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A simple acquisition strategy to avoid off-resonance blurring in spiral imaging with redundant spiral-in/out k-space trajectories

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Abstract

Purpose—The number hardle to widespread adoption of spiral trajectories has been their poor offresonance performance. Here we present a self-conjecting spiral k space trajectory that avoids much of the well-known spiral blurn ig during duta *e* equisition

Theory and Methods an comparison with a traditional spire?-out trajectory, the spiral-in/out trajectory has imploved off-resonance performance. Dy comparison were spiral-in/out acquisitions, one rotated 180° in 'c-space compared to the other, multi-chot spiral-in/out artifacts are eliminated. A phantom was can led with the center frequency manually runed 20 40, 80, and 160 Hz off-resonance with board a spiral-out gradient echo sequence and the reduction spiral-in/out sequence. The phantom was also imaged in an oblique orientation in couer to demonstrate improved concomitant gradient field performance or the sequence, and was additionally incorporated into a spiral turbo spin echo sequence or brain imaging.

Results—Phantom studies with manually-tuned off-resonal of agree well with theoretical calculations, showing that moderate on-resonance is well-corrected by this acquisition scheme. Blur due to concomitant fields is reduced, and good results are obtained in vivo.

Conclusion—The redundant spinal-in/out trajectory results in less image by race a given readout length than a traditional spiral-out scan. teducing the need for complet off-resonance correction algorithms.

Keywords

Spiral imaging; off-resonance correction

Introduction

Spiral k-space trajectories have many advantages over traditional rectilines, acquisitions, including better acquisition efficiency, less stringent hurdware requirements, and natural resilience to flow and motion (1). The major hurdle to wide-spread adoption of spiral trajectories has been their poor of irresonance perfermance (2). The blurring and disortion in spiral images in the presence of system non-idealities led to the two-promote strategy of

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mitigation and correction in spiral imaging. First, splitting the acquisition into multiple short interleaves minimizes artificate by resuring that an undue amount of undesirable phase does not accrue in a single rendout. Seconday, much effort has been expended to correct for offresonance effects in image reconcaruction algorithms. These techniques vary in complexity and computational cost, from a relatively simple center frequency correction and first-order trajectory warping method based on a least squares fit to an acquired field-map (3), to time-(4, 5), frequency- (6, 7), and polynomia l-approximation approaches (8), to automatic (9, 10) and secur-automatic (11) methods which demodulate the image at multiple frequencies in order to build a composite image free of blurring. Many of these algorithms have become large, requiring several seconds to reconstruct a single image.

Main field inhomogeneity is the primary source of off-resonance in MRI. However, offresonance may be deused by other system imporfections resides B₀ inhomogeneity. Particularly at lower field strengths and for off-center files, concomitant fields generated by no mail gradient operation can cause noticeable blur in spiral images. With some exceptions (12-14), the blur due to concomitant fields is largely ignored in the spiral literature, as it requires a more complete nodel to appropriately address deblurring.

The most commonly encountered spiral trajectories are "spiral-out". That is, the trajectory begins at the origin of k-space and moves outward slong a spiral. Another option is the spiral-in/out trajectory, which was first proposed for efficient sampling of spin-echoes for abdominal imaging (15) and has recently been showed to improve SNR and image quality for real-time spiral bCGrP carding imaging due to its natural ability to center TE within TR (16). Its major use, however, is in fMRI (17–19), where its GNR, special, and resistance to flow artifacts n ake it an attractive alternative to reculinear Eri methody.

For spiral in aging, the amount of undesired phase eccrued between when the center of kspace is sampled and when the edge of k-space is tampled determines the severity of the well-known spiral blue. Civen a desired resolution, spiral dat and spiral-in/out trajectories require readout lengths of meanly of same duration (within direw percent due to differing amounts of time spent near the center of k-space where the k-space velocity is small). For a given readout length they, a spiral-in/out trajectory requires doort half is king to move from the edge of k-space to the center compared to a spiral-out trajectory, and therefore accrues about half of the undesitable phase.

In this work, we take a closer look at spiral-in/out trajectories to address one method by which blurring due to of resonance can be availed through their use. The specific onlyin of off-resonance (B₀ or concomitant fields) is milliport int, and it will be shown that using a redundant spiral-in/out sampling scheme naturally removes the most schere off-resonance image artifacts during image acquisition, allowing for fast and simple correction methods in the image reconstruction step.

Theory

Figure 1 shows the readout gradients and trajectories for spiral-out and spiral-in out imaging. Conceptually, spiral-in/out trajectories only be implemented in two ways. In the firs method, termed here the "non-redundant" scheme, each spiral-out arm fills in the conjugate

k-space location of the opinal-in ann, requiring the same number of interleaves as a spiralout trajectory for equal k-spinor coverage (rig. 1d).

The second method, which we call the "redundant" scheme, acquires each interleaf twice, once in each direction "laroug". K-space (i.e., the redundant scheme consists of acquiring two i on-redundant trajectories with the second acquisition rotated 180° in k-space). Thus, the redundant scheme requires twice the number of excitations as its non-redundant counterpart. Despite this prolongation of scan time, this is a for more robust acquisition scheme, as each location in k-space is sampled twice: once with a spiral-in arm and once with a spiral-out arm. The data is then averaged, either before or after gridding the data onto a Cartesian matrix. In this way, amplitude and phase mismatches between the data acquired with spiralin and spiral-out arms of the trajectory are averaged, out. A simple illustration of the k-space weighting functions that occur when cach of these trajectories are used is provided in the Supplemental Materials.

Redundant trajectory response to system non-idealitios: Bo off-resonance

Is norms, relaxation and including B_0 in nomesceners, the classic demodulated signal equation in MixL is

$$s(t) = \int m(\mathbf{r}) e^{-j2\pi \nu(\mathbf{r})\mathbf{r}} e^{-j\omega(\mathbf{r},\tau(t))} d\mathbf{r}, \quad [1]$$

where $m(\mathbf{r})$ is the signal L(t) the k-space trajectory, and m(t) the off-resonance. In this subsection, the character phase-accrual time parameter $\tau(t) - t$ because phase accrues proportionally to time for B₀ off-resonance. Phase accrual due to concomitant fields will be addressed in the parameter $\tau(t)$ is more complex.

It can be shown (Appendix A, Supplemental Material) that the signal resulting from averaging the cata from a redundant spiral-in/out trainsform is

$$S(t) = \int m(\mathbf{r}) e^{-i^{\gamma}\pi k(t)\mathbf{r}} [\cos \omega(\mathbf{r})t] d\mathbf{r}$$
 [2]

For small-to-moderate off-resonance values, Eq. 2 shows that the signal eleperience a relatively benign cosine amplitude modulation rather than the more serious phase modulation that arises with spiral out trajectories in the presence of off-resonance. Point-spread-functions (PSFs) for a spiral-out, non-reglament in/out, and reductant in/or: trajectories with and without off-resonance were climulated and show that the P.T. is sharper in the presence of off-resonance when the redundant trajectory is used (Supplemental Material). However, the modulation transfer function (NTF) provides a more functive grasp of the situation in this case. Figure 2a shows simulated non-malized NTFs for various amounts of phase accumulated by the and of the readout for the redundant on al-in/out trajectory (corresponding to the off-resonance-time riodulation Eq. 2). There are two regimes under which the shape of the MTF may fall. In the first, the number of accumulated eveles is small, either because there is not much off-resonance present, or because the readout length.

is short. In this regime, the cosine modulation never reaches its first zero point during the readout, so the sonal experiences c windowing function that only slightly attenuates high-file juency components. In the second regime, when the value of the off-resonance-time product is high, the cosine modulation will begin nulling important frequencies as a function of k-space radius, resulting in image artifacts. This figure indicates the redundant spiral-in/out mound will work well as long as the number of cycles of phase accumulated during the readout remains lest man 0.5. This ymount of off-resonance is easily achieved during normal operation of clinical-strength scatners, so vie anticipate first-order correction should by performed on the data/trajectories prior to gridling in order to quickly correct gross off-resonance.

Concomitant gradient effects

when phase accrual is due to conconitant field effects, it can be shown that $\omega_c(\mathbf{r})$ is a complex function of the imaging gradients and opatial coordinates of the slice, the actual for n of which is unimportant for this case. As mentioned previously, in this case the phase-accrual time function takes a more complex form. Specifically,

$$\tau(t) = \frac{1}{g_m^2} \int_{0}^{1} \left[G_x^2(t') + G_y^2(t') \right] dt', \quad [3]$$

where q_m^2 is the maximum gradient strength required during the scan, and $G_x(t')$ and $G_y(t')$ are the spiral gradients on the two in-plane axes (9). Performed this time function depends on the gradients squared and because the spiral-in fout gradients are symmetric, $\tau(-t) = -\tau(t)$, and the signal in the presence of componitant gradient effects is

$$z(z) = \int m(r)e^{-j2\pi k(t)r} [\cos \omega_c(r)\tau(t)] r'. \quad [4]$$

The typical phase-accrual time function for spirals is, over all, less show than the linear function that governs B_0 off-resonance. Thus, the redundant involutions me is actually *more* robust to phase errors cause a by concomitant fields than it is to those caused by B_0 inhomogeneity and will perform well up to about 1 cycle cracerular phase (Fig. 2c).

Relaxation

In non-redundant multi-side spiral-in/out scanning, Γ_2^* relaxation during the readout results in stronger signal at one side of the periphery of k-space than the same, the set lt of which is artifacts that look strikingly similar to these caused by off-resonance. Adding Γ_2^* decay term to Eq. 1 and assuming that the readout time is short compared to T_2^* (Supplemental Material, Appendix B), the following expression to the signal in the presence of both offresonance and Γ_2^* decay can be shown to be

$$s(t) = \int m(r) e^{-j2\pi k(t)r} \cos \omega(r) t_{\perp}^{\gamma} dr + j \int m(r) e^{-j2\pi k(t)r} \frac{t}{T_{2}^{*}(r)} [\sin \omega(r)t] dr.$$
 [5]

This signal is complex, with γ real part corresponding to the previously described cosinebodulated signal equation and an imaginary part that varies in amplitude with time. This is potentially vorrisome, since phase can relation is required to remove off-resonance effects. Nowever, first note that that the maginary component in Eq. 5 becomes zero. It follows that, trajectories and that at t = 0, the imaginary component in Eq. 5 becomes zero. It follows that, at the conter of k-space where the majority of the image energy resides, there is little impact from the imaginary component of Eq. 5. Second, the ratio of $t/T_2^*(r)$ that controls the amplitude of the imaginary component will always be shall as long as T_2^* is larger than

It there is no off-resonance (or that off-resonance is corrected somehow), then $\omega(\mathbf{r}) = 0$, and the averaging operation in redundant sampling models to remove T_2^* -induced artifacts. In truth, even if the linear approximation utilized in the derivation were relaxed, as would be necessary for a chort T_2^* species, then there will be a symmetric emphasis on the outer regions or k-space, one result of which is nore beingn than the asymmetric T_2^* weighting that occurs for non-redundant spiral-in/out trajectories.

In simulations, the combination of T_2^* and on-resonance is not too different from either case alone. Figure 2c shows the performance of the redundant model in the presence of both inhomogeneity and T_2^* relaxation in terms of the MTF. As expected, strong shaping of the MTF only occurs for T_2^* on the order of reade at length (10 ms). Nowever, even at this short T_2^* , there is little degradation of the 2SF (not shown). The imaginary term in Eq. 5, and thus potentially damaging phase due to T_2^* relaxation, is negligible

Although the redundant <u>spiral in/cut trajectory</u> works for both gradient-echo and spin-echo imaging, spin-echoes (and <u>spin-echo</u> train)) provide a natural coung in which to apply them. Since the TE of spin-echo scane is generally longer, the <u>spiral in</u> po tion of the trajectory can be inserted with little or no increase in minimum TE. Second, the in/out trajectory aligns the gradient echo generated by the spiral gradient, with the upin echo formed by the RF pulses at the center of the gradient waveform, resulting in higher signal when the center of k-space is sampled.

One attractive application for the redundant spire! in/out trajector is in a slab-selective version of the 3D spiral TSE sequence (20). In 2D TSE sequences with slab-selective excitation pulses and nonselective refocusing pulses, spurious FID artifects arising from imperfect refocusing pulses are removed via an RF chopping technique, in think two averages are acquired with alt mating refocusing pulse phase (2). In the spiral TSE sequence, then, multiple averages are already being performed and since the origins of the spurious FID artifacts and the spiral-in out artifacts are different, the record, redundant acquisition can be performed with the RF-chopped requisition to acquire a fully redundant trajectory with no increase in scantime. For further discussion, see Supplemental Materials online.

Methods

A resolution phantom vias share d on a 1.5 T Siemens Avanto scanner with a GRE sequence using a spiral-cut trajectory and a redundant spiral-in/out trajectory. Acquisition parameters were: number of interleaves 14 Spiral duration 10 ms, in-plane FOV 300 mm, slice thickness 5 mm. To examine charesonarise performance, the sequences were run once with 5 good shim applied, and again with the receive frequency manually tuned 20, 40, 80, and 160 rdz off-resonance (corresponding to 0, 0, 2, 0.4, 0.8, and 1.6 cycles of off-resonance accumulated at the end of the readoul). All images viere acquired in the transverse plane, and more gridged and Fourier-transform-reconstructed with no off-resonance correction algorithm applied. The gridcing operation automatically sums the data at the proper k-space locations, given the redundant trajectories.

To investigate concornitant field performance, the resolution phantom was imaged again using 1 oth trajectories with 14 interleaves, spinal duration 6.4 ms, in-plane FOV 300 mm, and slice thickness 3 mm in a double-oblique orientation (($C \rightarrow S - 41.8^{\circ}$) $\rightarrow -27.8^{\circ}$) near the magnet isocenter (X -9.8 mm, Y -39.6 mm, Z -21.7 n.m), then moved 50 mm along the zvxis (X - 9.8 mm, Y -39.6 mm, Z -71.7 mm) and in aged again to ensure significant concordiant fields.

A comparison between the effectiveness of the reductant spiral-in/out trajectory in reducing blurring versus two post-processing techniques was performed by manually setting the second-order shim of the system, such that a choog non-linear variation in B₀ existed across the FOV in one direction. This ensures the simple, linear correction method will fail and represents a situation in which more advanced on-resonance correction algorithms are necessary. The same imaging parameters used the previous experiment were used, except that the phantom was set up in the coronal orientation so that Z^2 shim could be manually detuned by $\pm 14\%$ pT/m². The data was reconstructed first with no off-resonance correction applied; second, with a linear correction; and finally with a semi-antiomatic method that has been reported previously (12). For opiral-in/out data, the off resonance correction was applied to each component image prior to combining the data in image space. The reconstruction time for each method was recorded.

To test the redundant in/c at trajectory *in vi* v, a slab-selective version of the variable-flipangle 3D spiral TSE (spiral SPACE) sequence was used on a normal volunteer for 1_{2} weighted brain imaging. Sector parameters were as follows: TR/7 E 3000/200 spiral duration 6.4 ms, in-plane FOV 250 mm, number of slices = 64, slice thickness 1.0 mm. Forty-mine interleaves were used for both spiral-out and spiral-in/out acquisitions, nor spiral-in/out, the second, redundant interleaf scan was combined with the chopped scan so that the total acquisition time for both sequence variations was identical. No of the improved off-resonance correction algorithm was applied in reconstruction in order to better exhibit the improved off-resonance performance of the redundant in/or eservence.

Results

The redundant spiral-in/out trajectory shows execution robustness for off-resonance values ranging up to 0.5 cycles (Fig. 3a). Above this value, blurring appears and is comparable to a

spiral-out scan performed with han the about of off-resonance applied; consistent with the result, predicted by Fig. 2. Figure 5 compares the spiral-out and spiral-in/out trajectories for double-oblique imaging places near the magnet isocenter and off-center. The spiral-in/out trajectory results in image targely free from blurring due to concomitant fields.

In the presence of non-linear B₀ variation, an uncorrected redundant in/out trajectory natively performs better than an uncorrected spiral-out acquisition (Fig. 4a vs. 4d). Applying a linear correction to the data trajectorie during reconstruction cannot fully correct the spiral-out image because the underlying field is non-linear (Fig. 4b vs. 4a). With only a timeer correction, the redundant in/out trajectory can correct the residual blurring (Fig. 4e), whereas the optical-out data requires more advanced reconstruction algorithms, such as a semi-automatic correction (Fig. 4c), to correct for the non-linear field. The reconstruction times for each off resonance correction method was nearly identical when applied to spiralout and spiral in/out data and was as follows. No correction: 0.14 sec; Linear correction: 0.26 sec; Semi-artiomatic: 9.7 sec.

One slice from a 3D stack-of-spirals TSC *in vivo* dataset acquired with both spiral rajectories is shown in Fig. 5. Overall, the images acquired with the spiral-in/out trajectory are shown in the much improved off-resonance performance. An additional figure showing the trajectories' performance for non-contrast MPA is available in the Supplemental Materials.

Discussion and Conclusion

Although all of the data presented here was accorded with a single receive channel, we anticipate this technique will perform well in pereller imaging implementations, allowing fast reconstructions to take place at scan time. One major nurdle to the wide-spread adoption of non-Cartesian parallel reconstructions is their complexity due to the fact that they must address the off-resonance issue in addition to removing non-Cartesian aliasing artifacts.

For small values or on resonance, the *k*-space signal 'n coundant's impling experiences a cosine amplitude window, which is a different (and more benign) mechanism for resolution loss compared to the PST broadening observed in spiral out scarning. As in regular spiral imaging, this slight l turring will be space-variant depending on tocal off-resonance values. For larger off-resonance values (or for long readout length), the cosine modulation will begin nulling important frequencies in k-space as a function of the space radius, resulting in more severe image artifacts. Therefore, we ant organe the use of this trajectory inclusion and other methods where appropriate. The use of the reduction in/or trajectory allows a grouter notion of error for post-processing methods, resulting in better image quality and/or facilitating the use of less robust (but potentially simple, and faster) on resonance correction algorithms, as demonstrated in Figure 4.

The robustness of this method with regards to concordinate fields is potentially very important. Most current off-resonance correction, anethous do not incorporate concomitant gradient correction, and we believe that this is essential for our isocenter cans

Spirals are mostly lower and cliner for their data acquisition efficiency (speed), their robustness to metion for their shour rEs. A few general comments in regards to spiral-in/out trajectories should be made. The data acquisition efficiency and motion characteristics of in/out trajectories are ver a comprisable to spiral-out scanning. However, the time required to spiral in with these trajectories precludes short TE acquisitions. For all experiments in this work, the resulting in longer TEs that high hormally be used with spiral-in/out scans were identical, resulting in longer TEs that high hormally be used with spiral-in/out imaging has other issues that spiral-out scanning dries not, such as eddy currents causing trajectory warping leading to inconsistencies at the conter of k-space.

The fact that the redundant spiral-in/out trajectory necessarily requires twice the number of interleaves as a spire l-out trajectory to achieve a similar resolution cannot be overlooked. We have shown that for slab-selective 3D spiral TSE imaging, at least, the redundant accurstion may be combined with the RF-chopped second average for no penalty in scan time. For other cases where spiral trajectories are regularly used with multiple averages and long TFs (e.g. fMRI, A.Sil), or for time resolved second regularly used with multiple averages and long TFs (e.g. fMRI, A.Sil), or for time resolved second regularly used with multiple averages and long TFs (e.g. fMRI, A.Sil), or for time resolved second regularly used at a cquisition can be interleaved with the base acquisition in time and neighboring redundant in/out data can be combined with a view-sharing" approach. In this way, overall scan time is not affected (althoug in temporal resolution is). Recently, Jung, et al (21) reported an interleaved, high-resolution in/out trajectory for fMRI and showed and substituting with a conjugate phase reconstruction. There, we have shown this technique to have autity for other types of spiral imaging and have explored the limits of the approach.

The redundant spiral-in/out trajectory is an attractive accorrection-hased method to naturally mitigate blue due to off-resonance in spiral scanning. The technique produces excellent image quality with the simplest of post-acquisition correction methods, allowing simple, fast reconstructions

Supplementary Material

Refer to Web version or rubMr a Central for supply mentary material.

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Fig. re 1

Spiral gradients and train crories. a) Spiral-out read gradient (solid) with rewinding lobe (dashed), b) Spiral-ir, out read gradient, with prewinding and rewinding lobes. c,d) k-space trajectories for the gradients in (a,b), respectively. Arrows indicate direction of travel in k-space. c) k-space trajectory for second spiral-in/out trajectory, rotated 180° from the trajectory in (d). The redundant spiral-in/out schemes averages data acquired with trajectories (d) and (e)



Figure 2

Normalized MTFs of spilal in/cut trajectories MTFs are normalized to on-resonance (0 vycles) case, a) Center-frequency official Until ~ 0.5 cycles of phase are accrued, the cosine function never reaches its first zero, reculting in a relatively benign amplitude modulation in k-cpace: however, important frequencies are lost as the amount of off-resonance increases. b) Ch-resonance due to concomitant field's for off-resonance due to concomitant fields shows good performance control expected up to about 1 cycle of accrued phase. c) The effect of decay during the readout Phinering amounts of T₂ decay are simulated with 0.5 cycles of off-resonance phase accrual. Severe shapping of the MTF is only seen when the T₂ time constant approaches the readout length (10 ms)



Figure 3

Trajectory performance in a phantom. a) Phantom images at various amounts of offresonance. Images were acquired with a 10 ms readout, resulting in cycles of phase accumulated at the end of the readout: 0 int = 0 cycles, 20 Hz = 0.2 cycles, 40 Hz = 0.4 cycles, 8° Hz = 0.8 cycles, 16° int = 1.0 cycles. b) Trajectory performance in the presence of concomitant fields. Note blurring hear the top of the phantom due to concomitant fields in the spiral-out, offcenter image, and the absence of this blurring in the in/out image. Across the entire figure top row: spiral-out, bottom row, spiral-in/out. No off-resonance correction was applied in any reconstruction.

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Figure 4

Post processing techniques apriled to remove blurring due to off-resonance in spiral-out (a-·) and "coundant spiral-in/out (d-f) "equisition. A highly non-linear field was applied in the left-right direction with internationally prod shimming. a) No correction; spiral-out. b) Linear conection, spiral-out. c) Semi-autonatic correction; spiral-out. d) No correction; redundant spiral-in/out. e) Linear correction; redundant spiral-in/out. f) Semi-automatic correction; redundant spiral-in/out.



Figure 5

Compari, on of spiral-out and rodundant spiral-in/out in a volunteer with no off-resonance correction applied to either dataset. One slice is shown from a slab-selective 3D stack-ofspirals trajectory. a) Spiral-out. Significant blurring is observed in areas with poor B_0 homoger sity (insets). b) Spiral-in/out. Excellent off-resonance performance is obtained with no increase in scan time.