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A Method for Simultaneous Echo Planar Imaging of Hyperpolarized ¹³C Pyruvate and ¹³C Lactate

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Abstract

A rapid echo planar imaging sequence for dynamic imaging of [1-¹³C] lactate and [1-¹³C] pyruvate simultaneously was developed. Frequency-based separation of these metabolites was achieved by spatial shifting in the phase-encoded direction with the appropriate choice of echo spacing. Suppression of the pyruvate-hydrate and alanine resonances is achieved through an optimized spectral-spatial RF waveform. Signal sampling efficiency as a function of pyruvate and lactate excitation angle was simulated using two site exchange models. Dynamic imaging is demonstrated in a transgenic mouse model, and phantom validations of the RF pulse frequency selectivity were performed.

Keywords

EPI; Hyperpolarized; pyruvate; spectral spatial RF; C13; kinetics

1. Introduction

Recent studies have shown the great potential of hyperpolarized ¹³C pyruvate for the in vivo monitoring of cellular metabolism and the characterization of disease [1–11]. Dynamic nuclear polarization (DNP) of ¹³C labeled pyruvate followed by rapid dissolution generates an injectable contrast agent with a four order-of-magnitude signal enhancement [12] enabling the measurement of the spatial distribution of pyruvate and its conversion to lactate, alanine, and bicarbonate. The conversion of pyruvate to lactate catalyzed by the enzyme lactate dehydrogenase is of particular interest, as the kinetics of this process have been shown to be sensitive to the presence and severity of disease in preclinical models [1–8, 13–15].

Acquiring hyperpolarized ¹³C images presents numerous challenges. Typically, multiple frequency components must be imaged in a time period short compared with the spin-lattice

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Several novel techniques have been developed to rapidly acquire ¹³C images. A small-angle band-selective excitation radio frequency (RF) pulse [16, 17] followed by a train of rapid imaging gradient pulses such as echo planar imaging (EPI) [18] or spiral readouts [19] allows for rapid imaging of single frequencies. The transmitter frequency is offset periodically for each temporal phase, and dynamic images of multiple frequency components can be acquired in an interleaved manner. Gradient echo images acquired at multiple echo times can also be used to resolve multiple frequency components via least squares inversion of the image phase [20–22]. Although the long scan time from three dimensional phase encoding typically precludes three dimensional spatially-resolved dynamic spectroscopic imaging, echo planar spectroscopic imaging (EPSI) readouts [23] and compressive sensing acquisitions [24, 25] can sufficiently increase acquisition speed to enable dynamic imaging.

In this study, we developed an echo planar imaging sequence for dynamic imaging of [1-¹³C] lactate and [1-¹³C] pyruvate. Frequency separation was achieved by spatial misregistration in the phase-encoded direction by the appropriate choice of echo spacing. Suppression of the pyruvate-hydrate and alanine resonances is achieved through an optimized spectral-spatial RF waveform. The signal dynamics were simulated using two site exchange models for the pyruvate-to-lactate conversion. These models highlight the benefit of applying differential flip angles to pyruvate and lactate, and optimal combination of these angles was studied along with the signal tradeoffs. Dynamic imaging was demonstrated in a transgenic mouse model, and phantom validations of the RF pulse and imaging gradient frequency selectivity were performed. Given that this method separates frequencies using the low effective bandwidth in the phase-encoded EPI direction, imaging of resonances with small frequency separations would be challenging. This is due to the fact that echoes must be spaced far apart to generate an adequate spatial shift. However, this technique is more readily applicable to imaging two compounds with large chemical shift dispersions since the phase-encoded dimension can be aliased multiple times while still generating a small imagedomain shift. The only limitation is that the frequency-encoded field of view must large enough to encode the shifts in the frequency-encoded dimension.

2. Theory

2.1. Frequency Separation

Two noninteracting spins that are encoded by an echo-planar readout will be spatially shifted based on their difference in chemical shift. The spins have Larmor frequencies ω_1 and ω_2 with a chemical shift dispersion $\Delta \omega = \omega_2 - \omega_1$. The signal is spatially encoded after RF excitation using an flyback EPI gradient waveform with N_{int} interleaves and N_{etl} equally spaced echoes (with spacing t_{esp}). For $N_{int} > 1$, the echo train is shifted in subsequent repetitions by t_{esp}/N_{int} to mitigate ghosting errors [26]. With the transmitter and receiver set at the mean of the two frequencies and phase encoding along the y dimension, the twodimensional k space image is

$$F(k_x, k_y) = F_1(k_x, k_y) e^{-i\varphi(k_y)} e^{-i\Delta\omega T E_1/2} + F_2(k_x, k_y) e^{+i\varphi(k_y)} e^{+i\Delta\omega T E_1/2}$$
(1)

where $F_1(k_x, k_y)$ and $F_2(k_x, k_y)$ are the images of ω_1 and ω_2 , and

$$\varphi(k_y) = \frac{t_{esp} \Delta \omega k_y}{2N_{int} \Delta k_y}.$$
(2)

Here, $\Delta k_y = 1/FOV_y$ is the spacing of k_y lines. TE_1 is the distance from the midpoint of the RF excitation to the midpoint of the first gradient readout lobe. The phase factors $\exp(\pm i\Delta\omega TE_1/2)$ are independent of k_y and can be excluded from further analysis. The Fourier transformed image is

$$f(x, y) = f_1(x, y - y_s) + f_2(x, y + y_s),$$
(3)

where

$$y_s = \frac{\varphi(k_y)}{2\pi k_y} = \frac{t_{esp} \Delta \omega F O V_y}{4\pi N_{int}}.$$
(4)

We define the quantity $\Phi = y_s/FOV_y$, the fraction of the field of view the image is shifted. Since y is a phase-encoded dimension, this shift is periodic ($\Phi + N_{alias} = \Phi$, with N_{alias} an integer). Substituting into (4), this relation gives a compact form for the echo spacing in terms of the other imaging factors:

$$t_{esp} = \frac{4\pi (\Phi + N_{alias})N_{int}}{\Delta \omega}.$$
(5)

If the transmitter is at the true mean frequency, then the shift of each image will be equal and in opposite directions as shown by (3). If the resonance frequency if shifted, the images will be shifted asymmetrically about the center, but their relative shift with respect to each other will be the same.

With the appropriate choice of echo spacing, we can resolve two frequencies spatially in the phase-encoded direction. With $\Phi = 1/4$ and FOV_y set to at least twice the object's extent in y, equations (4) and (5) imply that images of ω_1 and ω_2 are generated side by side in y. N_{alias} can be treated as a design parameter: for large $\Delta \omega$, we can let $N_{alias} > 0$ to maintain a reasonably short t_{esp} while still generating side by side images. For pyruvate and lactate at 3T, $\Delta \omega \approx 2\pi \times 392$ Hz, so allowing for aliasing is not necessary.

Although the acquisition bandwidth of the frequency-encoded direction is much greater than that of the phase-encoded dimension, these shifts can be corrected for since the images of the two frequencies are not spatially overlapping. For larger echo spacings, the readout bandwidth is lowered to increase the sampling duty cycle and SNR [27], so characterization of the shift in the frequency direction is important. With the transmitter set at the mean of the two frequencies and the frequency encoding along *x*, the shift is

$$\Delta x = \frac{\Delta \omega FOV_x}{4\pi BW_{read}} \tag{6}$$

where FOV_x and BW_{read} are the field of view and the readout bandwidth in the x dimension respectively. Since a filter is applied in this dimension to eliminate aliasing, the field of view

must be large enough to account for the shifts. If L_x is the spatial extent of the object in x, then

$$FOV_x \ge L_x + 2\Delta x. \tag{7}$$

Provided this condition is met, the frequency-encoded shift can be accounted for by shifting the images by $\pm \Delta x$ from (6) similar to the technique used in [28].

This simultaneous encoding method differs from scanning individual frequencies one at a time in several important respects. Since FOV_y must be twice the object's spatial extent, the minimum number of EPI phase encodes per scan is doubled from acquiring a single frequency. The total scan time, however, is roughly constant since the minimum number of scans is halved. The number of RF excitations applied at each frequency is doubled with simultaneous acquisition suggesting an analysis of the effects of depolarization from RF saturation and cumulative signal encoded as a function of excitation angle imparted to each frequency.

2.2. Simulation of Excitation Angle Effects on Exchanging Hyperpolarized Spins

The sequence implemented uses a spectral spatial RF pulse [29] to suppress the alanine and pyruvate-hydrate signal while allowing for different excitation angles imparted to pyruvate and lactate. Qualitatively, a small flip angle on pyruvate is expected to yield a larger cumulative signal on lactate since depolarization is reduced; this effect is shown empirically in [17]. To quantify the tradeoffs in applying differential flip angles to lactate and pyruvate, a simulation was developed to study the depolarization effects from RF saturation as well as the net signal encoded as a function of flip angle. The simulation was based on a two-site exchange model for pyruvate and lactate after infusion into the animal [9, 11, 13, 14]. This model is determined by k_{PI} which characterizes the conversion rate from pyruvate to lactate, k_{LP} which determines the rate in the reverse direction, and $\rho_P = 1/T_{1,P}$ and $\rho_L = 1/T_{1,L}$, the spin lattice relaxation rates for pyruvate and lactate respectively. θ_P and θ_L are the flip angles for pyruvate and lactate, and their effect on the longitudinal polarization is accounted for by a quasi-continuous approximation to the actual discrete effect. Unlike the methods presented in [30–32] where an optimal flip angle schedule is derived to maximize the integrated signal intensity of hyperpolarized spins, we only analyze the time-independent flip angle scenario. This is due to the fact that a time-dependent signal weighting could preclude quantitative analysis of kinetic data.

Let $M_{z,P}(t)$ and $M_{z,L}(t)$ be the longitudinal magnetization components of lactate and pyruvate respectively. The state vector $M_z(t)$ is defined as

$$M_{z}(t) = \begin{pmatrix} M_{z,P} \\ M_{z,L} \end{pmatrix}.$$
(8)

Combining with a continuous model for RF saturation and T_1 decay with the chemical dynamics, the state space model can be represented as

$$\frac{d}{dt}M_z(t) = AM_z(t),\tag{9}$$

with

Reed et al.

$$A = \begin{pmatrix} -\rho_{P} - k_{PL} + \frac{1}{TR}(\cos\theta_{P} - 1) & k_{LP} \\ k_{PL} & -\rho_{L} - k_{LP} + \frac{1}{TR}(\cos\theta_{L} - 1) \end{pmatrix},$$
(10)

where *TR* is the repetition time. The solution to (9) is a matrix exponential that can be calculated using inverse Laplace transforms:

$$M_{z}(t) = e^{At} M_{z}(0)$$

= $\mathcal{L}^{-1}[sI - A]^{-1} M_{z}(0).$ (11)

Here, s is the complex frequency variable, I is the 2×2 identity matrix, and

$$M_z(0) = \begin{pmatrix} 1\\0 \end{pmatrix} \tag{12}$$

The cumulative signal for pyruvate and lactate can be modeled as

$$S_{p} = \sum_{k=1}^{N_{acqs}} \sin\theta_{p} \cdot M_{z,p}(k \cdot TR)$$
(13)

$$S_{L} = \sum_{k=1}^{N_{acqs}} \sin\theta_{L} \cdot M_{z,L}(k \cdot TR).$$
(14)

The initial conditions represented by equation (12) imply that the pulse sequence begins immediately after substrate infusion, and the function describing the time dependence of the substrate delivery is ignored. This could be accounted for as in [14] but was not performed here since we are only concerned with getting an estimate of signal intensity as a function of flip angle.

3. Methods

3.1. Simulations

The signal of the hyperpolarized pyruvate and lactate as acquired with this sequence were simulated in Mathematica (Wolfram Research, Champaign, II, USA). The solution (11) to equation (9) and the signal intensities (14) and (13) were computed as functions of both θ_P and θ_L at $k_{PL} = .01 \text{ s}^{-1}$, $.05 \text{ s}^{-1}$, and $.09 \text{ s}^{-1}$. These values of the forward rate constant represent the range of previously measured kinetic parameters for transgenic mice [14]. Backward conversion from lactate to pyruvate was assumed to be negligible ($k_{LP} = 0$). Relaxation times of $T_{1,P} = T_{1,L} = 25$ s were used for the simulations in Figure 2. These values were chosen as rough estimates and are similar to the previously quoted values of $T_{1,L} = 21$ s and $T_{1,P} = 32$ s measured from slices encompassing the tumors of TRAMP mice [14] and $T_{1,L} = T_{1,P} = 33$ s measured from spatially-resolved saturation transfer data in a mouse [9].

Since the $T_{1,L}$ and $T_{1,P}$ values are widely variable and cannot be measured beforehand, the non-optimality of the integrated signal as a function of flip angle and relaxation time was also simulated. The ratio of the integrated signal S_L acquired using a $\theta_L = 15^\circ$ flip angle to

that acquired using the T_1 -optimal flip angle was simulated in Figure 3a. The optimal flip angle is shown in Figure 3b.

The mean repetition time was used in simulations:

$$TR = \frac{N_{int}TR_{act} + t_{delay}}{N_{int}},$$
(15)

where TR_{act} is the actual repetition time used in the scan and t_{delay} is the delay period between dynamic acquisitions. N_{acqs} was set to $N_{int} \times N_{phases} = 4 \times 20$, where N_{phases} is the number of dynamic acquisitions.

3.2. Radio Frequency Excitation

A spectral-spatial pulse giving a nominal excitation flip angle of 7.5° to pyruvate, 15° to lactate, and 0° to both alanine and pyruvate hydrate was designed for this sequence (Figure 4). This choice of flip angles was based on the simulation results of the cumulative lactate signal as a function of both θ_P and θ_L given the repetition time for the sequence (see Figure 2). The spectral filter was designed using a spectral factorization/convex optimization for reduced peak B_1 and minimum sampling rate [17, 33]. Spatially-selective sublobe samples were played using a 0.5 VERSE fraction [34], and the irregular $t - k_z$ sampling was corrected for using the least-squares algorithm presented in [18]. The selectivity of the RF pulse was verified in a proton phantom using body coil for excitation (Figure 4).

3.3. Acquisition Parameters

This sequence was demonstrated in a single slice in the coronal plane. The phase encoding was set along the left/right dimension, and frequency was set along superior/inferior. The phase field-of-view FOV_y was set to at least twice the extent of the object. For mouse imaging, this was always possible using $FOV_y = FOV_x = 10$ cm. The transmitter resonance was set at the mean of the *in vivo* resonant frequencies of $[1^{-13}C]$ pyruvate and $[1^{-13}C]$ lactate, corresponding to $\omega/2\pi = 32.1315$ MHz on the 3T scanner. The echo spacing was set for a $\pm FOV_y/4$ shift ($\Phi = 1/4$) so that the sequence would generate pyruvate and lactate images side by side in *y*. The pyruvate/lactate frequency separation at 3T is $\Delta\omega/2\pi = 384$ Hz. A 32×32 grid was acquired in plane using $N_{etl} = 8$, $N_{int} = 4$, so from equation (5), $t_{esp} = 5.19$ ms. The readout bandwidth was then reduced to ± 5 kHz increase the sampling duty cycle and SNR. With the given resolution and readout bandwidth, a shift of 0.6 times the pixel size is expected in the readout (S/I) direction. This effect was neglected, but could be compensated for with an image domain shift in S/I. The total repetition time TR = 80 ms giving a single slice scan time of $N_{int}TR = 240$ ms. See Figure 1 for a diagram explaining the timing parameters.

3.4. Animal Experiments

Experiments were performed on a GE 3T Signa MRI system (GE Health-care, Waukesha, WI, USA). A custom dual-tuned ${}^{1}\text{H}/{}^{13}\text{C}$, transmit/receive mouse coil was used for carbon and proton imaging. Prior to infusion of the hyperpolarized contrast, the transmitter attenuation was calibrated by obtaining images at multiple attenuation settings from a syringe containing 0.7 mL of 8 M ${}^{13}\text{C}$ -urea placed next to the animal. 24 uL of 99 % [1- ${}^{13}\text{C}$] pyruvate was mixed with the OX063 radical (GE Healthcare, Oslo, Norway) and polarized in an Oxford Instruments Hyperpsense polarizer (Oxford Instruments Biotools, Oxford, UK) operating at 3.35 T and 1.2 K. After approximately one hour of 94.114 GHz microwave irradiation, the sample was rapidly dissolved in 4.25 mL of a heated TRIS/NaOH solution giving a final concentration of 80 mM. A transgenic prostate tumor bearing mouse [35] was injected with 350 μ L of the hyperpolarized solution over 8 seconds. Imaging was started 20

seconds after the beginning of injection. A 3 second delay period is inserted between slices giving a temporal resolution of $t_{delay} + N_{int}TR = 3.2$ seconds.

3.5. Reconstruction

The pulse sequence produces the pyruvate and lactate images side by side in the left/right dimension. The image was shifted by $+FOV_y/4$ and multiplied by an ideal filter with width $FOV_y/2$ to generate a pyruvate image with proper spatial localization. The same process was repeated for lactate images, but the shift was in the opposite direction.

4. Results and Discussion

Figure 2 shows plots summarizing the simulation results. The integrated lactate signal intensity from Figure 2a has a maximum at approximately $\theta_L \approx 15^\circ$ and is approximately independent of k_{PL} for the range of values simulated. As a function of θ_P , S_L is a monotonically decreasing function (Figures 2b, 2d) as expected since we assumed one-way conversion. Since the pyruvate signal is generally much stronger than lactate, θ_P was set to 7.5°. According to the model, this would sacrifice roughly 30% of S_P (2.5/3.5) versus encoding with the θ_P optimizing S_P ($\approx 15^\circ$ as shown in Figure 2c). However, S_L is expected to increase 30% (1.8/1.4) from encoding using $\theta_P = 7.5^\circ$ instead of 15° (Figure 2b). Figure 3 shows the simulated signal acquired if θ_L was set based on a $T_{1,L}$ differing from the actual value. The relative S_L is within 10% of the maximum for $T_{1,L}$ values down to 9 seconds and is even more stable for $T_{1,L} > 25$ s. Variations in $T_{1,P}$ had less then a 5% influence on the relative signal for the range of $T_{1,P}$ (5 s–50 s) simulated. Variations in k_{PL} did not strongly influence the non-optimality of the integrated signal nor the optimal θ_L for different relaxation values.

Figures 5, 6, and 7 show preliminary data acquired from infusion of hyperpolarized pyruvate into a TRAMP mouse. Figure 5 shows the coronal spin echo localizer (left) and a representative pyruvate/lactate image (right). Figure 6 shows anoverlay of the color-mapped SNR values of the hyperpolarized images. For ease of viewing, only every other dynamic repetition is displayed. Figure 7 shows the average SNR of all the pixels within two ROIs drawn around the prostate tumor and the thorax of the mouse. This figure highlights how multiband excitation with the EPI sequence generates comparable SNR for lactate and pyruvate (≈ 25 in the prostate tumor and ≈ 9 in the thorax).

Some limitations of the pulse sequence presented are discussed. Ghosting artifacts stemming from interleaving strongly degrades the images in this simultaneous encoding strategy since these artifacts appear along the same dimension in which the different frequencies are resolved. Although this effect was largely mitigated by the echo shifting strategy, some persistent ghosts were visible in early time points. This was likely due to the fact that in the early time points, the contrast is largely concentrated in the vasculature, and the echo shifting cannot fully correct the phase errors from flowing spins. Setting t_{esp} to a value larger than the minimum does make the sequence more susceptible to geometric distortion from B_0 homogeneity. However, the value used (5.2 ms) is on the order of the minimum achievable spacing by the clinical scanner using whole-body gradients (≈ 3 ms) for the same field-of-view and resolution. Combined with the low gyromagnetic ratio of carbon, we do not expect geometric distortions to be detrimental, and they could be compensated for using a B_0 field map and the methods presented in [36]. The transverse dynamics of this sequence including T_2^* effects were not investigated, and some signal loss by the long TE_1 due to the spectral-spatial RF pulse is to be expected. While using spectral-spatial excitation is necessary for encoding lactate and pyruvate, this pulse sequence should show the greatest advantage for encoding only two frequencies. In this case, a short pulse with a broad

excitation band could be used for excitation, and TE_1 could be made extremely short, allowing for a longer echo train and signal enhancement with reduced T_2^* loss.

5. Conclusion

Simultaneous encoding of pyruvate and lactate using a flyback echo planar sequence is demonstrated. This method spatially resolves the two frequencies by deliberate chemical shift misregistration based on proper selection of the echo spacing. Robust performance of the pulse sequence was observed in phantom and in vivo hyperpolarized experiments. Simulations of the combined exchange, excitation, saturation, and relaxation effects of hyperpolarized compounds were implemented. These simulations show promise in determining optimum flip angle for the metabolic product as well as elucidating the tradeoff in product signal encoded versus substrate flip angle. Possible drawbacks of the pulse sequence presented as well as some interesting extensions are discussed.

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Highlights

- EPI sequence for simultaneous imaging of two hyperpolarized compounds was developed
- Two site exchange models were used to optimize flip angles.
- In vivo validation was performed with 13C pyruvate in a transgenic mouse.

Reed et al.



Figure 1.

RF and gradient waveforms in a single repetition. $t_{esp} = 5.2ms$, $TE_1 = 13.5$ ms, TR = 80 ms. To increase the sampling duty cycle, the readout bandwidth was lowered to ± 5 kHz corresponding to a readout gradient strength of 1 G/cm.

Reed et al.



Figure 2.

Total encoded signal of pyruvate (S_P) and lactate (S_L) given by equations (13) and (14) respectively. Simulations were computed assuming $T_{1,L} = T_{1,P} = 25$ s. S_L and S_P both show maxima as functions of θ_L and θ_P , respectively (a,c). However, the lactate signal S_L decreases monotonically with increasing the pyruvate flip angle θ_P (b,d) suggesting the use of a small θ_P value as to not diminish S_L . For the given repetition times and assumed relaxation rates, the optimal θ_L was approximately 15° and was largely independent on the conversion times (a).



Figure 3.

a) the ratio of S_L calculated for a lactate flip angle of $\theta_L = 15^\circ$ to the S_L value computed with the optimal angle for the given lactate relaxation time $T_{1,L}$. The optimal flip angle is shown in b). Both plots assume $T_{1,P} = 25$ s. The acquired S_L using a 15° flip is within 10% of the optimum value for $T_{1,L}$ values down to 9 seconds.



Figure 4.

(a) RF waveform (top) and slice select gradient (bottom). The total waveform duration is 16.9 ms. (b) spectral/spatial excitation profile from Bloch simulation (top), measured excitation profile from a proton phantom (center), and simulated spectral profile at z = 0 shown on a semilog scale (bottom). The spectral locations of pyruvate, alanine, pyruvate-hydrate, and lactate are denoted. The minimum full width half max of the spatial profile is 1.5 cm.

Reed et al.



Figure 5.

Coronal FSE localizer (left) and a representative image produced by the pulse sequence (right).

Reed et al.



Figure 6.

Every other dynamic repetition overlaid on a coronal FSE localizer. Pyruvate and lactate SNR values represented as overlays. Concordant with prior studies in the TRAMP model, the tumor shows persistent lactate signal at the later time points.



Figure 7.

Coronal FSE localizer on the left shows the locations of the two ROIs. The mean SNR of all the pixels within each ROI is plotted at each repetition on the right. The comparable SNR values between pyruvate and lactate in the early time points highlight the benefit of the multiband excitation pulses.