Carbon-13 NMR relaxation studies of pre-melt structural dynamics in $[4^{13}C$ -uracil] labeled E. coli transfer RNAYal*

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ABSTRACT

We report 67.8 MHz carbon-13 spin-lattice relaxation studies.on [4-¹³C-
uracil] labeled tRNA^{Yal} purified from <u>E. coli</u> SO-187. Following ¹³C-enriched \mathtt{C}_4 carbonyl resonances from modified and unsubstituted uridines scattered throughout the polymer backbone enables us to determine dynamical features in both loop and helical stem regions. The experimental results have been analyzed in terms of a model of isotropic overall molecular reorientation. "Anomalous" residues for which the experimental data cannot be accounted for in terms of the model provide an assessment of local and regional properties. Thus, "native" tRNA{"' under physiological conditions of magnesium (10 mM) and temperature (20°-40 C), exhibits the following characteristics: 1) uridines held rigidly in helical stems and tertiary interactions display correlation times for rotational reorientation of 15-20 nsecs, typical for overall tRNA motion; 2) uridines in loops such as the wobble residue uridine-5-oxyacetic acid (V₃₄) are quite accessible to solvent; moreover V₃₄ and another loop
residue, D₁₇, exhibit local mobility; 3) the tertiary interactions involving
4-thio uridine (s⁴U₈) and A₁₄ and ribothymidine (rT₅₄) an temperature increases.

INTRODUCTION

Nuclear magnetic resonance techniques applied to transfer RNA structure, conformation and interaction in solution have been utilized very extensively and fruitfully during the past few years, as evidenced by the yearly reviews on the subject (1-6). A large fraction of the studies on individual tRNA's have dealt with the low field imino and amino protons involved in hydrogen bonding and the non-exchangeable high field protons from assorted modified nucleosides usually located in loop segments. Other investigations utilizing natural abundance $31p$ (7,8) and $13c$ -nmr on specifically enriched residues (9-11) have appeared. Despite the considerable effort in obtaining site specifically 13 C-enriched individual tRNA's, the use of 13 C-nmr offers certain advantages: 1) wide chemical shift range; 2) ease of resonance assignment; and 3) ability to relate relaxation data to dynamic features of tRNA itself and in complexes with proteins of interest.

In the case of E. coli tRNA $\{a\}$, Reid and co-workers (4, 6, 12) have shown that the low field imino proton spectrum can be interpreted in terms of a folded structure with \sim 7 tertiary hydrogen bond interactions and \sim 20 secondary base-pairs, two of the latter shown by Johnston and Redfield (13) to be due to the G_{50} -U₆₄ wobble base pair (Fig. 1). Solution structures of Class I tRNA's such as $tRNA₁²$ and $tRNA₂²$ probably are fairly close to the crystal structure (14) determined for $tRNA^{Phe}$ (15,16), shown in Fig. 3.

In the presence of 10 mM magnesium ion, the E. coli tRNA $\{a\}$ solution structure as depicted in Fig. 3 is maintained to 45-50°C (12) at which point the tertiary interactions begin to be disrupted. We have an interest in the dynamics of the native structure of $tRNA₁^{Ya1}$ at physiologically relevant temperatures around 30-40°C and how this structure might be altered in the presence of cognate aminoacyl synthetase. In this paper we describe 67.8 MHz 13 C-spin-lattice relaxation and NOE investigations on $[4-$ ¹³C-uracil] labeled $tRNA$ ²al of E. coli S0-187 as a function of temperature. We have monitored C₄ carbonyl resonances from eleven of the fourteen modified and unmodified uracil residues scattered throughout the loop and stem regions of this molecule (Fig. 1). Based upon a model of isotropic rotational reorientation, we have calculated motional correlation times at each of these residues. Residues which display anomalous behavior provide insights about local and regional dynamics. In addition, we have estimated the relative importance of chemical shift anisotropic (CSA), proton dipolar and nitrogen dipolar mechanisms to the spinlattice relaxation processes.

METHODS

Isolation and purification of $[4-13C-uracil]$ labeled tRNA $]$ ^{al} from the \underline{E} . coli uracil auxotroph SO-187 was accomplished essentially as outlined earlier (9), except that homogeneous $tRNA²$ could be obtained in two chromatographic steps: 1) BD cellulose (ambient temperature) and 2) Sepharose 4B at 4°C using reverse (NH₄)₂SO₄ gradients. 19 mg of the tRNA^{Yal} was dissolved in a 0.4 ml solution of double deionized charcoal filtered water containing 20% D₂0 (Stohler, 99.8% d), K₂HPO₄ (20 mM; pD 7.2), NaCl (50 mM), MgCl₂ (10 mM), EDTA (5 mM) , $\text{Na}_2\text{S}_2\text{O}_3$ (3 mM) , and 0.02% (w/v) Na_3 to retard bacterial growth. The sample was placed in a 5 mm tube previously cleaned with 50% nitric acid. Even though dissolved oxygen would probably be a minor influence on the relaxation rate in the $0.2 - 0.5 \text{ sec}^{-1}$ range found in this study, the tRNA $\{^{al}$ sample was not degassed.

Carbon-13 NMR measurements at 67.8 MHz were conducted on a JEOL FX-270,

using proton noise decoupling at 2 watts power to prevent dielectric heating in these aqueous salt solutions. For these carbonyls we used a spectral window of 3 KHz, centered at 177 ppm from TMS. T_1 values were determined using the fast inversion-recovery method (17) with the sequence $-(180^{\circ} - \tau -$ 90°(FID)-WT)n, where τ_{∞} equals 5T₁ (20 sec), waiting time, WT, between pulse trains equals $3T_1$ (12 sec) and n is 2000. Usually spectra were collected using 9 different t values. Using the manufacturer's software, standard least squares fitting of the line intensities, A_i , at τ_i was conducted and T_1 values determined from the equation $ln(A_{\infty} - A_i)/2A_0 = -i/\top_1$ where A_{∞} is the limiting intensity at $\tau = 5T_1$ and A_o is the extrapolated value for $\tau_i = 0$.

Nuclear Overhauser enhancements, 1n, were determined by comparing peak intensities in spectra taken with the decoupler on continuously and the decoupler gated off during a $5T_1$ delay between 90° observe pulses. Probe temperatures were controlled with a programmed variable temperature unit, previously calibrated with 1,3-propanediol.

RESULTS AND DISCUSSION

Based upon nmr studies of the imino and amino protons involved in secondary and tertiary hydrogen bonding interactions in E_0 coli tRNA Va1 (12), one may confidently assume that the general overall three-dimensional "native" structure of the molecule in solvent mileau containing 10 mM MgCl₂ is maintained over the 20-40°C temperature range under consideration here. Therefore the observed changes in relaxation parameters with temperature described below undoubtedly reflect segmental modulations of this basic molecular framework.

In Fig. 1 we display the primary sequence of E . coli tRNA 2d [18, 19] arranged in the usual cloverleaf pattern with the $[4-13c]$ uridines underlined and the corresponding structures of these nucleosides as shown.

The 67.8 MHz carbon-13 spectra for $[4-13C-uracil]$ tRNA $]$ ^{al} at the temperature extremes for this study, 20 and 40°C are presented in Fig. 2. These are the fully recovered " τ_{∞} spectra" of the inversion-recovery measurements, where $\tau_{\rm m}$ = 20 sec and 2000 transients were transformed. In addition to the four assigned resonances (assignment criteria were reported earlier (9) for $S^4 \cup_{8}$, D_{17} , and V₃₄] T₅₄ has subsequently been assigned via specific $13c\{1H\}$ decoupling) we have monitored unassigned resonances #3, 4, 5, 6, 7, 9 and 12 in the T, and NOE experiments.

Temperature increase generally results in line narrowing and shifts to lower field (Fig. 2). Of particular interest is the sharpening of $S^{4}U_{R}$ and the larger downfield shift of the V₃₄ resonance (0.5 ppm). The former may be a result of disruption of the tertiary $S^4\mathsf{U}_8$ - A_{14} interaction (12; see Fig. 3) while the latter undoubtedly comes about from flexibility in the anticodon loop enabling destacking of V_{34} from purine neighbors particularly A₃₅; (see Ref. 9).

Spin-lattice Relaxation

Hamill, et.al., (20) have previously described relaxation studies on [4-

¹³C-enriched Uracil Derivative in tRNA^{Yal}

FIGURE 1. Top: cloverleaf representation of tRNAi sequence (18,19) with uridines having ''C-enriched C. carbonyls underlined. The structures of these nucleosides are shown in the lower panel

FIGURE 2. Fully recovered (τ_{∞} = 20 sec.) spectra of 19 mg (4-'°C-ruacil)
labeled tRNAY^{al} in 0.4 ml 80/20 H₂O/D₂O system (see Methods). 2000 transients, 90° pulse width was 7.2 usec.

13C-uracil] labeled unfractionated tRNA of Salmonella typhimurium. In this work an assessment was made of the predominant relaxation mechanisms of the quaternary C_4 carbonyl carbon of the uracils as a function of magnetic field strength. The calculated contribution of chemical shift anisotropy to the spin-lattice relaxation increased from about 25% at 23.5 kG to about 80% at 84.6 kG, with less than $5-10\%$ 14 N dipolar contribution and the remainder due to proton dipolar mechanisms. Experimentally the relaxation rate was found to decrease with increase in magnetic field although not as drastically as theory would predict for a molecule such as tRNA tumbling outside the extreme narrowing limit.

 R_1 values obtained at 23.5 kG by Hamill, et. al, for S^4 U, D and several uridines in the S. typhimurium mixture, at 37°C, in water are respectively 0.7, 0.59 and 0.59 to 0.67 sec^{-1} . From Figs. 4 and 5, the values we obtained

FIGURE 3. Three dimensional framework as per X-ray structure of yeast tRNA^{phe} (15,16) with superimposed ¹³C-enriched uridines as in <u>E</u>. <u>COI1</u> tRNAI"'

at 63.5 kG, 34°C in 80/20 H₂0/D₂O are 0.47, 0.34 and 0.3 to 0.4 sec⁻¹. These significantly lowered values confirm that $tRNA^{2d}$ is undergoing molecular reorientation with a correlation time, τ_R , to the right of the R₁ vs τ_R maximum, e.g. outside the extreme narrowing region.

Spin-lattice relaxation times at several temperatures for the uracil C_{4} carbonyl resonances of $[4-13c-uracil]$ labeled tRNA $\{^{al}$ are presented in Table I. The corresponding relaxation rates, R_1 (sec⁻¹), for the samples in 80/20 H₂0/D₂0 solution are plotted in Figs. 4 and 5 versus temperature. From the different dependencies seen here, we designate three "classes" of uridines: Class I - large increases in R₁ vs. temperature, e.g., $S^{4}U_{8}$, #3, and #12 (Figs. 2 and 4); Class II - moderate increase in R_1 vs. temperature, e.g., #4, 5, 6, 7, 9 and rT_{54} (Fig. 2 and 5); Class III - very little net change in R₁, e.g., D_{17} and V_{34} (Figs. 2 and 5).

We consider that the uridines whose C_{Δ} carbonyl resonances exhibit a small but consistent increase in R_1 vs. temperature are those which are located in helical H-bonded stems and H-bonded tertiary folding and thus experience molecular reorientation similar to that of the tRNA molecule itself. Since we know from the field effect on R_1 mentioned above that the

FIGURE 4. Spin lattice relaxation rates, R_1 (sec $\tilde{\ }$) versus temperature for highly dependent uridines (Class I). Error estimate ± 0.02 sec-1.

<code>FIGURE 5. Spin</code> lattice relaxation rates, <code>R_{1</code> (sec $\texttt{--}^{\star}$) versus</code>} temperature for slightly dependent (Class II) and essentially independent (Class III) uridines. Error estimate ± 0.02 sec-1.

SPIN LATTICE RELAXATION TIMES FOR URACIL C₄ CARBONYLS IN E. COLI tRNAYal (a)

(a) Values are in seconds, measured at 67.8 MHz via fast inversion-recovery method. (Canet, et. al., 1975). 2000 transients transformed. Estimated error i 10% or less.

(b) Chemical shifts measured from internal dioxane in another tRNA^{Val} converted to TMS by adding 66.3 ppm. sample and

(c) Peak observed, but not processed in computer calculations.

(d) Peak not resolved
(e) This value may be underestimated since T_∞ = 30 sec was used**.**

correlation time, τ_R , for overall reorientation of the tRNA lies to the right of the R_1 vs. T_R maximum, it would be expected that as temperature increases and τ_R becomes smaller that R_1 should increase until the maximum is reached. Calculated Relaxation Parameters

In order to estimate the correlation times for rotational reorientation of the various uridines and modified uridines and to assess the contributions of proton-dipolar, nitrogen-dipolar and chemical shift anisotropy mechanisms, we have utilized the following treatment. We assume an isotropic tumbling model and consider only auto correlated internuclear spin pairs in the dipoledipole mechanism. We justify, in part, the neglect of cross-correlated interactions involving three or four nuclear spins on the basis that our data did not exhibit any pronounced deviations from single exponential behavior, and since initial slopes depend only on auto correlated terms there is insufficient information to characterize the importance of such terms. Finally, relaxation data on macromolecules such as tRNA unfortunately do not have

intrinsic accuracy sufficient to provide cross-correlation data in the inversion recovery type experiments on total carbon magnetization in the proton decoupled experiment. The following equations for isotropic tumbling were therefore used:

$$
R = \frac{1}{T_1} = R_{CSA} + \frac{\sum}{s} R_{IS}
$$
 (1)

$$
R_{CSA} = \frac{1}{T_{1CSA}(1)} = \frac{1}{15} (\gamma_C H_0)^2 \Delta \sigma^2 J(\omega_1)
$$
 (2)

$$
R_{IS} = \frac{1}{T_1(IS)} = \frac{N_S \hbar^2 \gamma_I^2 \gamma_S^2 S(S+1)}{15 \cdot r_{IS}^6} \chi_{IS}
$$
 (3)

where

$$
\chi_{IS} = J(\omega_I - \omega_S) + 3J(\omega_I) + 6J(\omega_I + \omega_S)
$$
 (4)

and

$$
J(\omega) = \frac{2\tau_R}{1 + \omega^2 \tau_R^2}
$$
 (5)

Where $\Delta\sigma$ is an effective chemical shift anisotropy (σ_{11} - σ_{12} for an axially symmetric case), N_c is the number of chemically, and structurally identical S spins (protons, nitrogens, etc.) and τ_R is the reorientational correlation time for effective isotropic-tumbling. The other symbols are standard.

Using the approach of Kuhlmann, Grant and Harris(21) for calculating NOE values in systems having three different magnetic nuclei $(^{13}C, ^{14}N$ and 1H in this study), one may calculate the NOE for the carbon-13 isotope, under proton decoupling only, from the following

$$
n_C = \frac{\gamma_H \gamma_C f_{CH} R_{1N} - \gamma_H \gamma_C f_{CN}}{\gamma_C^2 R_{1N} R_{1C} - \gamma_N \gamma_H f_{CN} f_{CH}}
$$
(6)

where

$$
f_{ij} = 6J(\omega_i + \omega_j) - J(\omega_i - \omega_j)
$$
 (7)

and
$$
R_{1N} = \sum_{S} R_{NS}
$$
 and $R_{1C} = \sum_{S} R_{CS}$

as defined above. The value for n_c differ slightly from the standard expression given for the carbon-proton NOE neglecting the presence of the nitrogen spin. Thus

$$
m_{\rm C} \approx \frac{\gamma_{\rm H} f_{\rm CH}}{\gamma_{\rm C} R_{\rm 1C}} \tag{8}
$$

Use of Equation 1 not only neglects the cross terms between various R_{IS} terms (the cross correlated power densities) but it also assumes a negligible cross term between the R_{CSA} and R_{IS} terms. Theoretical estimates (22) of this interference term indicates its importance in deviations from single exponential behavior, but so far an experimental verification of its importance is lacking. Furthermore, its effect on inversion recovery, proton-decoupled data should be rather small.

While the assumption of isotropic rigid body tumbling may approximate the motion for the stems and other rigid portions of the tRNA molecule, this assumption may break down seriously for segments of the molecule undergoing internal motion and for sections of the loops exposed to the aqueous protons in the water media. Attention to details of segmental motion go beyond the scope of our theoretical treatment, but experimental anomalies may suggest the importance of such terms, as noted below.

Hydrogen atoms at N_3 and C_5 were assumed to exert the primary influence in proton-carbon dipolar relaxation, and N_3 for nitrogen-carbon dipolar effects. Distances from the C₄ to these protons and N₃ were determined using bond lengths and angles from crystallographic data on uridine (23) and dihydrouridine (24). The following distances (in A) were used: C-H (1.08); C-^N (1.37); N-H (1.10); in uracil, C4-N3-H (2.03); C4-C5-H (2.16); C4-C5-C6-H (3.40); in dihydro uracil $C_4-C_5 < H$ $\{2.96\}$. The value for $\Delta\sigma$ was taken to be 160 ppm (25) for all residues except dihydrouracil, for which we studied the

TABLE II

Predicted Spin-Lattice Relaxation Parameters and Rotational Correlation Times for Uridine C₄ Carbons in E. Coli tRNA^{yal} at 34°C^(a)

^(a) T_l's in seconds, τ_R in nanoseconds. Values have also been predicted at
20, 26, and 40°C: they order with the 34° data above and the experimental data in Table I. First entry, 80% H_2O ; second, 100% D_2O .

effects of varying $\Delta\sigma$ in 20 ppm increments from 140 to 200 ppm.

Table II contains the calculated contributions to T_1 from the three mechanisms and the correlation times for isotropic reorientation at 34°C in 80% H₂0 and in 100% D₂0 solvent. Several interesting features emerge upon comparing the calculated and experimental T_1 's (Table I). We estimate that the CSA contribution accounts for 60-75% of the relaxation for these quaternary carbons at 6.3 Tesla. Of the 25-40% due to dipolar effects, all but 6-8% is due to nearby protons.

The range of τ_R , 12-30 nsec, is of similar magnitude reported by others, e. g., in 13 C studies of bulk tRNA (20, 26), and 31 P-nmr (8) and pulsed fluorometry (27) of yeast tRNA^{Phe}.

Applicability of the model

If the model is realistic, with respect to those atoms contributing to dipolar relaxation, the τ_R 's calculated at 34°C in H₂O and D₂O solvents should be similar. One can see that for "Class II" peaks #4, 5, 9 and 11 (rT_{54}), the τ_R 's are indeed quite comparable (due to coalescence in D₂0, peaks 6 & 7 could not be monitored). It is reasonable to suppose that the residues giving rise to these resonances are in stems or other relatively rigid regions of the molecule.

Several resonances displayed relaxation behavior which could not be reconciled with the model. These interesting anomalies, which could provide information about regional and local dynamics, include the following: Peaks #3, 12 and V_{34} - For Class I peaks #3 and 12 and V_{34} (Class III) the τ_R 's between H₂0 and D₂0 differ by nearly a factor of two; the large diminuition of relaxation rate in going from H_2O to D_2O might occur if solvent is more readily accessible, for example in loop segments. Thus peaks #3 and #13 may be due to uracil located in loops, as is the case for V_{34} . The very large solvent effect for this latter unusual uracil in the wobble position may be due in part to the fact that divalent metal binds to the carboxylate (De and Schweizer, unpublished), thus one is replacing approximately six solvent molecules about the cation, in this case Mg(II).

No better agreement between H₂0 and D₂0 solvents was found for peaks #3 and 12 if either were considered to arise from pseudouridine (Ψ_{55}) , therefore not having a proton at C₅.

 $_2$ ^{4U}₈ - It is of interest to note the considerably smaller T₁ (and T_R) and the large dependence of R₁ on temperature (Fig. 4) for $S^4\mathbb{U}_8$. This residue participates in a tertiary H-binding interaction with A_{14} (12) involving the N_3 -H. Even at 40° C ¹H-nmr studies show this H-bond is intact. However, motion of the proton dipole may exhibit a shift in frequency distribution ("breathing rate") toward the 13 C Larmor value with temperature increase, providing more efficient relaxation of C_4 . Thus it appears that gradual weakening of the tertiary interaction occurs from 20°-40°C.

 rT_{54} - Although the R₁ vs temperature behavior exhibited by this "Class II" residue from 20°-34°C (Fig. 5) is what we expect from a rigidly constrained monomer, it is interesting to note the large rise in R_1 between 34° and 40°C. What we are observing here is undoubtedly incipient melting of the tertiary intraloop interaction involving rT_{54} and A_{58} (28).

 D_{17} - Dihydrouridine (D₁₇), like V₃₄, displays very little net change in R₁

vs. temperature (class III behavior). This may indicate the presence of independent local motion. Previous results have been interpreted in terms of D-loop flexibility (20,26). From the comparisons in Table II, there is nothing to indicate anything unusual about D_{17} ; no peculiar solvent effects are exhibited. Characteristics of the tertiary folding between D- and TPCG loops (Fig. 3) may push D_{17} into the molecular interior so that solvent is inaccessible.

There was some concern about the use of 160 ppm for $\Delta\sigma$ determined for uridine (25) since we are dealing with the saturated C_5-C_6 segment in D_{17} . We studied the effect of varying $\Delta\sigma$ in 20 ppm increments from 140 to 200 ppm upon the calculated relaxation parameters (Table III). The closest match between parameters calculated for H₂0 and D₂0 solvents was obtained using $\Delta\sigma$ = 200 ppm. Employing the higher $\Delta\sigma$ also results in a change in estimated CSA contribution to the spin-lattice relaxation from 49 to 60%, the latter percentage equivalent to the results calculated for the various uridines (Table II). NOE Measurements

We have measured NOE's for the various C_4 peaks as a means of detecting potential flexibility in the D-loop (Table IV). The η values for D₁₇ are

TABLE III

Effect of Varying $\Delta\sigma$ on Predicted Relaxation Parameters for Dihydrouridine (D_{17}) at 34°C^(a)

 (a) T₁'s in seconds, TR in nanoseconds. First entry is for 80% H₂0, second for 100% D_2 0.

consistently higher (0.3 - 0.6) than for other uridines (0.1 to 0.3) which is to be expected for residues in more mobile segments. The larger enhancements for the D_{17} C₄ carbonyl were also observed in D_2O solvent in contrast to the observations of Hamill, et. al. (20), for the unfractionated Salmonella tRNA mixture. On the basis of these measured $n's$, we would say that the D_{17} undergoes local motion which is faster than the rate of overall molecular reorientation in agreement with previous studies (20, 26).

Also listed in Table III are n values calculated as per Equations 6-8 above. Since the model assumes overall isotropic molecular reorientation in a motional regime outside extreme narrowing and does not provide for segmental modulation, it is not surprising that the calculated n's are small and near 0.1, the expected limit for slowly tumbling molecules.

Although the η values for V₃₄ are not significantly larger than uridines,

TABLE IV

Measured and Calculated Nuclear Overhauser Enhancements, n, for Uracil C4 Carbonyls in E. coli tRNAYal (a)

(a) Peak intensities compared between spectra taken with decoupler on
continuously (+NOE) and with decoupler gated off (-NOE) during delay of 5T₁
(20 sec) between 90° observe pulses. 2000 transients transformed; solution conditions as in Figure 3. Estimated error limits *0.1. Calculated values in second column for each temperature.

(b) Peaks are numbered from low to high field (see Table ^I and Figure 3).

- (c) Difficult to measure due to low intensities for $S⁴U$
- (d) Unable to measure

in contrast to those for D_{17} , three findings may be interpreted in terms of flexibility in the anticodon loop segment containing V_{34} : 1) the relatively invariant response of R_1 to temperature cited above; 2) the low R_1 itself, $0.24 - 0.27 \text{ sec}^{-1}$ (Fig. 5); 3) the facile temperature induced destacking of V_{34} from A₃₅ and other purine bases in the anticodon loop as monitored by downfield shifting (Figs. ¹ & 2).

Conclusions

From these experimental and theoretical studies we can construct the following conceptualization of $tRNA^{2d}$ structural dynamics in solution over the physiologiclly relevant temperature range 20-40°C. The secondary and tertiary solution structure promulgated earlier (1, 6, 12) remains essentially intact over these temperatures. However, the tertiary interactions involving $s⁴U_R$ are gradually loosened as temperature rises and there is incipient melting of the rT_{54} -A₅₈ tertiary bond between 34° and 40°C.

We confirm that the loop segment containing dihydrouridine is more flexible than the overall molecular backbone. In addition, we find that the anticodon loop is indeed also quite flexible, as monitored by V_{34} in the wobble position of the anticodon.

The similar solvent effects on C₄ T₁'s observed for V₃₄, #3, and #12 indicates that the latter two uridines are probably also located in loops where solvent is readily accessible and tertiary interaction is not occurring. With reference to Fig. 1, the likely candidates are U_{33} , U_{59} , U_{49} or Ψ_{55} , although the latter supposedly undergoes tertiary H-bonding with G_{18} in the Dloop (28). Those residues more rigidly held in secondary (stems) and tertiary interactions with τ_R 's characteristic of the tRNA molecule itself may produce peaks #4, 5, 6, 7, and 9. Again with reference to Fig. 1, the corresponding residues may be U_4 , U_7 , U_{12} , U_{29} , U_{64} , or U_{67} .

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