{31}P images of the bone mineral of human wrists were obtained with the same protocol used in the imaging of the two-tube phantom, except that the TR was 275 ms due to the longer T1 of human bone mineral, which resulted in a longer acquisition time (37 min). The SNR is about 30 (Fig. 9).

Since the intrinsic pixel size is the dominating factor over the projection pixel size for the overall resolution, ^{31}P images with the available maximum gradient strength 30.11 mT/M were performed while keeping the FOV and the number of projections the same. The resulted wrist images were with higher resolution. According to Eq. 3, the Δx_i was reduced from 4.1 mm ($G = 28$ mT/M) to 3.7 mT/M. The overall resolution Δx_A was improved from 5.1 mm to 4.8 mm. (Fig. 10). However, the SNRs are reduced to 16–18 (Table 1).

DISCUSSION

The results demonstrate that despite the low intrinsic SNR, long T_1 and broad spectral linewidth of the ^{31}P isotope in bone, as well as the technical barriers to measuring short T_2 signals with clinical scanners, ^{31}P MRI of bone mineral in humans is feasible. Synthetic hydroxyapatite powder, which resembles bone mineral chemically, is a useful phantom material for setting up scans and for measuring the spatial resolution in the image. The resonance line width of hydroxyapatite and human bone mineral are similar (about 2 kHz FWHH), corresponding to a T_2 of 160 μs . The achievable SNR of the images of the HA tubes and human wrists demonstrated that the difficulty caused by the short T_2 of the ^{31}P signal from HA and the bone mineral is largely overcome by the utilization of a strong B_1 provided by the wrist coil and a shortened receiver dead time by using the proton channel of the scanner with a frequency converter.

The intrinsic pixel size for ^{31}P was 4.1 mm. The gyromagnetic ratio γ_P of ^{31}P is 2.5 times smaller than that of the proton, γ_H , which means that a $2.5\times$ stronger gradient strength is required for ^{31}P to achieve the same intrinsic pixel size as that of ^1H images. In this study, due to the limitation of the maximum gradient strength of 30 mT/m, the intrinsic pixel size (4.1 mm) is the dominating factor over the projection pixel size (3.0 mm) for the overall resolution ($\Delta x_{\text{all}} \sim 5.1$ mm). The measured PSF (FWHM ~ 5.0 mm) confirmed this. If the applied gradient strength could be higher, such as 40 mT/m (available in many 3 T clinical scanners), the intrinsic pixel size will be 2.9 mm. A stronger gradient strength would reduce the FOV from the current 150 mm to 105 mm (at 3.0 T), hence reducing the Δx_p from 3 mm to 2 mm and the overall resolution Δx_{all} from 5.1 mm to 3.5 mm.

In quantitative solid state MR imaging or spectroscopy measurements based on FIDs, it is critical to reduce the receiver dead time to minimize distortion or complete loss of the initial FID samples. The limiting initial value of the time or k-space signal corresponds to the total area under the spectrum or image, respectively. It is therefore essential to capture the early points with high fidelity to achieve accurate quantitative measurements. For example, a ^{31}P spectroscopic resonance (no gradient or inhomogeneity present) with a 2.0 kHz FWHH (typical of bone apatite) has a T_2^* of 160 μs if the lineshape is approximated by a Lorentzian function. A receiver dead time or sampling delay of 70 μs yields an amplitude of the first sampled FID value which is 65% of the true value of this signal in spectroscopy. During imaging, a gradient is present, and T_2^* is far shorter by a large factor, making the amplitude error far worse. Although it is possible to reconstruct an image of a short- T_2 substance acquired with a large dead time or sampling delay, the SNR and quantitative accuracy of that image will be severely degraded compared to the image obtained when the k-space data near the origin is properly acquired. Acquisition schemes based on echoes can in principle measure the data at the k-space origin, but in practice it is difficult for *in vivo* ^{31}P MRI of bone mineral by clinical scanners to achieve a true echo time short enough to capture the k-space values near the origin without severe distorted by T_2 and T_2^* decay. This decay arises from the inherent resonance offset distribution due to ^{31}P chemical shift anisotropy and heteronuclear proton coupling. The “echo time” of the UTE sequence is actually a sampling delay, and does not represent a time at which these dephasing processes are rephased.

Likewise, the use of slice selection or self-refocusing pulses, whose durations can be on the order of or longer than T_2 and T_2^* , further degrades the quantitative accuracy.

Following the end of an RF pulse, the coil rings down with stored RF energy, and the scanner receiver is briefly overloaded. There is also a finite time required for the transmit/receive (T/R) switch to change from transmit mode to receive mode. This requires modification of the T/R switch to permit faster switching, reducing the ring down time by possibly lowering of the coil Q, and reducing the software-imposed minimum recovery delay so that signals may be sampled sooner following the RF pulse than normally permitted.

Reconstructing the image with recovered data corrects the artifacts which most seriously affect quantitative measurements such as bone mineral or matrix density. The most common means of correcting for receiver dead time or sampling delay is to place sampled data points in their correct k-space positions and assign zeros to the k-space positions of missing data. The baseline distortion created by this procedure is less severe in ^{31}P SMRI than in ^1H SMRI, since the sinc broadening function decays as $1/k$ (17) where

$$k = \gamma G t_{\text{dw}} \quad [6]$$

and

$$\gamma_{\text{P}} = (1/2.5) \gamma_{\text{H}} \quad [7]$$

In UTE, the gradient is turned on after the RF excitation pulse. This timing scheme has a significant effect on spatial resolution and SNR because many or a majority of the data points are not acquired under the maximum gradient strength. For example, if the maximum gradient is 30 mT/m, and the slew rate is 100 mT/m/ms, it will take 75 μs to reach the maximum strength, which is roughly half the T_2 of bone mineral, and there is extra time (receiver dead time) before the gradient ramp begins. This will reduce the overall spatial resolution. Therefore we expect that the SMRI method will provide better spatial resolution and quantitative accuracy for short T_2 imaging than UTE.

Applying the RF excitation pulse while the readout gradient is on and after it has stabilized helps to capture early k-space points with high fidelity. The loss of these points is due to the electronic recovery time of the receiver, T/R switch, preamp and coil. These are the same irrespective of whether or not a gradient is present. Turning on the gradient following the end of the RF pulse imposes complexity, instability and consequent possibility of degradation and artifacts in the reconstruction because of the need for regridding or nonuniform sampling intervals. Sampling in the presence of a fully stable and constant gradient for all k-space points eliminates these potential sources of image degradation. Because the RF pulse amplitude is very high and the pulse is short, the excitation profile is not substantially affected from view to view.

In addition to the much lower SNR of ^{31}P compared to ^1H , the long ^{31}P T_1 of bone mineral (about 17–19 s at 3.0 T) represents a major challenge. Both factors result in long scan times. In addition to utilizing phase-array coils and parallel acquisition methods, there are two kinds of tradeoffs/optimizations that relate SNR, scan duration and spatial resolution one can consider. In the first tradeoff, TR could be made shorter while increasing the number of k-space radii acquired for the same overall scan duration, while keeping the flip angle at the Ernst optimum (18). However, the intrinsic spectral linewidth of the ^{31}P resonance limits the spatial resolution for a given gradient strength. Likewise, the number of k-space radii limits

the spatial resolution of the reconstruction. Because the gradient strength used in this study is close to the maximum stable gradient strength of the scanner, and since the number of k-space radii has been chosen to match the linewidth limitation, there is not a great deal to be gained by shortening TR in this situation. In the second tradeoff, if the goal is to measure the mean bone mineral density over a sizable volume, such as $6 \times 6 \times 6 \text{ cm}^3$, and the bone mineral is distributed across the entire volume, the resolution can be reduced to 6 mm, the number of projections reduced to 1964, resulting in a total acquisition time of 9 min for the same TR = 275 ms used in this study. Although DXA and CT have much higher spatial resolution than 6 mm, the results of these bone density measurements are typically summarized as a single number (the BMD score) averaged over a whole bone or large volume of bone. The implication for MR-based BMD measurement is that a very small number of pixels will suffice in order to measure a BMD value; for example, DXA or qCT images are ultimately reduced to a single score reflecting the average BMD over the spine or hip. Even in the limiting case of measuring the entire FOV as a single pixel (spectroscopic limit) (19), it is conceivable that clinically useful bone mineral density data of a large volume might be obtainable by solid state MRI with very short scan times. A detailed treatment of the choice of number of radii, intrinsic linewidth and spatial resolution for 3D radial acquisitions is found in (15).

In conclusion, visualizing human wrist bone mineral by ^{31}P SMRI with a clinical MR scanner is feasible. The short T_2 difficulty associated with ^{31}P signal of bone mineral can be overcome with the coil and T/R system meeting the requirements of strong B_1 field and short receiver recovery time. In the case of the scanner used in this study, a significant advantage in reduction of the receiver recovery time was gained by using a frequency converter to enable ^{31}P reception via the faster-recovering ^1H channel, although other means to reduce the receiver recovery time may be possible. A gradient strength of 40 mT/m is sufficient to achieve ^{31}P bone image resolution of about 3.5 mm.

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References

1. U.S. Department of Health and Human Services. Bone health and Osteoporosis: A Report of the Surgeon General. 2004.
2. Kanis JA, Delmas P, Burckhardt P, Cooper C, Torgerson D. Guidelines for diagnosis and management of osteoporosis. *Osteoporos Int.* 1997; 7:390–406. [PubMed: 9373575]
3. Engelke K, Adams J, Armbrrecht G, Augat P, Bogado C, Bouxsein M, Felsenberg D, Ito Masako, Prevrhal S, Hans DB, Lewiecki EM. Clinical use of quantitative computed tomography and peripheral quantitative computed tomography in the management of osteoporosis in adults: the 2007 ISCD official positions. *J Clin Densitom.* 2008; 11:123–62. [PubMed: 18442757]
4. Tyler DJ, Robson MD, Henkelman RM, Young IR, Bydder GM. Magnetic resonance imaging with ultrashort TE (UTE) pulse sequences: technical considerations. *J Magn Reson Imaging.* 2007; 25:279–289. [PubMed: 17260388]

5. Anumula S, Magland J, Wehrli SK, Zhang H, Ong H, Song HK, Wehrli FW. Measurement of phosphorus content in normal and osteomalacic rabbit bone by solid-state 3D radial imaging. *Magn Reson Med*. 2006; 56:946–52. [PubMed: 17041893]
6. Anumula S, Wehrli SL, Magland J, Wright AC, Wehrli FW. Ultra-short echo-time MRI detects changes in bone mineralization and water content in OVX rat bone in response to alendronate treatment. *Bone*. 2010; 46:1391–1399. [PubMed: 20096815]
7. Wu Y, Ackerman JL, Chesler DA, Li J, Neer RM, Wang J, Glimcher MJ. Evaluation of bone mineral density using three dimensional solid state phosphorus-31 NMR projection imaging. *Calcif Tissue Int*. 1998; 62:512–8. [PubMed: 9576979]
8. Wu Y, Chesler DA, Glimcher ML, Garrido L, Wang J, Jiang HJ, Ackerman JL. Multinuclear solid state three dimensional MRI of bone and synthetic calcium phosphates. *Proc Nat Acad Sci USA*. 1999; 96:1574–8. [PubMed: 9990066]
9. Wu Y, Ackerman JL, Chesler DA, Graham L, Wang Y, Glimcher MJ. Density of organic matrix of native mineralized bone measured by water and fat suppressed proton projection MRI. *Magn Reson Med*. 2003; 50:59–68. [PubMed: 12815679]
10. Wu Y, Hrovat MI, Ackerman J, Reese TA, Cao H, Ecklund K, Glimcher MJ. Bone matrix imaged in vivo by water and fat suppressed proton projection MRI (WASPI) of animal and human subjects. *J Magn Reson Imaging*. 2010; 31:954–963. [PubMed: 20373441]
11. Idiyatullin D, Corum C, Park JY, Garwood M. Fast and quiet MRI using a swept radiofrequency. *J Magn Reson*. 2006; 181:342–349. [PubMed: 16782371]
12. O'Connell, RD.; Moeller, S.; Idiyatullin, D.; Corum, C.; Garwood, M. theoretical sensitivities of SWIFT and the ideal sequence (delta pulse-acquire) for ultra-short T2. Abstracts, Eighteenth Scientific Meeting, International Society of Magnetic Resonance in Medicine; Stockholm, Sweden. May 1–7, 2010; p. 2975
13. Carl, M.; Chiang, J-T.; Bydder, G.; King, K. Bloch simulations of UTE, WASPI and SWIFT for imaging short T2 tissues. Abstracts, Eighteenth Scientific Meeting, International Society of Magnetic Resonance in Medicine; Stockholm, Sweden. May 1–7, 2010; p. 694
14. Wu, Y.; Ackerman, JL.; Chesler, DA.; Wang, J.; Glimcher, MJ. In vivo solid state ³¹P MRI of human tibia at 1.5 T. Seventh Scientific Meeting, International Society of Magnetic Resonance in Medicine; Philadelphia, PA. May 22–28, 1999;
15. Wu Y, Dai G, Ackerman JL, Hrovat MI, Glimcher MJ, Snyder BD, Nazarian A, Chesler DA. Water and fat suppression projection MRI (WASPI) of rat femur bone. *Magn Reson Med*. 2007; 57:554–67. [PubMed: 17326184]
16. Liang, Z-P.; Lauterbur, PC. Principles of Magnetic Resonance Imaging. IEEE Press; 2000. p. 91
17. Liang, Z-P.; Lauterbur, PC. Principles of Magnetic Resonance Imaging. IEEE Press; 2000. p. 15
18. Ernst, RR.; Bodenhausen, G.; Wokaun, A. Principles of nuclear magnetic resonance in one and two dimensions. Clarendon Press; 1987. p. 124
19. Cao H, Nazarian A, Ackerman JL, Snyder BD, Rosenberg AE, Nazarian RM, Hrovat MI, Dai G, Mintzopoulos D, Wu Y. Quantitative ³¹P NMR spectroscopy and ¹H MRI measurements of bone mineral and matrix density differentiate metabolic bone diseases in rat models. *Bone*. 2010; 46:1582–1590. [PubMed: 20188225]

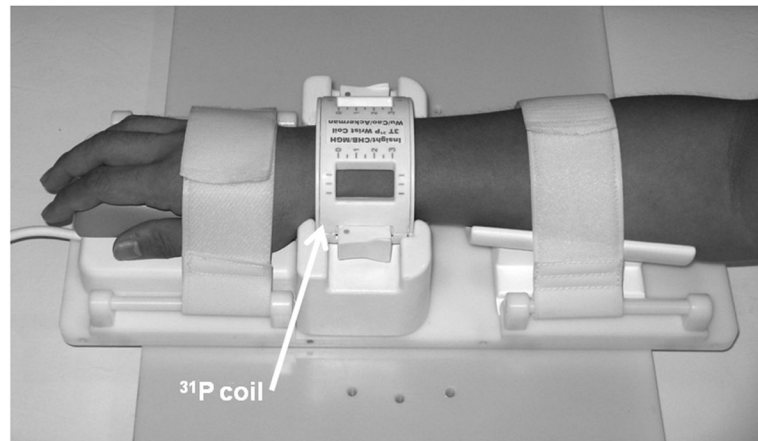


Figure 1. A ^{31}P quadrature low pass birdcage transmit/receive coil for human wrist imaging. The enclosure containing the passive T/R switch, quadrature hybrid, preamplifier and scanner interface is not shown.

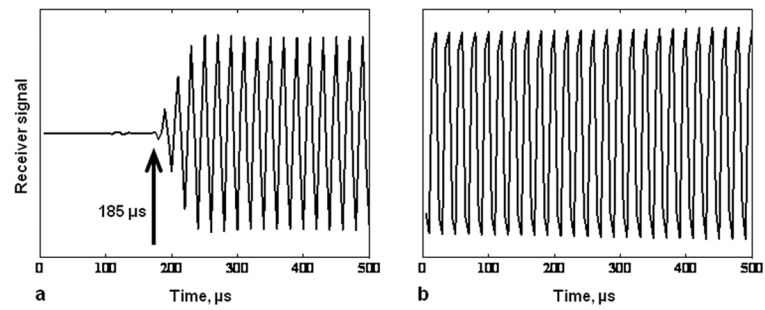


Figure 2.

Scanner receiver dynamic response to a continuous RF signal from a synthesizer. **a.** The response of the ^{31}P channel of the Siemens Trio 3T scanner receiver to which a continuous wave fixed amplitude signal was applied from a synthesizer. The attenuated output of the synthesizer was applied to the point at which the RF coil is connected, the transmitter was disconnected, and an acquisition was conducted. The receiver is completely blocked 185 μs after the receiver was turned on and is completely recovered after about 275 μs . **b.** Using the frequency converter scheme which detects the ^{31}P signal via the proton channel of the scanner (see the text and Figure 3), receiver blocking is reduced to about 5 μs .

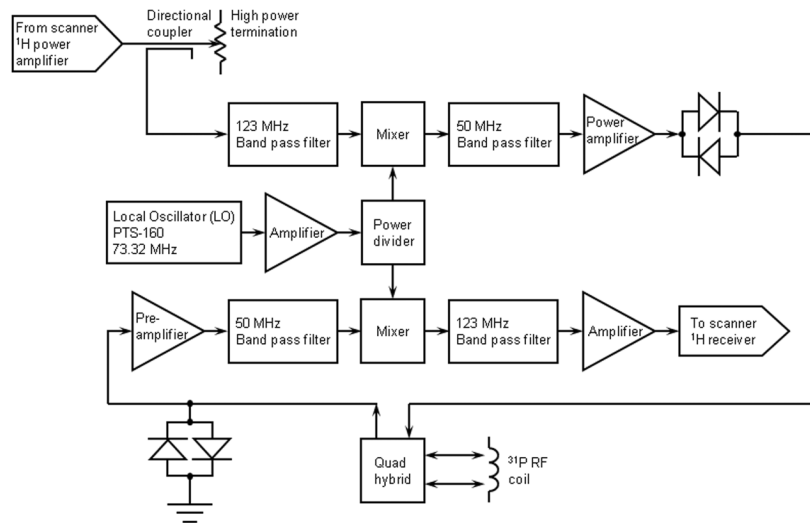


Figure 3. Simplified schematic block diagram of the frequency converter circuit for ^{31}P SMRI. A ^1H transmit pulse from the scanner is mixed with a local oscillator (LO) frequency from an external synthesizer to translate the pulse to the ^{31}P frequency, which is amplified by an external RF power amplifier and sent to the coil for ^{31}P excitation. The ^{31}P MR signal detected by the coil is amplified by a low noise narrow band preamplifier, mixed with the LO frequency to translated it to the ^1H frequency, and routed to the scanner ^1H receiver. Crossed diode symbols represent multiple diode pairs. Other amplifiers, attenuators, filters, bias tees, etc., are not shown.

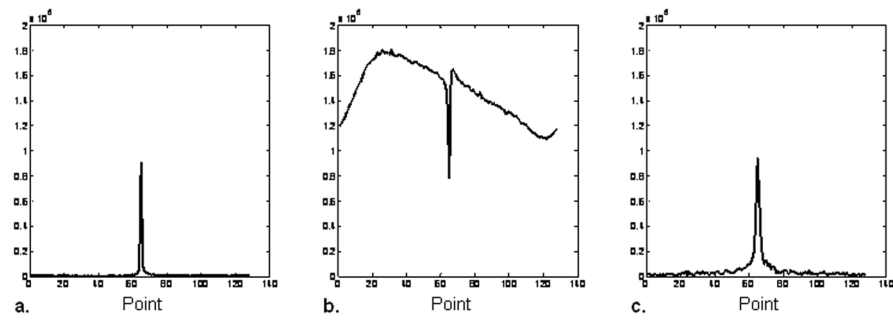


Figure 4.

Effect of coil loading on the receiver recovery time. **a.** ^{31}P magnitude spectrum of a vial of 85% phosphoric acid, which loads the coil severely, shortens the coil ring-down and reduces the Q to 37. The software receiver dead time delay was set to $10\ \mu\text{s}$. **b.** ^{31}P magnitude spectrum of a vial of hydroxyapatite, which does not load the coil and maintains the Q at 115. With the software receiver dead time delay set to $10\ \mu\text{s}$ the coil ring-down creates a strong artifact in the spectrum. **c.** Same as **b**, except that the receiver dead time delay is increased to $20\ \mu\text{s}$, which blocks the coil ring-down signal and permits acquisition of a distortion-free spectrum.

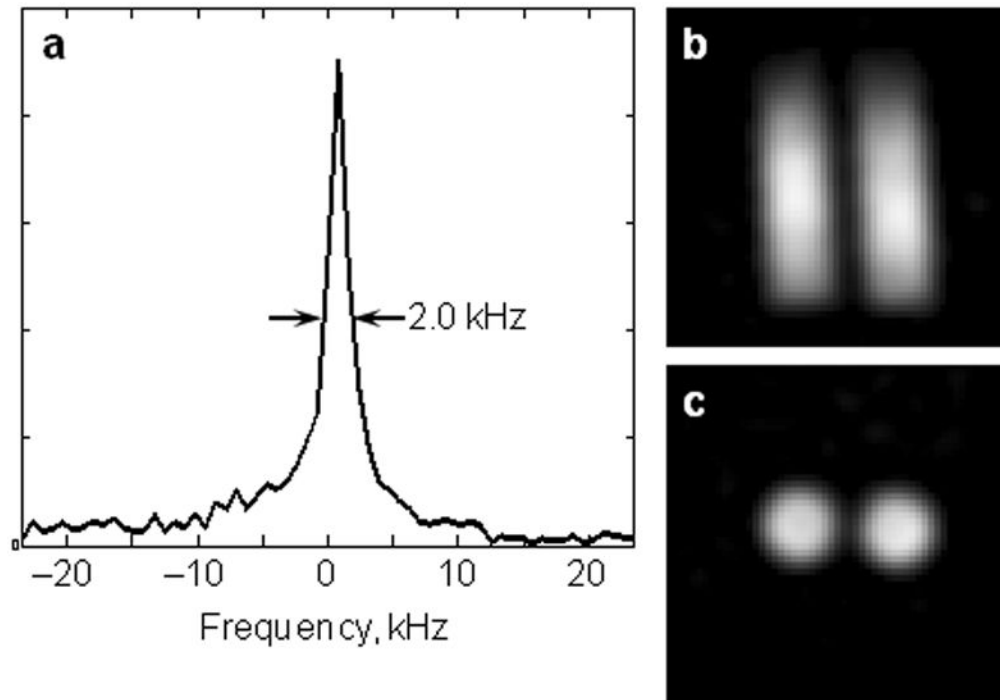


Figure 5. ^{31}P spectrum and solid state MRI (SMRI) of hydroxyapatite. **a.** ^{31}P magnitude spectrum of a vial of hydroxyapatite (HA). The full width at half maximum of the resonance is 2.0 kHz. **b.** Coronal slice of the ^{31}P image of the two-tube phantom. The inside diameter of each HA tube is 14 mm and the space between the two columns of HA is 2.9 mm. **c.** Transverse slice of the ^{31}P image of the two-tube phantom.

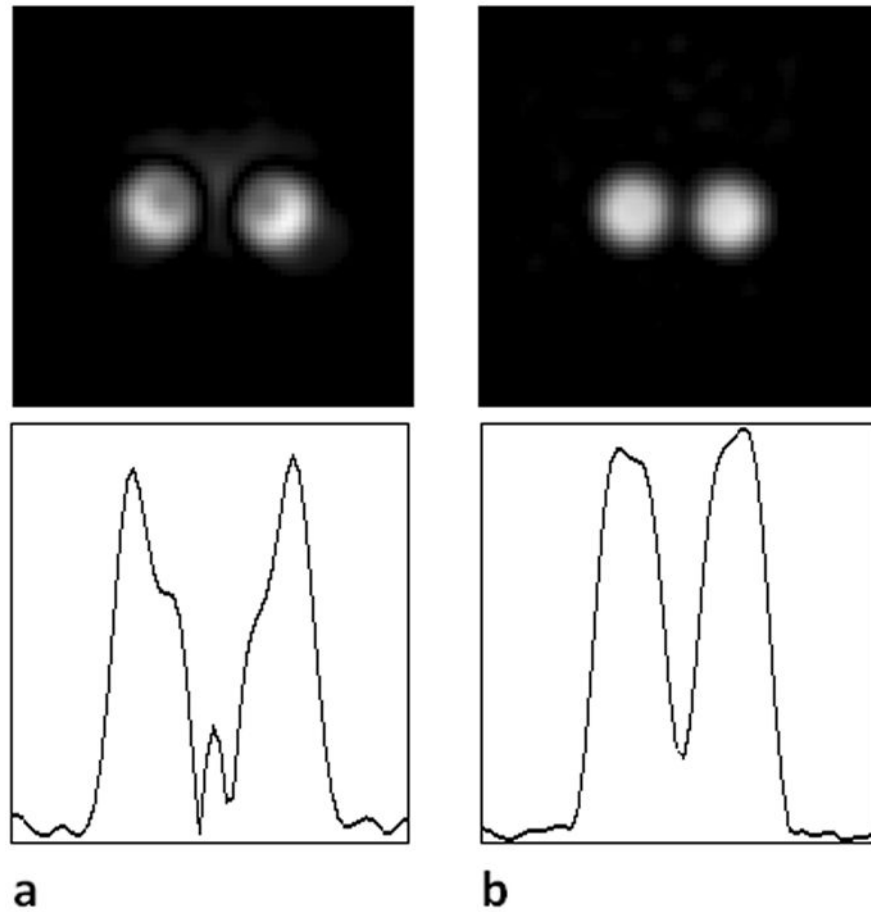


Figure 6. ^{31}P image reconstruction with and without the recovered data missed in the receiver recovery. The receiver dead time delay was $20\ \mu\text{s}$, and the dwell time was $7\ \mu\text{s}$ in both panels. Therefore, three data points were lost in the receiver dead time during the regular acquisition. **Upper row:** Transverse slice from the ^{31}P SMRI of the two-tube phantom. **Lower row:** Line profile across the corresponding image. **a.** Image reconstructed without any correction of time domain data or recovery of missing points; the first sampled data point is positioned at the origin of k-space. The image is severely distorted. **b.** Image reconstructed with the data in correct k-space positions and with missing points recovered (second data set). The artifacts in **a** disappear.

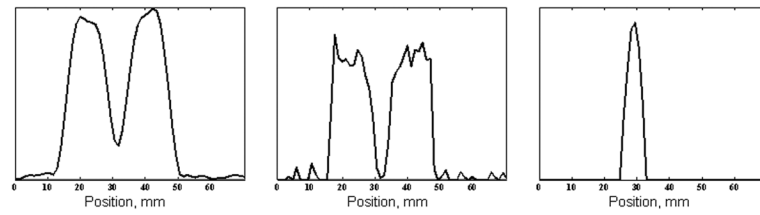


Figure 7. Point spread function of ^{31}P SMRI. **a.** Line profile through a transverse slice of the ^{31}P SMR image of the two-tube phantom. **b.** Deconvolved axial profile of the ^{31}P SMRI obtained using blind deconvolution. **c.** The combined point spread function of the system (including intrinsic spectral width of the HA) and reconstruction. The full width at half maximum is 5.0 mm.

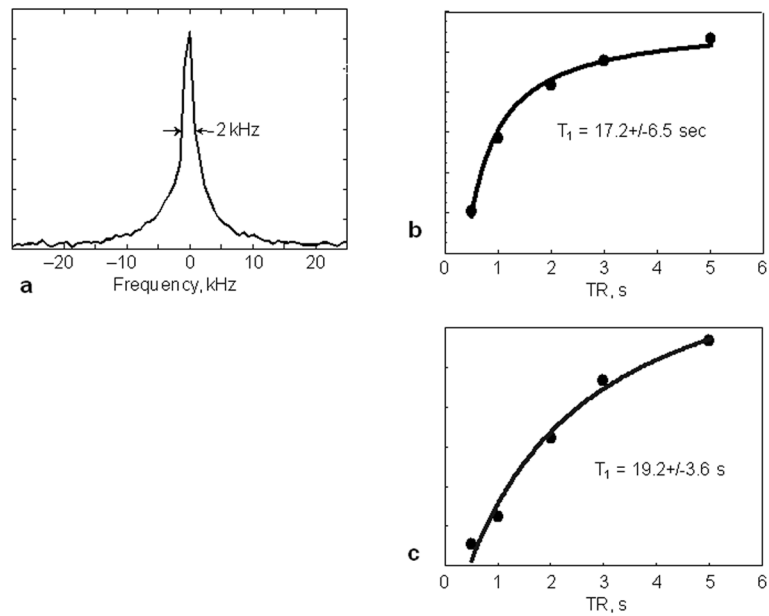


Figure 8.

^{31}P resonance line width and T_1 of human bone mineral measured *in vivo*.

- a.** Single pulse ^{31}P spectrum showing a full width half maximum resonance line width of 2 kHz from the wrist of a 28-year-old female volunteer.
- b.** Spectroscopy progressive saturation T_1 measurement from the same volunteer of a. yields an *in vivo* human bone mineral ^{31}P T_1 of 17.2 s.
- c.** SMRI progressive saturation T_1 measurement from a 27-year-old male volunteer yields ^{31}P T_1 of bone mineral of 19.2 s.

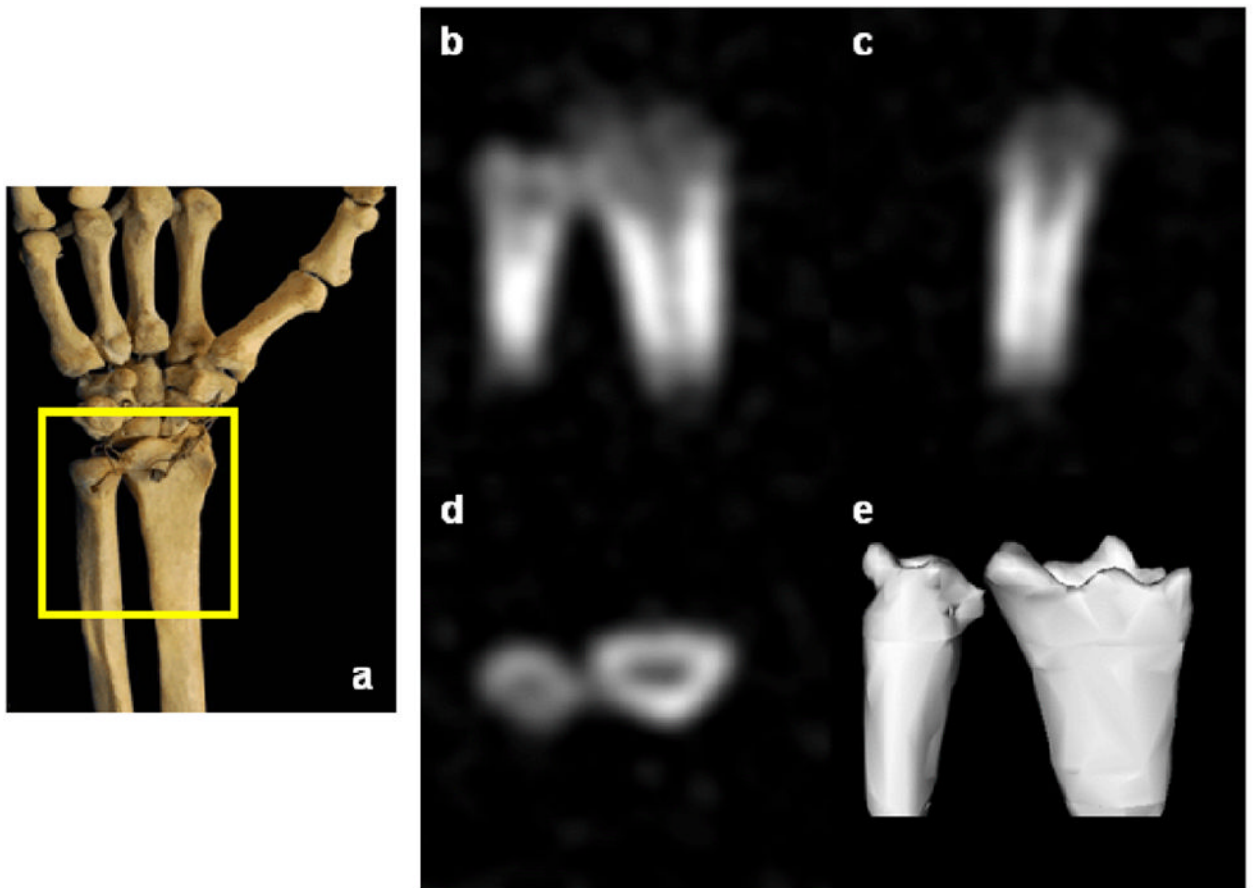


Figure 9.
In vivo ^{31}P SMRI of bone mineral of the wrist of a healthy 41-year-old male volunteer. **a.** Schematic view of the scanned region. **b–d.** Posterior-anterior, lateral, transverse image slices. **e.** Three-dimensional isosurface rendering of the full data set.

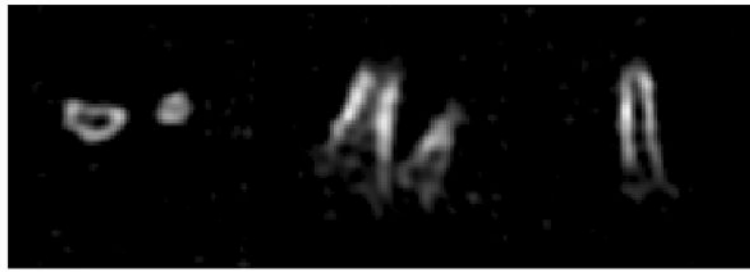


Figure 10.
In vivo ^{31}P SMRI of bone mineral of the wrist of a healthy 60 year old male volunteer, 30.11 mT/m gradient, 37 min scan, cortical SNR = 16, overall resolution: 4.8 mm.
Left: Transversal view.
Middle: Coronal view.
Right: Sagittal view.

Table 1

Details of ^3P images of healthy volunteers

Image	Age, yr	Gender	Gradient, mT/M	Dwell, μs	Number of Projections	Resolution, mm ^d	Duration, min:s	Cortical SNR	Note
1	27	Male	27.96	7.0	184	21	1:32-15:20	9-23	<i>b</i>
2	28	Female		5.0					<i>c</i>
3	47	Male	27.96	7.0	998	8.8	4:35	33	<i>d</i>
4	40	Male	27.96	7.0	8148	5.1	37:0	30	<i>e</i>
5	33	Male	30.11	6.5	8148	4.8	37:0	17	<i>f</i>
6	25	Female	30.11	6.5	8148	4.8	37:0	18	
7	60	Male	30.11	6.5	8148	4.8	37:0	16	

^aComputed from imaging parameters.^b T_1 measurements with TR: 0.5, 1, 2, 3, 5 s. Scan duration varied from 1 min 32 s (TR=0.5 s) to 15 min 32 s (TR=5 s), SNR varied from 9 (TR=0.5 s) to 23 (TR=5 s). $T_1 = 19.2$ s.^c T_1 measurements using ^3P spectroscopy. $T_1 = 17.2$ s (see text).^dSpatial resolution insufficient for usable image quality.^eShown in Fig. 9.^fShown in Fig. 10. Δx_j was improved to 3.7 mm due to stronger gradient strength, Δx_p remained at 3.0 mm, Δx_{all} was improved to 4.8 mm.