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# **Improving Chemical Shift Encoded Water-Fat Separation Using Object-Based Information of the Magnetic Field Inhomogeneity**

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# **Abstract**

**Purpose—** he purpose of this work was to improve the robustness of existing chemical shift encoded water-fat separation methods by incorporating object  $\mathcal{L}$  ased information of the B0 field inhomogeneity.

**Theory—The primary challenge in water-fat separation is the estimation of phase shifts that arise** from B0 field inhomogeneit, which is composed of the background field and susceptibilityinduced field. The susceptibility-induced rield can be estimated in the susceptibility distribution is known or can be approximated. In this work, the susceptibility distribution is approximated from the source images using the known susceptibility values of vater,  $f^2$ , and air. The field estimate is then demodulated from the source images prior to water-fat separation. **Philaded Is small education and EVALUATION**<br> *Magn Record 16:* 2015 February 17:(2), 59° 604 doi:10.1002/amm23163,<br> **EVALUATION CINEMEL CONSTRANT CONSTRANT CONSTRANT CONSTRANT CONSTRANT CONSTRANT CONSTRANT CONSTRANT CONST AHFORD SCIENCIAL CONSUMPTER**<br> **AHFORD SCIENCIAL CONSUMPTER**<br> **AHFORD SCIENCIAL CONSUMPTER**<br> **AHFORD SCIENCIAL CONSUMPTER**<br> **AHFORD OF THE CONSUMPTER CONSUMPTER**<br> **AHFORD SCIENCIAL CONSUMPTER**<br> **ALEX CONSUMPTER**<br> **ALEX CON** 

**Methods—**Chemical shift encoded source *i*mages were acquired in anatomical regions that are prone to water-fat swaps. The images were processed using algorithms  $\mu$  on the ISMRM Fat-Water Toolbox, with and without the object-based field map information. The estimates were compared to examine the benefit of using the object-based field map information.

**Results—**Multiple cases are shown in which water-rat swaps were avoided by using the objectbased information of the B0 field map.

**Conclusion—Object-based information of the B0 field may improve the robustness of existing** chemical shift encoded water-fat separation methods.

### **Keywords**

magnetic resonance imaging; chemical shift encoded water fat separation; Dixon imaging; susceptibility

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# **INTRODUCTION**

 $C<sup>L</sup>$ emical shift encoded techniques for water-fat separation have experienced considerable development and application in recent decades. Originally proposed by Dixon (1) and subsequently expanded by  $G^1$  ver et al. (2), these techniques have been adopted by applications that require improved visualization of water-based tissues as well as those that demand robust fat suppression in areas of severe B0 field inhomogeneity. In addition, the  $v$  se of  $c^1$  emical shift encoding in T<sub>1</sub>-we ghted contrast enhanced imaging is particularly im<sub>z</sub> ortant since alternatives such as short-tau inversion recovery (STIR) techniques (3) should be avoided with post-contrast T1-weighted imaging. A variety of water-fat separation techniques have been proposed, including a single-ech method (4), dual-echo methods  $(1,5–8)$ , as well as numerous  $3+$ -echo methods  $(2,9,11)$ 

The primary challenge in water-fat separation is the estimation of the time-dependent phase shifts that  $a^r$  se from B0 field inhomogeneity. It the B0 field map can be estimated accurately, then the water and fat signals can be uniquely separated using a straightforward line at inversion. However, inaccurate estimation of the B0 field map can lead to swaps of the water and  $f^*$  signals. This is commonly recognized as the main challenge for chemical suft encoded water-jat separation methods.

Robust estimation of the B0 field map has been a major focus of technical development in water-fat separation. To overcome the ambiguity when estimating the B0 field map, many techniques assume that the field is slowly varying (11–19). However, this assumption is empirically  $\frac{1}{4}$  ased and  $\frac{1}{4}$  oes not fully represent the underlying physics of the B0 field perturbations. Nevertheless, it is sufficiently valid in many cares, which explains the effectiveness of these techniques. Repeat work from Yu et al. has exploited the *material* properties  $\epsilon$ <sup>f</sup> tissue by using the spectral complexity of fat to reduce water-fat swaps (20). However, none of the previously developed methods use any *ana omical* or other geometrically based information to aid in the determination of the B0 field map. The use of such information may further improve the robustness of existing water-fat separation methods. **EVALUATION** Chemical shift encodes lead interest for water-fit separation have equally the problem of a proposition in product decomply proposed interest of states (Depending to the system of a system of a system of the Page 2<br>
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Past work in electromagnetic decay has shown that the component of the B0 magnetic field that is caused by susceptibility variations can be estimated efficiently and with high accuracy if the susceptibility distribution of the object bein, imaged is known  $(2^1-24)$ . Further, this susceptibility-induced field represents a significant portion of the field inhomogeneity that must be resolved for accuration contribution of the chemical species. Interestingly, previous techniques for chemical shift encoded water- $f^2$  separation do not exploit this information.

Therefore, the purpose of this work was to develop a framework for incorporating objectbased information of the B0 field map into chemical shift encoded water-fat separation. The proposed method is intended to augment, rather than to replace, existing  $\omega_c$  in iques for water-fat separation. This method may be used with any complex-based water-rat separation algorithm to improve the robustness of the  $\hat{a}$  superithm.

#### **THEORY**

The signal from a voxel containing water  $(\rho_w)$  and fat  $(\rho_f)$  experiencing a local B0 field inhomogeneity ( $\psi$ , in Hz), measured at echo time  $t_n$  can be written as

$$
s^{f+}_{n} = (\rho_{1} - c_{n}\rho_{f})e^{j2\pi \psi t_{n}} + N \quad [1]
$$

where  $c_n = \sum_{p=1}^{p} \alpha_p e^{j2\pi r^2 p^2 n}$  represents a known multi-peak fat spectrum (25,26),  $\alpha_p$  is the  $r<sub>other</sub>$  amplitude of the  $p<sup>th</sup>$  fat peak (such that  $\sum_{n=1}^{p} \alpha_p = 1$ ),  $\Delta f_p$  is the frequency shift (in  $H_{\epsilon}$ ) or the *p*<sup>*th*</sup> fat peak relative to the water peak, and *N* denotes complex additive white Gaussian noise (AWGN). The  $e^{ff}$  cus of  $T2^*$  have been ignored since they do not generally affect the quality of water-fat separation, howe ver the approach described in this paper could also be applied. to methods using  $T2*$  correction (27).

Separating water and fat requires the estimation of  $\mathbf{u}$  unknown parameters from the multiecho measurements. In the presence of AWGN, the maximum likelihood estimate is typically found by minimizing the least-squares cost. However, the least-squares cost contains multin<sup>1</sup>, iocal minima as a function of  $\psi(12)$ . Converging to a local minimum may result in a swap of the water and fat signals. Past works  $h_{\mu\nu}$  proposed to overcome the challenge of multiple local minima by constraining the estimate of the B0 field map to be slowly varying  $(5,7-9,11,-19)$ . However, the assumption of a slowly varying B0 field breaks down in regions where the susceptibility distribution changes rapidly (e.g. air-tissue interfaces) in analomical regions with irregular  $\sigma$  one try (e.g. brachial plexus and ankle), and in iso ated regions of an<sup> $\star$ </sup>  $\sigma$  and  $(\epsilon, g$ . liver dome  $\epsilon$  two legs imaged axially). Further, assumptions regarding the degree of field map smoothness generally have no physical basis and are chosen empirically. **EVALUATION To a signal from a vector cast case of**  $f(x) = dx$  **and fit (b) experience in<br>his magnetary (v, in Hz), massar 2.d are cho time**  $t_0$  **and fit (b) experience<br>inhomogeneity (v, in Hz), massar 2.d are cho time,**  $t_0$ The state of the state of  $(\rho_0)$  and fat  $(\rho)$  experiencing a local D0 field<br>
(w, in ILe), measured at echo time  $I_n$  can be written as<br>  $S(t_{n0}) = (p - t_{n0}t) e^{i2\pi i t} \pi + N$  [1]<br>  $\pm (t_{n0}t)^{2/2 + 1/2} \pi^{3/2}$  represents  $\frac{$ 

The B0 field in nomogeneity can be written as  $\psi = \psi_b + \psi_{fs}$ , where  $\psi_b$  represents the background field, which is caused by the shim fields and imperfections in the magnet system, and  $ψ<sub>s</sub>$  represents the susceptibility-induced field, which results from the interaction of the object being image d with the applied  $\alpha$  agnetic field. Because of the low order shim fields on MR systems, rapidly varying B0 fields can be at ributed to the susceptibilityinduced field. The susceptibility-induced field  $(\nu_s)$  can be estimated if the susceptibility distribution of the object ( $\chi$ ) is known (24), i.e.

$$
\psi_{S}(\mathbf{r}) = \frac{\gamma}{2\tau} R_{\mathsf{U}}(u(\mathbf{r})^* \chi(\mathbf{r})) \quad \text{for}
$$

In Eq. 2, \* denotes convolution,  $\gamma/2\pi$  is the gyromagnetic ratio (i.e. 4) 58 MHz/T for <sup>1</sup>H), *B*<sup>0</sup> is the main magnetic field strength  $(i_A \Gamma)$ , and  $\frac{d}{dx}i_B = \frac{3c_0^2}{\pi} \frac{2(\theta) - 1}{\theta}$  $\frac{(\theta) - 1}{4i |\mathbf{r}|^3}$  represents the dipole response kernel, where  $\theta$  is the angle with respect to the main magnetic field axis and **r** is the position vector. Eq. 2 can be equivalently represented in Fourier-space ( $\angle 1-23$ ) as:

$$
\psi_{s}(\mathbf{k}) = \frac{1}{2\pi} B_0 \left( \frac{1}{3} - \frac{k_z^2}{\mathbf{k}^H \mathbf{k}} \right) \chi(\mathbf{k}) \quad [3]
$$

where  $\mathbf{k} = [k_x, k_y, k_z]$  denotes the location in Ferrier-space, and  $\mathbf{k}^H$  represents the conjugate transpose of **k**. By convention, the z-axis is oriented along the superior-inferior direction (i.e.  $\alpha$  arection of the main magnetic field) and the x-y plane is the plane orthogonal to this axis. In a Cartesian acquisition, the points along the z-axis can be calculated as  $[-N_z/2:N_z/2]$ 2 1/FOV<sub>z</sub>, where N<sub>z</sub> (assumed to be an even v<sup>ol</sup><sub>u</sub>c) and FOV<sub>z</sub> are the number of acquired points and the field-of-view in the z-direction, respectively (similar for the x- and ydirections). The value of  $\psi_s$  is undefined  $\gamma$ ,  $\kappa = 0$ . Since this point defines the DC offset of the B0 field in image space, it  $m_{x,y}$  be reasonable to set it to zero (23), although other offsets based on the center frequency of the magnet could be chosen. The calculation of  $\psi_s$  is typically performed in Fourier-space, rather than in image-space, because it is more computationally efficient. The image domain representation,  $\psi_s(\mathbf{r})$ , can then be determined using the inverse Fourier transform of Eq. 3. Note that  $d(\mathbf{r})$  and  $\chi(\mathbf{r})$  must be appropriately zero-paded before Fourier transform to avoid the effects of circular convolution when  $e^{v}$ uating Eq. 3 (28). **EVALUATION**<br>
W<sub>2</sub>(**Ka**) =  $\frac{f_2}{2\pi}B_4\left[\frac{1}{3} - \frac{k_f^2}{k^2\pi}\right]$ <br>
where, **K**. [ $k_0, k_1, k_2$ ] censed a the location and the same parts on the same parts of the same of t **Pages**<br> *P<sub>2</sub>***, <b>A**<sub>2</sub> denots the location in Touries space, and  $k^2$  persons the conjugate<br> **A**<sub>N</sub>, **A**<sub>z</sub> denots the location in Touries space, and  $k^4$  represents the conjugate<br> **A**N<sub>2</sub> denots affective in Touries

The sus eptibility-induced field that is calculated using Eq. 3 serves as the object-based field map estimate. This object-based field map estimate is demodulated from the multi-echo sources in ages, as shown in Eq. 4, where *N*<sup>*r*</sup> represents complex AWGN.

$$
\hat{s}(z_n) = s(t_n)e^{-j2\pi\psi_s t_n} = (\rho_w + z_n \rho_f)e^{j2\pi(\psi_b - \psi_s - \psi_s)t_n} + N \quad [4]
$$

The resultant multi-ccho source images,  $\hat{s}(t_n)$ , are expected to have a majority of the B0 field demodulated by this supect-based field map estimate. The demodulation step is expected to simplify the task of the subsequent water-fat separation algorithm, which must now only correct for the slowly varying background field  $(\psi_b)$  and any remaining susceptibilityinduced field compone it  $(\psi_{\gamma} - \hat{\psi_{s}})$ .

Due to the nonlocal response  $\epsilon_1$  the dipole kernel in Eq. 2 (24), calculation of  $\hat{\psi_s}$  requires 3D information about the susceptibility distribution. However, approximating the susceptibility distribution in regions outside of the imaging field-of-view (FOV) is not possible without additional assumptions, and therefore an incomplete estimation of the susceptibility-induced field may occur near the edges of the imaging FOV. These effects were analyzed via the point-spread function (PS<sub>F</sub>) of the dipole kernel  $u(\mathbf{r})$ . In particular, we focused on the PSF along the slice-encoding direction because many water-fat separation algorithms impose field map constraints only in  $2\nu$ , and thus may be least robust to these artifacts along the slice-encoding direction.

#### **METHODS**

Experiments were conducted after obtaining informed consent and IRB approval using a clinical 1.5T scanner (HDxt, v16.0,  $CE$  Healthcare, Waukesha, WI) and a clinical 3 1 scanner (MR750 v22.0, GE Healthcare, Waukesha, W1). Cardi $\therefore$  datasets were acquired in

28 subjects using a 3D free-breathing, navigated, cardiac-gated four-echo SPGR sequence (29). Data were also acquired in the individual volunteers and in the brachial plexus  $\log n$  five volunteers (or e volunteer was scanned twice, on different days) using a 3D threeecho SPGR sequence. Lastly,  $d\sigma^2$  were acquired in the abdomen from ten volunteers using a 3D breath-held six-ec<sup>1</sup> o (six monopolar  $\sim$  noes per TR) SPGR sequence. We were unable to  $\therefore$  acquire the abdominal datasets using a true dual-echo acquisition because of the unavailability of the product reconstruction pipeline, which was needed to obtain the source images from the acquired data. To serve as a surrogate, a six-echo acquisition was used, from which the source images at two echoes  $cov^1$  do extracted. The source images at exhoes 4 and 5 were selected because they maximized the effective number of signal  $\alpha$  averages (NSA) over all possible echo combinations for this particular six-echo acquisition (8). Table 1 lists the acquisition  $r_{\text{in}}$  ameters for each of the datasets that are presented in this work. (20). Also were shown with  $\frac{1}{2}$  is consider a matter has an axiomheres and in the solutions of th the 2012 chromating manipund, curtina-gand from-celon SPGR sequences<br>because 2012 chromating manipund, curtina-gand four-celon SPGR sequences<br>because the sequence  $\frac{1}{2}$  change when exact the subsequence of the subsequ

Prior to any furth  $\alpha$  processing, the raw source images were corrected for the effects of gradient nonlinearity and were coil-combined using an adaptive phase preserving algorithm ( $3<sup>c</sup>$ ). Because the proposed method uses object-based in formation, it was important to correct for the gradient nonlinearity effects, which introduce image distortions, before generating the estimate of the B0 field. All  $\mathcal{L}_{\text{total}}$  for processing was then done in Matlab (The Ma'hworks Inc., Natick, MA) (64-bit Linux, 4 Octo Core AMD 6134, 128 GB RAM).

The coil-combined multi-exilos surce images were  $f_{\text{ext}}$  processed using algorithms from the ISMRI  $I_{\text{F-d}}$ -Water Toolbox (31). To demonstrate the general applicability of the proposed method, the cardiac datasets were processed using a  $\epsilon$  aphcut c' corithm (15), the brachial plexus datasets were processed  $\text{min}$  hierarchical IDFAL (19), the ankle datasets were processed using a region-merging algorithm  $(12)$ , and the absolution al datasets were processed using a dual-equority method with flexible value of  $\sim$ ing times (8). Note that the dual-echo algorithm that appears in the Fat-Water  $\sqrt{ }$  volbox is a voxel-independent method, which does not incorporate neighborhood information when estimating the phase shifts that are caused by  $B\circ \lim_{n \to \infty} \log$  neity.

For the dual-echo algorithm, the weighted smoothing of the  $\mu$ <sub>n</sub> sort was found to increase the occurrence of incomplete water-fat separation. Therefore, the smoothing was removed from the reconstruction. Other than this modification, no changes were made to the default settings of each algorithm in the  $Fr'$ . water Toolbox. For the sealgorithms that implemented water-fat separation in 2D, the 3D d<sup>3+</sup> usets we re processed on a slice-by-slice basis.

Each dataset was then processed a second time using the identical algorithm as described above, but after the object-based field map estimate had been first demodulated from the multi-echo source images. The object-based field map was estimated using the source images that were acquired at multiple  $\sim$  no times. Figure 1 shows a flowchart of the proposed approach. A maximum in ensity, projection (MIF) image was calculated from the multi-echo source images by projecting along the echo time dimension. A binary mask consisting of regions that contain either air or tissue was 'nen created from the MIP image. To create the binary mask, an air-tissue threshold was set at  $5\%$  of the maximum value in the MIP image. Those voxels in the MIP that were pelow the threshold were considered to

contain air while those above the threshold were considered to contain tissue. An estimated susce<sub>k</sub> tibility distribution  $\left(\hat{\zeta}, \cos \gamma\right)$  nerated from the binary mask, using the known  $\sin$  ceptibility values of vater,  $f_{at}$ , and air (32,33). Because water-fat separation had not yet  $\text{occ}$  an equal distribution of water and fat was assumed, to minimize the maximum error of the susceptibility estimate over <sup>1</sup> possible water-fat ratios. A susceptibility value of 8.42ppr<sub>a</sub> vas assumed, which is the average of the susceptibility of water (−9.05ppm) and fat  $(-7.79r\omega)$  (32,33). For the voxels containing air, a susceptibility value of 0.36ppm was used ( $2/2$ ). The estimated susceptibility-induced field  $(\hat{\psi_s})$  was then calculated via Eq. 3. To compensate for center frequency shifts during provided, a constant shift was applied to  $\hat{\psi_s}$ such that its mean value over the regions of tissue was zero. Finally, the susceptibilityinduced field was demodulated from the original multi-echo source images.

The demodulated multi-echo source images were processed with the same algorithm that was used to process the original source images. The two results (i.e. with and without de nod ulation of the object-based field map estimate) were visually compared to determine whether using the object-based field map information improved the quality of the water-fat separation. In addition, the total reconstruction time for each approach was computed using the Matlab *profile* function.

# **RESULTS**

Figure 2 shows the susceptibility-induced field for the center slice of a 3D acquisition that was estimated using all slices of the volume. The PSF of the original kernel along the sliceencoding ax is also shown. The PSF was normalized such that its maximum value equals one. It is important to note that the energy of the PSF is concentrated near the origin, which suggests that only a small subset of  $\text{L}$  neighboring slices, rather than the full 3D volume, may be used to sufficiently represent the susceptibility-induced field at the center slice. To determine the number of neighboring slices, a threshold on the magnitude of the normalized PSF was established. From the inset of the PSF in Figure 2, a sharp transition is observed as the coefficient magnitudes cross a value of 0.01. Using this threshold of 0.01, a total of only seven slices (i.e. three slices  $\sigma_A$  each side of the center slice) would be required to represent the susceptibility-induce I field  $\alpha$  the center slice. Figure 2 shows the susceptibility-induced fields that were estimated using only a subset of the total slices, as well as their corresponding differences. It is seen that  $t^{\mu}$  field estimate using only the seven center slices captures much of the susceptibility-induced variations. The PSF was computed assuming that the 3D acquisition y as done in the coronal plane. Similar results (not shown) were found for both sagittal and  $\alpha$  axial acquisitions. since this time particles on  $\zeta_{\text{max}}$ , a can come that from the binary mast.<br>
since particular surface for  $\zeta_{\text{max}}$  ( $\zeta_{\text{max}}$ ) ( $\zeta_{\text{max}}$ ) ( $\zeta_{\text{max}}$ )) ( $\zeta_{\text{max}}$ ) ( $\zeta_{\text{max}}$ ) ( $\zeta_{\text{max}}$ ) ( $\zeta_{\text{max}}$ )  $\alpha$ <sup>hm-2</sup>-love die mensan) were considered to contain itssue. An estimated<br>furtheron ( $\hat{\phi}$ -may a dietra of formation in the hinary mask, using the known<br>has of variety  $\hat{M}$ , and an (32,33). Because water-fat separa

Water-fat swaps were visually observed in six of the 28 cardiac datasets that were processed using graphcut. All observed swaps  $\sim$ curred  $\sim$ ar  $\sim$  dome of the liver. The proposed method resolved the swaps in all six cases and did not introduce any new swaps. Figure 3 shows the water and fat estimates  $f$  com  $\angle$  ne cardiac da aset, with and without the objectbased field map information. A swap was observed in the dome of the liver when using graphcut alone. By first demodula ing the object-based field map estimate from the original multi-echo source images, the graph out method was "ole to correctly separate water and fat. For reference, both the object-based field map estimate and the final field map estimate are

shown. It is seen that the object-based field map estimate provided an accurate estimate of the B0 field. The calculation and demodulation of the object-based field map estimate took  $2.3$  s for the entire 3D detase. For comparison, the average reconstruction time for each 2D slice using graphcut (with or without demodulation of the object-based field map estimate) was 18s,  $\alpha$  r approximately 15 minutes for the 50 slices.

Water rat swaps were observed in three of the six ankle datasets that were processed using region-merging. The proposed method resolved  $t_{\rm av}$  swaps in all three cases and did not introduce any new swaps. Figure 4 shows the water and fat estimates from one of the ankle datasets using region-merging, with and without the object-based field map information. The severe B0 field inhomogene ty in  $\sin$  anatomy caused a swap of the water and fat signals when using region-merging alone.  $P_y$  first demodulating the object-based field map estimate from the original source images, region-merging successibilly separated the water and fat signals. Calculation and demodulation of the 3L object Lased field map estimate required 1.9s of computation time for the entire 3D dataset. The average reconstruction time for each 2D lice using region-merging (with or with  $\omega$ u demodulation of the object-based field map estimate) was 31s, or  $\epsilon_{\rm r}$  oximately 25 minutes for the 48 slices.

Water- $f_{\mu\nu}$  swaps were observed in four of the six brachial plexus datasets that were processed using hierarchical IDEAL. The proposed method resolved the swaps in all four cases and it did not introduce any new swaps. Figure 5 shows the water and fat estimates of the brachial plexus using hierarchical IDEAL, with and without the object-based field map estimate. The water and fat estimates using only hierarchical IDEAL exhibited swaps of the water and f. signals. In contrast, by first demodulating the object-based field map estimate from the criginal source images, the archical IDEAL was able to correctly separate the water and fat signals. Calculation can demodulation of the object-based field map estimate took 2.1s for the entire  $3\sum$  dataset. The reconstruction time  $f_{\text{out}}$  is  $3\sum$  volume using hierarchical IDEAL (with or without demodulation  $\sigma_1$  the object-based field map estimate) was approximately 55s.

Water-fat swaps were observed in all of the abdominal datasets that were reconstructed using the dual-echo, voxel-independent method. This was expected due to the known limitation of voxel-independent methods to robustly estimate the B0  $\tilde{i}$ eld map in regions of severe field inhomogeneity. The proposed method was unable to resolve all of the swaps, however a marked improvement in the water- $f_{\alpha}$  separation vas observed especially at the air-tissue interfaces. Figure 6 shows the water, fat, and B0 field map estimates from  $0.2$  subject using the dual-echo, voxel-independent method and the proposed method. Water-fat swaps were largely resolved when the object-based  $f_1e_2$  map estimate was first demodulated, however swaps are still visible. The calculation and  $d$  modulation of the c bject-based field map estimate took 1.9s for the entire  $3\Gamma$  dataset. The reconstruction time for the 3D volume using the dual-echo, voxel-independent method (with or vithout demodulation of the c bjectbased field map estimate) was approximately  $1.2$ s. The Bell field the *melodiation* of the object-based from the object-based space (words in the object-based space). The consideration of the object-based space (words in the object-based space) was 185. Consider the and g **AHFORT**  $\frac{1}{2}$   $\$ 

#### **DISCUSSION**

In this work we have described  $\gamma$  novel technique that augments existing chemical shift encoded water-fat separation methods by incorporating object-based information of the B0 field map into the reconstruction. We have  $d_{\text{sub}}$  onstrated that using this information can improve the robustness of existing water-fat separation methods. We propose that this approach may be effective as a preprocessing of the multi-echo source images, and therefore  $s'$ ould  $\frac{1}{2}$  applicable to any complex-based chemical shift encoded water-fat separation method. The primary advantage of this approach is that it exploits a physical property of the  $\overline{1}$ <sup>1</sup> imaging volume (i.e. susceptibility distribution based on the anatomy) rather than relying solely on assumptions of field  $m_{\nu}$  smoothness. It is particularly applicable when imaging regions in which the B0 field map varies rapid, either due to abrupt changes in the susceptibility distribution or in anatomy with i regular geometry.

It is calculation of the object-based field map estimate was implemented efficiently using a point-wise multiplication in Fourier-space. In all examples, the calculation and demodulation of the 3D object-based field map took less than 3.5s seconds, representing a minimal increase to the overall reconstruction time while consistently improving the robustness of water-fat separation. No significant difference in the reconstruction time of the water-fat separation algorithm was observed after using the object-based field map estimate. However, because this estimate is a fairly accurate representation of the true B0 field, it may be possible to reduce the reconstruction time, particularly for iterative gradient-based methods with carefully chosen convergence criteria. In this work, the default criterion for each algorithm was used.

The main imitation of this approach is related to the assumption of the magnetic susceptibility distribution.  $T_{\text{def}}$  assumptions used in  $t_{\text{def}}$  work should be valid for most situations unless there is a foreign body with high susceptibility (e.g. metallic prosthesis). Further, iron is the only maturally occurring substance with high susceptibility that can occur in high concentration with  $\frac{1}{2}$  tissue. In the case of tissue iron overload, it may be possible to approximate the R2\* from the multi-cho source images,  $\gamma$  is used that value to estimate the susceptibility of those tissues  $(3<sup>2</sup>)$ . In addition, the air-tissue  $\mu$ <sub>1</sub> sk was determined by using a 5% threshold on the MIP image. The accuracy of the object-based field map estimate will be affected by the accuracy of the air-tissue mask, however we have found that setting a threshold value within a range of  $\frac{3}{2}$ % does not significantly affect the final water-fat separation. When using the sholds above 9%, the bone, which has a susceptibility value close to that of tissue (32),  $m\omega y$  be masked as air. This error in the air-tissue masking  $m\omega y$ introduce error in the object-based field  $m_{\mu\nu}$  estimate. It should be noted that  $\Box$  e a r-tissue masking algorithm that we have used was sufficient for the cases that were tested in this work, and that the modular nature of our proposed manework would all ow the use of more sophisticated air-tissue masking algorithms. **EVALUATION**<br>
For this work we have described in two fix-fits procedure that anymorite strategies that any interest that the external of the interest of the interest of the contribution of the contribution of the contribu The between the transmisterial material material material material shift<br>this separation methods by incorporating object-based information of the Bo<br>be recording to the skew of covered that using this information can<br>be re

The forward calculation of the field in homogeneity requires 3D information of the susceptibility distribution. Thus, the proposed method may not be suitable for  $2D^{\frac{1}{2}}$  imaging unless multiple closely spaced slices are acquired, providing an accurate  $\mathcal{I}_D$  representation of the tissue. Further, incomplete estimation of the susceptibility-induced field in 3D

acquisitions may occur at the edges of the imaging FOV. For the slices at either edge of the imaging volume, one may be able to synthetically extend the imaging volume by replicating  $\alpha$  extrapolating the edge slices. This would only serve as an approximation of the actual slices, but it may be adequate for improving the accuracy of the field map estimate for the edge slices. Based on our analysis of the point-spread function of the dipole kernel, we estimate  $\Delta t$  t errors in the estimate of the susceptibility-induced field map may occur for the three slices at each end  $\sigma_1$  a 3D volume. Further work is required to determine the performance of the proposed algorithm at the edges of volumetric acquisitions. **EVALUATION** (and the signal of this continuous control in the main space of the signal of the Passarian developes on the imaging FOV fee the sliess at either edge of the<br>
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the edge sliese. This would only serve as an approximation of the actua

 $\lim_{n \to \infty}$  the proposed method was not tested using data from a true dual-echo acquisition. This was not possible due to the unavailability at our site of the product reconstruction pipeline that was needed to obtain  $t^h$ . complex-valued, gradient nonlinearity-corrected source images from the acquired data. Indeed, the proposed method may provide greatest benefit to two-point v ater-fat separation algorithms, which are widely used for 3D vo'um tric imaging and are known to be especially vulnerable to swaps because of the  $\lim_{h \to \infty}$  red number of measurements.

In conclusion, we have described a novel approach to improve the robustness of existing water- $f_{\alpha}$  separation algorithms by incorporating object-based information of the B0 field inhomogeneity. This approach can be applied to any complex-based chemical shift encoded water-fat separation technique. Initial results are  $h^2$  guly promising for improving the robus ness of water-fat separation.

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#### **Figure 1.**

Flow chart of the proposed approach. A maximum intensity projection (MIP) image is valculated from the multi-ceno source images by projecting along the echo time dimension. An air-tissue mask is then created from the MIP image. The estimated susceptibility  $di$ r-tibution is generated from the air-tissue mask, using the known susceptibility value of air and the mean of the values of water and  $\Delta t$ . The estimated susceptibility-induced field is calculated and is then demodulated from the original source images. Finally, the demodulated source images serve as the input into any water-fat separation algorithm. **Example 18**<br> **EVALUATION**<br> **EVALUATION** Proposed a spectra che Al maximum intensity projection (MIP) image is<br>the radiation content in the MIP image. The estimated susceptibility<br>the values of water indicate in the MIP image. The set match is the March of the va



#### **Figure 2.**

(top 'eft). The susceptibility-induced field that was calculated using all slices of a 3D acquisition. (bottom left): The point-spread function (PSF) of the dipole kernel along the slice-encoding axis. Notice that the energy of the PSF is concentrated near the origin, which sugests that only a subset of the neighboring slices may be used to calculate the susceptibility-induced field at the center  $s_i$  is  $\infty$ . A threshold of 0.01 was established (dotted line), which corresponded to a total of seven slices (i.e. three slices on each side of the center slice). (right): The susceptibility-induced fields that were estimated using only a subset of the total slices, as well as the corresponding difference images. Using only the seven center slices captures much of the susceptibility-induced variations. **EVALUATION**<br> **EVALUATION Property of the second that the two schedulated using all sites of a 3D**<br>tuon sell), "The point c-seried fact the two calculated using all sites of a 3D<br>town sight, a major of the PSF is concentrated from the origin, whic

# **Figure 3.**

Water, fat, and B0 field  $\epsilon$  stimates using a graphcut algorithm, without and with the objectbased  $f_{\text{tot}}$  map information. The B<sup>0</sup> field estimate in the bottom row is shown as a sum of the object-based field estimated by graphcut. A swap is observed in the dome  $\mathcal{L}$  the liver when using graphcut alone (white arrow). When the object-based field map estimate is first demodulated from the original source images, the graphcut method correctly separates the water and fat signals. Further, notice that the objectbased field map estimate provides an accurate estimate of the B0 field map. In the liver dome, the mean / minimum / maximum differences between the object-based estimate and the final estimate were  $16.7 / -0.41 / 77$ .  $5$ Hz, respectively. **EVALUATION**<br>
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# **Figure 4.**

Water, fat, and B0 field  $\epsilon$  stimates in the ankle using a region-merging algorithm, without and with the object-based field map information. The B0 field estimate in the bottom row is shown as a sum of the object-based field estimate and the remaining field that was estimated by region-merging. The rapid<sup>1</sup>, varying B0 field in this anatomy caused a swap when the source images were processed using only  $r_{\leq 1}$  on-merging (white arrow). By first demodulating the object-based field map estimate from the source images, region-merging was able to correctly separate the water and fat signals. The object-based field map estimate provided an accurate estimate of the  $f^2$  in inhomogeneity, especially in the region of the toe.

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#### Fig. re 5.

Water, fat, and B0 field  $\epsilon$  stimates in the brachial plexus using hierarchical IDEAL, without and w<sup>14</sup> the object-based field map information. The B0 field estimate in the bottom row is shown as a sum of the object-based field estimate and the remaining field that was estimated by hierarchical IDEAL. Water-fat swaps in the head and neck (white arrow) are properly reso<sup>1</sup> ved by hierarchical IDEAL only when the object-based field map estimate is first  $d$  demodulated  $f$  on the source images. **EVALUATION**<br>
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#### **Figure 6.**

Water, fat, and B0 field map estimates in the abdomen using a voxel-independent algorithm, vithout and with the object-based field map information. Numerous swaps are seen throughout the abdomen, some of which are highlighted (solid arrows). By first demodulating the object-based field map estimate, the voxel-independent method dem nstrates an improvement in the water-fact separation. However, some swaps are still visible in the reconstructed images (dashed arrows). **EVALUATION**<br>
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**Table 1**

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Acquisition parameters for each of the datasets presented in this work. Each dataset was acquired using a 3D SPGR sequence. nTE denotes the number of echo times and ΔTE represents the echo time spacing. **Anatomy Matrix Size Imaging Plane B0 Field (T) nTE TE1 (ms) ΔTE (ms)** Cardiac 256×256×50 Axial 3 4 1.22 0.98 Ankle 256×256×48 Sagittal 1.5 3 1.984 1.588 Brachial Plexus 256×256×30 Coronal 1.5 3 1.984 1.588 Abdomen  $256\times256\times28$  Axial 1.5 2 *\** 7.42 *\** 2.06 <sup>(\*)</sup>The abdominal datasets were acquired using a six-echo acquisition, however or y echoes 4 and 5 were used for the chal-echo, voxel-independent reconstruction. *Magn Reson Med*. Author manuscript; available in PMC 2016 February 01. **EVALUATION** AHFORMATTER