



OPEN

# Preparation of Few-Layer Bismuth Selenide by Liquid-Phase-Exfoliation and Its Optical Absorption Properties

Liping Sun<sup>1</sup>, Zhiqin Lin<sup>1</sup>, Jian Peng<sup>1</sup>, Jian Weng<sup>1,3</sup>, Yizhong Huang<sup>2</sup> & Zhengqian Luo<sup>2</sup>

SUBJECT AREAS:

LASERS, LEDS AND LIGHT  
SOURCES

FIBRE LASERS

SYNTHESIS AND PROCESSING

Received  
4 November 2013Accepted  
8 April 2014Published  
25 April 2014

Correspondence and  
requests for materials  
should be addressed to  
J.W. (jweng@xmu.  
edu.cn) or Z.Q.L.  
(zqluo@xmu.edu.cn)

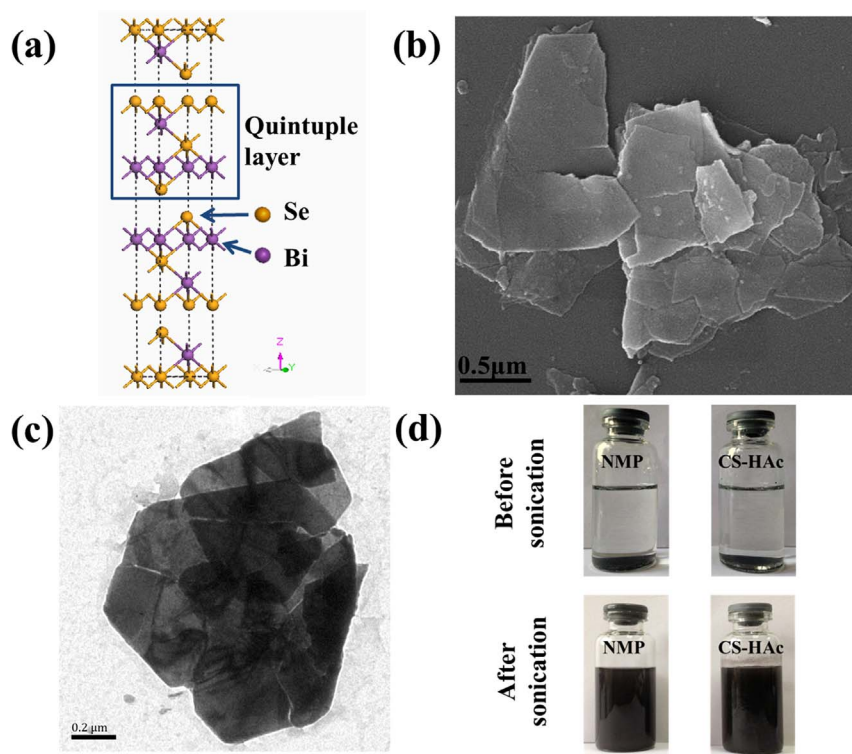
<sup>1</sup>Department of Biomaterials, College of Materials, Xiamen University, Xiamen 361005, China, <sup>2</sup>Institute of Optoelectronic Technology, Department of Electronic Engineering, Xiamen University, Xiamen 361005, China, <sup>3</sup>ShenZhen Research Institute of Xiamen University, Shenzhen 518057, China.

Bismuth selenide ( $\text{Bi}_2\text{Se}_3$ ), a new topological insulator, has attracted much attention in recent years owing to its relatively simple band structure and large bulk band gap. Compared to bulk, few-layer  $\text{Bi}_2\text{Se}_3$  is recently considered as a highly promising material. Here, we use a liquid-phase exfoliation method to prepare few-layer  $\text{Bi}_2\text{Se}_3$  in N-methyl-2-pyrrolidone or chitosan acetic solution. The resulted few-layer  $\text{Bi}_2\text{Se}_3$  dispersion demonstrates an interesting absorption in the visible light region, which is different from bulk  $\text{Bi}_2\text{Se}_3$  without any absorption in this region. The absorption spectrum of few-layer  $\text{Bi}_2\text{Se}_3$  depends on its size and layer number. At the same time, the nonlinear and saturable absorption of few-layer  $\text{Bi}_2\text{Se}_3$  thin film in near infrared is also characterized well and further exploited to generate laser pulses by a passive Q-switching technique. Stable Q-switched operation is achieved with a lower pump threshold of 9.3 mW at 974 nm, pulse energy of 39.8 nJ and a wide range of pulse-repetition-rate from 6.2 to 40.1 kHz. Therefore, the few-layer  $\text{Bi}_2\text{Se}_3$  may excite a potential applications in laser photonics and optoelectronic devices.

Topological insulators (TIs) as interesting insulators now have become the rising star in physics, chemistry and materials fields because they are insulating in the bulk phase but possess exotic metal surface state as a result of the combination of spin-orbit interactions and time-reversal symmetry<sup>1-3</sup>. In the past few years, some research groups<sup>4-6</sup> achieved great success in the prediction and experimental confirmation of TIs, including  $\text{Bi}_2\text{Se}_3$ ,  $\text{Bi}_2\text{Te}_3$  and  $\text{Sb}_2\text{Te}_3$ , which have a large band gap and a single Dirac cone. Especially, the remarkable band gap of  $\text{Bi}_2\text{Se}_3$  is approximately up to 0.3 eV (equivalent to 3600 K) that is much larger than the room temperature energy scale<sup>4</sup>. It means that  $\text{Bi}_2\text{Se}_3$  is able to exhibit topological insulator behavior at room temperature, which is considered as a promising topological system with a good application prospect<sup>7</sup>. Recently, most researchers paid attention to the physical basis<sup>8-10</sup>, synthesis method<sup>11,12</sup> and exploration of the nanostructure<sup>13-15</sup> of TIs. However, it is worth noting that topological properties of  $\text{Bi}_2\text{Se}_3$  as three-dimensional (3D) TIs are often covered up by the bulk state due to high carrier density<sup>5,16</sup>. Therefore, it is necessary to prepare two-dimensional (2D)  $\text{Bi}_2\text{Se}_3$  from its 3D bulk materials in order to acquire the superior performance for some potential applications.

$\text{Bi}_2\text{Se}_3$  possesses stacked layers of laminated structure that are held together by weak van der Waals interactions. Each layer is one quintuple layer (QL) and the five atoms are covalently bonded together along the z axis in the order of Se-Bi-Se-Bi-Se (Fig. 1a). The thickness of each layer is about 0.96 nm<sup>17</sup>. It is possible to exfoliate bulk  $\text{Bi}_2\text{Se}_3$  into few-layer nanosheets due to the weak interaction between layers. Up to date, bottom-up synthesis and top-down exfoliation are two main methods to prepare 2D nanomaterials<sup>3</sup>. Bottom-up synthesis approach is used to obtain single-layer or fewer layer 2D nanomaterials by a chemical reaction from the atomic or molecular scale synthesis<sup>18-21</sup>. 3D materials held together by weak van der Waals forces can be exfoliated into thin flakes by the methods of mechanical or chemical exfoliation<sup>22-24</sup>, which is a top-down process. Therefore, it is possible to obtain few-layer QLs from bulk  $\text{Bi}_2\text{Se}_3$  with “graphene-inspired” exfoliation methods because bulk  $\text{Bi}_2\text{Se}_3$  possesses the graphene-like layered structure. Liquid-phase exfoliation has been used to produce single-layer or few-layer graphene because it is easier and more convenient than other methods. Furthermore, the as-obtained graphene could form colloidal dispersions in solvents<sup>24,25</sup>. Therefore, we attempted to exfoliate bulk  $\text{Bi}_2\text{Se}_3$  by liquid-phase exfoliation method to prepare few-layer  $\text{Bi}_2\text{Se}_3$  in solutions.

As a new type of Dirac material, TIs with the unique energy-band structure can induce some fantastically electronic and optical properties<sup>26</sup>, opening up many new applications, such as superconductors<sup>27</sup> and ultrafast lasers<sup>28,29</sup>. Nowadays, these researches are focused on pulsed lasers due to their versatile applications in range



**Figure 1** | Preparation and exfoliation of as-synthesized bulk  $\text{Bi}_2\text{Se}_3$ . (a) Schematic of rhombohedral layer structure held together by weak van der Waals interactions in  $\text{Bi}_2\text{Se}_3$ . Each QL consists of five covalently bonded atomic sheets along the z axis in the order of Se-Bi-Se-Bi-Se. (b), (c) SEM and TEM images of as-synthesized bulk  $\text{Bi}_2\text{Se}_3$ . (d) Photographs of  $\text{Bi}_2\text{Se}_3$  dispersed in NMP and CS-HAc before and after sonication.

finding, medicine, laser processing, remote sensing and telecommunications<sup>30</sup>. In the field of pulsed lasers, passive Q-switched fiber laser for generating short and large-energy laser pulse is one of most effective ways because of their significant advantages of compactness, simplicity, and flexibility in design<sup>31</sup>. The key element in the passive Q-switched fiber laser is an excellently saturable absorber. Therefore, researchers have never stopped to seek for new saturable absorbers (e.g. semiconductor<sup>32</sup>, carbon nanotubes<sup>33</sup>, graphene<sup>33,34</sup>). Compared with bulk materials<sup>28,29</sup>, one can expect that the few-layer nanomaterials would possess the more excellent performance of saturable absorption, and could be a potentially saturable absorber. Therefore, we are strongly motivated to develop the pulsed fiber lasers Q-switched with few-layer  $\text{Bi}_2\text{Se}_3$  as the saturable absorber.

Here, N-methyl-2-pyrrolidone (NMP), the more promising organic solvent to exfoliate 2D layered materials<sup>24</sup>, is used to exfoliate bulk  $\text{Bi}_2\text{Se}_3$  for producing few-layer  $\text{Bi}_2\text{Se}_3$  (Supplementary Fig. S1). Another is chitosan acetic solution (CS-HAc), which possesses the low-toxic, good-biocompatible and environmentally friendly properties<sup>35</sup>. Meanwhile, we also investigated the optical absorption characterization of as-prepared few-layer  $\text{Bi}_2\text{Se}_3$  dispersed in solvents in visible light region, and saturable-absorption performance of few-layer  $\text{Bi}_2\text{Se}_3$  thin film in near infrared region. At last, few-layer  $\text{Bi}_2\text{Se}_3$  was successfully used as the new fiber-compatibly saturable absorber to attain passive Q-switched fiber laser at 1.53  $\mu\text{m}$  wavelength.

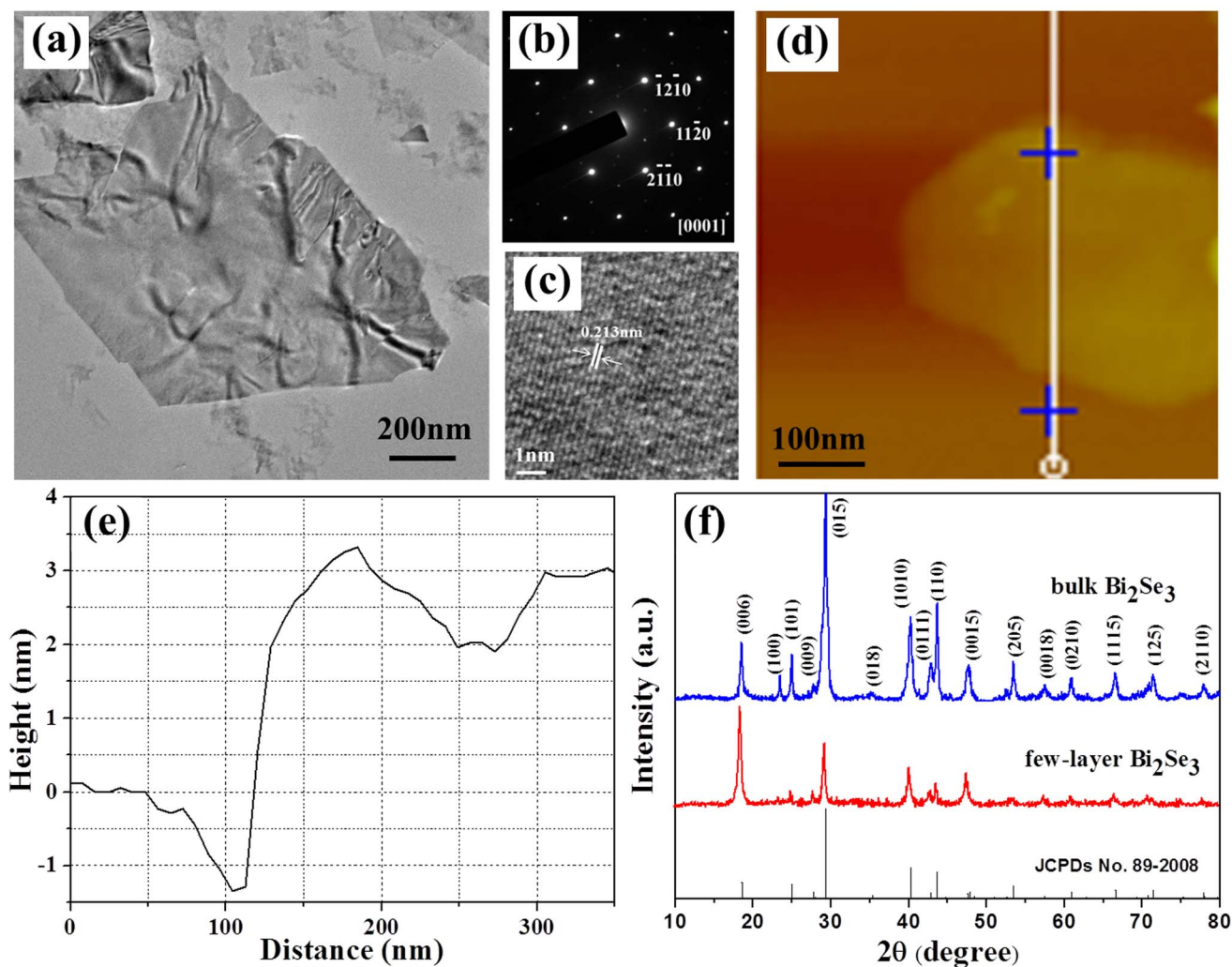
## Results

**Preparation and exfoliation of bulk  $\text{Bi}_2\text{Se}_3$ .** Bulk  $\text{Bi}_2\text{Se}_3$  was prepared by hydrothermal synthesis and characterized by X-ray diffraction (XRD, Supplementary Fig. S2a). All the labeled peaks can be readily indexed to rhombohedral  $\text{Bi}_2\text{Se}_3$  (JCPDS no. 89-2008). The scanning electron microscope (SEM) image in Fig. 1b and transmission electron microscope (TEM) image in Fig. 1c show that the as-synthesized bulk  $\text{Bi}_2\text{Se}_3$  exhibits sheet-like structure with a wide size distribution, and is easily to aggregate together.

The thickness of as-synthesized bulk  $\text{Bi}_2\text{Se}_3$  is about 40–100 nm determined by atomic force microscopy (AFM) (Supplementary Fig. S2c and d).

We further exfoliated as-synthesized bulk  $\text{Bi}_2\text{Se}_3$  with solution-phase exfoliation method, which is similar to the exfoliation of graphite in NMP and CS-HAc<sup>35</sup>. The as-synthesized  $\text{Bi}_2\text{Se}_3$  powders were insoluble in two solvents before sonication (Fig. 1d). After sonication of 30 h, the colors of two solutions were deepened, which means that the exfoliated  $\text{Bi}_2\text{Se}_3$  had been dispersed in these solvents. We also investigated the exfoliation of as-synthesized bulk  $\text{Bi}_2\text{Se}_3$  in other solvents (Supplementary Fig. S3). The result shows that NMP and CS-HAc are the optimal solvents to exfoliate as-synthesized bulk  $\text{Bi}_2\text{Se}_3$ . Therefore, NMP and CS-HAc are selected to investigate the exfoliation of bulk  $\text{Bi}_2\text{Se}_3$ . We further investigated the effect of sonication time on exfoliation of bulk  $\text{Bi}_2\text{Se}_3$  (Supplementary Fig. S4). With increasing ultrasonic time, the color of CS-HAc was deepened, but color is already deep dark in NMP at 2 h, which reveals a better exfoliating effect in NMP. Longer ultrasonic time should produce higher concentration of few-layer  $\text{Bi}_2\text{Se}_3$ . However, it needs more power and time. Therefore, we chose 30 h as the appropriately ultrasonic time because  $\text{Bi}_2\text{Se}_3$  has already been well dispersed in these two solvents, meeting the requirement of following experiments in this study. The exfoliated  $\text{Bi}_2\text{Se}_3$  also presented the Tyndall effect of the colloidal suspension (Supplementary Fig. S5). The result shows that the colloidal suspension of exfoliated  $\text{Bi}_2\text{Se}_3$  in the two solvents is stable.

**Characterization of few-layer  $\text{Bi}_2\text{Se}_3$ .** The TEM image (Fig. 2a) of exfoliated  $\text{Bi}_2\text{Se}_3$  showed that the as-obtained few-layer  $\text{Bi}_2\text{Se}_3$  was extremely thin 2D flake. According to the selected area electron diffraction (SAED) pattern (Fig. 2b), it could be indexed as a 6-fold symmetry [001] zone axis pattern, which is consistent with the layered structure along the z axis. Also, it revealed the single-crystalline nature of the thin 2D flake. Furthermore, the distance



**Figure 2 | Confirmation of few-layer  $\text{Bi}_2\text{Se}_3$  exfoliated in NMP.** (a) TEM image of few-layer  $\text{Bi}_2\text{Se}_3$ . (b) SAED pattern of few-layer  $\text{Bi}_2\text{Se}_3$ . (c) HRTEM image of few-layer  $\text{Bi}_2\text{Se}_3$ . (d), (e) AFM image and the corresponding height profile of few-layer  $\text{Bi}_2\text{Se}_3$ . (f) XRD patterns of few-layer  $\text{Bi}_2\text{Se}_3$  and bulk  $\text{Bi}_2\text{Se}_3$ .

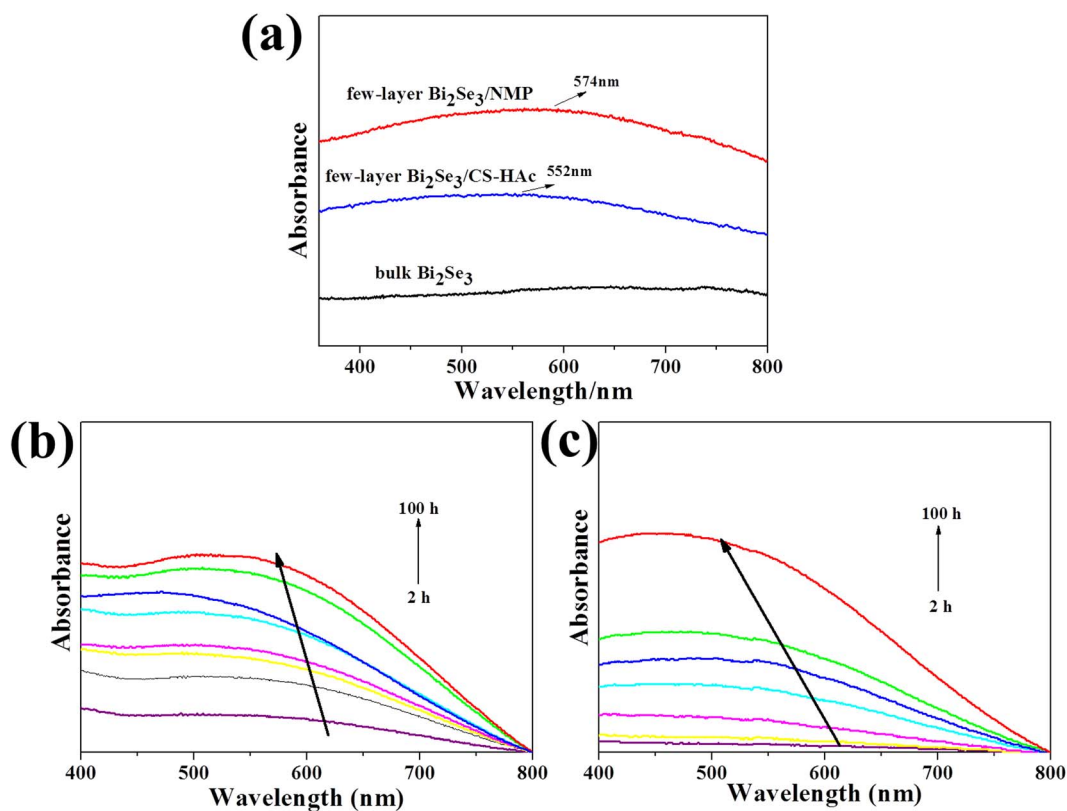
between the adjacent hexagonal lattice fringes investigated by the high-resolution TEM (HRTEM) is 0.213 nm for  $\text{Bi}_2\text{Se}_3$  (Fig. 2c), which is consistent with the lattice space of the (110) plane. The AFM image (Fig. 2d) also shows that the exfoliated  $\text{Bi}_2\text{Se}_3$  is a flake structure and its thickness is about 3–4 nm (Fig. 2e), which nearly equals to 4 layers of  $\text{Bi}_2\text{Se}_3$ <sup>17</sup>. The XRD pattern (Fig. 2f) of few-layer  $\text{Bi}_2\text{Se}_3$  showed a high [006] orientation and some characteristic peaks disappeared compared to bulk  $\text{Bi}_2\text{Se}_3$ , which indicates that bulk  $\text{Bi}_2\text{Se}_3$  had been successfully exfoliated as we expected. At the same time, the bulk  $\text{Bi}_2\text{Se}_3$  has successfully been exfoliated to few-layer  $\text{Bi}_2\text{Se}_3$  in CS-HAc (Supplementary Fig. S6). Besides, Raman spectrum was also used to further confirm the exfoliation of  $\text{Bi}_2\text{Se}_3$  (Supplementary Fig. S7). The A mode of few-layer  $\text{Bi}_2\text{Se}_3$  produced a red shift compared to that of bulk  $\text{Bi}_2\text{Se}_3$ , which could be attributed to the phonon softening<sup>36,37</sup>. Therefore, we successfully prepared few-layer  $\text{Bi}_2\text{Se}_3$  using the solution-phase exfoliation method.

**Optical absorption characterization of few-layer  $\text{Bi}_2\text{Se}_3$ .** The optical absorption properties of few-layer  $\text{Bi}_2\text{Se}_3$  were firstly investigated with ultraviolet-visible (UV-vis) spectra. Interestingly, we found that the dispersion solutions of few-layer  $\text{Bi}_2\text{Se}_3$  produced a broad absorption in the visible light region compared to as-synthesized bulk  $\text{Bi}_2\text{Se}_3$  (Fig. 3a). The UV-vis spectrum of as-synthesized bulk  $\text{Bi}_2\text{Se}_3$  showed a nearly straight line without any absorption peak in

the visible light region. However, few-layer  $\text{Bi}_2\text{Se}_3$  displayed an absorption band at about 552 nm in CS-HAc and 574 nm in NMP, respectively. The appearance of absorption band after exfoliation is remarkable, which might be due to the exfoliation of bulk  $\text{Bi}_2\text{Se}_3$  into nanosheets with a few nanometers thickness. The absorption also increases gradually as the sonication time extended (Fig. 3b and c), which reveals that more few-layer  $\text{Bi}_2\text{Se}_3$  would be obtained with increasing sonication times. The result further suggests that the absorption would be resulted from few-layer  $\text{Bi}_2\text{Se}_3$ .

We further investigate the effect of size and thickness on absorption property. After sonication in NMP, few-layer  $\text{Bi}_2\text{Se}_3$  was separated in different centrifugal speeds (Fig. 4a). The size distribution and corresponding height profile of few-layer  $\text{Bi}_2\text{Se}_3$  collected at three centrifugation speeds were distinguishing. With centrifugal speed increasing, the size of few-layer  $\text{Bi}_2\text{Se}_3$  was decreased from 500 to 100 nm and the thickness was also decreased from 10 to 2 nm, and the maximal absorption wavelength was blue-shifted from 613 to 459 nm (Fig. 4b). The similar result is also obtained in CS-HAc (Supplementary Fig. S8 and S9). The result further suggests that the broad absorption in the visible light region would be resulted from few-layer  $\text{Bi}_2\text{Se}_3$  but not from bulk  $\text{Bi}_2\text{Se}_3$ .

To further investigate the optical absorption properties of few-layer  $\text{Bi}_2\text{Se}_3$ , we used a spin-coating method to prepare  $\text{Bi}_2\text{Se}_3/\text{NMP}$  (few-layer  $\text{Bi}_2\text{Se}_3$  exfoliated in NMP) and  $\text{Bi}_2\text{Se}_3/\text{CS-HAc}$

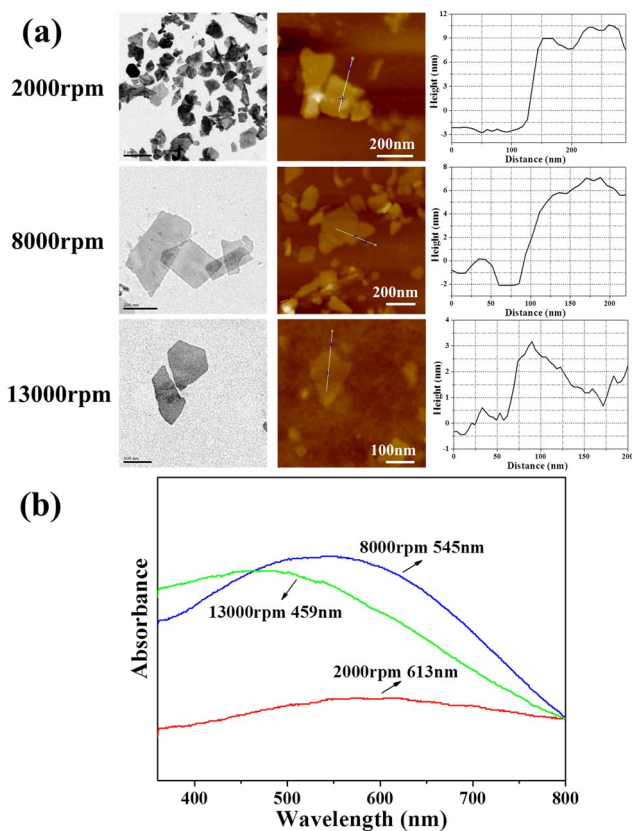


**Figure 3** | UV-vis absorption spectra of few-layer Bi<sub>2</sub>Se<sub>3</sub>. (a) UV-vis absorption spectra of as-synthesized bulk and few-layer Bi<sub>2</sub>Se<sub>3</sub> suspension. (b) UV-vis absorption spectra of few-layer Bi<sub>2</sub>Se<sub>3</sub> suspension prepared with different sonication times in NMP. (c) UV-vis absorption spectra of few-layer Bi<sub>2</sub>Se<sub>3</sub> suspension prepared with different sonication times in CS-HAc. The upper part of resulting suspension of each sample after sonication was collected, and then centrifuged for 30 min at 1000 rpm to receive the supernatant as the measurement solution.

(few-layer Bi<sub>2</sub>Se<sub>3</sub> exfoliated in CS-HAc) thin films on quartz plate, respectively. As shown in Fig. 5a, we measured the linear absorption spectra of the two films by a spectrophotometer scanning from 300 to 2000 nm. One can clearly see that both of the films have the relatively flat transmission curves in the UV-to-near infrared (NIR) region, e.g. the transmittance of Bi<sub>2</sub>Se<sub>3</sub>/NMP varies only from 0.67 to 0.84 in the broad wavelength range of 350 ~ 2000 nm. It indicates that the few-layer Bi<sub>2</sub>Se<sub>3</sub> would be a promising broadband optical material. In order to compare the nonlinear absorption of our few-layer Bi<sub>2</sub>Se<sub>3</sub> with that of bulk Bi<sub>2</sub>Se<sub>3</sub> (>50 layers) previously reported<sup>38,39</sup>, we also used the same Z-scan technique to measure the nonlinear transmission responses of the two few-layer Bi<sub>2</sub>Se<sub>3</sub> films. When the two samples were strongly excited by a femtosecond Ti: sapphire laser with the highest optical intensity of 2.6 GW/cm<sup>2</sup> (Fig. 5b and c), the open-aperture Z-scan transmission curves of Bi<sub>2</sub>Se<sub>3</sub>/NMP and Bi<sub>2</sub>Se<sub>3</sub>/CS-HAc were obtained, respectively. One can obviously see that the two samples possess the saturable absorption, i.e. the optical transmittance is different under differently optical intensity. The modulation depths ( $\delta T$ ) are 3.8% for Bi<sub>2</sub>Se<sub>3</sub>/NMP and 3.7% for Bi<sub>2</sub>Se<sub>3</sub>/CS-HAc, respectively, which is comparable to that of graphene<sup>40,41</sup>. Furthermore, by carefully fitting the curves in Fig. 5b and 5c, the produced saturable intensities ( $I_{sa}$ ) are 53 MW/cm<sup>2</sup> for Bi<sub>2</sub>Se<sub>3</sub>/NMP and 41 MW/cm<sup>2</sup> for Bi<sub>2</sub>Se<sub>3</sub>/CS-HAc, respectively. It is very interesting that the  $I_{sa}$  values are much less than that of bulk Bi<sub>2</sub>Se<sub>3</sub> reported previously<sup>38,39</sup>, mainly benefiting from the few-layer structure of exfoliated Bi<sub>2</sub>Se<sub>3</sub>. In the field of passive Q-switched or mode-locked lasers, the lower  $I_{sa}$  of saturable absorber, the easier the start of Q-switching/mode-locking operation is, implying that few-layer Bi<sub>2</sub>Se<sub>3</sub> might be very helpful for developing the low-threshold Q-switched/mode-locked lasers.

**Generation of Q-switched laser pulses using the saturable absorption of few-layer Bi<sub>2</sub>Se<sub>3</sub>.** As well as known, the optically saturable absorption can be used to efficiently generate the laser pulses by the passive Q-switching or mode-locking techniques<sup>42,43</sup>. The lower  $I_{sa}$  of few-layer Bi<sub>2</sub>Se<sub>3</sub> may be very helpful for developing the low-threshold Q-switched/mode-locked lasers. In this section, to testify the performance of few-layer Bi<sub>2</sub>Se<sub>3</sub>, we will exploit the saturable absorption of few-layer Bi<sub>2</sub>Se<sub>3</sub> to passive Q-switch erbium-doped fiber laser (EDFL) for generating laser pulses. Supplementary Fig. S10 shows the experimental setup of Q-switched EDFL using few-layer Bi<sub>2</sub>Se<sub>3</sub> as a saturable absorber. In order to clearly evaluate the significance of few-layer Bi<sub>2</sub>Se<sub>3</sub> to Q-switching operation, we purposely performed the following control experiments. At first, when as-synthesized bulk Bi<sub>2</sub>Se<sub>3</sub> was deliberately inserted into the laser cavity, we found that the Q-switching operation at 1530.2 nm was extremely unstable with a large pulse-intensity and repetition-rate fluctuation (see the Supplementary Fig. S11 for more details). Moreover, the Q-switching operation has a high pump threshold of 22.1 mW, a broad pulse duration of 22.8  $\mu$ s and a small operating range of pump power (22.1 ~ 67.5 mW). In contrast, a very stable Q-switching operation was produced when the few-layer Bi<sub>2</sub>Se<sub>3</sub> was placed in the laser cavity to replace as-synthesized bulk Bi<sub>2</sub>Se<sub>3</sub> as followed.

As increasing the pump power, we found that the laser with few-layer Bi<sub>2</sub>Se<sub>3</sub> as saturable absorber reached its threshold at the pump power of 9.3 mW only, and the stable Q-switching operation occurred simultaneously. The pump threshold for Q-switching is much lower than that of as-synthesized bulk Bi<sub>2</sub>Se<sub>3</sub> (22.1 mW), and other saturable absorber-based pulsed EDFLs reported previously<sup>29,38,39</sup>, mainly benefiting from the lower  $I_{sa}$  of few-layer



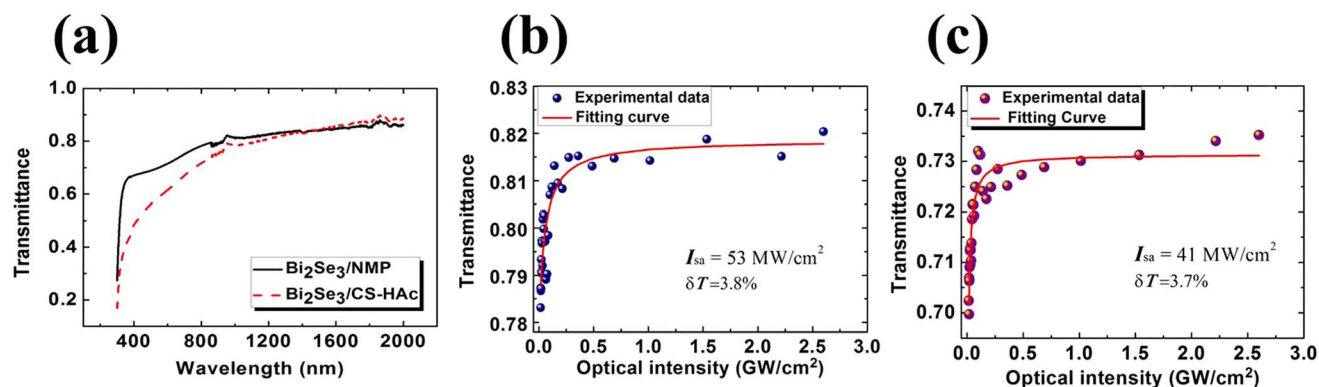
**Figure 4** | Few-layer  $\text{Bi}_2\text{Se}_3$  in NMP collected in different centrifugal speeds. (a) TEM and AFM images, and the corresponding height profiles of few-layer  $\text{Bi}_2\text{Se}_3$  in NMP collected in different centrifugal speeds. Firstly, the stock solution was centrifuged at 2000 rpm for 30 min, and the precipitate was collected as sample one (top). Then, the remaining supernatant was centrifuged at 8000 rpm for 20 min, and the precipitate was collected as sample two (middle). At last, the supernatant collected in second step was further centrifuged at 13000 rpm for 12 min, and the precipitate was collected as sample three (bottom). (b) UV-vis absorption spectra of few-layer  $\text{Bi}_2\text{Se}_3$  in NMP. The absorption band was blue-shifted with decreasing thickness and size of few-layer  $\text{Bi}_2\text{Se}_3$ .

$\text{Bi}_2\text{Se}_3$ . Fig. 6 summarizes the output characteristics of the Q-switched pulses. Fig. 6a and 6b give the typical oscilloscope trace of the Q-switched pulse trains and the single pulse envelope at the pump power of 87.2 mW, respectively. The Q-switching pulse output with the repetition rate of 23.8 kHz was stable, and no significant pulse jitter was observed on the oscilloscope. The measured

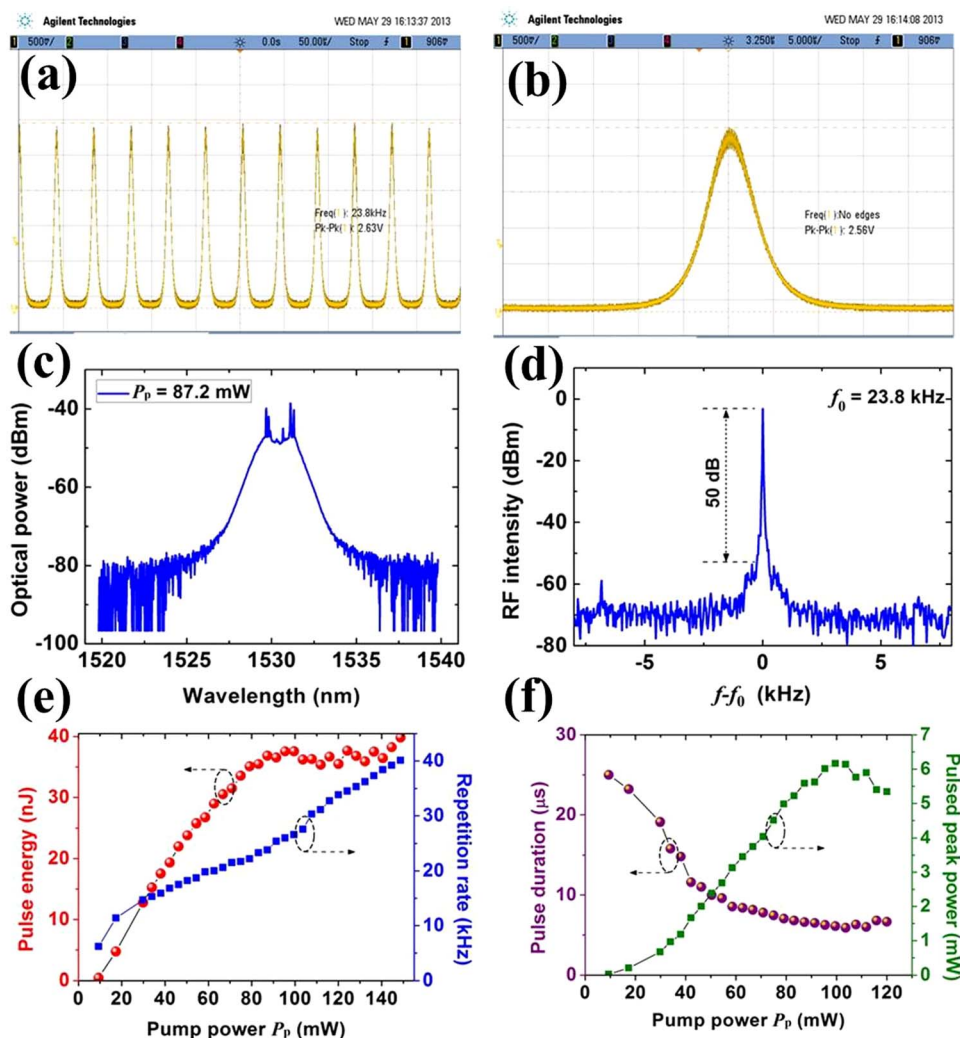
single-pulse envelope shows the good symmetry and has the pulse duration of 5.4  $\mu\text{s}$ . The typical laser spectrum of the Q-switching operation depicted in Fig. 6c has the central wavelength of 1530.3 nm with the 10-dB bandwidth of 2.2 nm. As usually observed in Q-switched fiber lasers<sup>44</sup>, the sideband structure appeared in the optical spectrum was due to the multimode oscillation and the cavity perturbations of Q-switching<sup>44</sup>. As shown in Fig. 6d, we also measured the RF output spectrum of Q-switching pulses at the same pump power of 87.2 mW. The pulse repetition rate is 23.8 kHz. The RF signal-to-noise ratio is more than 50 dB, and the 20-dB RF linewidth is less than 10 Hz (limited by the RF resolution bandwidth of 10 Hz), further indicating the good stability of the Q-switching operation. Moreover, the stability of the Q-switching is excellent in our testing period of 4 h, and the stable Q-switching is available in the large range of pump power (9.3 ~ 150.1 mW), which is superior to that of bulk  $\text{Bi}_2\text{Se}_3$  (22.1 ~ 67.5 mW). Fig. 6e plots the pulse repetition rate and the pulse energy as a function of the pump power. As increasing the pump power from 9.3 to 150.1 mW, one can see that: 1) the pulse repetition rate linearly increases from 6.2 to 40.1 kHz; and 2) the pulse energy monotonically increases in the lower pump power, but slightly saturates after exceeding the pump power of 100 mW. The maximum pulse energy obtained in our experiment is 39.8 nJ, corresponding to the average output power of 1.6 mW at the pump power of 150.1 mW. In addition, we also recorded the evolution of pulse duration in different pump powers. As shown in Fig. 6f, the pulse duration can be significantly narrowed from 24.0  $\mu\text{s}$  to 4.9  $\mu\text{s}$  with the increase of pump power. The pulse duration might be further reduced by shortening the cavity length and optimizing the cavity loss<sup>45</sup>.

## Discussion

In this work, we attempted to exfoliate as-synthesized bulk  $\text{Bi}_2\text{Se}_3$  for preparing few-layer  $\text{Bi}_2\text{Se}_3$  by liquid-phase exfoliation method, and the result shows that it is viable. In the process of preparation, ten solvents were used to exfoliate  $\text{Bi}_2\text{Se}_3$  with same ultrasonic time and concentration in order to find the optimal solvents to exfoliate  $\text{Bi}_2\text{Se}_3$ . With the aid of ultrasound wave, few-layer  $\text{Bi}_2\text{Se}_3$  has successfully been prepared in NMP and CS-HAc. The exfoliation of as-synthesized bulk  $\text{Bi}_2\text{Se}_3$  is attributed to the energy provided by the ultrasound wave which overcomes the van der Waals force between  $\text{Bi}_2\text{Se}_3$  QLs. With the increasing of ultrasonic time, higher concentration of few-layer  $\text{Bi}_2\text{Se}_3$  was produced. However, the increasing amount of few-layer  $\text{Bi}_2\text{Se}_3$  is not obvious after 30 h. Considering the efficiency of preparation, 30 h is selected as the appropriate ultrasonic time to prepare few-layer  $\text{Bi}_2\text{Se}_3$ . In NMP solvent, many materials held by van der Waals forces could be exfoliated to produce 2D nanosheets due to its appropriate surface tension<sup>46</sup>.  $\text{Bi}_2\text{Se}_3$  has a similar structure held together via van der Waals forces between QLs, so it is possible



**Figure 5** | Optical absorption of few-layer  $\text{Bi}_2\text{Se}_3$ . (a) The linear absorption of few-layer  $\text{Bi}_2\text{Se}_3$ . (b) The nonlinear optical absorption (i.e. saturable absorption) of  $\text{Bi}_2\text{Se}_3/\text{NMP}$ . (c) The nonlinear optical absorption (i.e. saturable absorption) of  $\text{Bi}_2\text{Se}_3/\text{CS-HAc}$ .



**Figure 6** | Performance of few-layer Bi<sub>2</sub>Se<sub>3</sub> to passively Q-switch erbium-doped fiber laser. (a) The typical oscilloscope trace of Q-switched pulses at the pump power of 87.2 mW. (b) The single pulse envelope. (c) The typical optical spectrum of Q-switching operation. (d) The RF output spectrum. (e) The pulse repetition rate and the pulse energy vs the pump power. (f) The pulse duration as a function of the pump power.

to obtain few-layer Bi<sub>2</sub>Se<sub>3</sub> after sonication in NMP. Really, the expected results have been obtained as we suppose so. Another aqueous surfactant solution, CS-HAc, was also used to prepare few-layer Bi<sub>2</sub>Se<sub>3</sub> through hydrophobic interaction of main chains of chitosan and the surface of Bi<sub>2</sub>Se<sub>3</sub>. Few-layer Bi<sub>2</sub>Se<sub>3</sub> can stably be dispersed in chitosan dispersion, which is due to the electrostatic repulsion<sup>35</sup> between NH<sub>4</sub><sup>+</sup> in the side chains of chitosan adsorbed on the surface of Bi<sub>2</sub>Se<sub>3</sub>. Therefore, liquid-phase-exfoliation of as-synthesized bulk Bi<sub>2</sub>Se<sub>3</sub> allows production of few-layer Bi<sub>2</sub>Se<sub>3</sub> suspensions in NMP or CS-HAc, which might be a simple and convenient method to prepare few-layer Bi<sub>2</sub>Se<sub>3</sub> for further investigating its properties and exploring the promising applications.

The optical absorption spectrum of few-layer Bi<sub>2</sub>Se<sub>3</sub> in solution exhibits a strong absorption band in the visible light region, which is different from as-synthesized bulk Bi<sub>2</sub>Se<sub>3</sub> without any absorption peak in this region. The few-layer Bi<sub>2</sub>Se<sub>3</sub> with a small size can be considered as a “quantum dot” that would result in quantum confinement, leading to resonance that can be tuned with size. Therefore, the optical absorption of few-layer Bi<sub>2</sub>Se<sub>3</sub> is size-dependent (Fig. 4). The products in this study consist of nanosheets with different sizes and thicknesses, and it is difficult to prepare the sample of fixed size with different thickness or fixed thickness with different sizes. Therefore, for systematic and in-depth investigating the effect of thickness and size on the optical absorption of few-layer Bi<sub>2</sub>Se<sub>3</sub>, it

is necessary to synthesize nanosheets that are monodispersed in both size and thickness, but it remains a challenge by liquid-phase-exfoliation method at this stage.

The blue-shift of UV-vis absorption of few-layer Bi<sub>2</sub>Se<sub>3</sub> with its size decreasing is similar to the results for many semiconductor nanoparticles<sup>47</sup>, which the small dimensions result in differently physical properties compared with their corresponding bulk materials. Therefore, we also studied the optical band gap ( $E_g$ ) of few-layer Bi<sub>2</sub>Se<sub>3</sub> according to their optical absorption spectra in solution. The optical absorption properties of few-layer Bi<sub>2</sub>Se<sub>3</sub> with different size and thickness in solutions were further investigated in UV-Vis-NIR spectral region (Supplementary Fig. S12 and 13). There is not obvious absorption peak in NIR region. The recorded absorption spectra were mathematically processed to acquire the values of  $E_g$ <sup>48</sup>. The optical absorption is calculated using the following equation:

$$(\alpha h\nu)^n = B(h\nu - E_g) \quad (1)$$

where  $\alpha$  is the absorption coefficient,  $h$  is Planck constant,  $\nu$  is the frequency of photon,  $E_g$  is the band gap and  $B$  is a constant. For the direct band gap semiconductor Bi<sub>2</sub>Se<sub>3</sub>,  $n$  is 2. The  $(\alpha h\nu)^2$  vs.  $h\nu$  curves for all samples were shown in Supplementary Fig. S12 and 13. The  $E_g$  of few-layer Bi<sub>2</sub>Se<sub>3</sub> with different sizes and thicknesses obtained by different centrifugation speeds in NMP were determined by extrapolating the straight portion of the plot to the energy axis. The  $E_g$



were 1.22 eV for 2000 rpm, 1.39 eV for 8000 rpm and 1.50 eV for 13000 rpm (more details in supplementary Table S1 and S2), respectively, which is higher than that of as-synthesized bulk  $\text{Bi}_2\text{Se}_3$  (1.08 eV). Meanwhile, the  $E_g$  of few-layer  $\text{Bi}_2\text{Se}_3$  increases with the centrifugation speed increasing (size and thickness of few-layer  $\text{Bi}_2\text{Se}_3$  decreasing), which indicates a blue-shift phenomenon. The reason for this larger  $E_g$  of few-layer  $\text{Bi}_2\text{Se}_3$  might be due to the well-known quantum confinement effect by shifting the conduction and valence band edges in opposite directions<sup>49–51</sup>. It is worth noting that the  $E_g$  of as-synthesized bulk  $\text{Bi}_2\text{Se}_3$  is larger than the theoretical value (0.3 eV) calculated by first-principle electronic structure. The reason is that the as-synthesized “bulk”  $\text{Bi}_2\text{Se}_3$  with a thickness of 40–100 nm and a size of 100–500 nm (Supplementary Figure S2) is smaller than those of the generally bulk  $\text{Bi}_2\text{Se}_3$  (thickness and size  $\geq 10 \mu\text{m}$ ). That is to say, the as-synthesized “bulk”  $\text{Bi}_2\text{Se}_3$  is nanosheet and not the real bulk  $\text{Bi}_2\text{Se}_3$ . The higher  $E_g$  of as-synthesized “bulk”  $\text{Bi}_2\text{Se}_3$  (1.08 eV) is attributed to the quantum size effect as also considered by Goror and Hodes<sup>52</sup>. Therefore, it is reasonable that the theoretical  $E_g$  of  $\text{Bi}_2\text{Se}_3$  is smaller than the experimental  $E_g$  because the  $E_g$  growing is nearly inversely proportional to the lateral size<sup>53</sup>.

For a 2D crystallite, the band gap shift,  $\Delta E_g$ , is described by the equation<sup>54,55</sup>

$$\Delta E_g = \frac{h^2}{4\mu_{xy}L_{xy}^2} + \frac{h^2}{8\mu_zL_z^2} \quad (2)$$

where  $\mu_{xy}$  and  $\mu_z$  are the reduced effective masses of electron-hole pairs in parallel (xy) and perpendicular (z) directions, respectively, and  $L_{xy}$  and  $L_z$  are the corresponding dimensions of the crystallite. For the ideally thin nanosheets,  $L_{xy}$  (0.1–1  $\mu\text{m}$ ) is much larger than  $L_z$  (0.96 nm for  $\text{Bi}_2\text{Se}_3$  QL), so the first term in eq. 2 can be neglected. Consequently, the band gap shift depends only on  $L_z$ . As shown in Supplementary Fig. S12, we can get an approximate  $\Delta E_g = E_g$  (13000 rpm) –  $E_g$ (bulk) = 0.42 eV. Therefore, the calculated  $\mu_z$  is 0.24  $m_e$  ( $m_e$ : electron mass). The Bohr radius R of exciton can be calculated by the following equation<sup>56</sup>

$$R = \frac{\epsilon h^2}{\mu_z \pi e^2} \quad (3)$$

where  $\epsilon$  is the dielectric constant at optical frequencies. The dielectric constant for  $\text{Bi}_2\text{Se}_3$  can be typically set to be  $100\epsilon_0$ <sup>57,58</sup>, where  $\epsilon_0$  is the vacuum permittivity.  $e$  is electron charge,  $1.602 \times 10^{-19}$  C. The calculated R from eq. 3 is about 21.79 nm. Therefore, the calculated R is much larger than the thickness of few-layer  $\text{Bi}_2\text{Se}_3$  at 13000 rpm ( $0.96 \times 2 = 1.92$  nm), suggesting that electron-hole pairs would be physically confined in few-layer  $\text{Bi}_2\text{Se}_3$ . The calculated R values for other few-layer  $\text{Bi}_2\text{Se}_3$  were listed in Supplement Table S1 and S2. It is well-known that semiconductors perform dramatic quantization effect when charge carriers (electrons and holes) are confined by potential barriers to small regions of space<sup>59</sup>. Or equivalently, the thickness of few-layer  $\text{Bi}_2\text{Se}_3$  is less than twice the Bohr radius of excitons in the bulk material. In a word, the blue-shift phenomenon implies that the  $E_g$  would increase with decreasing thickness, especially for the molecularly thin nanosheets by quantum size effect.

Under strong light excitation (Supplement Fig. S14), the electrons in the valence band become depleted while the final state in the conduction band is partially occupied, and further excitation from the valence band is blocked and no further absorption is induced, leading to a saturable absorption effect. The saturable intensity of few-layer  $\text{Bi}_2\text{Se}_3$  thin film is much less than that of bulk  $\text{Bi}_2\text{Se}_3$  (>50 layers). Therefore, the saturable absorption of few-layer  $\text{Bi}_2\text{Se}_3$  is further exploited to Q-switch fiber laser, experimentally confirming the advantage of few-layer  $\text{Bi}_2\text{Se}_3$  as a broadband saturable absorber because the Q-switching operation of as-synthesized bulk  $\text{Bi}_2\text{Se}_3$  was extremely unstable with the large pulse-intensity and repetition-rate

fluctuation. Further exploiting the few-layer  $\text{Bi}_2\text{Se}_3$  with the saturable absorption, we have successfully obtained the few-layer  $\text{Bi}_2\text{Se}_3$ -based passive Q-switched EDFL. Compared with as-synthesized bulk  $\text{Bi}_2\text{Se}_3$ , we have revealed that few-layer  $\text{Bi}_2\text{Se}_3$  is more favorable for stable Q-switching. The reason why few-layer  $\text{Bi}_2\text{Se}_3$  for Q-switched pulsed laser is superior to bulk  $\text{Bi}_2\text{Se}_3$  could be explained as follows. It is well known that most of unique characteristics of topological insulator (including optical and electrical ones) originate from the metallic states on the surfaces or edges. As illustrated in Supplement Fig. S15, because bulk  $\text{Bi}_2\text{Se}_3$  can be exfoliated to many few-layer  $\text{Bi}_2\text{Se}_3$  sheets, in this process the surfaces/edges can be sharply increased. Therefore, one can think that the metallic states of few-layer  $\text{Bi}_2\text{Se}_3$  should be stronger than that of bulk  $\text{Bi}_2\text{Se}_3$ . As is well known, the electrons at the surface, such as metals, are very active with very low surface energy, and they are readily excited by externally electromagnetic (e.g. lightwave) or thermal fields. According to this way, one can easily understand that under light excitation, the surface electrons of few-layer  $\text{Bi}_2\text{Se}_3$  can be transitioned more readily, because few-layer  $\text{Bi}_2\text{Se}_3$  possesses more metallic surfaces/edges in comparison with the bulk one. Thus, the optically saturable absorption of few-layer  $\text{Bi}_2\text{Se}_3$  is more excellent than that of bulk  $\text{Bi}_2\text{Se}_3$ . Also, the few-layer  $\text{Bi}_2\text{Se}_3$  can significantly enlarge the surface-to-volume ratio, and can be considered as a “quantum dot” that would result in quantum confinement. This could lead to the easier occurrence of the saturable absorption which has been partially verified by the lower saturable intensity (53 and 41 MW/cm<sup>2</sup>) in Fig. 5. Therefore, few-layer  $\text{Bi}_2\text{Se}_3$  can generate the stable Q-switching operation compared to the unstable operation with bulk  $\text{Bi}_2\text{Se}_3$ . The Q-switched laser based on few-layer  $\text{Bi}_2\text{Se}_3$  has the low pump threshold of 9.3 mW, the pulse energy of 39.8 nJ, the pulse duration of 4.9  $\mu\text{s}$  and the wide range of pulse-repetition-rate from 6.2 to 40.1 kHz, comparable to those reported fiber lasers Q-switched by other saturable absorbers (e.g. graphene<sup>60</sup>, carbon nanotubes<sup>44</sup>, and semiconductor<sup>38,39</sup>). The promising results might have been due to the unique energy-band structure of few-layer  $\text{Bi}_2\text{Se}_3$ . This performance of the Q-switched laser shows good prospects of few-layer  $\text{Bi}_2\text{Se}_3$  as an excellently saturable absorber in the future.

## Methods

**Synthesis of bulk  $\text{Bi}_2\text{Se}_3$ .** Polyvinyl pyrrolidone (0.9 g) was dissolved in ethylene glycol (EG, 36 mL). Then bismuth oxide powder ( $\text{Bi}_2\text{O}_3$ , 1 mmol), selenium powder (Se, 3 mmol) and ethylenediamine tetraacetic acid powder (4 mmol) were added into above-mentioned EG solution. The resulting suspension was stirred vigorously and subsequently sealed in a steel autoclave. The autoclave was heated to 200 °C in 30 min and maintained this temperature for 20 h. The as-obtained product was collected by high-speed centrifugation, washed several times with deionized water and absolute ethanol, and finally dried at 60 °C for 96 h in an oven.

**Preparation of few-layer of  $\text{Bi}_2\text{Se}_3$ .** The as-synthesized bulk  $\text{Bi}_2\text{Se}_3$  was dispersed in NMP or stock solution of chitosan (0.2 mg·mL<sup>-1</sup>) that was prepared in 0.5% acetic acid aqueous solution at a concentration of 1 mg·mL<sup>-1</sup> by sonication in a sonic bath for 30 h (KQ–250 DB). The upper part of the resulting suspension after leaving to stand for 24 h was collected and centrifuged for 30 min at 1000 rpm. Subsequently, the supernatant was decanted to another centrifuge tube. After centrifuging the supernatant at 10000 rpm for 10 min, the as-obtained product was collected into phials and dispersed in the solvent used above for further characterization.

**Characterization.** Powder X-ray diffraction system (Rigaku Ultima IV XRD) equipped with Cu K $\alpha$  radiation ( $\lambda = 1.542 \text{ \AA}$ ) over the  $2\theta$  range of 10–80° was used to characterize the crystal structure of as-synthesized bulk and few-layer  $\text{Bi}_2\text{Se}_3$ . The sample was prepared by dropping the dispersive solutions on the surface of glass slide which had been etched a groove, then drying with an infrared lamp. Again and again to depositing a film on the fluted glass was named a continuous drop-dry process. SEM images were obtained on LEO-1530 operated at 20 kV. SEM samples were prepared by depositing a small drop of solution on small pieces of silicon wafer and then dried at room temperature. Energy dispersive X-ray spectrum pattern was acquired through spreading as-synthesized  $\text{Bi}_2\text{Se}_3$  powders on sample stage directly. The micrographs of samples were taken using a transmission electron microscope (JEOL JEM-1400, JEM-2100) at an accelerating voltage of 200 kV. To prepare the TEM samples, a small drop of sample was deposited onto copper grids coating with lacey carbon film and then dried under room temperature at atmospheric pressure. AFM images were obtained in the tapping mode in air using an Agilent 5500 atomic force microscope. The samples were prepared by dropping their dispersions on mica



- substrates. Raman spectra (XploRA, Jobin-Yvon) were recorded with a diode laser at the excitation wavelength of 532 nm. The UV-vis absorption spectrum was measured on UV-vis spectrometer (UV-2550, Shimadzu). The UV-vis-NIR absorption spectrum was recorded on a Varian Cary 5000. The linear absorption spectra were measured by a spectrophotometer (Perkinelmer Lambda 7500) scanning from 300 to 2000 nm. The output laser spectrum was monitored by an optically spectral analyzer (Advantest Q8384) with the spectral resolution of 0.01 nm. The pulsed characteristics of this laser were detected by a 10 GHz photodetector (Nortel PP-10G) together with a digital storage oscilloscope (DSO, Agilent MSO7104A) and a radio-frequency (RF) spectrum analyzer (Gwinstek GSP-930).
- Moore, J. Topological insulators: the next generation. *Nat. Phys.* **5**, 378–380 (2009).
  - Moore, J. E. The birth of topological insulators. *Nature* **464**, 194–198 (2010).
  - Kong, D. & Cui, Y. Opportunities in chemistry and materials science for topological insulators and their nanostructures. *Nat. Chem.* **3**, 845–849 (2011).
  - Zhang, H. *et al.* Topological insulators in Bi<sub>2</sub>Se<sub>3</sub>, Bi<sub>2</sub>Te<sub>3</sub> and Sb<sub>2</sub>Te<sub>3</sub> with a single Dirac cone on the surface. *Nat. Phys.* **5**, 438–442 (2009).
  - Xia, Y. *et al.* Observation of a large-gap topological-insulator class with a single Dirac cone on the surface. *Nat. Phys.* **5**, 398–402 (2009).
  - Chen, Y. L. *et al.* Experimental realization of a three-dimensional topological insulator, Bi<sub>2</sub>Te<sub>3</sub>. *Science* **325**, 178–181 (2009).
  - Hsieh, D. *et al.* A tunable topological insulator in the spin helical Dirac transport regime. *Nature* **460**, 1101–1105 (2009).
  - Moore, J. & Balents, L. Topological invariants of time-reversal-invariant band structures. *Phys. Rev. B* **75**, 121306(R) (2007).
  - Hsieh, D. *et al.* Observation of unconventional quantum spin textures in topological insulators. *Science* **323**, 919–922 (2009).
  - Cao, H. L. *et al.* Quantized Hall effect and Shubnikov-de Haas oscillations in highly doped Bi<sub>2</sub>Se<sub>3</sub>: evidence for layered transport of bulk carriers. *Phys. Rev. Lett.* **108**, 216803 (2012).
  - Chen, X., Ma, X. C., He, K., Jia, J. F. & Xue, Q. K. Molecular beam epitaxial growth of topological insulators. *Adv. Mater.* **23**, 1162–1165 (2011).
  - Liu, H. T., Dai, J., Zhang, J. J. & Xiang, W. D. Solvothermal synthesis of Bi<sub>2</sub>Se<sub>3</sub> hexagonal nanosheet crystals. *Adv. Mater. Res.* **236–238**, 1712–1716 (2011).
  - Xiu, F. *et al.* Manipulating surface states in topological insulator nanoribbons. *Nat. Nanotechnol.* **6**, 216–221 (2011).
  - Kong, D. *et al.* Topological insulator nanowires and nanoribbons. *Nano Lett.* **10**, 329–333 (2010).
  - Min, Y. *et al.* Surfactant-free scalable synthesis of Bi<sub>2</sub>Te<sub>3</sub> and Bi<sub>2</sub>Se<sub>3</sub> nanoflakes and enhanced thermoelectric properties of their nanocomposites. *Adv. Mater.* **25**, 1425–1429 (2013).
  - Qi, X. L. & Zhang, S. C. The quantum spin Hall effect and topological insulators. *Phys. Today* **63**, 33–38 (2010).
  - Sun, Y. *et al.* Atomically thick bismuth selenide freestanding single layers achieving enhanced thermoelectric energy harvesting. *J. Am. Chem. Soc.* **134**, 20294–20297 (2012).
  - Li, Y. Y. *et al.* Intrinsic topological insulator Bi<sub>2</sub>Te<sub>3</sub> thin films on Si and their thickness limit. *Adv. Mater.* **22**, 4002–4007 (2010).
  - Peng, H. *et al.* Aharonov-Bohm interference in topological insulator nanoribbons. *Nat. Mater.* **9**, 225–229 (2010).
  - Li, H. *et al.* Controlled synthesis of topological insulator nanoplate arrays on mica. *J. Am. Chem. Soc.* **134**, 6132–6135 (2012).
  - Min, Y. *et al.* Quick, controlled synthesis of ultrathin Bi<sub>2</sub>Se<sub>3</sub> nanodiscs and nanosheets. *J. Am. Chem. Soc.* **134**, 2872–2875 (2012).
  - Shahil, K. M. F., Hossain, M. Z., Goyal, V. & Balandin, A. A. Micro-Raman spectroscopy of mechanically exfoliated few-quintuple layers of Bi<sub>2</sub>Te<sub>3</sub>, Bi<sub>2</sub>Se<sub>3</sub>, and Sb<sub>2</sub>Te<sub>3</sub> materials. *J. Appl. Phys.* **111**, 054305 (2012).
  - Novoselov, K. S. *et al.* Electric field effect in atomically thin carbon films. *Science* **306**, 666–669 (2004).
  - Hernandez, Y. *et al.* High-yield production of graphene by liquid-phase exfoliation of graphite. *Nat. Nanotechnol.* **3**, 563–568 (2008).
  - Lotya, M. *et al.* Liquid phase production of graphene by exfoliation of graphite in surfactant/water solutions. *J. Am. Chem. Soc.* **131**, 3611–3620 (2009).
  - Hasan, M. & Kane, C. Colloquium: topological insulators. *Rev. Mod. Phys.* **82**, 3045–3067 (2010).
  - Zarepour, P. *et al.* Proximity-induced high-temperature superconductivity in the topological insulators Bi<sub>2</sub>Se<sub>3</sub> and Bi<sub>2</sub>Te<sub>3</sub>. *Nat. Commun.* **3**, 1056 (2012).
  - Bernard, F., Zhang, H., Gorza, S. P. & Emplit, P. Towards mode-locked fiber laser using topological insulators. *Nonlinear Photonics; 2012: Optical Society of America; 2012.*
  - Zhao, C. *et al.* Ultra-short pulse generation by a topological insulator based saturable absorber. *Appl. Phys. Lett.* **101**, 211106 (2012).
  - Tsai, T. Y., Fang, Y. C. & Hung, S. H. Passively Q-switched erbium all-fiber lasers by use of thulium-doped saturable-absorber fibers. *Opt. Express* **18**, 10049–10054 (2010).
  - Petropoulos, P. & Offerhaus, H. L. Passive Q-switching of fiber lasers using a broadband liquefying gallium mirror. *Appl. Phys. Lett.* **74**, 3619 (1999).
  - Keller, U. *et al.* Solid-state low-loss intracavity saturable absorber for Nd: YLF lasers: an antiresonant semiconductor Fabry-Perot saturable absorber. *Opt. Lett.* **17**, 505 (1992).
  - Hasan, T. *et al.* Nanotube-polymer composites for ultrafast photonics. *Adv. Mater.* **21**, 3874–3899 (2009).
  - Bao, Q. *et al.* Atomic-layer graphene as a saturable absorber for ultrafast pulsed lasers. *Adv. Funct. Mater.* **19**, 3077–3083 (2009).
  - Wang, Z., Lv, X. & Weng, J. High peroxidase catalytic activity of exfoliated few-layer graphene. *Carbon* **62**, 51–60 (2013).
  - Richter, W. & Becker, C. R. A Raman and far-infrared investigation of phonons in the rhombohedral V<sub>2</sub>-VI<sub>3</sub> compounds Bi<sub>2</sub>Te<sub>3</sub>, Bi<sub>2</sub>Se<sub>3</sub>, Sb<sub>2</sub>Te<sub>3</sub> and Bi<sub>2</sub>(Te<sub>1-x</sub>Se<sub>x</sub>)<sub>3</sub> (0 < x < 1), (Bi<sub>1-y</sub>Sb<sub>y</sub>)<sub>2</sub>Te<sub>3</sub> (0 < y < 1). *Phys. Status Solidi B* **84**, 619–628 (1977).
  - Zhang, J. *et al.* Raman spectroscopy of few-quintuple layer topological insulator Bi<sub>2</sub>Se<sub>3</sub> nanoplatelets. *Nano Lett.* **11**, 2407–2414 (2011).
  - Zhao, C. *et al.* Wavelength-tunable picosecond soliton fiber laser with topological insulator: Bi<sub>2</sub>Se<sub>3</sub> as a mode locker. *Opt. Express* **20**, 27888–27895 (2012).
  - Lu, S. *et al.* Third order nonlinear optical property of Bi<sub>2</sub>Se<sub>3</sub>. *Opt. Express* **21**, 2072–2082 (2013).
  - Sun, Z. *et al.* Graphene mode-locked ultrafast laser. *ACS Nano* **4**, 803–810 (2010).
  - Luo, Z. *et al.* Multiwavelength dissipative-soliton generation in Yb-fiber laser using graphene-deposited fiber-taper. *IEEE Photonics Technol. Lett.* **24**, 1539–1542 (2012).
  - Kelleher, E. *et al.* Nanosecond-pulse fiber lasers mode-locked with nanotubes. *Appl. Phys. Lett.* **95**, 111108 (2009).
  - Yamashita, S. *et al.* Saturable absorbers incorporating carbon nanotubes directly synthesized onto substrates and fibers and their application to mode-locked fiber lasers. *Opt. Lett.* **29**, 1581 (2004).
  - Zhou, D. P., Wei, L., Dong, B. & Liu, W. K. Tunable passively Q-switched erbium-doped fiber laser with carbon nanotubes as a saturable absorber. *IEEE Photonics Technol. Lett.* **22**, 9–11 (2010).
  - Zayhowski, J. & Kelley, P. Optimization of Q-switched lasers. *IEEE J. Quantum Electron.* **27**, 2220–2225 (1991).
  - Coleman, J. N. *et al.* Two-dimensional nanosheets produced by liquid exfoliation of layered materials. *Science* **331**, 568–571 (2011).
  - Trindade, T., O'Brien, P. & Pickett, N. Nanocrystalline semiconductors: synthesis, properties and perspectives. *Chem. Mater.* **13**, 3843–3858 (2001).
  - Ueda, K., Tabata, H. & Kawai, T. Magnetic and electric properties of transition-metal-doped ZnO films. *Appl. Phys. Lett.* **79**, 988 (2001).
  - Brus, L. A simple model for the ionization potential, electron affinity, and aqueous redox potentials of small semiconductor crystallites. *J. Chem. Phys.* **79**, 5566 (1983).
  - Henglein, A. Small-particle research: physicochemical properties of extremely small colloidal metal and semiconductor particles. *Chem. Rev.* **89**, 1861–1873 (1989).
  - Leutwyler, W., Bürgi, S. & Burg, H. Semiconductor clusters, nanocrystals, and quantum dots. *Science* **271**, 933–937 (1996).
  - Gorer, S. & Hodes, G. Quantum size effects in the study of chemical solution deposition mechanisms of semiconductor films. *J. Phys. Chem.* **98**, 5338–5346 (1994).
  - Vargas, A. *et al.* The Changing Colors of a Quantum-Confined Topological Insulator. *ACS Nano* **8**, 1222–1230 (2014).
  - Sandhoff, C., Hwang, D. & Chung, W. Carrier confinement and special crystallite dimensions in layered semiconductor colloids. *Phys. Rev. B* **33**, 5953–5955 (1986).
  - Smotkin, E. *et al.* Size quantization effects in cadmium sulfide layers formed by a Langmuir-Blodgett technique. *Chem. Phys. Lett.* **152**, 265–268 (1988).
  - Sakai, N., Ebina, Y., Takada, K. & Sasaki, T. Electronic band structure of titania semiconductor nanosheets revealed by electrochemical and photoelectrochemical studies. *J. Am. Chem. Soc.* **126**, 5851–5858 (2004).
  - Kim, D. *et al.* Surface conduction of topological Dirac electrons in bulk insulating Bi<sub>2</sub>Se<sub>3</sub>. *Nat. Phys.* **8**, 459–463 (2012).
  - Culcer, D., Hwang, E., Stanescu, T. & Sarma, S. Two-dimensional surface charge transport in topological insulators. *Phys. Rev. B* **82**, 155457 (2010).
  - Nozik, A. & Memming, R. Physical chemistry of semiconductor-liquid interfaces. *J. Phys. Chem.* **100**, 13061–13078 (1996).
  - Luo, Z. *et al.* Graphene-based passively Q-switched dual-wavelength erbium-doped fiber laser. *Opt. Lett.* **35**, 3709–3711 (2010).

## Acknowledgments

This work is supported by the National Basic Research 973 Project (2014CB932004), National Natural Science Foundation of China (31371005, 81171453, 61107044), the Knowledge Innovation Program of Shenzhen City (JCYJ20130327150937484), the Fundamental Research Funds for the Central Universities, Program for New Century Excellent Talents in University, the Ministry of Education.

## Author contributions

J.W., Z.Q.L. and L.P.S. conceived the work. Z.Q.L. performed the preparation experiments. J.P. assisted with characterization of materials. Y.Z.H. performed the optical experiment of fiber laser. All authors analyzed the data and prepared the manuscript.

## Additional information

Supplementary information accompanies this paper at <http://www.nature.com/scientificreports>





**Competing financial interests:** The authors declare no competing financial interests.

**How to cite this article:** Sun, L.P. *et al.* Preparation of Few-Layer Bismuth Selenide by Liquid-Phase-Exfoliation and Its Optical Absorption Properties. *Sci. Rep.* 4, 4794; DOI:10.1038/srep04794 (2014).



This work is licensed under a Creative Commons Attribution-NonCommercial-NoDerivs 3.0 Unported License. The images in this article are included in the article's Creative Commons license, unless indicated otherwise in the image credit; if the image is not included under the Creative Commons license, users will need to obtain permission from the license holder in order to reproduce the image. To view a copy of this license, visit <http://creativecommons.org/licenses/by-nc-nd/3.0/>