

NIH Public Access

Author Manuscript

J Chem Inf Model. Author manuscript; available in PMC 2014 May 24.

Published in final edited form as:

J Chem Inf Model. 2013 May 24; 53(5): 1073–1083. doi:10.1021/ci300619x.

Development of the Knowledge-based & Empirical Combined Scoring Algorithm (KECSA) to Score Protein-Ligand Interactions

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Abstract

We describe a novel knowledge-based protein-ligand scoring function that employs a new definition for the reference state, allowing us to relate a statistical potential to a Lennard-Jones (LJ) potential. In this way, the LJ potential parameters were generated from protein-ligand complex structural data contained in the PDB. Forty-nine types of atomic pairwise interactions were derived using this method, which we call the knowledge-based and empirical combined scoring algorithm (KECSA). Two validation benchmarks were introduced to test the performance of KECSA. The first validation benchmark included two test sets that address the training-set and enthalpy/entropy of KECSA The second validation benchmark suite included two large-scale and five small-scale test sets to compare the reproducibility of KECSA with respect to two empirical score functions previously developed in our laboratory (LISA and LISA+), as well as to other well-known scoring methods. Validation results illustrate that KECSA shows improved performance in all test sets when compared with other scoring methods especially in its ability to minimize the RMSE. LISA and LISA+ displayed similar performance using the correlation coefficient and Kendall τ as the metric of quality for some of the small test sets. Further pathways for improvement are discussed which would KECSA more sensitive to subtle changes in ligand structure.

Keywords

Knowledge-based scoring function; Empirical scoring function; Potential of mean force

Introduction

Knowledge-based protein-ligand scoring functions $1-18$, building on the idea of potential of mean force (PMF),¹⁹ are derived from structural information regarding protein-ligand complexes. Their pairwise interaction parameters are directly converted from the frequency of occurrence of given atom pairs contained in a large database of complexes. The concept of the potential of mean force can be illustrated by a simple fluid system of N particles whose positions are $r_1 \ldots r_N$.

The average potential $\omega^{(n)}(r_1... r_N)$ is expressed as:

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Supporting Information

Supporting Information is available including the Lennard-Jones potential models parameter η values and Predicted pKd or pKi vs. experimental p K_d or p K_i for our test sets. This information is available free of charge via the Internet at<http://pubs.acs.org/>.

$$
\omega^{(n)}(r_1 \cdots r_n) = -\frac{1}{\beta} \ln(g^{(n)}(r_1 \cdots r_n)) \quad (1)
$$

where $g^{(n)}$ is called a correlation function. β= $1/k_BT$ and k_B is the Boltzmann constant and T is the system temperature. Hence the mean potential of the system with N particles is strictly the potential that gives the average force over all the configurations of the $n+1...N$ particles acting on a particle at any fixed configuration keeping the 1…n particles fixed. The mean potential can be described as follows:

$$
-\nabla_j \omega^{(n)} = \frac{\int \cdots \int e^{-\beta U} (\nabla_j U) dr_{n+1} \cdots dr_N}{\int \cdots \int e^{-\beta U} dr_{n+1} \cdots dr_N}, j = 1, 2, \cdots, n \quad (2)
$$

where U is the total potential energy of the system. Described by Sippl and others,1–5 the average potential is expressed as Equation 3 for the special case of a system with an observed particle number of n=2, as is the case herein (pairwise atoms from the protein and ligand).

$$
\omega_{ij}^{(2)}(r_{12}) = -\frac{1}{\beta} \ln(g^{(2)}(r_{12})) = -\frac{1}{\beta} \ln(\frac{\rho_{ij}(r_{12})}{\rho_{ij}^*(r_{12})})
$$
(3)

where $g^{(2)}(r)$ is the pair distribution function, $\rho_{ij}(r)$ is the number density for the atom pairs of types *i* and *j* observed in the known protein structures and $\rho^*_{ij}(r)$ is the number density of the corresponding pair in a reference state. In order to obtain the pure interaction potential between atoms, a reference state is required to remove the contribution of the ideal-gas state potential. So, in the reference state, the system of particles is like an ideal-gas state defined by fundamental statistical mechanics, in which particles would be evenly distributed in the binding site. Equation 3 can also be expressed as:

$$
\omega_{ij}^{(2)}(r_{12}) = -\frac{1}{\beta} \ln(g^{(2)}(r_{12})) = -\frac{1}{\beta} \ln(\frac{n_{ij}(r_{12})}{n_{ij}^*(r_{12})})
$$
 (4)

where $n_{ij}(r)$ and $n^*_{ij}(r)$ are numbers of atom pairs of type *i* and *j*, respectively, at distance *r* for the observed structures and the reference state.

In potential of mean force methods, the number of the corresponding pairs in the reference state cannot be exactly obtained for protein-ligand systems due to the effects of connectivity, excluded volume, composition, etc.⁶ Therefore, the pairwise interaction potential cannot be accurately calculated. Nonetheless, this idea of PMF scoring has advantages over empirical scoring, because it directly relates pairwise interaction to structural data instead of fitting to known binding affinity data. Additionally, the PMF is more efficient than force field scoring due to the avoidance of higher expense computations. Our intent is to introduce a new concept of the reference state, in order to relate the statistical potential to atomic pairwise interaction potential. Hence the atomic pairwise interaction model can be parameterized exclusively from structural data instead of binding data or quantum calculations.

Methods and Results

Construction of traditional statistical potentials starts by collecting structural information from large numbers of protein-ligand complexes, in order to simulate a "mean force" state in which the protein-ligand atomic pairwise radial distribution arises from all possible interactions in the binding site. Various reference states have been designed to remove the

non-interacting energy from the "mean force" state in order to correlate the pairwise radial distribution to the interaction potential between selected atoms of a specific atom type \dot{i} , \dot{j} with all other atoms in the protein-ligand binding site.

Our goal is to equate the statistical potential to the Lennard-Jones potential for each pairwise interaction. However, the LJ potential reflects pairwise interactions between two types of atoms, while a statistical potential is an average potential contributed by all atoms within the binding region. In this case, when trying to equate the statistical potential to a pairwise interaction potential, we need to remove all interactions except the pairwise interaction between atoms of type \vec{i} and \vec{j} in the binding region by defining a new reference state (denominator) in the PMF model. Unlike the traditional reference state, in which the selected atom pairs *i* and *j* are at an infinite separation where the interaction energy is zero (as in the ideal gas state), within this new reference state (which we will call reference state II), a system of particles is under an average force contributed by all atoms in the binding region excluding the interaction force between the selected atom pairs i and j . In other words, the only difference between the mean force state and the reference state II is that the latter state does not contain the pairwise interaction potential between the selected atoms of type *i* and *j*. Figure 1 provides a graphic illustration of the KECSA statistical potential model.

When equated to the LJ potential, the statistical potential can be expressed as:

$$
E_{ij}(r) = -RT \ln \left[\frac{n_{ij}(r)}{n_{ij}^{**}(r)} \right] = RT \left(\ln \left[n_{ij}^{**}(r) \right] - \ln \left[n_{ij}(r) \right] \right) = \frac{-1}{\left(\frac{\beta}{\alpha} \right)^{\frac{\alpha}{\alpha-\beta}} - \left(\frac{\beta}{\alpha} \right)^{\frac{\beta}{\alpha-\beta}}} \varepsilon \left[\left(\frac{\sigma}{r_{ij}} \right)^{\alpha} - \left(\frac{\sigma}{r_{ij}} \right)^{\beta} \right] \tag{5}
$$

where σ is the distance at which the inter-particle potential is zero and ε is the well depth. The exponents for the repulsive term and attractive term are α and β , respectively. We derive the exponents instead of assigning "typical" exponent values $(i.e., 12–6)$, because (1) the repulsive and attractive forces change with different types of pairwise interactions and (2) $E_{ij}(r)$ in Equation 5 includes both van der Waals and electrostatic interactions, which means the LJ-potential formula on the right hand side of Equation 5 accounts for two components:

$$
\frac{-1}{\left(\frac{\beta}{\alpha}\right)^{\frac{\alpha}{\alpha-\beta}}-\left(\frac{\beta}{\alpha}\right)^{\frac{\beta}{\alpha-\beta}}}\varepsilon\left[\left(\frac{\sigma}{r_{ij}}\right)^{\alpha}-\left(\frac{\sigma}{r_{ij}}\right)^{\beta}\right]\approx 4\varepsilon_0\left[\left(\frac{\sigma_0}{r_{ij}}\right)^{12}-\left(\frac{\sigma_0}{r_{ij}}\right)^{6}\right]+\frac{q_1q_2}{\varepsilon_1r_{ij}}\quad(6)
$$

The reason we use the LJ formula on the left hand side of Equation 6 instead of partitioning them into van der Waals and electrostatic potentials is that the LJ potential reaches 0 at σ

and R, while reaching its minimum value when r is $\left(\frac{\alpha}{\beta}\right)^{\frac{1}{\alpha-\beta}}\sigma$. Based on these properties, equations can be derived in order to determine the unknown parameters.

In Equation 5, $n_{ij}^{**}(r)$ and $n_{ij}(r)$ are the number of protein-ligand atomic pairwise interactions in the bin (r, r+ Δ r), with the volume $4\pi r^a \Delta r$ in reference state II and in the training set that mimics the mean force state, respectively. Δr is defined as 0.005Å. We introduce a to-bedetermined parameter a for the shell volume because of the inaccessible volume present in protein-ligand systems, and because of the deviation of $n_{ij}(r)$ in the training set from the "perfect" pairwise number. Hence, the expectation is the parameter a will adopt values other than 2.

The central issue in the KECSA model construction is to build up the radial distribution function of the selected atom pairs in reference state II. A way in which to do this is to measure the similarity of the reference state II with two known states: the mean force state, and ideal gas state. Then build the radial distribution function with information collected from these two states. In reference state II, the radial distribution of a certain atom pair i and j is associated with a certain "background interaction" which is related to the total number of selected atom pairs N_{ij} . Because the "background interaction" potential contains all atom pairwise interactions in the binding site excluding the "selected interactions" between atom pairs i and j, the difference in energy between the mean force state and reference state II for each atom pair type depends on the total number of the selected atom pairs N_{ij} and the total number of atom pairwise interactions N found in the binding site. The "background

interaction" energy approaches the "mean force" state energy as $\frac{N_{ij}}{N}$ becomes smaller, while

the energy difference increases when $\frac{1}{N}$ becomes larger. Hence we decide to make the "background interaction" potential, as well as the radial distribution function, a function of

 $\frac{N_{ij}}{N}$. The modeling of the atom pairwise number distribution function in reference state II

starts from the two extreme situations for $\frac{1}{N}$: (1) When N_{ij} approaches zero $(N_{ij} \rightarrow 0)$, the background energy \approx the "mean force" state energy, resulting in $n_{ij}^{**}(r) \approx n_{ij}(r)$. (2) When N_{ij} approaches to $N(N_{ij} \rightarrow N)$, the background energy \approx the ideal-gas energy, resulting in $n_{ij}^{**}(r)$ resembling an ideal gas state radial distribution. The radial distribution function for an ideal

gas state is defined as $\left(\overline{V}^{4\pi r \Delta r}\right)$ implying that the number of the selected atom pairs i and j is evenly distributed in the binding site which has an average volume V . The average V of

protein-ligand binding site is given as $\frac{\pi R^{a+1}}{a+1}$, with the same to-be-determined parameter ^a as introduced above. Hence, the two extreme situations for reference state II can be defined as:

$$
n_{ij}^{**}(r) = \left(\frac{N_{ij}}{V} 4\pi r^a \Delta r\right), N_{ij} \to N \quad (7)
$$

$$
n_{ij}^{**}(r) = n_{ij}(r), N_{ij} \to 0 \quad (8)
$$

At certain distance r, $n_{ij}^{**}(r)$ is a function of $\frac{1}{N}$ with a range from $\left(\frac{1}{V}4\pi r \Delta r\right)$ to $n_{ij}(r)$. In addition, with N_{ij} tending towards 0 or N, the reference state II would be more similar to the mean force state or the ideal gas state, respectively. Hence, $n_{ij}^{**}(r)$ is defined as a weighted combination of both the ideal gas state and the mean force state radial distribution functions. Due to the fact that the integral of $n_{ij}^{**}(r)$ from 0 to R (cutoff distance where the atomic interaction is regarded as zero) is N_{ij} , a linear combination (Equation 9) of the weighted radial distribution functions for both the ideal gas and mean force state meets all the necessary conditions:

$$
n_{ij}^{**}(r) = \left(\frac{N_{ij}}{V} 4\pi r^a \Delta r\right) \frac{N_{ij}}{N} + \left(n_{ij}(r)\right) \left(1 - \frac{N_{ij}}{N}\right) \tag{9}
$$

In this way the new reference state is designed as state intermediate between the ideal gas and the mean force state. At a certain distance between the atom pairs of type i and j , the total energy of reference state II $(E_{ij}^*(r))$, or the "background interaction" energy is:

$$
E_{ij}^*(r) = -\frac{1}{\beta} \ln(n_{ij}^*(r)) = -\frac{1}{\beta} \ln\left(\frac{N_{ij}}{V} 4\pi r^a \Delta r\right) \frac{N_{ij}}{N} + (n_{ij}(r)) \left(1 - \frac{N_{ij}}{N}\right) \tag{10}
$$

A plot illustrating the relationship between the number fraction $\frac{N_{ij}}{N}$ and the new reference state potential $E_{ij}^*(r)$ is shown in Figure 2, to better illustrate the differences in energy between the different states.

Combining Equations 5 and 9 we obtain:

$$
\ln\left[\left(\frac{N_{ij}}{V}4\pi r^a \Delta r\right)\frac{N_{ij}}{N} + (n_{ij}(r))\left(1 - \frac{N_{ij}}{N}\right)\right] - \ln[n_{ij}(r)] = \frac{-1}{\left(\frac{\beta}{\alpha}\right)^{\frac{\alpha}{\alpha-\beta}} - \left(\frac{\beta}{\alpha}\right)^{\frac{\beta}{\alpha-\beta}}} \frac{\varepsilon}{RT} \left[\left(\frac{\sigma}{r_{ij}}\right)^{\alpha} - \left(\frac{\sigma}{r_{ij}}\right)^{\beta}\right] \tag{11}
$$

Using $V = \frac{1}{a+1} \pi R^{a+1}$ and using the fact that the Lennard-Jones potential is zero at $r_{ij} = \sigma$ and at $r_{ij} = R$, we arrive at:

$$
\ln\left[\frac{N_{ij}}{N}\left(\frac{N_{ij}(a+1)\sigma^a\Delta r}{R^{a+1}n_{ij}(\sigma)}\right)+(1-\frac{N_{ij}}{N})\right]=0, \text{ and } (12)
$$

$$
\ln\left[\frac{N_{ij}}{N}\left(\frac{N_{ij}(a+1)R^a\Delta r}{R^{a+1}n_{ij}(R)}\right)+(1-\frac{N_{ij}}{N})\right]=0. (13)
$$

In addition, the LJ potential reaches its minimum value when r is $\sigma\left(\frac{\alpha}{\beta}\right)^{\frac{1}{\alpha-\beta}}$. So the first derivative (D) of the statistical energy term with respect to r is zero at $r = \sigma \left(\frac{\alpha}{\beta}\right)^{\frac{1}{\alpha-\beta}}$.

$$
D\left[-\frac{1}{RT}\ln\left(\frac{n_{ij}(r)}{n_{ij}^{**}(r)}\right),\,r=\sigma\left(\frac{\alpha}{\beta}\right)^{\frac{1}{\alpha-\beta}}\right]=0\quad(14)
$$

To simplify the resultant expressions the factor $\left(\frac{\alpha}{\beta}\right)^{\frac{1}{\alpha-\beta}}$ is given as η .

$$
D[-\frac{1}{RT}\ln(\frac{n_{ij}(r)}{n_{ij}^{**}(r)}), r=\eta\sigma] = \frac{a\frac{N_{ij}}{N}\left(\frac{N_{ij}(a+1)(\eta\sigma)^{a-1}\Delta r}{R^{a+1}n_{ij}(\eta\sigma)}\right) - \frac{N_{ij}}{N}\left(\frac{N_{ij}(a+1)(\eta\sigma)^{a-1}\Delta r}{R^{a+1}n_{ij}^2(\eta\sigma)}\right)D(n_{ij}(\eta\sigma))}{\frac{N_{ij}}{N}\left(\frac{N_{ij}(a+1)(\eta\sigma)^{a}\Delta r}{R^{a+1}n_{ij}(\eta\sigma)}\right) + (1-\frac{N_{ij}}{N})} = 0 \quad (15)
$$

Simplifying Equations 12, 13 and 15 yields:

$$
N_{ij}(a+1)\sigma^a \Delta r = R^{a+1} n_{ij}(\sigma) \quad (16)
$$

$$
N_{ij}(a+1)R^a \Delta r = R^{a+1} n_{ij}(R) \quad (17)
$$

and

$$
an_{ij}(\eta\sigma)=D(n_{ij}(\eta\sigma))
$$
 (18)

Although we don't know the values of α and β yet, we do know that the value of η is unique for each combination of α and β. Supplementary Table 1 lists all η values for each integer combination of α and β from 2-1 to 15-14. Different η values will be chosen for every

pairwise interaction, to satisfy the well depth distance at $\eta \sigma$ (e.g., $\left(\frac{\alpha}{\beta}\right)^{\frac{1}{\alpha-\beta}} \sigma$ is $2^{1/6}\sigma$ or r_{ij} * (well-depth at the minimum) for the 12-6 potential).

In order to find the a , σ and η values within Equations 16–18, we still need to determine the value of R , the cutoff distance. We introduce a nonlinear programming method to find a reasonable R for each pairwise interaction type instead of assigning a fixed R value. Ideally, *should be as large as possible since the LJ potential approaches 0 when the distance* approaches infinity. Meanwhile, for any r between σ and R, the potential value is below 0. Here we use the following inequality constraint in our nonlinear programming approach:

$$
\frac{1}{RT}\ln\left[\frac{N_{ij}}{N}\left(\frac{N_{ij}(a+1)r^a\Delta r}{R^{a+1}n_{ij}(r)}\right)+(1-\frac{N_{ij}}{N})\right]<0,\,\sigma
$$

which can be simplified as:

$$
N_{ij}(a+1)r^a \Delta r \langle R^{a+1}n_{ij}(r), \sigma \langle r \langle R \quad (20)
$$

With the goal of maximizing the value of R , coupled with the three constraint equations (Equations 16–18) and an inequality constraint (Equation 20), a , σ and η can be determined. The values of η obtained in this way can then be compared with the η values in supplementary Table 1, in order to determine the closest α and β pair. Inserting these values into Equation 11 we can calculate all of the corresponding ε values.

One important issue in the parameterization of KECSA is that the LJ potential parameters for each type of pairwise interaction should be independent of the other types of interactions,

instead of being $\frac{d}{d}$ dependent. Derivation of a, σ , α , β and R comes from Equations 16–18 and 20, none of which contains the total interaction number N. This indicates that for each

type of pairwise interaction, the average volume $(\frac{4}{a+1}\pi R^{a+1})$, the distances at which the LJ

potential reaches zero and has a minimum $(\sigma \text{and} \left(\frac{\alpha}{\beta}\right)^{\frac{1}{\alpha-\beta}}\sigma)$, the relative strength of the repulsive and attractive forces in the LJ potential $(α'$ and $β)$, and the long-range cutoff distance (R) are independently derived in KECSA. The only issue lies in the derivation of

the e values from Equation 11, where the probability of occurrence $\left(\frac{N_{ij}}{N}\right)$ is included in the

calculation. In order to avoid relative energies generated for each interaction type based on

their probability of occurrence in protein-ligand binding sites, we used a normalized $\frac{N_{ij}}{N}$ for each interaction type. Thereby the number fractions for all interaction types are identical.

In the present work, all pairwise interactions among 18 atom types (listed in Table 1) were examined resulting in 49 significant interaction types being identified. The remaining interaction types were abandoned or merged into similar interaction types due of the paucity of data to fit to or because they are randomly distributed across the observed distance range. The chosen interaction types included 38 van der Waals and 11 hydrogen bonding interaction types. In this case, all interaction types share the same probability of occurrence

 in protein-ligand binding sites. Equation 11 can be rewritten as follows in order to generate the ε values. All derived parameters are listed in Table 2.

$$
\ln\left[\left(\frac{N_{ij}}{V}4\pi r^a \Delta r\right)\frac{1}{49} + (n_{ij}(r))\left(1 - \frac{1}{49}\right)\right] - \ln[n_{ij}(r)] = \frac{-1}{\left(\frac{\beta}{\alpha}\right)^{\frac{\alpha}{\alpha-\beta}} - \left(\frac{\beta}{\alpha}\right)^{\frac{\beta}{\alpha-\beta}}} \frac{\varepsilon}{RT}\left[\left(\frac{\sigma}{r_{ij}}\right)^{\alpha} - \left(\frac{\sigma}{r_{ij}}\right)^{\beta}\right] \tag{21}
$$

With all of the enthalpy terms determined in the analytical manner described above, the entropy terms are then decided upon in an empirical manner. Structural information such as the number of rotatable bonds, number of double and aromatic bonds, molecular mass, counts of carbon/oxygen/nitrogen atoms, buried surface area, etc. were collected for all ligands contained in the training set. The selection of entropy terms is based on their contribution to our linear regression model, whose 95% confidence interval should not include 0. Finally, 9 entropy terms are selected: number of rotatable bonds in the ligand, the molecular mass of the ligand, number of aromatic bonds in the ligand, number of oxygen atoms in the ligand, number of nitrogen atoms in the ligand, the nonpolar buried surface area, total buried surface area, the ratio of the nonpolar buried surface and total ligand surface area and, finally, the ratio of the total buried surface area and the total ligand surface area. The PDB bind v2010 data set $2^{1,22}$ including 5054 protein-ligand complexes is chosen as the training set for the parameterization of the enthalpy terms. We chose 1982 proteinligand complexes found in the PDBbind v2011 refined data set as the training set for the selection and parameterization of the entropy terms.

Model Validation and Discussion

I) Validation to Detect Over-Fitting and Ligand-Size Dependence

Several validation benchmarks were introduced to test the performance of KECSA. The first benchmark included two test sets in order to examine the dependence of KECSA on the training set. Because the entropy term was obtained by fitting to experimental binding free energies care was taken to ensure that the resultant model was not over-fit. First, a leaveone-out cross validation was used against the training set, which includes 1982 proteinligand complexes used in the KECSA entropy term parameterization. Comparison of the Pearson correlation coefficient r, RMSE (root-mean-square error) and Kendall τ between the training set and the leave-one-out cross validation are shown in Table 4. The three statistical measures all showed small differences between the training set and the leave-one-out prediction, indicating that the KECSA entropy model was properly built.

Second, we introduced a test set including 1934 protein-ligand complexes chosen from the PDBbind v2011 dataset, with no overlap with the training set used above. From the total of 6051 protein-ligand complexes with binding affinity data in the PDBbind v2011 dataset,

complexes forthis test set were selected following four criteria: (1) Crystal structures of all selected complexes had X-ray resolutions $\;$ 2.5Å. (2) Only complexes with p $K_{\rm i}$ or p $K_{\rm d}$ values distributed between 2 to 8 were selected, mimicking what might be found in a virtual screening database or a pharmaceutically relevant ligand database. (3) Only complexes with molecular weights (MWs) distributed from 80 to 800 were selected, to avoid ligand sizedependent prediction results. (4) Complexes used in the KECSA entropy term training set were excluded.

The second test set was used to verify the robustness of KECSA against and external dataset as well as investigating the contributions of the enthalpy and entropy terms in KECSA's binding affinity prediction. In addition, because the entropy term in KECSA was modeled as a linear combination of several ligand properties including ligand size/mass information this test demonstrated that KECSA was not ligand size-dependent. We split the KECSA scoring function and used both LJ potential and entropy terms for binding affinity prediction and then compared with the full KECSA scoring function. We also calculated the correlation coefficient between the experimental pK_i or pK_d with ligand MW. The predictions and ranking results are listed in Table 5.

KECSA produces a Pearson's r of 0.590 and a Kendall τ of 0.404. Comparison to the training set result (a Pearson's r of 0.610 and a Kendall τ of 0.442) indicates that KECSA gives robust predictions against an unknown binding affinity dataset. When just using the ligand MW for binding affinity prediction, the Pearson's r drops by 0.128 and the Kendall τ drops by 0.08 when compared to the LJ potential only prediction, while Pearson's r drops by 0.140 and the Kendall τ drops by 0.08 when compared to the entropy term prediction. The drop in prediction is even greater when compared to the full KECSA score function. This result suggests that the KECSA prediction is minimally affected by MW considerations. The LJ potential and entropy only predictions are quite similar, but lower than the full KECSA prediction (see Table 5). None of the two independent parts, enthalpy or entropy, shows a significant performance over the other, suggesting that both enthalpy and entropy play important roles in the KECSA scoring function. RMSE is not listed for comparison because the LJ potential is generated from a statistical potential while entropy is derived by fitting to pK_d or pK_i values, which results in different scales making comparison difficult.

II) Validation Benchmark for Comparison with LISA, LISA+ and Other Scoring Methods

KECSA is the second-generation scoring function we have developed after $LISA^{23}$ and $LISA+²⁴$ The latter two were developed for the fast calculation of protein-ligand binding affinity (pK_d and pK_i), and were successful in the SAMPL3 challenge (first rank for all scoring methods), which was a blind test for docking and scoring methods.²⁴ Comparisons between KECSA and other scoring methods including LISA and LISA+ are necessary for further understanding its performance. The second validation benchmark consists of two large-scale test sets and four smaller test sets for comparison of KECSA with LISA/LISA+, as well as one small scale test set (Wang's test set²⁵) with 100 diverse protein-ligand complexes for comparison of KECSA and several other well-known scoring methods.

First, two large scale test sets both with more than 1000 complexes were introduced to KECSA, LISA and LISA+. The first test set contained 1399 complexes from the PDBbind v2010 database, which was previously used for LISA validation. KECSA reproduces a Pearson correlation coefficient r of 0.553, an RMSE of 2.46kcal/mol and a Kendall τ of 0.401, while LISA reproduces a Pearson correlation coefficient r of 0.534, an RMSE of 2.65kcal/mol and a Kendall τ of 0.378. LISA+ was trained based on this data set, so it was excluded from this validation benchmark. A larger test set was applied for all three scoring functions, including 2456 protein-ligand complexes from the PDBbind v2011 refined data set, of out which 290 complexes had Zn-ligand binding. For those 2166 non-metal

containing complexes, KECSA gets an r of 0.589, an RMSE of 2.31kcal/mol and a Kendall τ of 0.429, while LISA gets an r of 0.542, an RMSE of 3.06kcal/mol and a Kendall τ of 0.397, LISA+ yields an r of 0.572, an RMSE of 2.81kcal/mol and a Kendall τ of 0.419. For those complexes with Zn-ligand binding, KECSA has an r of 0.415, an RMSE of 2.33kcal/ mol and a Kendall τ of 0.267, LISA has an r of 0.409, an RMSE of 3.08kcal/mol and a Kendall τ of 0.252, and LISA+ has an r of 0.420, an RMSE of 3.00kcal/mol and a Kendall τ of 0.257. Calculation and statistical results of KECSA, LISA+ and LISA for all large-scale validation studies are shown in Table 6 and Figure 3.

In the large-scale test, KECSA yields a better prediction than our two first-generation scoring functions. It produces better predicted results based on RMSE and more reliable binding affinity ranking based on Kendall τ compared with the other two scoring methods. LISA+ can compete with KECSA based on correlation coefficient, and even achieves better r in the subset of complexes with Zn-ligand binding. We believe that the improvement of LISA+ compared with LISA is because the complexes in the training set are categorized based on ligand properties (mass and hydrophobicity) and different models are trained for each category. This proves that a multi-model scheme can improve the predictive ability of empirical scoring functions.

Our validation studies indicate an improvement from LISA to KECSA. Introducing PMF theory for non-bonding interaction modeling shows its advantages over simply fitting to binding affinity data. However, metal-ligand binding prediction remains a challenge for classical-mechanics based or statistical scoring methods. KECSA improves the binding affinity predicting ability mostly in RMSE for this subgroup of complexes. However, although KECSA shows improvement in both correlation coefficient r and RMSE, predictions in the low and high binding regions are poor. Seen from the linear regression functions in Figure 3 the slopes of LISA and LISA+ generated data vs. experimental data are 0.66 and 0.64, while that for KECSA is 0.41; the intercepts for LISA and LISA+ are 2.02 and 2.06, while that for KECSA result is 3.72. The reason is that 3314 complexes, which comprised 65.3% of the whole training set, are in the mid-binding region (pK_d or pK_i) between 4 and 8). Hence the scoring function tends to overestimate binding affinity of the low-binding region while underestimating that of the high-binding region if there is no significant decrease or increase in contact number for the protein-ligand complexes from these two regions. We used a scaling procedure to help improve the prediction of binding affinity in the high and low binding regions. The KECSA generated results were fit to a linear model to reproduce the PDBbind v2011 refined data set pK_d values. So we get:

 $pK_d = 1.889 \times (pK_d)_{\text{KECSA}} - 5.539$ (22)

Next, four test sets containing 427 protein-ligand complexes from four protein families was examined. The list of complexes, their families and binding constants are given in Supplementary Table 2. Statistical and calculation results are shown in Table 7 and Figure 4. KECSA improves the RMSE for all test sets, indicating that KECSA makes a more robust prediction with respect to experimental binding affinity data than do LISA and LISA+. However, LISA and LISA+ both give a better correlation coefficient and Kendall τ for the serine protease, endothiapepsin and HIV-1 protease test sets, showing that they perform better in binding affinity ranking in these small test sets. Results for each test set were carefully examined. The serine protease test set contains many complexes with low-mass ligands, while most ligands are relatively larger in the endothiapepsin test set. Statistical results of the three scoring methods' predictions are similar in these two test sets: KECSA generates a better RMSE, while LISA and LISA+ yielded both a better correlation coefficient and Kendall τ . In the serine protease test set, 11 out of 96 protein-ligand

complexes had small ligands with molecular mass lower than 200 Daltons. KECSA prediction overestimates most of their pK_d or pK_i values (Supplementary Table 2), while LISA and LISA+ to some degree underestimates these values. Hence, KECSA decreases the binding affinity differences between these low-binding complexes and other high-binding complexes and LISA/LISA+ increases these differences. The implication is that LISA and especially LISA+ differentiate the binding affinity values of the complexes better in this test set and give higher r and τ values, but they have worse RMSEs. In the endothiapepsin test set, on the other hand, all complexes have ligands with molecular mass higher than 500 Daltons. LISA and LISA+ overestimated binding affinities with larger errors compared with KECSA, while distinguishing the complexes better, generating a better linear correlation towards the experimental data. These test results suggest that LISA and LISA+ are more sensitive to ligand mass changes, while KECSA makes more precise prediction with smaller error, but has difficulty in ranking complexes with similar pK_d s or pK_i s. The HIV-1 protease test set is a significant challenge for all three scoring methods. Most of the complexes have high-mass ligands and high binding affinity. Because the ligands in this test set have similar structures binding affinity predictions from all three scoring methods are not able to rank these complexes. LISA+ does better in correlation coefficient and ranking than others, which indicates that training scoring functions for different complex categories based on ligand or binding pocket properties may help scoring methods improve their ability to identify subtle changes in ligand structure. The carbonic anhydrase II test set includes 100 out of 110 complexes with Zn chelation, and contains more polar and charged interactions. KECSA demonstrates better performance in the correlation coefficient, RMSE and Kendall τ, pointing to its advantage over LISA and LISA+ on reproducing binding affinity data of complexes with more hydrophilic interactions and metal chelation. For the reproduction of the binding affinity for the four combined test sets all three scoring methods give a similar correlation coefficient r, while LISA+ has a small advantage. KECSA does better in both RMSE and Kendall τ. For all four test sets, scaling of the KECSA score does help to improve the slope of calculated data vs. experimental data, but does not improve the statistical tests. Overall, KECSA give better RMSE values and better binding affinity ranking of the complexes belonging to different protein families.

For the last test set, we introduced Wang's test set²⁵ with 100 diverse protein-ligand complexes. The purpose was to compare binding affinity prediction ability of KECSA not only with LISA and LISA+, but also with other well-known scoring functions. We obtain Pearson's $r = 0.69$, RMSE = 2.25 kcal/mol using KECSA, compared with Pearson's $r = 0.72$, RMSE = 2.32 kcal/mol using LISA, Pearson's $r = 0.67$, RMSE = 2.80 kcal/mol using LISA +. This result coincides with the conclusion gained from the second validation benchmark, that among these three scoring functions, KECSA prediction has the smallest RMSE. Pearson correlation coefficients of KECSA together with other score functions are presented in Figure 5, showing its performance on this test set.

Conclusion and Outlook

Based on atom pairwise interactions, interaction enthalpy terms in KECSA were parameterized by combining PMF theory with the Lennard-Jones potential, without fitting to any binding affinity data. This procedure parameterizes the LJ potential with neither QM calculations nor binding affinity data, hence lowering the computational expense while improving the prediction accuracy relative to empirical scoring functions. Generally, KECSA improves the binding affinity RMSE, when compared to LISA and LISA+, especially for complexes dominated by polar and charged interactions. With respect to ranking predictions, KECSA better distinguishes complexes in the large-scale test sets. KECSA yields the lowest RMSE values illustrated by its superior performance in all test sets for this measure of quality. It is less responsive, however, to minor structural changes of

the ligand or binding pocket, reducing its ability to rank complexes from the same or similar protein families. In the KECSA model, the solvent accessible surface area is introduced to describe the desolvation effect and entropy terms were empirically modeled. Since we have formulated and parameterized the enthalpy component of non-covalent interactions with this alternative method, an interesting possibility is to use similar procedures to build a force field model solely based on experimental structural data. In this way, the desolvation and entropy terms can also be included instead of using empirical models. We believe that more accurate and effective scoring methods can be developed using this concept.

Our group has constructed an in-house docking program where KECSA (as well as LISA and LISA+) is employed as the scoring module in this program. KECSA's ability to score docking poses and distinguish native poses from decoys will be further evaluated and refined in future work using this docking module.

Supplementary Material

Refer to Web version on PubMed Central for supplementary material.

Acknowledgments

We would like to thank the NIH (GM044974 and GM066859) for supporting the present research. Dr. Michael Weaver is also acknowledged for numerous helpful discussions.

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 $\omega_{ij}^{(2)}(r_{12}) = -$

a. Mean Force State

c. the State with only the potential of interaction between selected atomic type pair *i* and *j*.

b. Reference State II

Figure 1.

A protein-ligand structural illustration (using PDBID 1xbc) of how the KECSA statistical potential is modeled. The protein binding site is shown as a grey surface with the ligand located within the binding site surrounded by protein residues which it makes contacts with. The pink dashed lines indicate interactions between certain atom pair types i and j , (*i.e.*) carbonyl oxygen with amine nitrogens in this example) which are defined as "selected interactions" in this manuscript. Green dashed lines indicate all other non-covalent interactions between the protein and ligand atoms in the binding pocket, defined as "background interactions". (a) In the mean force state, the system is filled with all types of interactions. (b) The reference state II contains all the background interactions. (c) Removing all the background interactions from total interactions results in a state with only the selected interactions for each i and j combination.

Figure 2.

Ratio of the observed atom pairs to the total interacting atom pairs $\frac{N_{ij}}{N}$ vs. the new reference state II potential.

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Figure 3.

Plot of KECSA, LISA+ and LISA calculated pK_d or pK_i vs. Experimental pK_d or pK_i in obtained in validation studies.

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Figure 4.

Plot of KECSA, scaled KECSA, LISA+ and LISA calculated pK_d or pK_i vs. Experimental pK_d or pK_i in four small test sets.

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Figure 5.

With the test set built by Wang,²⁵ binding affinity comparison was done for KECSA, LISA, LISA+ and several other well-known scoring functions, ITScore/SE,¹⁸ ITScore,¹⁷ X-Score,²⁶ DFIRE,²⁷ DrugScoreCSD,¹² DrugScorePDB,¹¹ Cerius2/PLP,^{28,29} SYBYL/G-Score, ³⁰ SYBYL/D-Score, ³¹ SYBYL/ChemScore, ³² Cerius2/PMF, ⁸ DOCK/FF, ³¹ Cerius2/ LUDI,^{33,34} Cerius2/LigScore,³⁵ SYBYL/F-Score,³⁶ AutoDock.³⁷

List of selected atom types.

interaction interaction

σaRε

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Entropy parameters and their 95% confidence intervals

Leave-one-out cross validation of KECSA.

Statistical results for KECSA, KECSA LJ, KECSA entropy and Ligand MW correlated with experimental binding affinity.

Statistical results of KECSA, LISA+ and LISA from large-scale validation studies. Statistical results of KECSA, LISA+ and LISA from large-scale validation studies.

Test set 2 contains 2166 non-metal complexes from PDBbind v2011 refined data set.

 c Test set 3 contains 290 metalloprotein-ligand complexes from PDB
bind v2011 refined data set. c . Test set 3 contains 290 metalloprotein-ligand complexes from PDBbind v2011 refined data set.

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