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Excitation Wavelength Dependent O₂ Release from Copper(II)-Superoxide Compounds: Laser Flash-Photolysis Experiments and Theoretical Studies

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Abstract

Irradiation of the copper(II)-superoxide synthetic complexes [(TMG₃tren)Cu^{II}(O₂)]⁺ (**1**) and [(PV-TMPA)Cu^{II}(O₂)]⁺ (**2**) with visible light resulted in direct photo-generation of O₂ gas at low temperature (from –40 °C to –70 °C for **1** and from –125 °C to –135 °C for **2**) in 2-methyltetrahydrofuran (MeTHF) solvent. The yield of O₂ release was wavelength dependent: λ_{exc} = 436 nm, φ = 0.29 (for **1**), φ = 0.11 (for **2**), and λ_{exc} = 683 nm, φ = 0.035 (for **1**), φ = 0.078 (for **2**), which was followed by fast O₂-recombination with [(TMG₃tren)Cu^I]⁺ (**3**) and [(PV-TMPA)Cu^I]⁺ (**4**). Enthalpic barriers for O₂ re-binding to the copper(I) center (~ 10 kJ mol⁻¹) and for O₂ dissociation from the superoxide compound **1** (45 kJ mol⁻¹) were determined. TD-DFT studies, carried out for **1**, support the experimental results confirming the dissociative character of the excited states formed upon blue or red light laser excitation.

Copper-containing proteins play a major role in O₂ transport and activation in biology. Thus, Cu^I/O₂ reactions and subsequent transformations are critical in this setting as well as in practical systems.¹ Initial O₂ adducts of copper(I) must form in all cases, including in O₂-carriers, oxygenases (oxygen transfer to the substrate) and oxidases (substrate oxidized by O₂), but these first formed species often further react with other electron/proton sources (which may be the substrate) to give Cu_n-peroxo, Cu^{II}-hydroperoxo^{2,3} or perhaps Cu_n-oxyl^{1b,4} active species or intermediates. In peptidylglycine α-hydroxylating monooxygenase⁵ and dopamine β-monooxygenase,⁶ such O₂ activation occurs at a single copper center. An X-ray structure of a pre-catalytic complex along with chemical⁷ and computational studies,^{4a,8} suggested an end-on bound Cu^{II}-superoxide species serves as the enzyme reactive intermediate effecting substrate hydrogen abstraction, further implicating the (bio)chemical importance of initially formed Cu^I/O₂ 1:1 adducts, i.e., Cu^{II}-superoxide species.

Here, for the first time, we show that O₂ can be photo-ejected directly from the 1:1 mononuclear copper/O₂ compounds [(TMG₃tren)Cu^{II}(O₂)]⁺ (**1**) and [(PV-TMPA)Cu^{II}(O₂)]⁺ (**2**) using either 436 nm or 683 nm pulsed laser light (Scheme 1).

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ASSOCIATED CONTENT

 Supporting Information. Experimental procedures, spectra, explanations, DFT calculations and supporting diagrams. This material is available free of charge via the Internet at <http://pubs.acs.org>.

Interestingly, a different yield for O₂ release was observed with these two excitation wavelengths which is different if compared to the O₂ photo-release found in heme systems, like myoglobin.⁹ Temperature-dependent kinetic and thermodynamic studies have been carried out to elucidate the nature of the barriers and the stability of the species involved in the O₂ binding and dissociation processes. Data are corroborated by DFT calculations that help to a) explain why O₂ photo-release is observed and b) interpret the experimentally observed excitation wavelength dependent quantum yield for the O₂ photo-release through new insights into the evolution of the excited states along the copper-oxygen reaction coordinate. To the best of our knowledge, this is the first time that a direct O₂ photo-ejection from 1:1 copper-superoxide adducts has been shown to occur.

Oxygenation of **3** at low temperature in MeTHF was accompanied by a drastic color change of the solution, from colorless to green, forming the previously well characterized compound **1**¹⁰ and leading to the red spectrum shown in Figure 1A. Oxygenation of **4** at low temperature also yielded the previously characterized mononuclear copper/O₂ species **2** (see Supporting Information for UV-visible spectra).¹¹ Cleavage of the copper-oxygen bond was, then, induced upon laser excitation of **1** and **2** ($\lambda_{\text{exc}} = 436$ or 683 nm) as shown by the transient absorption spectral data collected after laser excitation, for **1**. These spectra were in complete agreement with that expected for O₂ photo-release from **1** to yield **3** (Figure 1B) and from **2** to yield **4** (see Supporting Information). The products of the reaction (**3**, O₂ and **4**, O₂, respectively) were excitation wavelength independent, although the quantum yields differed markedly: $\phi = 0.29$ for **1** and $\phi = 0.11$ for **2** ($\lambda_{\text{exc}} = 436$ nm), $\phi = 0.035$ for **1** and $\phi = 0.078$ for **2** ($\lambda_{\text{exc}} = 683$ nm). The appearance of the products, **3** and **4**, occurred within the instrument response time indicating an O₂ time release of less than 10 ns.

The follow-up thermal reaction of [(TMG₃tren)Cu]¹⁺ (**3**) with O₂ led to the formation of the initial compound **1** as is shown in Figure 1B. Kinetic parameters for O₂ coordination to **3** were quantified based on microsecond time scale data. Thus, a plot of the observed rate constants versus the O₂ concentration under pseudo-first-order conditions (excess of O₂) revealed a linear correlation that allowed the determination of the second-order rate constants for O₂ coordinating to **3**, i.e., $k_{\text{O}_2} = 2.1 \times 10^6 \text{ M}^{-1}\text{s}^{-1}$ at -80 °C. For the same temperature, this compares to $k_{\text{O}_2} = 6.6 \times 10^5 \text{ M}^{-1}\text{s}^{-1}$ for **4** (Table 1).

The linear plots of k_{obs} vs [O₂] had a positive intercept that was indicative of the presence of an equilibrium between the reacting species, O₂ and **3** (see Supporting Information). Such a positive intercept was not observed for the coordination of O₂ to **4**, instead, indicating a quantitative formation of **2** from **4** and O₂. Consequently, rate constants for O₂ dissociation from **2** and equilibrium constants for the reaction between **4** and O₂ to give **2** could not be determined here. However, we were able to determine the equilibrium constant at several temperatures in MeTHF solvent through benchtop titration experiments for the binding of O₂ to **3**, to give **1** (Table 1 and Supporting Information). Equilibrium constant values were also determined from laser experiments as follows. In pseudo-first-order conditions, the rate law for O₂ binding to **3** is expressed by the equation $k_{\text{obs}} = k_{\text{O}_2} [\text{O}_2] + k_{-\text{O}_2}$ where k_{obs} is the observed rate constant, k_{O_2} is the second-order rate constant for the binding between **3** and O₂, and $k_{-\text{O}_2}$ is the first-order rate constant for the dissociation reaction of O₂ from **1** (see section 6 of Supporting Information). The values of $k_{-\text{O}_2}$ and k_{O_2} were determined from laser experiments, as a function of temperature through which the equilibrium constants were determined from the ratios $k_{\text{O}_2} / k_{-\text{O}_2}$. Van't Hoff analysis (see Section 5 of the SI) of the equilibrium constants determined with the two different methods (titration experiments and laser experiments) led to the same thermodynamic parameters within the experimental errors and are consistent with values found in a previous report by Roth and co-workers (Table S1 and Figures S2 and S5). Furthermore, equilibrium constants found in this work follow a trend with solvent dielectric constant (ϵ) that was previously established.¹² The

equilibrium constant should favor the superoxide adduct as ϵ increases because of the stabilization of the charge separation present in **1**. In fact, the equilibrium constant for the formation of **1** (K_{O_2}) determined here at -60 °C fits well into a linear correlation together with the previously determined K_{O_2} values in DMF (3030 and 4340)¹² and in chlorobenzene (216)¹³ at -60 °C as a function of ϵ (see Supporting Information).

A comparison of activation and thermodynamic parameters determined in this study with those previously reported for the [(TMPA)Cu^{II}(O₂)]⁺ adduct in MeTHF using [(TMPA)Cu^I(CO)]⁺ and the “flash-and-trap” method are also given in Table 1 (see Chart 1 for structure of ligands). This complex has been very well studied and it is the ‘parent’ ligand of PV-TMPA.^{1a,13,14} The “flash-and-trap” experiments, previously employed for [(L)Cu^I(CO)]⁺ (L = ligand) compounds, allowed characterization of O₂ binding to copper(I) after CO photo-release through competitive coordination of CO and O₂.^{13,14} The kinetic data obtained through the direct O₂ photo-ejection method described here are more straightforward to analyze compared to that of the “flash-and-trap” method where the competitive binding of CO needs to be taken into account. Furthermore, in fast time scale studies of heme-copper oxidases, it has been shown that the presence of CO and starting with a metal-CO adduct may interfere or alter the mechanism or rate of O₂ binding.¹⁵ The activation parameters found for the compounds studied here are quite similar to those previously determined by the flash-and-trap method, providing strong evidence for the reliability of the new method we have employed here to study the reactivity of mononuclear copper compounds with O₂.

TMG₃tren, PV-TMPA, and TMPA offer an analogous coordination sphere to the copper ion, all being tetradentate chelating ligands to study under the same experimental conditions (solvent, temperature, etc). The activation enthalpy found for the binding of O₂ to **3** and **4** falls within the same range (~ 10 kJ mol⁻¹). On the basis of the crystal structure of the starting compound [(PV-TMPA)Cu^I]⁺,¹¹ the coordination within **4** mostly likely also includes an interaction between the copper(I) ion and the O-atom of the pivalamido group. As a consequence, one would expect a higher activation enthalpy for the reaction between O₂ and **4** compared with that between O₂ and **3** as the Cu(I)-O interaction needs to be “disrupted” by O₂ coordination to **4** but not for **3**. Since the ΔH^\ddagger values for the binding of O₂ to **3** and **4** fall, instead, into the same range, this suggests a quite weak interaction for the Cu^I-O_(carbonyl) coordination in **4**. The activation entropy estimated for the reaction involving O₂ coordination to **3** and to **4** is, instead, smaller for the latter. This suggests a mechanism where O₂ coordination to **4** leads to a “highly ordered” transition state where both O₂ and the pivalamido O-atom are interacting with the copper center; for O₂ reacting with **3**, there is of course no pivalamido group present.

The activation enthalpy and entropy for O₂ coordination to [(TMPA)Cu^I]⁺ previously determined (Table 1) are smaller and less negative, respectively, compared with those found for **3** and **4**. This can be interpreted on the basis of a stronger Cu-O₂ interaction in the transition state for [(TMPA)Cu^I]⁺ compared to that for **3** and **4** due to an “easier” spatial approach of O₂ to the copper(I) in [(TMPA)Cu^I]⁺. In fact, the presence of guanidino groups which extend out away from the copper and its ligands in **3**, and of the Cu^I-O_(carbonyl) coordination, in **4**, would support this hypothesis. The less negative activation entropy found for the coordination of O₂ to [(TMPA)Cu^I]⁺ could reflect a smaller molecular reorganization occurring upon O₂ binding to [(TMPA)Cu^I]⁺ due to the absence of guanidino groups or specific Cu(I)-O interactions in [(TMPA)Cu^I]⁺ compared with **3** and **4**. Similar arguments can be used to interpret the difference between the activation enthalpy found for the O₂ dissociation from [(TMPA)Cu^{II}(O₂)]⁺ with the “flash-and-trap” method and those found, here, for **1** and **2**, although the large activation entropy found for O₂ dissociation from [(TMPA)Cu^{II}(O₂)]⁺ seems unclear.

TD-DFT calculations are in line with the previously assigned electronic ground state for **1**.^{10,12,16} In this rather peculiar electronic structure, the central copper ion is in a d^9 configuration and coordinated to a superoxide ligand. The singly occupied MOs are of copper $3d_{z^2}$ and $O_2^{\cdot-}-\pi^*_v$ character. The orthogonality of these two orbitals leads to a $S = 1$ ground state multiplicity in which the spin in both SOMOs are aligned parallel (see Supporting Information). In a spin-unrestricted description, the highest occupied spin-down orbital has mainly oxygen π^*_σ -character and it is bonding with respect to the Cu^{II} -superoxide Cu-O bond. The lowest unoccupied orbitals in the spin-down manifold are the empty partner orbitals of the two SOMOs. Importantly, the unoccupied $3d_{z^2}$ orbital is strongly σ -antibonding with respect to the Cu-O bond. Excitation from the bonding π^*_σ -based orbital to the antibonding d_{z^2} orbital corresponds to a ligand-to-metal charge transfer (LMCT) excitation that formally leads to a $Cu(I)-^3O_2$ electronic configuration. Importantly, this excitation leads to a dramatic weakening of the Cu-O bond to the point that the excited state becomes dissociative (Figure 2).

As it is evident from Figure 3, there is an avoided crossing of the d_{z^2} and π^* orbitals upon elongation of the Cu-O bond, resulting in a change from a triplet $Cu(II)$ superoxide ground state [d^8] $d_{z^2}^1\pi^{*3}$ to a triplet $Cu(I)-^3O_2$ ($d^{10}\pi^{*2}$) state at the dissociation limit. The triplet ground state potential energy surface of **1** (Figure 2, green line) shows a minimum at a Cu-O distance of about 1.9 Å. The calculated excited state energy at the same Cu-O bond distance was 18,843 cm^{-1} (530 nm) for the d-d and 21,635 cm^{-1} (462 nm) for the LMCT transition, consistent with the experimentally observed electronic transitions for these states. Moreover, the character of both excited states at a Cu-O bond distance of 1.9 Å is dissociative (Figure 2). The LMCT excited state (blue line) crosses the ground state at a Cu-O distance of about 4.5 Å. As the dissociative LMCT state crosses the d-d excited states (shown in red) there is an opportunity for the system to cross from one of the d-d excited surfaces to the dissociative LMCT surface. Hence, there can also be O_2 dissociation following d-d excitation, provided that these states live long enough to reach the crossing regime. The exact crossing probability will depend on many details the discussion of which is outside the scope of this work. Given the smaller oscillator strengths for the d-d absorptions and the finite probability for surface hopping, much lower quantum yields are theoretically predicted for d-d excitations. This is in agreement with the observations for O_2 photo-release observed experimentally following excitation of **1** with either red ($\phi_{683} = 0.035$) or blue light ($\phi_{436} = 0.29$). The theoretical results are also consistent with the activation enthalpy for the O_2 dissociation from **1** observed experimentally ($\Delta H^\ddagger_{\text{experim}} = 45 \text{ kJ mol}^{-1}$ vs $\Delta H^\ddagger_{\text{comput}} = 67 \text{ kJ mol}^{-1}$).

Finally, the crossing between the ground state (green) and LMCT (blue) surfaces explains the fact that an association barrier is observed for O_2 -rebinding. The calculated barrier from TD-DFT ($\sim 24 \text{ kJ mol}^{-1}$) is in the right ballpark but slightly overestimates the experimentally measured barrier ($\sim 10 \text{ kcal mol}^{-1}$).

Summarizing, we report here the first example of a photodissociation of molecular oxygen from cupric-superoxide complexes, thus also representing a new approach to study the kinetics and the thermodynamics of formation of 1:1 ligand-copper(I)/ O_2 compounds. Copper-oxygen bond breaking is induced in [based orbital to the antibonding TMG_3tren based orbital to the antibonding Cu^{II} based orbital to the antibonding O_2 based orbital to the antibonding] $^+$ and [based orbital to the antibonding PV-TMPA based orbital to the antibonding Cu^{II} based orbital to the antibonding O_2 based orbital to the antibonding] $^+$ through laser excitation either into the LMCT band, using 436 nm light, or into the d-d electronic transition, using 683 nm light. Interestingly, the quantum yield for O_2 release was wavelength dependent. TD-DFT studies elucidated the O_2 photo-release event occurring upon irradiation with red light on the basis of a) population of a molecular orbital ($3d_{z^2}$) that

has strong σ -antibonding character along the Cu-O bond and b) energy surface crossing between the d-d and the LMCT excited states to lead to O₂ release. Such findings add new insights into the observed wavelength dependent Cu-O₂ photochemistry which differs markedly from that observed with hemes, where for example, the O₂ adduct of myoglobin releases O₂ with a quantum yield of 0.3 following Soret ($\lambda_{\text{exc}} = 488$ nm) or Q ($\lambda_{\text{exc}} = 580$ nm) band excitation.⁹ Formation and decay of [(TMG₃tren)Cu^I]⁺ and [(PV-TMPA)Cu^I]⁺ formed *in situ* have been observed and both activation and thermodynamic parameters for the Cu/O₂ reactions have been determined. Additional experimental studies are on their way to further characterize the excited states involved in the copper-oxygen bond breaking process using ultra-fast laser spectroscopy.

Supplementary Material

Refer to Web version on PubMed Central for supplementary material.

Acknowledgments

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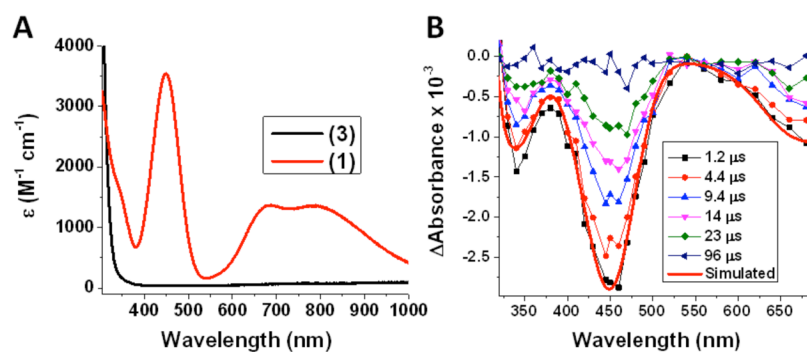


Figure 1.

(A) Absorption spectrum of **1** (red line) obtained from oxygenation of **3** (black line) at 218 K in MeTHF. (B) Transient absorption difference spectra collected at the indicated delay times after 436 nm laser excitation (15 mJ/pulse, 8–10 ns fwhm) of **1** in MeTHF at 218 K. Overlaid in red on the experimental data is a simulated spectrum ($\text{Abs}(\mathbf{3}) - \text{Abs}(\mathbf{1})$).

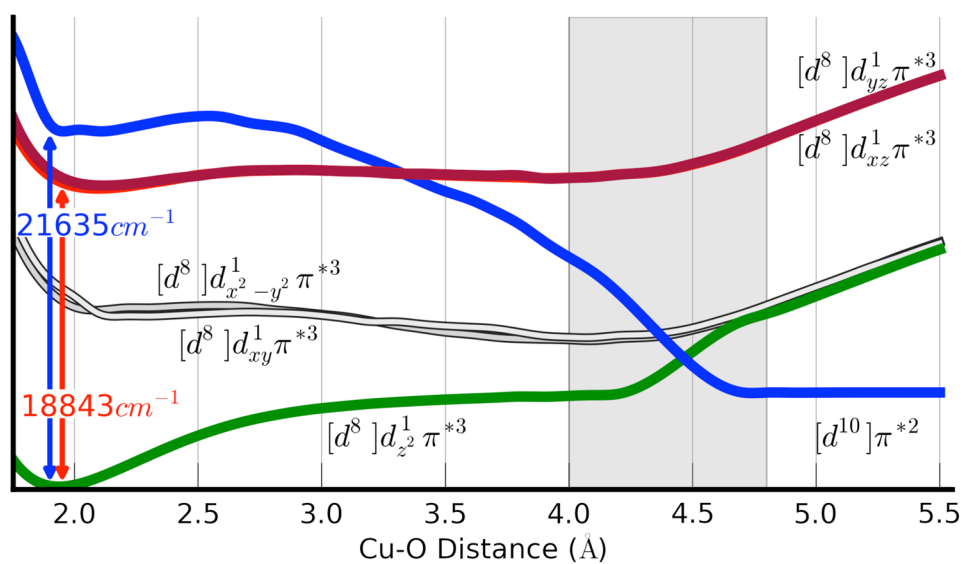


Figure 2. TD-DFT calculated excited state potential energy surfaces (PESs) as a function of copper-oxygen bond distance.

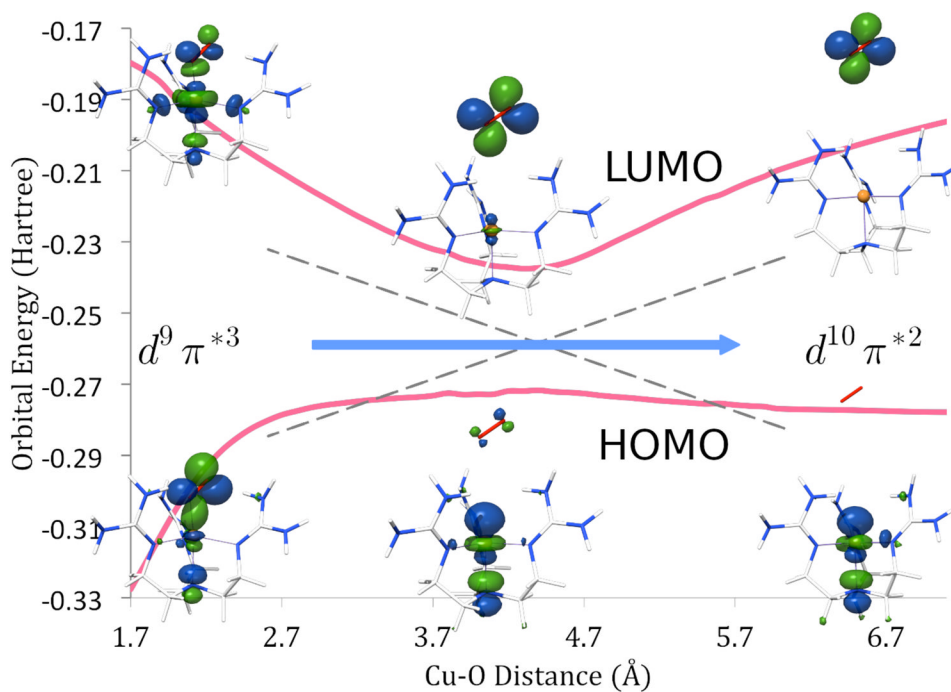
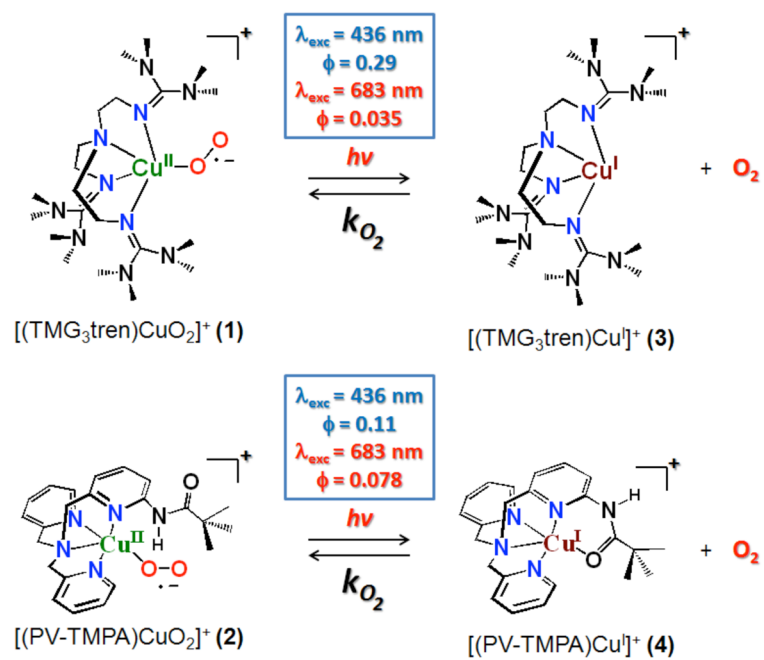


Figure 3. TD-DFT calculated energy and shape of the beta HOMO and beta LUMO orbitals as a function of copper-oxygen bond distance.



Scheme 1.

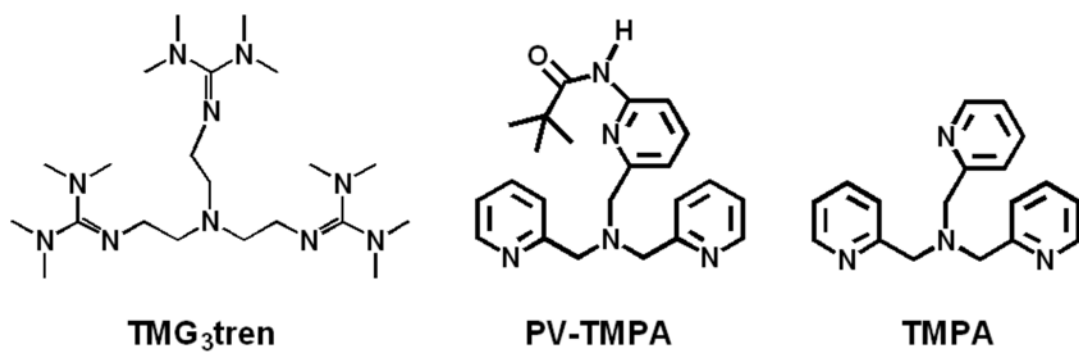


Chart 1.

Table 1Comparison of Kinetic and Thermodynamic Parameters for O₂ binding and dissociation for [(L)Cu]⁺ adducts.

TMG ₃ tren ^c			
	<i>k</i> _{O₂}	<i>k</i> _{-O₂}	K _{O₂}
ΔH^\ddagger or ΔH^0 <i>a</i>	10 ± 6	45 ± 7	-40 ± 2
ΔS^\ddagger or ΔS^0 <i>b</i>	-70 ± 26	42 ± 34	-134 ± 11
<i>k</i> or K 25°C	(2.7 ± 1.2) · 10 ⁷	(1.5 ± 0.8) · 10 ⁷	≈ 1
<i>k</i> or K -80°C	(2.1 ± 1.0) · 10 ⁶	(5.2 ± 2.0) · 10 ²	(6.3 ± 1.9) · 10 ³

PV-TMPA ^c			
	<i>k</i> _{O₂}	<i>k</i> _{-O₂}	K _{O₂}
ΔH^\ddagger or ΔH^0 <i>a</i>	9 ± 1	-	-
ΔS^\ddagger or ΔS^0 <i>b</i>	-97 ± 7	-	-
<i>k</i> or K 25°C	(4.8 ± 2.8) · 10 ⁷	-	-
<i>k</i> or K -80°C	(6.6 ± 3.5) · 10 ⁵	-	-

TMPA ^d			
	<i>k</i> _{O₂}	<i>k</i> _{-O₂}	K _{O₂}
ΔH^\ddagger or ΔH^0 <i>a</i>	7.62	58.0	-48.5
ΔS^\ddagger or ΔS^0 <i>b</i>	-45.1	105	-140
<i>k</i> or K 25°C	1.3 · 10 ⁹	1.3 · 10 ⁸	15.4
<i>k</i> or K -80°C	1.4–1.6 · 10 ⁸	240	6.5 · 10 ⁵

^a ΔH , kJ mol⁻¹^b ΔS , J K⁻¹ mol⁻¹^c In MeTHF, this work^d In THF, determined through flash-and-trap method.¹³ Values for [(TMPA)-Cu] in MeTHF have been found to be the same as in THF within experimental errors.