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# Different dissociation pathways. and observation of an excited deoxy state in picosecond photolysis of oxy- and carboxymyoglobin

(picosecond spectroscopy/myoglobin/hemoglobin/picosecond dynamics)

P. A. CORNELIUS, A. W. STEELE, D. A. CHERNOFF<sup>†</sup>, AND R. M. HOCHSTRASSER<sup>‡</sup>

Department of Chemistry, University of Pennsylvania, Philadelphia, Pennsylvania 19104

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ABSTRACT Picosecond transient absorption spectra of Mb, MbCO, and  $MbO<sub>2</sub>$  have been studied at time delays of up to 10 ns after excitation at 353 nm. Particular attention has been paid to the rapid spectral changes that occur in the Soret region during the first 50 ps in MbCO and  $MbO<sub>2</sub>$ . In MbCO both the bleaching of the Soret peak (feature I) and the appearance of new deoxy-like absorption (feature II) occur instantaneously, whereas in  $MbO<sub>2</sub>$ feature II is delayed with respect to feature I. A short-lived  $(\approx 12)$ ps) feature near 455 nm (feature III) was much more intense in MbO<sub>2</sub> than in MbCO and was also identified in the transient spectrum of Mb. No evidence of subnanosecond geminate recombination was found in either MbCO or MbO<sub>2</sub>. These observations are consistent with a scheme in which  $MbO<sub>2</sub>$  photodissociates through an excited state of Mb, whereas  $MbCO$  under the same conditions produces ground state Mb directly. The results and conclusions are compared with those of previous picosecond studies on these molecules and related hemoglobin derivatives.

Picosecond transient absorption spectroscopy is a valuable method for the study of heme proteins and their reactions with  $O<sub>2</sub>$  and CO. The unique feature of this technique is its ability to probe changes occurring on a time scale inwhich any largescale movements of the protein are unlikely. In the case of hemoglobin, this permits the observation of the primary processes that initiate the protein conformational changes associated with cooperativity. Understanding these primary processes requires a detailed knowledge of the electronic states of this system and the couplings and interactions between them. To this end, a great deal of experimental and theoretical work has been devoted to establishing an accurate mapping of the states ofvarious hemoglobin derivatives (1). We recently reported the results of a series of picosecond transient measurements on  $HbCO$  and  $HbO<sub>2</sub>$ , in which we showed that relaxation pathways derived from spin-orbit coupling arguments could be used to explain the observed spectral changes (2). That work showed how such considerations could account for the difference in the behavior of these two molecules:  $HbO<sub>2</sub>$  exhibited a fast transient (<90 ps) and subnanosecond geminate recombination, whereas in HbCO both of these effects were absent.

In order to examine further the scheme proposed in these hemoglobin studies, we have used our picosecond transient absorption spectrometer to study a closely related set of myoglobin derivatives: MbCO,  $MbO<sub>2</sub>$ , and unliganded Mb. Myoglobin has a similar electronic structure to hemoglobin but small spectral shifts result from the differences in protein environment. Because of its resemblance to a single hemoglobin subunit, myoglobin provides a natural model for the tetrameric protein, and it has been the object of previous kinetic studies in both picosecond (3, 4) and longer time regimes (5, 6). The

The publication costs of this article were defrayed in part by page charge payment. This article must therefore be hereby marked "advertisement" in accordance with 18 U. S. C. §1734 solely to indicate this fact. spectral and temporal information provided by our picosecond spectrometer makes possible a more detailed investigation of the primary photolytic processes in these molecules. We report here the observation and interpretation of spectra taken at time delays of up to 10 ns after excitation at 353 nm, with emphasis on the rapid spectral changes that occur during the first 50 ps.

## MATERIALS AND METHODS

Lyophilized sperm whale myoglobin (type II) was purchased from Sigma and used without purification. Reduction and oxygenation of the largely ferric form was accomplished with sodium dithionite on a Sephadex G-25 column (7). The thawed MbO<sub>2</sub> was diluted with  $0.025$  M Tris HCl buffer, pH 8.6 at 25°C, to produce samples with a Soret peak absorbance of  $\approx 0.9$ at 418 nm in <sup>a</sup> 1-mm quartz cell. For each of two independent preparations and runs, data collection was accomplished in a single day, the continuous wave (CW) spectra showing no evidence of MbO<sub>2</sub> degradation over this time. MbCO was prepared by passing CO over a rotating solution of freshly prepared MbO2 for <sup>1</sup> hr. The concentration was adjusted by.addition of the pH 8.6 Tris buffer saturated with CO to yield MbCO samples with absorbances of 0.81 and 1.2 at the Soret maximum  $(423 \text{ nm})$  as measured in a 1-mm quartz cell. Deoxymvoglobin (70  $\mu$ M) was prepared by adding a few grains of sodium dithionite to a 1-mm optical cell containing a freshly prepared  $MbO<sub>2</sub>$ sample. Human oxyhemoglobin, prepared by a standard procedure (8), was similarly treated to yield 70  $\mu$ M deoxyhemoglobin. Despite the dithionite absorbance in the near UV, blanks made in parallel with the deoxy samples were found to have no transient absorption signal upon excitation with 353-nm pulses. All samples were characterized in terms of purity and concentration by comparing CW spectra (Cary 14) with published spectra (9).

Our picosecond double beam transient absorption apparatus, designed to acquire transient measurements at delays ranging up to 10 ns after excitation, has been described in detail (2). Very briefly, a single 8-ps pulse is switched out of the pulse train generated by a passively mode-locked Nd/phosphate glass laser. This pulse is amplified to 20 mJ and then used for second and third harmonic generation; in this set- of experiments the UV pulse at 353 nm is used to excite the sample. The second harmonic can be delayed optically before it is focused into a cell of CC14, generating a weak continuum pulse of 8-ps duration. A beamsplitter selects <sup>a</sup> fraction of the continuum light as <sup>a</sup> reference beam, which traverses an unexcited region of the sample. The main continuum pulse propagates collinearly with the excitation pulse, both beams passing through a  $400-\mu m$ -diam-

Abbreviation: CW, continuous wave.

<sup>t</sup> Present address: The Standard Oil Co. (Ohio), 4440 Warrensville Ctr. Rd., Cleveland, OH 44128.

<sup>t</sup> To whom reprint requests should be addressed.

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eter pinhole located directly before the sample cell. The excitation pulse at 353 nm contains approximately 10  $\mu$ J of energy on entering the sample cell. The transmitted continuum, both "pinhole" and "reference" beams, provides two tracks of spectral data when imaged onto the slit of a low-dispersion spectrograph fitted with an optical multichannel analyzer vidicon detector.

In order to minimize the effects of shot-to-shot fluctuations in the continuum spectrum, we have employed the following data collection procedure throughout these studies: With the excitation pulse blocked, we collect a series of two-track continuum spectra. From this group we select and average five shots whose continuum spectra are similar. The reference track of each subsequent shot (with excitation present) is compared with the reference track of this average of five shots. If the spectral content differs beyond a certain tolerance, as determined by least squares analysis, the shot is rejected. In addition, we reject laser pulse trains containing satellite pulses and shots whose energy at <sup>353</sup> nm falls outside the typical range. We find that this data collection procedure results in reliable and reproducible spectra. All of the spectra presented here are averages of at least five laser shots and have an average noise level of  $\approx 0.01$  absorbance units in the region 430-500 nm, and  $\approx 0.03$ between 400 and 430 nm.

## RESULTS

 $MbO<sub>2</sub>$ . Excitation of  $MbO<sub>2</sub>$  results in the bleaching of the ground state Soret band and the production ofa new absorption feature near 440 nm. The bleaching feature is fully formed within 3 ps, whereas the new absorption does not achieve its maximal amplitude until a time delay of 15 ps. This difference cannot be accounted for by the effect of group velocity dispersion on the probing continuum, which might contribute a delay of at most 3 ps over this small wavelength region. Fig. <sup>1</sup> illustrates the evolution of the difference spectrum immediately after the initial buildup. At 13 ps the absorption shows a tail extending into the red, which has disappeared by 40 ps. Within our experimental error, this spectral change is not accompanied by recovery of the bleaching feature or increase of the main absorption peak at 440 nm. Fig. <sup>1</sup> suggests that the rapidly relaxing component is centered around 455 nm, and from our data we estimate the lifetime of this peak to be  $12 \pm 3$  ps. From 40 ps to 9.5 ns the shape of the transient difference spectrum undergoes no further change, while the amplitude decreases by at most 15%. While this is in accordance with recent studies



FIG. 1.  $MbO<sub>2</sub>$  transient difference spectra, Soret region, after ps excitation at 353 nm. These spectra show different rates of growth for excitation at 353 nm. These spectra show different rates of growth for<br>the absorption and bleaching features:  $\cdots$ , 0 ps;  $-\cdots$ , 3.3 ps;<br> $-\cdots$ , 6.6 ps;  $\cdots$ , 13.2 ps;  $-\cdots$ , 40.0 ps. Sample  $A_{418} = 0.90$  in 0.025 6.6 ps;  $\cdots$ , 13.2 ps; ——, 40.0 ps. Sample  $A_{418} = 0.90$  in 0.025 M Tris buffer, pH 8.6 at 23°C, 1-mm cell.



FIG. 2. MbCO transient difference spectra at  $6.6$  ps  $(......), 13.2$  ps (----), and 100 ps ( ). Absorption and bleaching features develop simultaneously. Sample  $A_{423} = 0.81$ ; excitation at 353 nm.

of  $HbO<sub>2</sub>$  recombination in the nanosecond regime (10), it indicates an absence of the significant (40%) amount of subnanosecond geminate recombination observed in HbO<sub>2</sub> excited with 530-nm ps pulses (2).

MbCO. Fig. <sup>2</sup> illustrates our finding that excitation of MbCO also leads to bleaching of the ground state Soret band and the appearance of new absorption near 440 nm. Spectra obtained at 0, 3.3, 6.6, and 13.2 ps indicate a simultaneous development of both bleaching and absorption intensities; these features are fully formed within the pulsewidth of the laser. During the first 100 ps after the buildup of the signal the new absorption band exhibits a spectral change similar to that in  $MbO<sub>2</sub>$ , as depicted in Fig. 2. The difference spectrum at 6.6 ps has absorption extending to the red that disappears rapidly, and the change in shape can be attributed to a broad, weak absorption centered at 455 nm. Because this effect is much smaller here than in  $MbO<sub>2</sub>$ , we were unable to estimate a lifetime for this band. The qualitative resemblance to  $\mathrm{MbO}_{2}$  is nevertheless clear from the figures. The shape that the MbCO transient difference spectrum assumes after the relaxation of the 455-nm component is subsequently unchanged up to 9.5 ns after excitation. The difference spectrum remains constant in amplitude, in contrast to MbO2, which suggests that recombination ofCO does not occur on this time scale.



FIG. 3. Short-lived difference spectrum of deoxy Mb  $(......)$  and difference-of-difference spectra (13 ps  $-$  50 ps) of MbO<sub>2</sub> (----) and<br>MbCO (-----) illustrating a common absorption feature near 455 nm -) illustrating a common absorption feature near 455 nm. The different amplitudes of these curves, which are normalized for  $\varepsilon_{353}$ , and their similarity in shape and lifetime ( $\tau = 12 \pm 3$  ps) suggest that an excited state of Mb is formed through different pathways in the photolysis of  $MbO<sub>2</sub>$  and  $MbCO$ .



FIG. 4. Diagram defining possible photolysis and relaxation pathways in MbCO and MbO<sub>2</sub>. MbX<sup>\*</sup> and MbX<sup>\*\*</sup> could represent bound, excited electronic states; Mb\* could be an electronically excited, unliganded myoglobin. Straight arrows represent transitions involving at least one photon; wavy lines are radiationless processes, which may include intermediates not shown. Check marks beneath each species indicate which spectral features are associated with it in our experiments. Our results suggest that in our experiments the dominant pathway in MbO<sub>2</sub> involves  $k_4$ ,  $k_5$ , or both and produces Mb<sup>\*</sup>, whereas in MbCO the dissociation is through  $k_3$  to ground state Mb.

Mb. The excitation of deoxygenated Mb produces <sup>a</sup> weak, short-lived transient difference spectrum. A bleaching feature at 433 nm and <sup>a</sup> new absorption peak at 451 nm both attain their maximal amplitudes ( $\pm \Delta A = 0.05$ ) at a delay of 10 ps (see Fig. 3). The relative prominence of the absorption peak relative to the bleaching in this spectrum indicates that the transient state produced has a strong Soret absorption. The transient decays rapidly and cannot be observed at 20 ps.

Summary of Results. The photolysis of MbX  $(X = CO \text{ or } O<sub>2</sub>)$ can be described in terms of three spectral features listed in Fig. 4. The bleaching at 420 nm (I) results from the disappearance of ground-state MbX. The absorption at 440 nm (II) is characteristic of deoxy Mb and monitors the production of that species. The rapidly relaxing absorption at 455 nm (III) may originate from an excited state of the deoxygenated molecule, denoted Mb\*. The basis for this assignment is discussed in the next section.

Fig. 4 also presents a diagram representing some possible pathways for excitation and relaxation in the MbX system. It should be emphasized that this diagram is shown here for didactic purposes and is not intended to exclude more complex or detailed descriptions. Transitions from one species to another in this diagram may involve many steps and include intermediates not shown; in some cases transitions may require absorption of at least one additional photon.

### DISCUSSION

In this section we examine the above experimental findings, together with previous work on myoglobin and hemoglobin, with a view to unifying the description of the available information. As a first step, we show in Fig. 5 the photoproduct spectra for both  $MbO<sub>2</sub>$  and MbCO photolysis. These have been constructed by using the transient difference spectrum at relatively long times (>50 ps) and the absorption spectrum of the starting material in each case, using the fractional photoconversion as an adjustable parameter. The photoproduct spectra are similar to each other though somewhat different from the CW spectrum of deoxygenated Mb. With respect to the spectrum of Mb the photoproduct spectra are slightly broader and weaker, and they show a 3-nm shift to the red. These results suggest that the photolysis products are the same in MbCO and  $MbO<sub>2</sub>$ , although different from stable deoxygenated Mb.



FIG. 5. Calculated photoproduct absorption spectra at  $t > 50$  ps for MbCO  $(\cdots \cdots)$  and MbO<sub>2</sub> (----) compared with the CW Soret absorption band of Mb (-).

From the figures it can be seen that all three of the molecules studied show a common spectral feature, denoted III in Fig. 4. To isolate this rapidly relaxing peak in the difference spectra of MbO<sub>2</sub> and MbCO, we subtracted the 50-ps spectrum from the 13-ps spectrum in each case. These curves are plotted in Fig. 3 on the same scale as the induced transient in Mb, with all three curves normalized for equal ground state absorbance at 353 nm. Because these curves show approximately the same peak position and lifetime, we have assumed that all three represent the same state, Mb\*.

The difference in amplitude among the curves in Fig. 3 is significant and must be incorporated into any model used to explain the data. In  $MbO<sub>2</sub>$  feature III is larger than when  $Mb$ is pumped directly, suggesting that the photolysis results in a very efficient production of Mb\*. In the terms of Fig. 4 this would entail a mechanism involving either  $k_4$  or  $k_5$ , and it would imply that the MbO<sub>2</sub> photoproduct rapidly reaches the excited state Mb\*. In MbCO feature III is weaker than it is in the direct pumping of Mb, which could be explained by a pathway through  $\bar{k}_3$  and  $k'_6$ . In this scheme the photoproduct appears in its ground state and is subsequently pumped to Mb\* by the absorption of a second photon from the excitation pulse.

The electronic structure of these molecules has been described in detail in the earlier hemoglobin work (1). In MbCO a lowest excited state is dissociative and can provide a simple pathway  $(k_3)$  for the photolytic production of ground-state Mb. In MbO<sub>2</sub> the presence of low-lying charge-transfer states complicates the kinetic picture, because these states could be weakly dissociative or even nondissociative. Therefore a major dissociative pathway in  $MbO<sub>2</sub>$  might not pass through the lowest-lying excited state. We suggest that in  $MbO<sub>2</sub>$  a multiplephoton mechanism operates that provides a significant amount of product through repeated excitation of a short-lived dissociative state MbX\*\* not reached by one-photon absorption. In Fig. 4 MbX\*\* is reached through  $k'_2$  and dissociates through  $k_5$ to produce a large amount of photoproduct in an intense short pulse even if  $k_2$  is relatively fast.

This interpretation would account for: (i) the low quantum yield for dissociation of  $MbO<sub>2</sub>$  under lower intensity illumination  $(11)$ ;  $(ii)$  the delayed appearance of the 440-nm absorption feature mentioned above, which implies the presence of a second step in the photolysis pathway; and *(iii)* the appearance of the initial photoproduct Mb in an excited state. This scheme suggests that the short pulse quantum yield of photolysis in MbO<sub>2</sub> would be wavelength dependent through the step  $k_2$ , which could account for the negligible quantum yield obtained by Shank et al.  $(12)$  in  $HbO_2$ , using 500-fs pulses at 615 nm.

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We first compare our present results with the previous picosecond studies of myoglobins by Rentzepis and coworkers (3, 4). In their experiments the excitation characteristics were similar to ours, whereas the probing process relied principally on kinetic measurements of the ground state Soret bleaching and transient Soret absorption at 420 nm and 440 nm, respectively. For this reason their experiments did not detect feature III at 455 nm; nonetheless it is possible to compare the three sets of results with regard to features <sup>I</sup> and II. In their earlier study the photodissociation of both  $MbO_2$  and  $MbCO$  was deduced to occur in less than 6 ps, and a reported  $125 \pm 50$  ps decay exhibited only by MbCO apparently was preparation related (3). In their more recent report (4), Reynolds et al focused more on early-time kinetics ( $t < 280$  ps) in samples at  $4^{\circ}$ C that were pumped at <sup>355</sup> nm and at <sup>530</sup> nm. They found both MbCO and  $\overline{M}bO_2$  exhibited an instantaneous bleaching, corresponding to feature I, and a new absorption feature II, delayed in its appearance by 3.5-5 ps. The authors accounted for this by proposing an intermediate state ( $\tau = 7{\text -}10$  ps) in the dissociation pathway of both MbCO and MbO<sub>2</sub>. The recent results and proposed explanations of Reynolds  $et$  al. (4) are consistent with our  $MbO<sub>2</sub>$  results reported here, but their MbCO results are at variance with ours. Thus their model accommodates neither our MbCO results nor our information regarding optically pumped Mb and feature III. Further work is required to examine quantitatively the possible effects of pulse width, pulse energy, experimental geometry, sample preparation, and sample conditions in order to explain the discrepancies between these studies. We note that Reynolds et al. found that, while their kinetic results were independent of excitation wavelength (530 or 355 nm), their quantum yield measurements implied that  $\phi_{355,530}^{\text{MbCO}}$   $\phi_{355}^{\text{MbO}_2}$   $\approx$  3 $\phi_{530}^{\text{MbO}_2}$ . Further careful determination of the wavelength dependence of the quantum yield of  $\mathrm{MbO}_2$  is necessary.

The results of the present studies also may be compared with the previous picosecond work on  $HbO_2$  and  $HbCO$  (1, 2). The following similarities exist: the photoproduct spectra for HbCO and  $HbO<sub>2</sub>$  resembled the spectrum of deoxygenated Hb but were broader and weaker. A fast transient similar to feature III in  $MbO<sub>2</sub>$  was seen in the  $HbO<sub>2</sub>$  difference spectrum. Recently we observed <sup>a</sup> rapidly relaxing transient difference spectrum, resembling feature III of Mb, in Hb pumped directly at 353 nm. The one major difference in the behavior of the Mb and Hb systems is the absence of subnanosecond recovery of the Soret bleaching in  $MbO<sub>2</sub>$ , which was observed in  $HbO<sub>2</sub>$ . In previous work with HbO<sub>2</sub> and HbCO we attributed the recovery of the Soret bleaching for  $HbO_2$  to geminate recombination of  $O_2$ . We suggested electronic reasons why CO might not be expected to recombine so rapidly, thereby improving its probability of escape from the protein. We have shown here that there is no significant recovery of the Soret bleaching in the cases of  $MbO<sub>2</sub>$ and MbCO. On the other hand, the electronic structures of the hemes in Hb are not expected to be significantly different from the structure in Mb. It is obvious then that arguments based on differences of the electronic structures of Hb,  $O_2$ , and CO cannot provide a complete explanation of the photokinetic behavior of  $HbO<sub>2</sub>$  and  $HbCO$ . These important parallels indicate that the scheme of Fig. 4 may also describe the photokinetic behavior of Hb and its derivatives. Our earlier results therefore provide evidence for a multiple-photon dissociative pathway to an electronically excited deoxy-like photo product in HbO<sub>2</sub> and for the relative unimportance of this path in HbCO.

A comparison can also be drawn between our work and recent time-resolved Raman studies of HbCO (13, 14), in which <sup>a</sup> flowing sample was irradiated with 30-ps pulses and the resulting spectra were recorded for different degrees of focusing. The

observed shift of an oxidation marker band was attributed to high-spin Hb, with the iron not fully relaxed out of the plane. This frequency shift persisted when the pulses were lengthened to <sup>20</sup> ns. This is consistent with our results on HbCO (2), in which a transient is generated within our pulsewidth and exhibits no subsequent spectral changes. However, the similarity in the peak position of feature III in MbO<sub>2</sub>, MbCO, and Mb suggests that the iron is out of the plane at the time this peak appears  $(<10 \text{ ps}$ ). If the motion of the iron had not yet occurred, <sup>a</sup> shift in peak position relative to equilibrium Mb would be expected. Our work has detected quite different behavior for  $H\dot{b}O_2$ , in which the dissociation apparently produces a form of Hb also as yet undetected by vibrational spectroscopy. Time resolution better than 12 ps will be needed to see these prompt species.

In conclusion, we consider first the dynamics immediately after excitation and then the results at longer delay times. The observation offeature III and its identification as an excited state of the deoxy species was made possible through careful study of transient spectra closely grouped around time zero. This feature seems to be characteristic of directly pumped Mb and was found to have an analog in Hb. It seems reasonable, then, to assign the fast-relaxing component  $(\tau < 90 \text{ ps})$  observed in HbO<sub>2</sub> to an excited state of Hb (2). It appears firmly established now that MbCO undergoes photolysis with high quantum yield and exhibits no geminate recombination phase in either the picosecond or the nanosecond time regime  $(15)$ . For MbO<sub>2</sub> there is no clear information available regarding the nanosecond time domain dynamics. Eisert et al. (3) found no relaxation in the first 450 ps for MbO<sub>2</sub> pumped at 530 nm, and we report here an upper limit of 15% recombination in 10 ns after 353-nm excitation. Further experiments are necessary to explore the excitation wavelength and heme environment dependence of subnanosecond geminate recombination.

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