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Exchange-Mediated Contrast in CEST and Spin-Lock Imaging

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Abstract

PURPOSE—Magnetic resonance images of biological media based on chemical exchange saturation transfer (CEST) show contrast that depends on chemical exchange between water and other protons. In addition, spin-lattice relaxation rates in the rotating frame $(R_{1\rho})$ are also affected by exchange, especially at high fields, and can be exploited to provide novel, exchange-dependent contrast. Here, we evaluate and compare the factors that modulate the exchange contrast for these methods using simulations and experiments on simple, biologically relevant samples.

METHODS—Simulations and experimental measurements at 9.4T of rotating frame relaxation rate dispersion and CEST contrast were performed on solutions of macromolecules containing amide and hydroxyl exchanging protons.

RESULTS—The simulations and experimental measurements confirm that both CEST and R_{1p} measurements depend on similar exchange parameters, but they manifest themselves differently in their effects on contrast. CEST contrast may be larger in the slow and intermediate exchange regimes for protons with large resonant frequency offsets (e.g. > 2ppm). Spin-locking techniques can produce larger contrast enhancement when resonant frequency offsets are small (< 2 ppm) and exchange is in the intermediate to fast regime. The image contrasts scale differently with field strength, exchange rate and concentration.

CONCLUSION—CEST and R_{1p} measurements provide different and somewhat complementary information about exchange in tissues. Whereas CEST can depict exchange of protons with specific chemical shifts, appropriate $R_{1\rho}$ dependent acquisitions can be employed to selectively portray protons of specific exchange rates.

Keywords

CEST; spin lock; $R_{1\rho}$; chemical exchange

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INTRODUCTION

There is continuing interest in developing and exploiting novel contrast mechanisms in proton magnetic resonance imaging (MRI) to characterize tissues. Proton exchange between water and exchangeable protons in other molecules provides one such mechanism that reports the presence of specific chemical components within a mixture. Methods such as chemical exchange saturation transfer (CEST) depend explicitly on exchange between water and protons with specific chemical shifts, but other approaches may also be affected by exchange on appropriate time scales, especially at high magnetic fields. In particular, exchange between sites of different chemical shifts contributes directly to both transverse relaxation rates (R₂) and spin-lattice relaxation rates in the rotating frame ($R_{1\rho}$). Here we compare and contrast the sensitivity and selectivity for producing contrast based on exchange processes of CEST with a novel approach based on measurements of R_{1p} .

The dynamics of exchange with amide and hydroxyl sites in particular have been extensively exploited to generate contrast in CEST imaging (1,2). To interpret CEST data, protons are considered to comprise at least two pools, the solvent water and the exchangeable protons in the solute, as shown in Figure 1. Each pool is characterized by its own relaxation times and chemical shift, but they communicate via chemical exchange at specific rates. Where the exchanging species is largely derived from a singular tissue constituent, such as glycogen (glycoCEST) or gycosaminoglycans (gagCEST), the endogenous CEST contrast potentially reports on the concentrations of specific molecules in tissues (3,4). In addition, exogenous agents such as paramagnetic chelates have been developed to shift proton resonance frequencies to increase exchange effects (paraCEST) (5,6), while common x-ray contrast agents which contain exchanging amide and hydroxyl groups have also been shown to produce significant CEST effects (7,8) and therefore are also potential MRI contrast agents. However, in practice, CEST signal changes may be contaminated by non-specific magnetization transfer and nuclear Overhauser effects, and are sensitive to the effects of direct saturation as well as field inhomogeneities. In addition, the derivation of explicit information on exchange rates or molecular concentration requires the acquisition of multiple images and fitting the data to a model.

Measurements of $R_{1\rho}$ using spin-locking sequences are also sensitive to exchange effects at high fields but are less affected by some of the above considerations and are differently influenced by other factors, including the field strength and exchange rates. Here, we present simulations and experimental data to illustrate the influences of factors that affect signals in both CEST and spin lock imaging and to quantify how specific parameters such as exchange rate and field strength modify contrast. The sensitivities of CEST and spin locking (SL) techniques to chemical exchange effects in specific experimental regimes are predicted theoretically and measured experimentally and used to quantify contrast in polypeptide and sugar systems of biologic interest.

CEST methods include radiofrequency (RF) saturation of an exchanging species, which is then transferred to water, reducing its longitudinal magnetization. This change does not, in the ideal case of perfectly selective RF saturation, explicitly depend on the chemical shift of the irradiated protons. In practice, RF saturation is applied at a series of frequencies (offset δω relative to the water peak), and the acquired images yield a "z-spectrum" of the signal intensity at each voxel in which peaks correspond to specific exchanging species. However, the RF pulses may also alter the water signal by direct saturation or non-specific magnetization transfer (9) with other broad resonances. One strategy to correct for these effects is to acquire images at the opposite frequency offset(s) for metabolites of interest. The difference in the normalized signals from opposite sides of the water peak is the magnetization transfer ratio asymmetry (MTR_{asym}) :

$$
MTR_{asym} = \frac{S(\omega_0 - \delta\omega) - S(\omega_0 + \delta\omega)}{S_0}
$$
 [1]

where $S(\omega)$ is the signal when the RF pulse is at frequency ω , S_0 is the signal without RF saturation, and ω_0 is the resonance frequency of water.

Spin-locking techniques may also generate exchange-dependent contrast and can be implemented for imaging experiments (10,11). Typically, a 90-degree adiabatic half-passage (AHP) pulse nutates longitudinal magnetization to the transverse plane, followed by a spinlocking pulse (B_1) on resonance for some duration. Another 90-degree reverse half-passage (RHP) pulse returns magnetization to the longitudinal axis, residual transverse magnetization is spoiled, and an imaging sequence may then acquire $R_{1\rho}$ -weighted images. Relaxation rates in the rotating frame are sensitive to molecular interactions on the time scale defined by the locking field amplitude, B_1 , about which magnetization would precess at frequency $\omega_1 = \gamma B_1$ (12,13). Variations in $R_{1\rho}$ with locking-field strength ($R_{1\rho}$ dispersion) can provide quantitative information on the relevant parameters that describe chemical and diffusive exchange (14-16). In tissues, values of $R_{1\rho}$ are affected by slow molecular motions which modulate dipole-dipole interactions, but at high fields the increased separation of resonance frequencies between water and other chemical species gives rise to large contributions from chemical exchange whose magnitudes depend on the locking field amplitude and exchange rate.

CEST and SL techniques are both sensitive to chemical exchange effects, but have different constraints and sensitivities. Two important parameters are the chemical shift resonant frequency difference (δω) and the rate of proton exchange between the metabolite and free water (r_{BA}) (3,17). As in conventional spectroscopy, exchange regimes in CEST can be divided into slow, intermediate, and fast exchange relative to the chemical shift between water and the exchanging species. In high resolution spectroscopy, the chemical shift of an exchanging species must be greater than the exchange rate ($\delta \omega / r_{BA} > 1$) in order for individual peaks to be well resolved. As r_{BA} increases into the intermediate ($\delta \omega / r_{BA} \approx 1$) and fast $(\delta\omega/r_{BA} << 1)$ exchange regimes, line broadening reduces the ability to resolve individual resonances. At moderate saturation power levels, CEST contrast exhibits similar exchange regimes and constraints due to the spectral blurring effects of the applied irradiation field. Faster exchange requires larger saturating field strengths in order to achieve solute saturation, and creates correspondingly larger blurring of z-spectra. However, if saturation requirements are reduced, the conventional exchange rate constraints can be stretched, allowing, for example, CEST peaks to be resolved when conventional spectroscopy peaks are blurred (1). Nonetheless, CEST peaks are easiest to resolve in the slow exchange regime.

By contrast, fast exchange can be advantageous for affecting $R_{1\rho}$ values. Chemical exchange between sites of different chemical shift causes spin dephasing, which increases the difference between longitudinal (R_1) and transverse (R_2) relaxation rates, and thereby the dispersion of $R_{1\rho}$ with locking field, to a degree that increases with chemical shift and exchange frequencies. The values of $R_{1\rho}$ observed for different media (and therefore potential image contrast) may be selectively modulated by adjusting the amplitude of spinlocking pulses to reduce the contributions to relaxation of a specific range of exchange rates, up to the point where the rates become so large that the necessary locking field becomes impractical. In practice these rates can be much higher than those detectable by CEST.

Here we study exchange in solutions of biologically relevant macromolecules (polypeptides and sugars) in order to better understand the factors that modulate exchange-based image

contrast for CEST and $R_{1\rho}$ imaging, and examine the limits of each technique. We show, using theory, simulations and experimental measurements, how CEST contrast and a novel metric based on $R_{1\rho}$ dispersion may be used to emphasize the presence of protons characterized either by chemical shifts or by specific exchange rates.

METHODS

(a) COMPUTER SIMULATIONS

Numerical simulations were used to quantify the contributions of exchange to $R_{1\rho}$ and CEST measurements for the two-pool and three-pool models shown in Figure 1. The effects of chemical exchange on $R_{1\rho}$ and CEST contrast were calculated using the Bloch equations modified with exchange terms for each species (18,19) and solved by numerical methods as described previously (6,13).

Chemical exchange effects were simulated using parameters appropriate for poly-L-lysine (PLK) and dextran (DXT) as a representative polypeptide and sugar, respectively. Simulation parameters were adapted from literature values of CEST z-spectra. PLK was modeled with the following parameters (the subscript A refers in each case to the solvent pool, whereas B refers to the solute pool): $T_{1A} = 3 \text{ sec}, T_{2A} = 2 \text{ sec}, T_{1B} = 1 \text{ sec}, T_{2B} = 30$ msec, fractional populations $p_A = 0.99$, $p_B = 0.01$, resonant offset frequencies $\delta \omega_A = 0$ ppm $\delta \omega_B = 3.5$ ppm, exchange rate r_{BA} = 140 Hz (1,22). DXT was modeled with the following parameters: $T_{1A} = 3$ sec, $T_{2A} = 2$ sec, $T_{1B} = 1$ sec, $T_{2B} = 30$ msec, $p_A = 0.99$, $p_B = 0.01$, $\delta\omega_A = 0$ ppm $\delta\omega_B = 1.2$ ppm, $r_{BA} = 1$ kHz (17,21). For spin-locking simulations, the locking

field amplitude (SLA, denoted in $\text{Hz} = \frac{1}{2\pi} = \frac{1}{2\pi}$) was varied over a range achievable experimentally, corresponding to frequencies of 13.5 Hz to 10 kHz. To calculate CEST zspectra, the frequency offset was varied from −10 to +10 ppm and assumed an RF amplitude of 1μ T and duration 8 sec.

Selected parameters of interest were varied, including chemical exchange rate, chemical shift, concentration, and main field strength, to assess their effects. The effect of main field $(B₀)$ was varied from 0.54 Tesla to 9.4 Tesla (a resonant frequency range of 23 to 400 MHz). The effects of increasing chemical exchange rate were modeled by varying r_{BA} from 0.1 to 50 kHz. The effects of chemical shift were explored by comparing a typical hydroxyl chemical shift of 1.2 ppm at $r_{BA} = 1$ kHz to a typical amide shift and rate ($\delta \omega_B = 3.5$ ppm and $r_{BA} = 140$ Hz). The effects of concentration of the labile species were evaluated by varying the values of the pool size (p_B) from 0.001 to 0.05.

A three-pool model was also evaluated to illustrate the effects of multiple exchanging sites for substances such as poly-L-threonine (PLT) or chondroitin sulfate (CS) that carry both NH and OH groups. Two exchanging pools at two different offset frequencies and exchange rates were modeled in equal concentrations (each 0.01 as above, using the same offset and exchange rates). Exchange between pools B and C was ignored.

(b) ANALYSES

CEST contrast was quantified in terms of MTR_{asym} calculated from the predicted z-spectra using Equation 1. To analyze and compare these with $R_{1\rho}$ effects, a novel metric was devised. Previously, Chopra et al. (23) derived an expression for exchange effects on R_{10} that in our case can be simplified to

$$
R_{1\rho} = R_{2A} + p_B \left[R_{2B} + \frac{r_{BA} \delta \omega_B^2}{r_{BA}^2 + \delta \omega_B^2 + \omega_1^2} \right] \tag{2}
$$

where R_{2A} , R_{2B} are the transverse relaxation rates of the A (water) and B (exchanging) pools, respectively. Parameters p_B , r_{BA} , $\delta \omega_B$ denote the exchanging pool fractional population, exchange rate from the B to the A pool, and chemical shift respectively, and ω_1 is the spin-locking field amplitude. This is valid when $r_{BA}>>R_{2B}>R_{1B}$ as is always the case here. If we consider three measurements of $R_{1\rho}$ taken at three different values of ω_1 (low (\approx 0), high ($\approx \infty$), and a selected value ω_1), we can define an Exchange Rate Contrast (*ERC*) as

$$
ERC\left(\omega_{1}\right)=4^{*}\frac{\left(R_{1\rho}\left(\omega_{1}=0\right)-R_{1\rho}\left(\omega_{1}\right)\right)^{*}\left(R_{1\rho}\left(\omega_{1}\right)-R_{1\rho}\left(\omega_{1}-\infty\right)\right)}{\left(R_{1\rho}\left(\omega_{1}=0\right)-R_{1\rho}\left(\omega_{1}=\infty\right)\right)^{2}}\tag{3}
$$

Using [2] and [3] we obtain the following simple expression

$$
ERC\left(\omega_{1}\right)=\frac{4\alpha^{2}}{\left(1+\alpha^{2}\right)^{2}}\quad\text{[4]}
$$

where $\alpha^2 = \frac{\omega_1^2}{3\omega_0^2}$ and

$$
3\omega_e^2{=}r_{\scriptscriptstyle BA}^2{+}\delta\omega_{\scriptscriptstyle B}^2\quad \text{[5]}
$$

The locking frequency ω_e is characteristic of the compound and corresponds to the field at which the dispersion of $R_{1\rho}$ with ω_1 shows an inflection point so the rate of change of $R_{1\rho}$ with locking field takes on its maximum negative value. Equation 4 has a straightforward interpretation that provides the theoretical basis for modulating the contrast in exchange imaging. The ERC is smooth, has a single peak, and by including the factor 4 in equations [3] and [4], it has a maximum value of 1 when $\alpha = 1$, i.e. when the selected locking field satisfies $\omega_1 = \sqrt{r_{BA}^2 + \delta \omega_B^2}$. The ERC decreases rapidly at higher and lower locking fields and the width of the overall response is again determined by the ratio α. For amides, even at moderate field, the chemical shift dominates the expression $r_{BA}^2 + \delta \omega_{B}^2$. Typically $r_{BA}^2 \approx (140 \text{ Hz})^2$ or less while a shift of 3.5ppm corresponds to $\delta \omega_B^2 \approx (6598 \text{ rad} \cdot \text{sec}^{-1})^2$ at 7T (1,15). For hydroxyls, at 7T the exchange rate is much more significant (typically $r_{BA}^2 \approx (1000 \text{ Hz})^2$ or more (27) and the chemical shift is smaller (\approx 1.2 ppm). The locking frequencies to obtain maximum ERC of 1 at 7T are then 394 and 1050 Hz for hydroxyl and amide protons respectively. When one of these locking frequencies is selected, the ERC for the other species is much reduced (≈ 0.43 for hydroxyls when amide frequency selected, and for amides when hydroxide frequency is selected). Thus by choosing the locking field we can emphasize different species according to their exchange rates.

Actual measurements of $R_{1\rho}$ are not essential to this analysis. A simpler quantity can be defined, similar to the ERC, but based on a comparison of image intensities rather than explicitly on relaxation rates, which behaves in a similar fashion. We name this the Exchange-Weighted Image Contrast or EWIC

$$
EWIC\left(\omega_{1}\right) = 4^{*} \frac{\left(\left(S\left(\omega_{1}=0\right)-S\left(\omega_{1}\right)\right)^{*}\left(S\left(\omega_{1}\right)-S\left(\omega_{1}=\infty\right)\right)\right)}{\left(S\left(\omega_{1}=0\right)-S\left(\omega_{1}=\infty\right)\right)^{2}} \tag{6}
$$

Here, *S* is the spin lock signal acquired for a fixed locking time. EWIC requires only 3 images and by comparing signals at different locking fields, different exchanging populations may be emphasized. Below, values of ERC and EWIC are computed and compared to MTR_{asym} for both experimental and simulated data.

(c) Experiments

Solutions of varied concentrations of polypeptides and sugars were created in 0.6 ml plastic tubes in $1\times$ phosphate-buffered saline (PBS) and titrated to pH 7.4. All chemicals were purchased from Sigma-Aldrich (St. Louis, MO).

Samples of poly-L-threonine (PLT, P8077), poly-L-arginine (PLR, P4663), and poly-Llysine (PLK, P7890) at 10 mg/mL of molecular weight ~15 kDa were prepared. These three polypeptides have previously been identified in a study of a wide range of compounds as having particularly large CEST effects by McMahon et al. (2). Their structures are shown in Figure 2 and were generated from the National Institutes of Health PubChem database [\(http://pubchem.ncbi.nlm.nih.gov/](http://pubchem.ncbi.nlm.nih.gov/)). PLT contains one backbone NH₂ and a single side-chain OH exchangeable group per sub-unit at chemical shifts of \sim 3.5 and \sim 0.8 ppm respectively. PLR contains two side-chain guanidyl NH2 groups and one backbone NH per sub-unit at \sim 1.8 and \sim 3.7 ppm respectively. PLK contains one side-chain NH₂ and a backbone NH group at ~3.6 ppm per lysine sub-unit.

Samples of 40% (wt/wt) glucose (G8270), and dextran (D9260) were created both in distilled water and in $1 \times PBS$ at pH 7.4. Samples of 10% (wt/wt) chondroitin sulfate (CS, C4384) were also created in water and PBS. Samples in PBS were diluted by half four times to create different concentrations. Samples in water were titrated to a pH range of 3 through 11 in 5 steps.

 R_{1p} and CEST data and images were acquired at 9.4T (Varian Medical Systems, Palo Alto, CA). Sample temperature was monitored and maintained at 37 °C. R_{1p} values at each locking field amplitude were measured by varying the time of the locking pulse (SLT) between 20 ms and 1 sec, and these data were then fit to a monoexponential decay. R_{10} dispersion was evaluated by varying the locking field amplitude (SLA) between 150 Hz to 10 kHz (corresponding to ω₁ between 943 and 62,840 rad.sec⁻¹). CEST data were obtained using an 8 second 1μ T saturation pulse with TR of 20 sec. The RF frequency offset was varied between +6 and −6 ppm in 0.1 ppm increments.

Magnetic resonance images were also acquired of the samples in a water bath. The R_{10} imaging sequence consisted of a SL pulse as described by Witschey et al. followed by a fast spin-echo acquisition (12). Images at each of 10 values of SLT were acquired as before at twelve spin lock amplitudes between 150 Hz and 10 kHz. Other imaging parameters included: $FOV = 25 \times 25 \times 1$ mm, matrix = $64 \times 128 \times 1$, TR = 4 sec, TE = 10 ms, ETL = 8, $NEX = 1$.

RESULTS

Figures 3.a and b show the effects of field strength on $R_{1\rho}$ dispersion obtained from simulations of the coupled Bloch equations assuming other relaxation rates (R_1, R_2) remain fixed. The degree of dispersion of $R_{1\rho}$ increases strongly with field, especially when exchange is moderately fast. Figures 3.c and d show the ERC values calculated from the

dispersion curves using Eq. 3. The locking frequency at which maximum contrast occurs scales with B_0 as the resonant frequency differences of the exchanging species increase.

Figure 3.e shows the dramatic effect of field on the separation of the water and metabolite peak of the polypeptide in CEST. The plots of MTR_{asym} for the peptide (Fig. 3.g) are very similar above 64 MHz with MTR_{asym} \approx 0.75. Figure 3.f demonstrates how the small chemical shift and rapid exchange rate of the sugar hydroxyl result in a broad coalesced CEST water peak below 200 MHz. Above 200 MHz, the OH shoulder becomes more pronounced, leading to a more distinct MTR_{asym} as shown in Figure 3.h. The MTR_{asym} peak approaches the actual chemical shift only as the main field increases above 2T. The variations with field reflect the direct effects of saturation on the water, which decrease as the chemical shift frequency difference increases.

Figure 4 demonstrates the predicted effects of increasing chemical exchange rates at 9.4T. $R_{1\rho}$ at low ω_1 increases dramatically as the exchange rate increases until a maximum is reached, after which $R_{1\rho}$ decreases. Figure 4 also shows that the ERC is maximized when ω_1 (SLA) is chosen appropriately. The locking field at which maximum contrast occurs increases with the chemical exchange rate to a degree that depends on the chemical shift.

Figures 4.e and 4.f show simulated z-spectra and MTR_{asym} for the polypeptide and sugar for increasing exchange rates. As the exchange rate increases from the slow to the intermediate regimes, the two peaks coalesce and result in a much-reduced MTR_{asym} (at this 1μ T irradiation strength). When the chemical shift is reduced from 3.5 to 1.2 ppm (polypeptide vs. sugar), this drop occurs at even lower exchange rates.

Figures 4.g and 4.h show trends for ERC and CEST contrasts, respectively. Spin lock contrast was simulated for DXT and PLK using a fixed locking field amplitude of 5 kHz, shown as the red vertical line in 4.c and 4.d, and exchange rates of between 100 Hz and 10 kHz. The results show spin lock contrast increasing with the exchange rate, whereas CEST contrast for the same materials and range of exchange rates decreases with increasing exchange rate. This simulation demonstrates simply that ERC may be best suited for examining materials undergoing rapid exchange compared to CEST methods that provide greater contrast for the same materials in slower exchange regimes.

The effects of increasing the concentration of the exchanging pool are shown in Figure 5. The $R_{1\rho}$ values increase monotonically with the concentration, whereas the ERC plots show complete insensitivity to concentration as expected. The magnitude of the CEST contrast varies non-linearly with the amount of solute present under the chosen conditions of irradiation. At moderate concentrations, MTR_{asym} changes little even for large changes in the pool size fraction.

Figure 6 highlights the relative contributions of two different chemically exchanging species. $R_{1\rho}$ dispersion curves of the individual and combined contributions of NH and OH at typical exchange rates confirm that OH exchange may dominate $R_{1\rho}$ dispersion in mixtures at typical rates and fields. Figure 6 also shows that by selecting appropriate values of ω1, contrast may be dominated by −OH exchange, but the relative contribution of peptide exchange can be manipulated. The simulated CEST z-spectrum and MTR_{asym} for the threepool system clearly show separate sugar and polypeptide exchanging peaks with only a mild conflation near 2.2 ppm on the MTR_{asym} plot.

Figure 7 shows measured CEST z-spectra and MTR_{asym} for glucose (GLU), dextran (DXT), and chondroitin sulfate (CS) at 5% (wt/wt) concentration. Dextran shows a dominant peak at \sim 1.2 ppm with a smaller peak near 2.8 ppm, attributed to residual sites on the molecule that are not cross-linked to other glucose subunits. This 2.8 ppm peak is much larger in the GLU

z-spectrum, which features three distinct exchangeable OH sites. The CS z-spectrum and MTR_{asym} show a dominant peak at ~ 0.8 ppm that is attributed to exchangeable OH sites and a smaller peak near ~3.5 ppm visible only at high concentrations that is attributed to the single NH⁺ site on the CS molecule. Figure 7 also shows $R_{1\rho}$ dispersion profiles for the same samples. Fitted values for R_2 , $R_{1\rho}(\omega_1 = \infty)$, and exchange rate (r_{BA}) for the Chopra model (Eq. 2) are given in Table 1.

Figure 8 shows the measured effects of varying the concentration of dextran. The dominant CEST peak at 1.2 ppm becomes more distinct as the concentration decreases but MTRasym does not vary linearly in proportion to concentration. By contrast, $R_{1\rho}$ shows a nearly linear dependence on concentration over the range of locking fields in which exchange contributions are significant. Exchange rates and other fitted parameters from the Chopra model are given in Table 1.

Figure 9 shows the effects of pH on the dispersion of $R_{1\rho}$ for CS. The values of $R_{1\rho}$ at low locking field were clustered together but at higher fields the curves separate, with the higher pH samples showing less rapid dispersion. The inflection point of the dispersion curve moves to lower frequencies as pH drops along with exchange rates. The fitted exchange rates, other model parameter values, and confidence intervals are given in Table 1 and show a sharp increase in exchange rate above pH 7.

Maps obtained by calculating EWIC using Eq. 6 at each pixel are shown in Figure 9.b. An initial short-TE, T₂-weighted image was recorded at SLA =150 Hz (ω_1 = 942.4 rad. sec⁻¹) and $SLT = 20$ msec. Four sample tubes are shown in clockwise order from $pH=5$ at the top, pH=7 at right, pH=9 at bottom and pH=11 at left. The phantoms appear equally bright in this initial image. Next, EWIC maps were calculated starting with the nearest acquired locking field value to the midpoint of the pH 5 dispersion curve in Figure 9.a for $S(\omega_1)$ in equation 6. The next sub-figure uses the next larger acquired value of ω_1 for $S(\omega_1)$ and so on, to produce the results in Figure 9.b. The sample with maximum contrast cycles from the low pH sample to the high pH sample as ω_1 varies. Artifacts are visible that may be attributed to variations in ω_1 and B_0 across the samples. Note how the pH = 5 and 7 samples are brighter at low ω_1 and the pH = 9 and 11 samples are emphasized at high ω_1 . Thus, by using Eqn. 6 and varying ω_1 exchange rate selective images are produced. However, whether pH changes of practical significance can be detected by such means remains to be established.

Figure 10 shows CEST z-spectra of the solutions of PLK, PLR, and PLT. PLK has a zspectrum "peak" apparent at ~3.5 ppm. The PLR polypeptide carries a guanidyl amine $(gNH₂)$ group in addition to backbone groups that resonate between 1.8 and 3.7 ppm. The PLT group carries an amide group resonating near 3.5 ppm and also a hydroxyl exchange site at 1 ppm, giving a wide spread in saturation effects as shown in the MTR_{asym} plot in Figure 10.b.

Figure 10 also shows the corresponding $R_{1\rho}$ dispersions for the polypeptides. PLR showed the largest dispersion and a fitted exchange rate to the Chopra model of ~930 Hz. PLK and PLT showed much smaller dispersions with estimated exchange rates ~70 Hz and ~3200 Hz respectively. These and other fitted parameters are given in Table 1. Note that no attempt was made to separate the individual contributions of the multiple exchanging species in PLR and PLT so the fitted rates are apparent average exchange rates (24,25).

Figure 10.d shows EWIC maps for the three polypeptides. The image acquired with a locking field of 150 Hz and $SLT = 20$ msec shows PLK, PLT, and PLR with nearly equal brightness. The image acquired with locking field 300 Hz shows PLK clearly, while the other samples are suppressed. The color bar indicates how closely the selected frequency

matches the apparent chemical exchange rate, with unity representing a good match. PLK has an average value of about 0.55 (if the banding artifact at the top of the phantom is left out of the measured region of interest). When the locking field is 1200 Hz, only the PLR sample is visible with an average value of 0.5. At 7 kHz, only the PLT with average image intensity of 0.95 is clearly seen.

Figure 10 also shows the corresponding CEST images resulting from selective irradiation at \pm 3.69 ppm. All peptide species studied here have significant MTR_{asym} near 3.69 ppm with PLK and PLT showing greater saturation effects and MTR_{asym} than PLR. When the saturation occurs at ± 1.8 ppm, only the PLR sample has significant contrast. At ± 0.8 ppm, the PLR and PLT samples show greater contrast than PLK, consistent with the z-spectra.

DISCUSSION

The studies described above illustrate several distinct features of both spin locking and CEST techniques and provide a basis for understanding some of the factors that affect contrast in exchange-dependent imaging. They confirm that both methods are sensitive to chemical exchange effects, especially at high fields. Additionally, some characteristic limitations of each technique are demonstrated, so that each may provide a complementary approach to obtaining chemical exchange-based contrast. The chemical shift frequency difference and exchange rate are the most important factors affecting measured signals and their relative values determine the precise behavior of both SL and CEST contrast.

Imaging is being performed at ever-greater magnetic field strengths both clinically and in pre-clinical research, and field strength is clearly of great importance for exchange imaging, but affects each method differently. Variations in the values of $R_{1\rho}$ with locking field depend explicitly on B_0 because spin locking is sensitive to the time scales of the local fields that promote relaxation, and in an exchanging system these depend on the frequency separation of the exchanging species and the lifetimes spent in each location. Increasing B_0 in SL experiments results in larger magnitude of dispersive effects that make the effects of exchange relatively more important, depending on the exchange rate. Conversely, when the saturating RF pulse in CEST has an ideal narrow-band frequency content and avoids direct water saturation, the CEST signal change does not depend explicitly on the chemical shift. Increasing B_0 clarifies where the exchanging peak occurs, but not the magnitude of the effects on the water resonance. Such effects are seen only indirectly as the resulting zspectrum suffers from the effects of pulse bandwidth and direct saturation that must be optimized by varying the irradiation power. Thus whereas spin lock dispersion may increase with field, CEST contrast becomes relatively field independent.

Our simulations and experiments confirm that each technique shows distinct chemical exchange contrast that is dependent on the exchange regime in which the experiment is performed. Each technique has highest sensitivity in different regimes. The specific effects of increasing chemical exchange rates on $R_{1\rho}$ and CEST contrasts are also distinct. For example, once an appropriate value for ω_1 is chosen there is a monotonic increase in SL contrast with exchange up to very high rates (-10 kHz) , whereas CEST contrast decreases significantly over the same exchange range. The ERC and EWIC are maximized when ω_1 is chosen appropriately which depends on the chemical exchange rate and shift as predicted by Equation 5. This implies also that more distinct contrast may be obtained from substances with smaller chemical shifts if exchange rates are sufficiently separated and sufficient RF power is available.

CEST contrast decreases as exchange rates approach the intermediate and fast exchange regimes. For example, our simulation data for hydroxyl exchange in dextran show that

MTRasym varies little with increasing exchange rate at rates below 1 kHz, but contrast dramatically reduces at rates above 5 kHz. When exchange is very fast compared to the chemical shift frequency difference, line broadening adversely affects the ability to produce difference images with CEST as the exchanging peak coalesces with the water resonance. For example, at rates greater than 1 kHz, the OH peak broadens into the water peak and the corresponding MTR_{asym} reduces to near zero. Thus, an increase in exchange rate reduces contrast in CEST. For spin locking techniques, fast exchange between the water and metabolite site promotes signal dephasing that may be reduced with sufficient locking field ω_1 . The larger this dephasing effect is, the greater the contrast enhancement that may be obtained. However, at extremely high rates (>10 kHz) $R_{1\rho}$ also loses sensitivity to ω_1 .

In a mixture of more than one type of exchanging species, the relative contributions to R_{10} dispersion will again depend on the locking field ω_1 . The combined OH and NH mixture considered above shows that OH exchange will dominate $R_{1\rho}$ dispersion in mixtures, consistent with the findings of Reddy et al. who noted a larger contribution from OH exchange to CEST than from NH sites in studies on chondroitin sulfate (4). The influence of exchange rate in SL compared to chemical shift can be understood by reference to Eqs. 3 through 5 and the relative sizes of r_{BA} and $\delta\omega_B$. These expressions demonstrate that exchange rate may be much more important for contrast than chemical shift when the exchange rate is >> chemical shift.

Figures 9 and 10 demonstrate novel images derived from spin-locking acquisitions in which the contrast emphasizes the presence of exchanging protons with specific exchange rates, as opposed to CEST techniques that emphasize contrast from particular chemical shifts. By appropriate selection of the spin lock field amplitude, maximal contrast can be derived from nuclei whose exchange rates occur around the values predicted by the theory given in Equations 3-5 where $\alpha = 1$. When the R_{1p} dispersion is different for different species, so that the critical locking frequency varies, the contrast may be manipulated by selecting different locking fields.

In conclusion, we have demonstrated with simulations and experiments on model systems that CEST and spin lock images are sensitive to chemical exchange and the influence of factors including the size of the metabolite pool, exchange rate and pH, and field strength. CEST may be particularly suited to exploring contrast from systems in the slow-tointermediate exchange regime and with low metabolite concentrations, but CEST contrast does not provide a clear indication of the concentration of a species under the irradiation conditions considered here. Spin-lock acquisitions and dispersion are well suited to the intermediate-to-fast regime where other approaches, such as CEST or CPMG, may be technically difficult. $R_{1\rho}$ -weighted images show contrast that scales with concentration, but novel metrics are suggested that could highlight disparate species by a combination of exchange rate and offset frequency rather than chemical shift alone. Each technique demonstrates strengths in a particular exchange regime, and therefore the two techniques may be thought of as providing complementary information about chemically exchanging systems.

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Exchanging Pool (B)

Figure 1.

Model of chemical exchange between a large pool of free water protons (A) and smaller pools of exchangeable protons B and C. The rate r_{AB} represents the exchange rate from free water to the exchangeable proton site, and r_{BA} is the reverse rate. The relaxation rates R_1 and R_2 and resonant frequency offsets $\delta\omega$ are the assumed independent parameters for each site and are distinguished by an appropriate subscript. The total magnetization $M_{0A} + M_{0B}$ + M_{0C} =1. There is negligible presumed communication between pools B and C, and a twopool model is obtained by simply removing pool C.

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Figure 2.

Structures of units of molecules in which proton chemical exchange occurs. 2.a. Poly-L-lysine with exchangeable amide (RC(O)NR'R) groups that resonate collectively near 3.69 ppm from water.

2.b. Poly-L-arginine with exchanging guadinyl NH^+ (gNH₂) groups.

2.c. Poly-L-threonine with exchanging NH⁺ and [−]OH groups.

2.d. Glucose molecule (CID: 5793) with 3 distinct hydroxyl (−OH) exchanging sites.

2.e. Chondroitin Sulfate (CID: 24766), a sulfated glycosaminoglycan, that carries multiple exchanging [−]OH and a single NH⁺ site.

2.f. Dextran (CID: 4125253), a poly-glucose molecule that is formed from multiple glycosidic linkages at the α-1,6 or α-1,3 sites on the glucose sub-unit and carries exchanging −OH sites that resonate near 1.2 ppm.

Figure 3.

Simulated effect of moving to higher main field on chemical exchange mediated R_{10} dispersion for peptides (PLK) and sugars (DXT), assuming other relaxation rates do not vary. (T_{1A} = 3 sec, T_{2A} = 2 sec, T_{1B} = 1 sec, T_{2B} = 30 msec, p_A = 0.99, p_B = 0.01, Δω_B = 3.5 ppm or 1.2 ppm, r_{BA} = 140 Hz for PLK or 1kHz for DXT)

3.a. Peptide R_{1p} dispersion vs. applied locking field amplitude SLA $\left(\frac{m}{2\pi} = \frac{m}{2\pi} \right)$ with increasing main field/frequency.

3.b. Sugar $R_{1\rho}$ dispersion.

3.c & 3.d. Simulated Exchange Rate Contrast (ERC) using Eq. 3. Note that the contrast enhancement scales with main field.

3.e & 3.f. CEST z-spectra for PLK (left) and DXT (right).

3.g & 3.h. CEST MTR_{asym} for PLK and DXT, respectively. The field dependence reflects direct saturation effects on water that reduce as the chemical shift frequency difference increases.

Figure 4.

Simulated effect of changes in chemical exchange rate (r_{BA}) on $R_{1\rho}$ dispersion and CEST contrast. (T_{1A} = 3 sec T_{2A} = 2 sec, T_{1B} = 1 sec, T_{2B} = 30 msec, p_A = 0.99, p_B = 0.01, $\Delta \omega_B$ = 3.5 ppm or 1.2 ppm, $\omega_0 = 2\pi * 400$ MHz)

4.a. Peptide $R_{1\rho}$ dispersion vs applied locking field amplitude SLA for increasing exchange rates.

4.b. Sugar $R_{1\rho}$ dispersion with increasing exchange rate.

4.c & 4.d. Exchange Rate Contrast (ERC) using Eq. 3. The DXT data in 4.d shift more with exchange rate than PLK because the chemical shift of the exchanging proton is smaller. The red vertical line at $2\pi * 5$ kHz is used to generate figure 4.g.

4.e & 4.f. CEST z-spectra with inset MTR_{asym} for PLK (left) and DXT (right). Note how CEST contrast sensitivity decreases rapidly as line broadening causes the metabolite peak to coalesce into the water peak above 10 kHz exchange rates.

4.g & 4.h. ERC contrast at fixed ω_1 (SLA=5 kHz) (See Fig. 4.c and 4.d) compared to CEST contrast for PLK and DXT. Note that SL contrast generally increases with increasing exchange rate in 4.g, and decreases for CEST contrast as shown in 4.h.

Figure 5.

Simulated effect of increasing concentration of metabolite pool on $R_{1\rho}$ dispersion and CEST. (T_{1A} = 3 sec T_{2A} = 2 sec, T_{1B} = 1 sec, T_{2B} = 30 msec, $\Delta \omega_B$ = 3.5 ppm or 1.2 ppm, r_{BA} = 140 Hz or 1 kHz, ω₀ = 2π*400 MHz)

5.a. $R_{1\rho}$ dispersion for different SLA for PLK. Note how dispersion is very small at low metabolite ($p_B < 0.01$) concentrations and the low simulated exchange rate (140 Hz). 5.b. $R_{1\rho}$ dispersion for DXT. Note the greater contribution of exchange to $R_{1\rho}$ dispersion attributed to the higher exchange rate (1 kHz), as compared to the peptide in 5.a.

5.c & 5.d. Simulated Exchange Rate Contrast (ERC) using Eq. 3 for both PLK and DXT. The potential contrast becomes insensitive to pool size in contrast with the CEST method shown in 5.g and 5.h.

5.e & 5.f. CEST z-spectra for PLK (left) and DXT (right). Note how CEST contrast is relatively insensitive to pool size, and large contrast enhancement is available at low concentrations and at low exchange rates.

5.g & 5.h. The corresponding CEST MTR_{asym} plots for PLK and DXT, respectively.

Figure 6.

Demonstration of a 3-pool model of proton relaxation on $R_{1\rho}$ dispersion and CEST contrast. This figure gives the simulated effects of multiple exchanging species and demonstrates the relative effects of slowly exchanging NH+ groups vs. typically faster −OH groups at typical chemical shifts and exchange rates under physiologic conditions.

6.a. The individual and combined predicted contribution of exchange at different sites to R_{10} dispersion from −OH and NH+ sites at typical rates. This demonstrates that for a substance with an equal amount of peptide and sugar, the sugar dispersion contribution may dominate the overall $R_{1\rho}$ dispersion curve, but by suitable selection of locking field, the relative peptide contribution can be modulated.

6.b. Calculated Exchange Rate Contrast (ERC) using Eq. 3.

6.c. & 6.d. CEST z-spectrum and MTRasym for combined 3-pool model. Note how, unlike

 $R_{1\rho}$ dispersion, the effects of each species may be separated by chemical shift.

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Figure 7.

CEST z-spectra, MTR_{asym}, and R_{1p} dispersion for 5% (wt/wt) concentrations of the sugars glucose (GLU), chondroitin sulfate (CS), and dextran (DXT). 7.a CEST z-spectra of CS, GLU, and DXT. 7.b Corresponding MTRasym of the three sugars.

7.c $R_{1\rho}$ dispersion of sugars.

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Figure 8.

Experimental results for solutions of dextran with increasing concentration. 8.a. CEST data for solutions of dextran of concentration 5-40% (wt/wt). Note the dominant peak at 1.2 ppm.

8.b. Corresponding MTR_{asym} for figure 8.a.

8.c. $R_{1\rho}$ dispersion results for the same dextran samples used in 8.a.

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6 pH₅ pH₇ ı 5 pH 9 4 $R_{\gamma\rho}$ [s $^{-1}$ α ['] 3 \overline{c} \mathbf{I} \mathbf{I} \mathbf{I} \mathbf{I} \blacksquare \mathbf{I} \mathbf{I} $\overline{}$ ¹ \mathbf{I} п 1 $\frac{1}{10^2}$ $\frac{10^3}{\gamma B_1^1[Hz]}$ 10^4

Figure 9.

Exchange rate selective imaging and R_{1p} dispersion measurements of 10% CS (wt/wt) with varied pH at 400 MHz.

9.a. Reference "T2-weighted" image (SLT = 20 msec, SLA = 150 Hz). Top phantom = pH

5, Right = pH 7, Bottom = pH 9, Left = pH 11.

9.b. Exchange Weighted Image Contrast (EWIC) with SLA = 1200 Hz.

- 9.c. EWIC with $SLA = 3450$ Hz.
- 9.d. EWIC with SLA = 5250 Hz.
- 9.e. EWIC with $SLA = 8400$ Hz.

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Figure 10.

CEST and exchange rate selective imaging and $R_{1\rho}$ dispersion for three 10 mM polypeptides: PLK, PLR, PLT.

10.a. CEST z-spectra of the three peptides. The PLR, with its gNH2 exchanging groups, has the greatest dispersion of R_{1p} from 3.3 s⁻¹ to 0.4 s⁻¹ and carries a fitted mean exchange rate near 980 Hz. The PLK (NH₂) has much slower fitted exchange rates near 70 Hz and shows a smaller dispersion profile. The PLT has both [−]OH and NH⁺ exchanging sites and shows a small dispersion at much higher exchange rates > 3 kHz.

10.b. Corresponding CEST MTR_{asym}

10.c. $R_{1\rho}$ dispersion profiles for each peptide.

10.d. Exchange rate selective image contrast (EWIC) using Eq. 6.

- **a.** Original image of PLK (left), PLT (right), and PLR (bottom)
- **b.** EWIC with SLA= 300 Hz.
- **c.** EWIC with SLA = 1200 Hz.
- **d.** EWIC with $SLA = 4300$ Hz.

10.e. CEST imaging using Eq. 1.

- **a.** Chemical shift selective image with saturation pulse offset = 3.69 ppm.
- **b.** Chemical shift selective image with pulse offset = 1.8 ppm.
- **c.** Chemical shift selective image with pulse offset = 0.8 ppm.

Table 1

Sugar Exchange Rates and Fitted Parameters:

Table 2

Peptide Exchange Rates and Fitted Parameters:

