

Gating transitions in the selectivity filter region of a sodium channel are coupled to the domain IV voltage sensor

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Voltage-dependent ion channels are crucial for generation and propagation of electrical activity in biological systems. The primary mechanism for voltage transduction in these proteins involves the movement of a voltage-sensing domain (D), which opens a gate located on the cytoplasmic side. A distinct conformational change in the selectivity filter near the extracellular side has been implicated in slow inactivation gating, which is important for spike frequency adaptation in neural circuits. However, it remains an open question whether gating transitions in the selectivity filter region are also actuated by voltage sensors. Here, we examine conformational coupling between each of the four voltage sensors and the outer pore of a eukaryotic voltage-dependent sodium channel. The voltage sensors of these sodium channels are not structurally symmetric and exhibit functional specialization. To track the conformational rearrangements of individual voltage-sensing domains, we recorded domain-specific gating pore currents. Our data show that, of the four voltage sensors, only the domain IV voltage sensor is coupled to the conformation of the selectivity filter region of the sodium channel. Trapping the outer pore in a particular conformation with a high-affinity toxin or disulphide crossbridge impedes the return of this voltage sensor to its resting conformation. Our findings directly establish that, in addition to the canonical electromechanical coupling between voltage sensor and inner pore gates of a sodium channel, gating transitions in the selectivity filter region are also coupled to the movement of a voltage sensor. Furthermore, our results also imply that the voltage sensor of domain IV is unique in this linkage and in the ability to initiate slow inactivation in sodium channels.

electrophysiology | Nav1.4 | outer pore conformation

Voltage-gated sodium channels, like the other constituents of the voltage-gated ion channel superfamily, undergo gating transitions that result in channel opening and inactivation. Structural studies mainly on potassium selective ion channels have established that the opening of the channel primarily involves splaying of the bundle crossing formed on the intracellular side by the S6 helices (1–3). C-type inactivation gating, which follows channel opening and is mechanistically distinct from fast N-type inactivation, involves conformational rearrangements in the selectivity filter region (4–8).

According to our current thinking, inactivation gating at the selectivity filter region arises as a result of structural changes that follow pore opening and are not due to the movement of voltage sensors (9–11). Indeed, gating currents, which are a measure of voltage-sensing charge movements, are recorded under conditions with altered C-type inactivation (12) or with blockers that interact intimately with groups at the ion binding site (13, 14). For instance, the gating currents in exemplar Shaker potassium channels are typically measured in the background of the W434F mutation which results in a permanently C-type inactivated channel (15, 16). Although it has been suggested that the slow C-type inactivation may be coupled to voltage sensor movement,

the evidence to date has been indirect and somewhat confounding. Charge–voltage ($Q-V$) curves of Shaker potassium channel were found to undergo time and voltage-dependent shifts, which correlated with development of C-type inactivation (17). Signals from fluorescent probes attached to the S3–S4 loop of the Shaker potassium channel develop an additional component, which also correlates with C-type inactivation (18). However, this may not reflect coupling between voltage sensor and outer pore because these fluorescent probes are bulky and, in all likelihood, directly sense the conformation of the outer pore. Recent studies have suggested that the holding-potential–dependent shift in the $Q-V$ curve is an intrinsic property of a voltage sensor and may not necessarily reflect coupling to the outer pore. Charge–voltage curves of an isolated voltage sensor lacking a pore domain (D) were found to exhibit these charge shifts (“mode shifts”) upon changing the holding potential (19). Moreover, a noninactivating mutant of the human ether-a-go-go related gene (HERG) channel (20) and noninactivating hyperpolarization-activated cyclic nucleotide-gated (HCN) channels (21) both show holding-potential–dependent shifts in $Q-V$ curves. Finally, this correlation between charge shift and C-type inactivation has been reexamined by other groups who have concluded that these two processes are not coupled in the Shaker potassium channel (22).

The goal of this study is to test whether, in a voltage-gated sodium channel, the outer pore conformation is energetically coupled to the movement of a voltage sensor. As in the tetrameric potassium channels, the voltage-gated sodium channels also have two gates that control ion flux through the central pore. Studies with accessory subunits (23, 24) and local anesthetics (25, 26) provide reassuring evidence that the primary pore gate is on the cytoplasmic side of the channel. The cytoplasmic region of the $\beta 4$ subunit behaves as a classical open pore blocker (1). It blocks the conductance of the open pore in a state-dependent manner and prevents the deactivation of the sodium channel gates (23). The recently solved landmark structure of a prokaryotic sodium channel, NaVab also shows rather strikingly that the intracellular pore helices can gate the ion conducting pathway (27). Several groups have proposed that the selectivity filter region of the sodium channel is the site for a second gate, which is involved in slow and ultraslow inactivation gating (28–36). In many respects, slow inactivation gating in the sodium channels is similar to C-type inactivation (4, 5, 28, 34, 35). Although residues

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magnitude of current was reduced by as much as 65%. Moreover, a further increase in toxin concentration (up to 10-fold) did not increase inhibition of gating pore currents, suggesting that this effect is saturated (Fig. S2). We can eliminate the possibility that the hyperpolarization activated currents through DIV-CN are due to ion flux through the central pore as opposed to being gating pore currents because the central pore currents remain completely blocked by TTX (Fig. S1). Therefore, we consider the possibility that TTX is somehow modifying the currents fluxing through gating pore within the DIV voltage sensor.

First, we examined the possibility that TTX binds directly to the voltage sensor in the DIV-CN mutant. We introduced a pore mutation at a site known to be critical for TTX binding (Y401S) in the background of the DIV-CN mutant. The Y401S mutation, located next to the selectivity determining DEKA (D400, E755, K1237, A1529) locus (48–52), yields a channel that is nearly insensitive to block by TTX (53, 54). In this background (Y401S DIV-CN), TTX had no effect on gating pore currents through the DIV voltage sensor (Fig. S3). Therefore, TTX likely influences the gating pore currents through the DIV voltage sensor by binding to the outer pore of the sodium channel.

Next, we considered the possibility that TTX interacts directly with the charged voltage sensor while being docked to its binding site in the outer pore. The highly charged pore blocking peptide, μ -conotoxin (μ -CTX), has been previously suggested to shift the voltage dependence of activation by a small degree by modifying the surface charge (55). Our prediction will be that μ -CTX, which has a nominal net charge of +6 at neutral pH, will have a much larger effect on gating pore currents if these interactions are determined solely by electrostatics (56). We observed that the saturating concentrations of μ -CTX cause a much smaller reduction (16% at most) in the DIV-CN gating pore currents (Fig. 3 and Fig. S1). These findings and others (57) support the notion that TTX perturbs the DIV voltage sensor through a mechanism that is unlikely to involve direct electrostatic interaction.

A plot of residual gating pore currents versus voltage reveals that μ -CTX inhibition of the gating pore currents is voltage independent, whereas TTX block is pronounced at depolarized potentials (Fig. S4). The voltage dependence of the block is not due to differences in voltage-dependent binding because these measurements were obtained at saturating concentrations of TTX (central pore currents remain fully blocked). Therefore, this voltage dependence reflects the effect of TTX on the gating transition of the voltage sensor and shows that TTX binding favors the activated state of the DIV voltage sensor.

TTX Modulates Off-Gating Currents. Whereas gating pore current data suggest that TTX specifically modulates the gating pore currents and by extension the voltage sensor movement, we sought to validate this observation independently. Alternate approaches to monitor conformational rearrangement of the voltage sensors involve measuring gating currents or fluorescence signals from site-specific probes. Gating currents are the most direct measure of voltage sensor movements but recording them in absence of pore blockers remains technically demanding.

To eliminate contamination due to ionic currents, we measured the off-gating currents in the wild-type channels after a depolarizing pulse sufficient to induce fast inactivation. Permeant ions in the external solution were substituted by *N*-methyl *D*-glucamine (NMG) to eliminate any ionic current contamination of tail currents. Robust nonlinear charge movements were recorded from the wild-type channel without pore blockers (Fig. S5) and, as expected, the charge–voltage curves obtained from these measurements were saturated at extreme potentials. A comparison of the off-gating currents before and after addition of saturating concentrations of TTX shows that the total off-

gating charge is reduced by 24% (Fig. 4). Note that a substantial shift in the Q – V curve to a more hyperpolarized potential in the presence of TTX would result in subtraction of some of the nonlinear charge component and will be manifested as a reduced gating charge. As opposed to TTX, μ -CTX caused no reduction in the total gating charge, consistent with its effect on gating pore currents through the DIV voltage sensor (Fig. 4).

The above differences in perturbation of voltage sensor movements by the two toxins may reflect differences in their binding site or their state dependence. Remarkably, both TTX and μ -CTX can be simultaneously accommodated in the pore and, furthermore, TTX can be trapped in its binding site by μ -CTX (58, 59). The mutagenesis data also suggest that the critical residues implicated in TTX binding is in the narrow ion access pathway in the selectivity filter region (53, 54, 60). Because TTX causes a reduction in the total charge of the off-gating currents and influences the behavior of the DIV voltage sensor to a much greater extent than μ -CTX, we suggest that μ -CTX may be a more appropriate tool for collecting sodium channel gating currents. In summary, these gating current experiments establish that the TTX binding in the selectivity filter region of the sodium channel influences the movements of one or more voltage sensors of the sodium channel.

Cross-Bridging in the Selectivity Filter Affects the DIV Voltage Sensor.

One possible explanation for such cross-talk is that the central pore toxins introduce a nonnative interaction that perturbs the conformation of the DIV voltage sensor. Alternatively, this long-range interaction may reflect an inherent allosteric coupling between the outer pore and this particular voltage sensor (61).

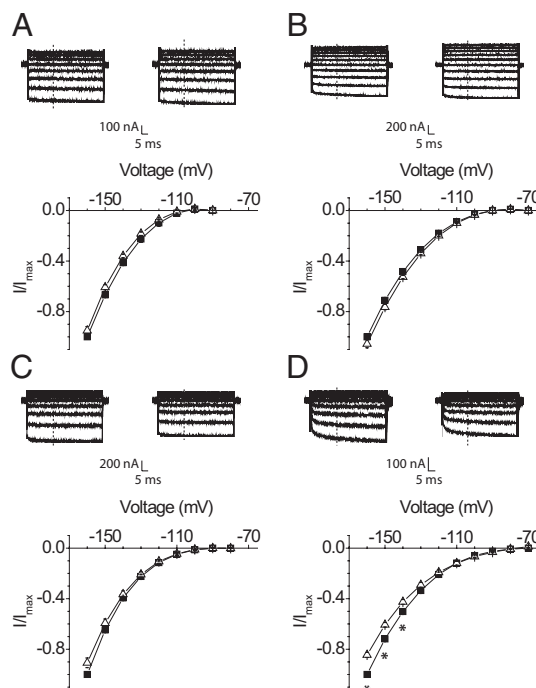


Fig. 3. Effect of μ -CTX on gating pore currents through the individual voltage sensors. A family of hyperpolarization activated currents before (Left; filled square) and after (Right; unfilled triangle) addition of μ -CTX from the DI-CN (A), DII-CN (B), DIII-CN (C), and DIV-CN (D) mutants (Upper). Normalized current–voltage plots of the same (Lower). Currents were recorded and normalized following the protocol described in Methods. Currents at each voltage were measured at the end of 20 ms, marked as a dashed line. Each plot represents the mean \pm SE of at least three independent experiments. * P value < 0.05 , statistical significance.

concurrently with canonical voltage sensor pore coupling. These mechanistic differences in coupling may underpin functional specialization of sodium channel domains (43, 69–72) and may be a universal feature of other nonsymmetric voltage-dependent ion channels.

Methods

Molecular Biology and Expression. Mutations were made using the QuikChange mutagenesis kit (Stratagene). cRNA was made using the mMessage mMachinE T7 kit (Ambion) after digesting the plasmids containing the sodium channel gene with NotI enzyme. For heterologous expression, 50 nL of cRNA containing a mixture of α - and β 1-subunit in 1:1 stoichiometry was injected into stage IV or V *Xenopus laevis* oocytes. After injection, the oocytes were kept at 18 °C in 1 × standard oocyte solution with Ca^{2+} and gentamicin for 2–5 d before recording.

Electrophysiology. Recordings were made using the cut-open oocyte voltage clamp configuration with CA-1B amplifier (Dagan) as described previously (73). Gating and ionic currents were sampled at 250 kHz with a Digidata 1400 interface (MDS Analytical Technologies) and low-pass filtered at 20 kHz. The oocytes were clamped at a holding potential of –80 mV for at least 5 min before recording. Current–voltage curves for DI, DIV-CN, and DIV-CN Y401S were obtained in 105 mM Na^+ -Mes (methanesulfonic acid) external solution, whereas those for DII-CN and DIII-CN mutants were obtained in an external solution containing either 52.5 mM Na^+ -Mes or 105 mM K^+ -Mes, respectively. Following a 50-ms prepulse to –120 mV, the currents were elicited by test pulses ranging from –80 to +70 mV for a duration of 60 ms. Gating pore currents were recorded in 105 mM Na^+ -Mes for both DI and DIV-CN, whereas the gating pore currents for DII-CN and DIII-CN were recorded in 105 mM K^+ -Mes external solution with 10 mM Hepes and 2 mM $\text{Ca}(\text{OH})_2$ at pH 7.4. Unless stated otherwise, these recordings used NMG-Mes internal

solution with 105 mM NMG, 10 mM Hepes, and 2 mM EGTA at pH 7.4. Gating pore currents were elicited by test pulses from –160 to –70 mV for 60 ms following a 50-ms prepulse to –120 mV. The currents were normalized with respect to the maximum current measured before toxin, typically at –160 mV. Channels were blocked using 0.6–1.1 μM of TTX or 6–25.7 μM of $\mu\text{-CTX}$ in the external solutions. Gating pore currents were recorded without a P/4 subtraction protocol, which means that the activation kinetics were likely to be contaminated by capacity transients. Offline leak subtraction was used to eliminate the linear leak component in experiments where it was significant (74).

To record the off-gating currents, we applied a prepulse to +50 mV for 30 ms to inactivate the channels. This was followed by a test pulse to potentials ranging from –130 mV to +50 mV for 20-ms durations. These experiments were performed in NMG-Mes internal and external (105 mM NMG-Mes with 10 mM Hepes, and 2 mM $\text{Ca}(\text{OH})_2$) solutions. To obtain gating currents in presence of toxin, 0.6–1.1 μM TTX or 6–25.7 μM $\mu\text{-CTX}$ was added to the external solution. At these concentrations, ionic currents were fully blocked. The charge–voltage curves for the off-gating currents were fit to a single Boltzmann equation: $Q/Q_{\text{max}}(V) = 1/(1 + \exp(-ze(V - V_{1/2})/KT))$ to estimate the $V_{1/2}$ (voltage for half-maximal activation) and z (valence) values. Linear capacity transients were eliminated by P/–4 subtraction at a holding potential of –130 mV.

For cross-linking experiments, the oocytes expressing mutants were treated with 2 mM H_2O_2 to maximize the disulfide formation. To measure currents under reducing conditions, 1 mM of freshly made DTT was added to the external solution containing 105 mM Na^+ -Mes.

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