

Dynamics of inherently bounded histone modification domains

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Contributed by Gerald R. Crabtree, July 2, 2012 (sent for review April 8, 2012)

A central goal of chromatin biology is to reveal how posttranslational histone marks modulate gene expression; however, relatively little is known about the spatial or temporal dynamics of these marks. We previously showed that a dynamic model of histone mark nucleation, propagation, and turnover fits the mean enrichment profiles from 99% of noncentromeric histone H3 lysine 9 trimethylation (H3K9me3) domains in mouse embryonic stem cells without the need for boundary or insulator elements. Here we report the full details of this “inherently bounded” model of histone modification dynamics and describe several dynamic features of the model using H3K9me3 as a paradigm. By analyzing the kinetic and structural constraints that drive formation of inherently bounded domains, we find that such domains are optimized when the rates of marking and turnover are comparable. Additionally, we find that to establish such domains, propagation of the histone marks must occur primarily through local contacts.

epigenetics | heterochromatin | methylation

Chromatin is a dynamic environment subject to spatial and temporal regulation through a number of factors. A central focus of this regulation is the posttranslational modification of histone proteins, whose modification states are associated with a variety of transcription states at a given genetic locus. Unfortunately, there are few quantitative models regarding the dynamics of histone marking and therefore little basis for describing the structure, stability, or fluctuations of histone modification domains.

One particular posttranslational histone modification, histone H3 lysine 9 trimethylation (H3K9me3), has served as a paradigm to study how histone marks are propagated in live cells to induce local silencing. In mammalian cells, H3K9me3 is associated with heterochromatic and transcriptionally silent regions (1–4). H3K9me3-based repression requires Heterochromatin Protein 1 (HP1), which mediates *cis*-spreading of the H3K9me3 mark by binding existing H3K9me3 sites (5–7), oligomerizing to bind neighboring nucleosomes (8, 9), and recruiting H3K9-specific histone methyltransferases (HMTs) (10–13) to induce outward spreading of the mark.

Unless opposed, this positive feedback would expand H3K9me3-containing domains without bounds. Several groups have described “insulator elements” in flies, yeast, and other organisms, which prevent further expansion of H3K9me3 marking (14). The presence of these insulator elements gives rise to distinct chromatin boundaries, which suggests chromatin may be packaged into modular, independent structural domains near these elements (15). Additionally, in fission yeast, the mating-type locus is silenced through expansion of H3K9me3 over an ~20-kbp domain with relatively sharp borders (16, 17). The distinct structure of this histone modification domain, with a large plateau and definite borders, indicates that the expansion of H3K9me3 marks is blocked or subject to increased turnover at the border of the domain, consistent with a boundary element.

In mammals, however, such expansive H3K9me3 plateaus in coding regions are rare. In mouse embryonic stem cells, the vast majority of noncentromeric H3K9me3 domains span less than 10 kbp, with most less than 5 kbp (18, 19). These domains also typically show peaked patterns of enrichment, rather than broad,

flat plateaus (20). We recently showed that small-molecule-mediated recruitment of HP1 α to the Oct4 (Pou5f1) locus gives rise to a symmetric and peaked H3K9me3 domain in both mouse ES cells and fibroblasts (20). This domain grew continuously over time to between 2 and 10 kbp while maintaining its peaked shape, allowing us to obtain kinetic measurements on the dynamics of H3K9me3 propagation *in vivo*. Additionally, the enrichment profile and dynamics of this domain were consistent with those of H3K9me3 domains throughout the mouse genome. In our previous work, we showed that this synthetic H3K9me3 domain and >99% of all euchromatic H3K9me3 domains were well described by an “inherently bounded” model of histone modification dynamics (20), where propagation of the H3K9me3 mark was opposed by random turnover.

In this model, histone modification domains like H3K9me3 do not require insulator or boundary elements because the spreading of the mark is inherently constrained by mark and/or histone lifetime. Because the average enrichment profiles of >99% of mammalian H3K9me3 domains are well described by this model, we suggested that these domains are in fact inherently bounded and that our model provides quantitative predictions into the kinetics and dynamics of H3K9me3 marking throughout the genome. Here, we focus on the kinetic and structural constraints that lead to the characteristic pattern of peaked mark enrichment within the domains. We also report the full details of our dynamic model, using H3K9me3 as a paradigm.

Results

Reaction Scheme and Assumptions. We model the dynamics of an individual genetic locus as a discrete one-dimensional lattice. We consider that each lattice position corresponds to an individual nucleosome, analogous to the “beads on a string” picture of chromatin. The lattice $\{I_j\}$ spans from $j = -128$ to $j = +128$, which corresponds to $256 + 1$ nucleosomes or 51.4 kbp of DNA if nucleosomal spacing is 200 bp. To make our model as general as possible, we consider that each position j along the lattice may occupy one of two states: an unmodified ($I_j = 0$) state or a modified ($I_j = 1$) state. Briefly, three features define the standard model:

- i) Nucleation: The origin (position 0) corresponds to a target site that is uniquely able to nucleate H3K9me3 modification. This nucleation occurs at rate k_+ .
- ii) Propagation: Our general kinetic scheme integrates the series of processes that propagate H3K9me3 into a single net propagation rate (k_+). In our kinetic model, k_+ describes the net rate of H3K9me3 addition at nucleosomes immediately adjacent to H3K9me3-marked sites. In this way, spreading of H3K9me3 along the chromosome occurs through linear propagation of the mark to neighboring nucleosomes, consistent

Author contributions: C.H. and G.R.C. designed research; C.H. performed research; C.H. and G.R.C. analyzed data; and C.H. and G.R.C. wrote the paper.

The authors declare no conflict of interest.

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This article contains supporting information online at www.pnas.org/lookup/suppl/doi:10.1073/pnas.1211172109/-DCSupplemental.

with proposed spreading via HP1 oligomerization (5, 6, 8, 21, 22). Mechanistically, this propagation likely requires multiple steps (e.g., recruitment of HP1, oligomerization, recruitment of HMT, enzymatic methylation, etc.); however, for simplicity, we treat propagation as subject to a single rate-limiting step with rate k_+ .

- iii) Turnover: Unlike propagation of the mark, turnover is equally likely everywhere, and thus k_- describes the stochastic turnover of H3K9me3 at any marked nucleosome. Mechanisms of mark turnover include enzymatic catalysis (e.g., demethylation), histone turnover, or any other process that results in removal of the mark. For simplicity, we treat turnover as subject to a single rate-limiting step with rate k_- .

We developed a simple probabilistic description of histone marking based on the kinetics of nucleation, propagation, and turnover. Monte Carlo simulations based on this model allowed us to determine the kinetic and structural constraints that promote formation of inherently bounded domains.

Implementation. The stochastic trajectory of a given genetic locus is carried out by iteration of the following processes at each time step (Fig. 1):

- Nucleation:** If the target site is unmarked ($I_0 = 0$), it is converted to a marked state ($I_0 = 1$) with probability α_+ (where $\alpha_+ = k_+ \Delta t$).
- Propagation:** For each marked nucleosome j on the lattice, if the $j-1$ position is unmarked ($I_{j-1} = 0$), it is converted to a marked state ($I_{j-1} = 1$) with probability α_+ . Similarly, if the $j+1$ position is unmarked ($I_{j+1} = 0$), it is converted to a marked state ($I_{j+1} = 1$) with probability α_+ . Mark propagation thus proceeds outward in both directions from each marked site.
- Turnover:** For each marked nucleosome j on the lattice, conversion to an unmarked state ($I_j = 0$) occurs with probability α_- (where $\alpha_- = k_- \Delta t$).
- Time evolution:** Simulation time t is incremented by Δt .

The above procedure treats propagation and turnover as homogeneous Poisson processes with exponentially distributed

lifetimes. We also allow all simulations to evolve using periodic boundary conditions; therefore, there is no boundary implicit in the simulations.

Constraints on the Establishment of an Inherently Bounded Domain.

When these processes are allowed to evolve, a stochastic domain of histone marks is established at the target site (Fig. 1D). To characterize which regions of the parameter space support the establishment of an inherently bounded H3K9me3 domain, we held k_- constant and varied k_+ between $0.1k_-$ and $3.0k_-$. We chose values of k_+ , k_- , and Δt such that $\alpha_- = 0.05$, and α_+ varied between 0.005 and 0.15. We found it convenient to define the ratio

$$\kappa = \frac{k_+}{k_-}$$

to convey the relative rate of mark propagation and turnover. By allowing the model to evolve constrained by different values of κ , we discovered that an inherently bounded domain was established only when $\kappa \leq 1.5$ (Fig. 2). At values of $\kappa > 1.5$, H3K9me3 marks spread without bounds (Fig. 2D and H and Fig. S1). Here the simulations undergo a second-order phase transition (Fig. 3A) to a new state where marks occupy the entire locus.

Intensity and Specificity of Marking. Although the H3K9me3 domains established according to our model are, on the average, symmetric and centered around the target site, the relative propagation rate κ plays an important role in determining the overall mark density and spatial specificity of the domain. At low values of κ , the average domain is not very densely marked but is extremely specific, because the only marking takes place at or very near to the target site. As κ increases, the target site is more and more likely to be marked (Fig. 3B).

Conversely, the overall specificity of the marking decreases when κ increases, because the domain spreads outward from the target site. To quantify the effect of κ on the overall specificity of the domain, we introduce a specificity score S ,

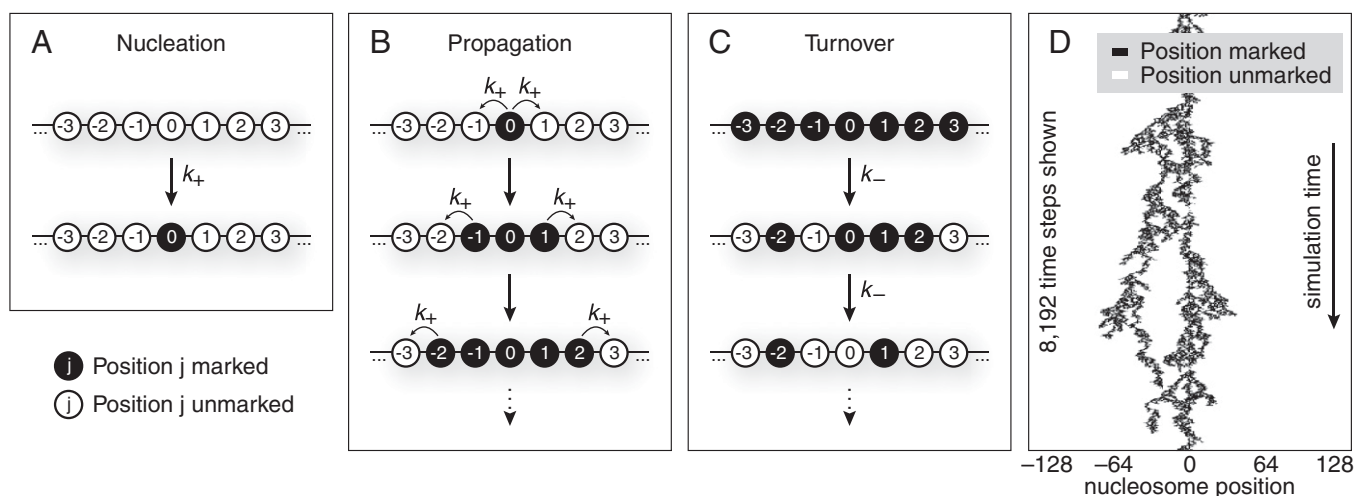


Fig. 1. An inherently bounded histone modification domain arises from three processes. (A) A target site is marked at rate k_+ , which nucleates the histone modification domain. (B) Neighboring unmarked nucleosomes are marked at rate k_+ to propagate the domain. (C) Random turnover occurs at rate k_- and is equally likely everywhere. Turnover reverses the marks placed during nucleation and propagation. When all three processes occur simultaneously, an inherently bounded steady-state histone modification domain arises. (D) Simulation by iterating these three processes results in a stochastic, bounded steady-state modification domain. Shown here are 8,192 time steps from a simulation where the relative propagation rate $\kappa = k_+/k_- = 1.5$. In an individual simulation (analogous to an individual locus in a cell), the marking can transiently bifurcate to form multiple smaller domains. Accumulation of marks is limited by turnover.

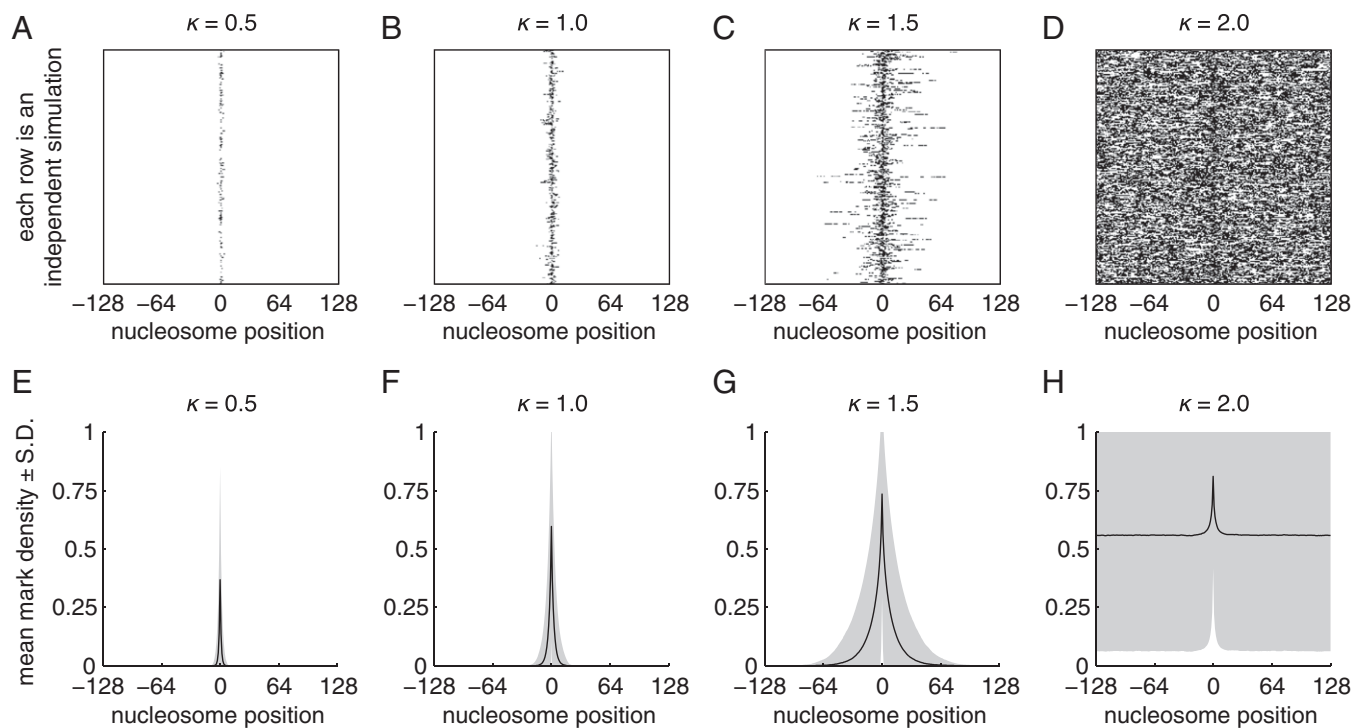


Fig. 2. The relative propagation rate κ determines the spatial extent of inherently bounded domains. (A–D) Snapshots from simulations obtained using different values of κ . Each row represents the state of an individual trajectory at a single time step. Black pixels indicate that the position is marked, and white pixels indicate the position is unmarked. When $\kappa \leq 1.5$, a localized steady-state H3K9me3 domain arises that is inherently limited by mark and/or histone turnover and does not spread. At values of $\kappa > 1.5$, marking spreads to occupy the entire simulated space. (E–H) For each value of κ , the mean mark density is plotted at each position (\pm SD in gray). For A–H, the value of κ is given above each plot. See also Fig. S1.

$$S = \frac{\langle I_{\text{target}} \rangle - \langle I \rangle}{\langle I_{\text{target}} \rangle},$$

where $\langle I_{\text{target}} \rangle$ is the mean mark density at the target site, and $\langle I \rangle$ is the mean mark density over the entire lattice. Because the

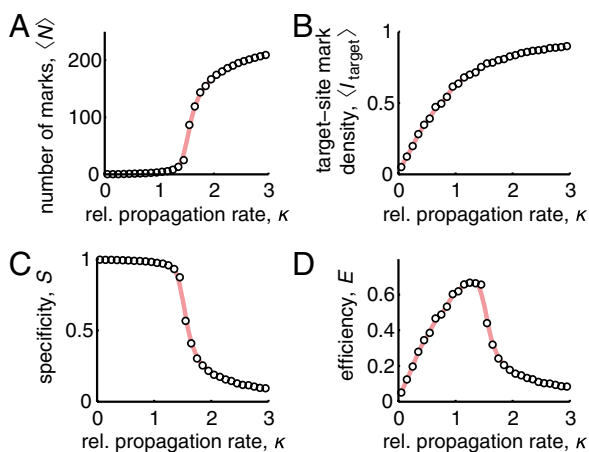


Fig. 3. Structural features vary with the relative propagation rate. (A) As the relative propagation rate κ increases, the mean number of marks increases. Above $\kappa > 1.5$, the number of marks undergoes a second-order phase transition and marks occupy the entire locus. (B) The mean mark density at the target site monotonically increases with the relative propagation rate κ . (C) The specificity of marking S (defined in the main text) is sharply reduced when the relative propagation rate $\kappa > 1.5$. (D) The efficiency score, which is a product of target-site mark density and specificity (see main text), shows a peak between $1 \leq \kappa \leq 1.5$.

target site always has the maximum mean density throughout the lattice, this score returns a value of 0 when marking is perfectly uniform across the locus (no specificity) and returns a value approaching 1 when marking occurs exclusively at the target site (perfect specificity). When $\kappa \leq 1.5$, the domain shows a high degree of specificity (Fig. 3C). The simulations undergo a phase transition for values of $\kappa > 1.5$, into a new state that shows very poor specificity. In this state, H3K9me3 marks spread without bounds.

We next considered an “efficiency” measure E , which weighs the contribution of target site-specific histone mark density and specificity by multiplying these two factors together:

$$E = \langle I_{\text{target}} \rangle S = \langle I_{\text{target}} \rangle - \langle I \rangle.$$

Conveniently, this measure also corresponds to the contrast between the mean mark density at the target site and the mean mark density over the entire locus. We found that marking efficiency is maximized when $1.0 \leq \kappa \leq 1.5$ (Fig. 3D). Thus, histone mark density and specificity are jointly optimized when the rates of mark propagation and turnover are coupled such that κ falls between 1.0 and 1.5.

Dynamics and Stability of Inherently Bounded Domains. An important aspect of our model is that each simulation undergoes fluctuations in the number and distribution of marks throughout its trajectory. To characterize these fluctuations and to examine the stability of inherently bounded domains, we investigated several dynamic properties of the domains under varying values of κ . To focus on the dynamics of bounded domains, we restrict the remainder of our analysis to values of $\kappa \leq 1.5$. To address the role of κ on the dynamics of the domains, we examined how κ affected fluctuations in the total number of marks $N(t)$:

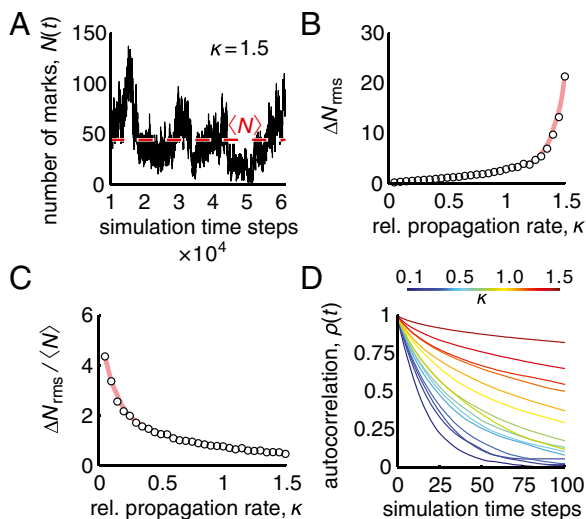


Fig. 4. Dynamic features vary with the relative propagation rate. (A) The total number of marks $N(t)$ fluctuates throughout each individual simulation about the average value $\langle N \rangle$. (B) The absolute root mean-squared (rms) fluctuations in $N(t)$ increase monotonically as κ increases. (C) When rms fluctuations are normalized by the average number of marks $\langle N \rangle$, the relative rms fluctuations decrease, indicating that higher values of κ lead to greater stability on a relative basis. (D) Autocorrelation of $N(t)$ shows that domains with higher values of κ maintain correlation over longer times, providing stability and memory of the domain over a longer time.

$$N(t) = \sum_j I_j(t).$$

The total number of marks $N(t)$ is not static, but constantly fluctuates about its mean value $\langle N \rangle$ (Fig. 4A). We characterized the scale of domain size fluctuations by calculating the root mean-squared (rms) fluctuation of $N(t)$ under different values of κ :

$$\Delta N_{\text{rms}} = \sqrt{\frac{1}{T} \sum_{t=1}^T (N(t) - \langle N \rangle)^2}.$$

The absolute scale of these fluctuations increases as κ rises (Fig. 4B). However, when normalized by the mean number of marks $\langle N \rangle$, these relative rms fluctuations instead decrease with increasing values of κ (Fig. 4C). Therefore, although an increased relative propagation rate κ induces larger fluctuations in absolute numbers of marks, these fluctuations are disproportionately smaller when normalized by the overall size of the domain.

To examine the stability and memory of the domain, we calculated the autocorrelation function of $N(t)$ in trajectories with different values of κ . The autocorrelation function $\rho(t)$ was calculated using the definition (23)

$$\rho(t) = \frac{\sum_{k=1}^{T-t} (N(k) - \langle N \rangle)(N(k+t) - \langle N \rangle)}{\sum_{k=1}^T (N(k) - \langle N \rangle)^2},$$

where t is in units of time steps and T is the total number of time steps used for calculating autocorrelation. We found autocorrelation rose monotonically with increasing values of κ (Fig. 4D); thus, raising the relative propagation rate causes the domains to maintain correlation with their current states over a longer time, providing a degree of stability and memory. We conclude that increasing the relative propagation rate contributes to the stability of the domain by reducing the relative scale of the fluctuations, as well as contributes to

memory of the domain over a longer time by increasing autocorrelation.

Model Extension: Varying Nucleation Rate. An extension of the standard model above is to allow the nucleation rate to vary independently of the propagation rate (i.e., allowing nucleation to occur at rate k_v rather than k_+). To examine this, we extended the standard model by replacing step i with alternative step ia .

ia) Nucleation: If the target site is unmarked ($I_0 = 0$), it is converted to a marked state ($I_0 = 1$) with probability α_v (where $\alpha_v = k_v \Delta t$).

Similar to κ , we defined $\kappa_v = k_v/k_+$. By allowing κ and κ_v to vary independently, we found that bounded domains are established across a wide range of nucleation rates (Fig. S2). However, when mark nucleation is much slower than turnover (i.e., near $\kappa_v \approx 0$), domains are not established with high efficiency. Additionally, because nucleation of the mark is a rate-limiting step for establishing a modification domain, the nucleation rate also controls the rate at which the domain is established de novo (Fig. S3). Thus, the primary effects of the nucleation rate are to control the marking density at the target site and the timescale over which inherently bounded domains may be established.

Model Extension: Marking Beyond Nearest Neighbors. Although marks may propagate only through nearest-neighbor contacts in the standard model above, several lines of experimental evidence suggest that discontinuous spreading may also occur (reviewed in ref. 24). To address whether inherently bounded domains may arise through marking at a distance, we have extended the standard model to allow one of two different modes of histone marking beyond nearest-neighbor contacts: a geometric dependence with respect to distance or a power-law dependence with respect to distance. We focus here on the geometric dependence, which may occur, for example, if HP1 oligomers behave as a simple one-stranded equilibrium polymer (25, 26) when extended beyond H3K9me3 marked sites and histone methyltransferases were recruited to the ends of HP1 oligomers. In our implementation, this dependence is treated by replacing step ii with alternative steps ii_a and ii_b .

ii_a) Propagation (left): For each marked nucleosome j on the lattice, integer distance $d_a \in [1, 256]$ is drawn from a geometric distribution with probability $P(d_a) = (1-p)^{d_a-1}p/D$. If the $j-d_a$ position is unmarked ($I_{j-d_a} = 0$), it is converted to a marked state ($I_{j-d_a} = 1$) with probability α_+ .

ii_b) Propagation (right): Similarly, for each marked nucleosome j on the lattice, integer distance $d_b \in [1, 256]$ is drawn from a geometric distribution with probability $P(d_b) = (1-p)^{d_b-1}p/D$. If the $j+d_b$ position is unmarked ($I_{j+d_b} = 0$), it is converted to a marked state ($I_{j+d_b} = 1$) with probability α_+ .

Note that in both *ii_a* and *ii_b*, the normalization constant

$$D = \sum_{i=1}^{256} (1-p)^{i-1} p$$

normalizes probability such that the sum over the entire locus is unity. Here, p describes the probability of marking the nearest-neighbor site, and the probability of mark propagation decays geometrically with increasing distance d . In this model, there are two floating parameters, κ and p . To describe the stability of inherently bounded domains using these two parameters, we constructed heat maps of mean target-site mark density (I_{target}), specificity S , and the efficiency score E as a function of both κ and p (Fig. 5).

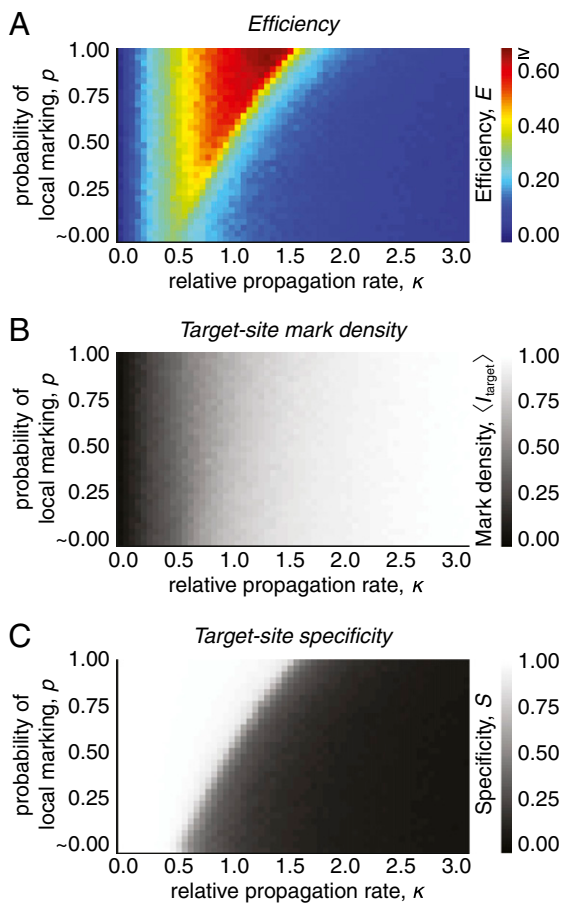


Fig. 5. Local marking increases the robustness of inherently bounded domains. (A) We extended our standard model of marking to include marking at a distance. The probability of marking follows a geometric dependence as described in the main text. Inherently bounded domains are formed most efficiently when marking occurs primarily at the nearest neighbor (near $p = 1$). As p is decreased, the efficiency of the domain is decreased and the rates that support domain formation are sharply reduced. (B) Average mark density of the target site depends on the relative propagation rate κ , but does not show a strong dependence on local marking. (C) Specificity of marking shows a strong dependence on local marking. As marking becomes more local (near $p = 1$), specificity is maintained over a broader range of κ . Because of this dependence, domains are established with high target-site mark density and specificity only when marking is primarily local.

Importantly, the efficiency score E shows that only certain regions of κ and p can support the formation of densely marked domains with high specificity (Fig. 5A). At one extreme, at the limit when $p \rightarrow 0$ and marking is equally likely everywhere along the locus, inherently bound domains largely do not form under any values of κ because marking is too diffuse. At the other extreme, when $p = 1$, the behavior of the system is identical to that described in the standard model above with propagation restricted only to nearest neighbors. Here as before, a specific and densely marked domain is established between $1.0 \leq \kappa \leq 1.5$ (Fig. 5A). Between these two extremes, we find that the values of κ that support formation of a densely marked domain with high specificity are gradually broadened as p is increased (i.e., as marking becomes more local).

To sort out the effects of marking at a distance on target-site mark density and specificity, we separately examined the dependence of these parameters on κ and p . Target-site mark density was only weakly dependent on the probability of nearest-neighbor marking, but displayed a strong dependence on the

relative propagation rate κ (Fig. 5B). On the other hand, specificity showed a mixed dependence on both κ and p (Fig. 5C): At low values of p , when marking is diffuse over the entire region, a high level of specificity was maintained up to $\kappa \approx 1$; above this value specificity was sharply reduced. As p increased, specificity was maintained over a broader range of up to $\kappa \approx 1.5$ as marks remained largely constrained to the narrow area around the target site. Thus, local marking positively contributes to the overall specificity of the domain. These results indicate that, although inherently bounded domains may form when marking is extended beyond nearest-neighbor histones, it is primarily local marking that drives formation of these domains.

Similar results were obtained when using a power-law dependence of marking at a distance (*SI Text* and Fig. S4). We conclude that, regardless of the mathematical details regarding the distance dependence of marking, inherently bounded domains are established primarily through local contacts. In particular, local marking contributes to the specificity of these domains by constraining histone marks to the vicinity of the target site. Nevertheless, our model provides a mechanism for the appearance of discontinuous marking. In an individual cell, discontinuities transiently arise even when propagation occurs exclusively through local contacts (Fig. 1D). Such discontinuities arise when turnover occurs in an area between two marked regions. Under these circumstances, the mark may continue to be propagated as two separate domains for some time. Additionally, in our simulations, “local” means nearest-neighbor nucleosomes; however, in chromatin, where physical proximity may be maintained through stable looping or other structures, marking may be transmitted across a loop, giving rise to the discontinuities previously described (24).

Discussion

Despite abundant experimental data regarding the genomic distributions of histone marks, there are surprisingly few physical models that predict the spatial distribution and dynamics of these marks. One earlier model of H3K9me3 marking at the yeast mating-type locus rationalizes how memory of the mark may be maintained over cell generations (27); however, the spatial distribution anticipated by that model is inconsistent with the genomic distribution of noncentromeric H3K9me3 marks in mammalian cells (20). We speculate that this discrepancy may reflect important differences in the way that H3K9me3 marks are propagated in yeast and mammals within nonrepetitive coding regions.

In our previous work, we showed that our standard model of inherently bounded histone marking—presented in detail here—is consistent with the local distribution of H3K9me3 marking at >99% of noncentromeric H3K9me3 domains throughout the genome in mouse embryonic stem cells (20). Additionally, when the accumulation of mark density predicted by the model is fitted to experimental results, we obtained specific values for k_+ and k_- , both rates on the order of $0.1\text{--}0.2\text{ h}^{-1}$. The values of k_- that we obtained in mouse embryonic fibroblasts and embryonic stem cells fall between two different experimental measures of histone turnover in live cells (20, 28, 29). On the basis of the reasonableness of these values, the degree of similarity in the shapes of H3K9me3 enrichment profiles, and the striking predictive power throughout the genome, we propose that most mammalian H3K9me3 domains arise through the general mechanisms described by our model. The simple processes of nucleation, local propagation, and histone turnover are sufficient to drive the formation of dynamic H3K9me3 domains that are peaked in their mean genomic profiles and are of the appropriate length scales. Importantly, our model predicts that inherently bounded domains have continuously decaying H3K9me3 profiles when averaged over an ensemble of cells, rather than a plateau with sharp edges that would arise through a well-positioned boundary

or insulator element. Although boundary elements have been described in several specific instances (30), we propose that the majority of mammalian H3K9me3 domains may not require such elements because the propagation of the domains is inherently limited by the rate of mark and/or histone turnover. In this context, turnover of the mark may occur through a variety of mechanisms. We expect that enzymatic demethylation and displacement of core histones are the primary contributors to mark turnover, although the dynamics described by our model are general and consistent with these and other mechanisms.

Multiple lines of experimental evidence suggest an ongoing competition between enzymatic methylation and demethylation throughout the genome. The H3K4me3 and H3K9me3 demethylase Lsd1 appears at both active and silenced genes and plays an important role during differentiation by decommissioning enhancer sites that have accumulated these marks (31). In addition, depletion of the H3K9-specific demethylases Jmjd1a and Jmjd2c leads to accumulation of H3K9me3 and differentiation of ES cells (32). Both of these results indicate that demethylases are continually opposing the continued placement of H3K9me3. Thus, H3K9me3 marks are governed not simply by the enzymatic placement of the mark, but instead by the continual steady-state balance between the placement and removal of the marks. Reduction of demethylase activity pushes the balance toward propagation of the mark, resulting in increased accumulation of H3K9me3. Importantly, this steady-state balancing act is integral to the maintenance of cell identity and pluripotency.

This balance is reflected in our model by the relative propagation rate κ . We find that when propagation of a histone mark is limited to local *cis*-spreading, formation of a robust, bounded steady-state histone modification domain is optimized when κ is between 1.0 and 1.5. In fact, this is the range of κ that also best fits experimental data from mouse embryonic stem cells and fibroblasts (20). This relatively narrow range suggests that such inherently bounded domains arise only when the rates of marking (k_+) and turnover (k_-) are coupled. It should not be surprising that these processes have comparable rates: No histone modification

domains would ever arise if turnover were much faster than marking, but on the other hand, the genome would quickly saturate with histone modifications if marking were much faster than turnover. Similarly in position-effect variegation of the *Drosophila* eye, the addition or loss of a single allele of a euchromatic or heterochromatic factor can shift the balance from a variegated phenotype to either a fully expressed or a silenced state at the *white* locus (33). This exquisite sensitivity indicates that heterochromatin dynamics are finely tuned to the dosage/activity balance of chromatin factors. On the basis of this argument, we speculate that the factors responsible for k_+ and k_- likely coevolved such that κ is maintained in this range.

This general model abstracts over much of the complexity that is already known or that remains an outstanding question about H3K9me3 biology. In particular, our model does not address the mechanisms by which these marks are transmitted during or after replication, although we have previously shown that straightforward variations of the model can provide the means for epigenetic memory of the mark (20). Other potential mechanisms for epigenetic memory include factors that promote continued recruitment of marking factors to the target site, such as antisense RNA-based targeting of H3K9me3 silencing factors (34, 35).

Our model of inherently bounded histone modification domains presents a challenge for direct observation of the chromatin states in individual cells over time. In vivo, the accumulation and removal of marks like H3K9me3 are relatively slow (20), indicating that single-cell assays should require observations over long periods of time with temporal resolution on the order of hours. As single-cell chromatin assays mature, quantitative models like the one presented here will provide a biophysical framework for studying chromatin regulation.

ACKNOWLEDGMENTS. We thank members of the G.R.C. laboratory, in particular O. Bell and N. Hathaway, for advice and helpful discussions. Research in the G.R.C. laboratory is supported by National Institutes of Health Grants HD55391 and NS046789 and the Howard Hughes Medical Institute. C.H. is supported by Eunice Kennedy Shriver National Institute of Child Health and Human Development Fellowship F32HD072627.

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