# Fate and Transport of Molybdenum Disulfide Nanomaterials in Sand Columns

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## Abstract

Research and development of two-dimensional transition metal dichalcogenides (TMDC) (e.g., molybdenum disulfide  $[MoS<sub>2</sub>]$ ) in electronic, optical, and catalytic applications has been growing rapidly. However, there is little known regarding the behavior of these particles once released into aquatic environments. Therefore, an indepth study regarding the fate and transport of two popular types of  $MoS<sub>2</sub>$  nanomaterials, lithiated ( $MoS<sub>2</sub>-Li$ ) and Pluronic PF-87 dispersed ( $MoS<sub>2</sub>-PL$ ), was conducted in saturated porous media (quartz sand) to identify which form would be least mobile in aquatic environments. The electrokinetic properties and hydrodynamic diameters of  $MoS<sub>2</sub>$  as a function of ionic strength and pH were determined using a zeta potential analyzer and dynamic light scattering techniques. Results suggest that the stability is significantly decreased beginning at 10 and 31.6 mM KCl, for  $MoS<sub>2</sub>-PL$  and  $MoS<sub>2</sub>-Li$ , respectively. Transport study results from breakthrough curves, column dissections, and release experiments suggest that  $MoS<sub>2</sub>$ -PL exhibits a greater affinity to be irreversibly bound to quartz surfaces as compared with the  $MoS<sub>2</sub>-Li$  at a similar ionic strength. Derjaguin–Landau–Verwey– Overbeek theory was used to help explain the unique interactions between the  $Mo_{2}$ -PL and  $Mo_{2}$ -Li surfaces between particles and with the quartz collectors. Overall, the results suggest that the fate and transport of  $MoS<sub>2</sub>$ is dependent on the type of  $MoS<sub>2</sub>$  that enters the environment, where  $MoS<sub>2</sub>-PL$  will be least mobile and more likely be deposited in porous media from pluronic–quartz interactions, whereas MoS<sub>2</sub>-Li will travel greater distances and have a greater tendency to be remobilized in sand columns.

Key words: effects of pH and ionic strength; molybdenum disulfide  $(MoS<sub>2</sub>)$ ; nanomaterials; pluronic coating; transport in packed bed column; stability; straining

## Introduction

**TWO-DIMENSIONAL NANOMATERIALS have gained a sig-**<br>nificant amount of attention recently (Cooper *et al.*, 2013). The most common two-dimensional material has traditionally been graphene (Lotya *et al.*, 2009) or its oxidized derivative graphene oxide (Compton and Nguyen, 2010). However, research and development of molybdenum disulfide  $(MoS<sub>2</sub>)$ , a two-dimensional transition metal dichalcogenide (TMDC) has been increasing over the last few years (Gong *et al.*, 2013) and is beginning to rival the popularity of graphene for monolayer materials (Liu *et al.*, 2014). In 2012, an editorial in the journal

are gaining a tremendous amount of attention with tungsten disulfide  $(WS_2)$  and  $MoS_2$  leading the way for future optical and electronic applications (Wang  $et$   $al$ , 2012b). MoS<sub>2</sub> is a layered TMDC, which includes two planes of sulfur (S) atoms with a single plane of molybdenum (Mo) atom in the center of the structure with the Mo and S atoms being bound together by covalent bonds (Silbernagel, 1975). The rise in popularity with  $MoS<sub>2</sub>$  is due to its many novel

*Nature* noted that ''Graphene is Not Alone'' and that TMDCs

electrical and optical properties (Najmaei *et al.*, 2014), including a tunable band gap (Lopez-Sanchez *et al.*, 2013), exceptional photo-response (Lopez-Sanchez *et al.*, 2013), and high charge carrier mobility (Wang *et al.*, 2012a). MoS<sub>2</sub> is also a direct band gap semiconductor (Zhou *et al.*, 2013), which could make it an ideal material for transistors or optoelectronics devices (Mak *et al.*, 2010). Other applications for MoS<sub>2</sub> include lithium ion batteries (Stephenson *et al.*, 2014), electrochemical charge storage (Winchester *et al.*,

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2014), lubricants (Khare and Burris, 2014), and ultra-fast photonics (Zhang *et al.*, 2014). With the widespread research and development of engineered nanomaterials, the implementation of  $MoS<sub>2</sub>$  nanomaterials has also gained a tremendous amount of interest for use in nanotechnological applications (e.g., catalysts, and nanotubes) (Lauritsen *et al.*, 2007). Whereas the production and development of  $MoS<sub>2</sub>$ NPs has a tremendous technological potential, there is a lack of knowledge and research regarding its behavior if released into the environment. To the authors' knowledge, there are no reports that investigate the fate or transport of  $MoS<sub>2</sub>$  in aquatic environments.

In the past, it has been reported that other two-dimensional nanomaterials (e.g., graphene oxide) can exhibit toxic effects on living organisms (e.g., increased mitochondrial respiration in mice) (Duch *et al.*, 2011). Currently there is little information regarding the potential toxicity of  $MoS<sub>2</sub>$  nanomaterials. In 2011, Wu *et al.* (2011) reported that inorganic fullerene-like  $MoS<sub>2</sub>$  nanomaterials prepared using pulsed laser ablation were biocompatible and provided a green method to synthesize nanomaterials. However, it has been reported in the past that elevated levels of molybdenum (Mo) in ruminate intake can result in a physiological copper (Cu) deficiency, which may have negative consequences (Ward, 1978).

In 2013, Keller *et al.* (2013) reported that it is highly likely that some portion of engineered nanomaterials that reach a landfill will make their way into the surrounding environment. In 2014, Conway *et al.* (2014) reported that phytoplankton sorbed engineered nanomaterials  $(CeO<sub>2</sub>)$  in less than 1 h and had negative impacts on marine mussels in seawater environments. This suggests that if the careful implementation of nanomaterials is not employed, then their accidental release into the environment may occur and have long-lasting effects on aquatic life. Therefore, it is crucial that the science community and policy makers have a comprehensive set of research data to help predict the fate and transport of  $MoS<sub>2</sub>$ nanomaterials once introduced into aquatic environments.

This study is set forth to bridge the gap of knowledge regarding the influence that environmentally relevant parameters (e.g., ionic strength and pH) will have on the fate, transport, and aggregate morphology of the two forms of  $MoS<sub>2</sub>$  nanomaterials: Lithiated  $(MoS_2-Li)$  and Pluronic  $(MoS_2-PL)$ .  $MoS<sub>2</sub>-Li$  is exfoliated through a lithiation process, during which lithium ions intercalate into the layered bulk material and enable the exfoliation of the material using gentle sonication. A dialysis step following the exfoliation removes the lithium ions as well as any residual organics, leaving only trace amounts of lithium present on the surface of the nanomaterial. The exfoliation process produces two-dimensional flakes of  $MoS<sub>2</sub>$  that are overall negatively charged and thus remain dispersed in aqueous environments without the addition of surfactants. The  $MoS_2$ -Li was chosen since this represents the most scalable and controllable synthesis method for monolayer  $MoS<sub>2</sub>$  that will likely be used in the future, therefore, representing one of the most abundant forms of two-dimensional  $MoS<sub>2</sub>$  that would be found in the environment. On the other hand,  $MoS_2-PL$  is exfoliated by ultrasonication with the aid of a biocompatible triblock copolymer, commercially available as Pluronic, which acts as a surfactant and stabilizes the  $MoS<sub>2</sub>$  flakes in aqueous solution. The Pluronic triblock copolymer is a linear molecule containing a central hydrophobic chain of polypropylene oxide (PPO) functional groups and is bordered by two hydrophilic chains of polyethylene oxide (PEO) of equal length (Seo *et al.*, 2011). It is predicted that for industrial applications, a stable and monodisperse suspension of  $MoS<sub>2</sub>$  will be required and, therefore, surfactants will play an increasing role in  $MoS<sub>2</sub>$  nanotechnology development (Vaisman *et al.*, 2006). Recent studies have also reported using Pluronic surfactants to stabilize carbon nanotubes (Vaisman *et al.*, 2006), and other two-dimensional nanomaterials such as graphene (Seo *et al.*, 2011). Therefore, in this study, Pluronic PF-87 was chosen as a representative surfactant to produce an additional form of dispersed  $M_0S_2$  to investigate its behavior in aquatic environments.

This study is the first to report the effects of two important environmentally relevant parameters (i.e., ionic strength and  $pH$ ) on  $MoS<sub>2</sub>$  nanomaterials stability and transport in a packed bed column and provides critical information regarding how interactions between quartz surfaces influence the transport of  $MoS_2-PL$  and  $MoS_2-Li$ . This report also helps to answer the question of whether traditional colloid filtration theory can apply to the planar TMDCs. The results from this report help to explain the behavior of  $MoS<sub>2</sub>$  in engineered or natural systems where they may settle or deposit on sediments in the subsurface and sheds light on the chemical and physical mechanisms responsible for this behavior.

### Experimental Protocols

### Molybdenum disulfide

The  $MoS<sub>2</sub>$  nanomaterials used in this study were synthesized by the following process. To synthesize the lithiated  $MoS<sub>2</sub> (MoS<sub>2</sub>-Li)$ , lithium intercalation was achieved by combining  $MoS<sub>2</sub>$  powder with butyllithium in a low vapor and oxygen environment for 2 days while stirring. The  $MoS<sub>2</sub>$ -Li was retrieved from the solution by filtration and immersed in deionized water. To avoid deintercalation, the exfoliation was carried out immediately using a bath sonicator. A centrifugation step was performed to remove unexfoliated material. The resulting supernatant was dialyzed for 7 days in a bath of deionized water to remove any excess lithium and hexane. To synthesize the pluronic  $MoS<sub>2</sub> (MoS<sub>2</sub>-PL)$ ,  $MoS<sub>2</sub>$ powder was exfoliated by ultrasonication in a solution of 2% w/v Pluronic F87 in deionized water.

Following ultrasonication, the slurry was centrifuged to remove any unexfoliated material and aggregates. Only the top 80% of the supernatant was retained. The lithiated  $MoS<sub>2</sub>$  $(MoS<sub>2</sub>-Li)$  used in this study had a mean square root surface area of  $(187.5 \pm 126.9 \text{ nm})^2$ , a mean height of  $3.9 \pm 0.6 \text{ nm}$ , and a mean perimeter of  $689.1 \pm 433.0$  nm. Whereas the pluronic  $MoS_2 (MoS_2-PL)$  used in this study had a mean square root surface area of  $(27.6 \pm 15.6 \text{ nm})^2$ , an average height of  $5.2 \pm 1.1$  nm, and an average perimeter of  $100.8 \pm$ 57.3 nm. The  $MoS_2-PL$  and  $MoS_2-Li$  were previously dispersed in deionized (DI) water at stock concentrations of 0.125 and  $0.038$  mg/mL, respectively. To create  $10$  mg/L MoS<sub>2</sub> dilutions at each ionic strength (IS) used in this study (0.1–  $100 \text{ mM KCl}$ ,  $MoS<sub>2</sub>$  nanomaterials were introduced to the electrolyte and either vortexed (Fisher Scientific, Mini Vortexer) or lightly shaken for  $\geq 5$  s. Potassium chloride (KCl) was used as the background electrolyte (ACS Research Grade, Fisher Scientific, Pittsburgh, PA) in all experiments. The unadjusted pH of the  $MoS_2$ -Li and  $MoS_2$ -PL stock suspensions were  $3.9 \pm 0.1$  and  $5.3 \pm 0.1$ , respectively. Potassium hydroxide (KOH) and hydrochloric acid (HCl) were used to adjust the pH of the  $MoS<sub>2</sub>$  suspensions. All  $MoS<sub>2</sub>$  suspensions were created using DI water (nanopure water at  $> 18.2 M\Omega$  cm at 25°C).

# General characterization of nanomaterials and porous media

Electrokinetic properties and hydrodynamic diameter of  $MoS<sub>2</sub>$ . The influence of pH and ionic strength (IS) on the  $MoS<sub>2</sub>-PL$  and  $MoS<sub>2</sub>-Li$  was determined by measuring the effective hydrodynamic diameter and electrophoretic mobility (EPM) across an environmentally relevant pH range of 4–10 and an IS range of 0.1–100 mM KCl, respectively. Dynamic light scattering (Brookhaven model BI-9000, Holtsville, NY) measurements taken at a wavelength of  $661$  nm and a scattering angle of  $90^\circ$  were used to determine the effective hydrodynamic diameter of the  $MoS<sub>2</sub>$  suspensions. The EPM was measured using a ZetaPALS analyzer (Brookhaven Instruments, Holtsville, NY). Triplicate measurements were taken at room temperature  $(23^{\circ}C \pm 1^{\circ}C)$  for all of these electrokinetic and size measurements.

Elektrokinetic characterization of quartz. Additional characterization of the quartz sand collectors was performed to help explain any potential interactions between the  $MoS<sub>2</sub>$ particles and the quartz surfaces during transport in the packed bed column and was used in the application of DLVO theory (Tables 2 and 3). The streaming potential of the quartz sand (used as the porous media in the column experiments) was measured using a streaming potential analyzer (sur-PASS, Anton Paar, Graz, Austria). Triplicate measurements were taken in deionized water at pH 5 and across the same range of ionic strength (1–31.6 mM KCl) conditions used in the transport experiments.

Scanning electron microscope images of  $MoS<sub>2</sub>$  and quartz. Scanning Electron Microscopy (SEM) images were taken of the dried aggregated  $MoS<sub>2</sub>$  nanoparticles on quartz slides. MoS<sub>2</sub> samples were prepared by adding  $\sim 15 \mu L$  of stock  $MoS_2-PL$  and  $MoS_2-Li$  at an unadjusted pH in deionized water to quartz slides and were allowed to dry overnight before taking SEM images. The SEM images of the quartz sand used as porous media in the column experiments were also taken. Quartz samples were prepared by adding a small amount  $\sim$  30 sand grains in deionized water to a quartz slide and left to dry overnight.

#### Packed bed column transport experiments

The effect of ionic strength (1–31.6 mM KCl) on the transport of  $MoS_2$ -Li and  $MoS_2$ -PL in saturated porous media was conducted using a packed bed column at pH 5. Transport experiments were conducted using an inverted borosilicate glass column (Omnifit, Boonton, NJ) with an inner diameter of 1.5 cm and a packing depth of 5 cm. Details regarding the packed bed column transport methods, conditions, and sand cleaning procedures can be found elsewhere (Lanphere *et al.*, 2013). Briefly, quartz sand was sieved to obtain an average diameter  $(d_{50})$  of 275  $\mu$ m and purified before being wet packed at a porosity of  $0.47 \pm 0.01$  in the column before transport experiments.  $MoS<sub>2</sub>$  suspensions were injected into the column at a flow rate of 2 mL/min to represent similar flow conditions found in the engineered or natural systems (Hemond, 2009) using peristaltic pumps and the effluent was collected every minute using a fraction collector (CF 1 Fraction Collector, Spectrum Chromatography, Houston, TX). The effluent was used to determine the concentration at each time interval using a spectrophotometer (DU 800 Beckman Coulter, Fullerton, CA) to correlate the absorbance with concentration. Calibration curves for the  $MoS<sub>2</sub>-Li$  and  $MoS<sub>2</sub>-PL$  were created by measuring the absorbance as a function of  $MoS<sub>2</sub>$  concentration and an  $R<sup>2</sup>$  value of 0.999 was obtained for both types of  $MoS<sub>2</sub>$ .

Release experiments were also conducted following select transport experiment at each IS for  $MoS_2-PL$  and  $MoS_2-Li$  to investigate the release of particles as a function of the decrease in IS. To facilitate the release experiments, an additional 12 pore volumes (PV) of DI water was injected into the column after the background absorbance dropped to  $\sim 0$  and the effluent was collected and the absorbance measured as mentioned earlier.

## Column dissection

To evaluate the distribution by which  $MoS<sub>2</sub>$  was retained in the column, select columns were dissected following a transport experiment. Deposition profiles of  $MoS<sub>2</sub>$  at each centimeter in the column were developed from the column dissections to interpret the distribution of  $MoS<sub>2</sub>$  following an experiment at each ionic strength. The procedure for the column dissection can be found elsewhere (Lanphere *et al.*, 2013) and was based upon a previously developed protocol (Tufenkji and Elimelech, 2004). Briefly, once the experiments were completed, the bottom fitting of the column was removed and the background electrolyte was injected into the column. Next, the sand slowly flowed out of the column by gravity and approximately one centimeter portions of the sand were collected in centrifuge tubes as they came out of the column. The collected samples–comprising of quartz, KCl, and retained  $M_0S_2$ particles–were rinsed with 25 mL of DI water and vortexed for 30 s. The supernatant from this step was decanted into a centrifuge tube and saved for further analysis. This step was repeated twice to achieve a total volume (50 mL) of reversibly bound  $MoS<sub>2</sub>$ . These  $MoS<sub>2</sub>$ -containing supernatant suspensions were then taken to a spectrophotometer to measure the absorbance of the  $MoS_2-PL$  and  $MoS_2-Li$  suspensions obtained for each centimeter at wavelengths of 258 and 247 nm, respectively. The absorbance values for the supernatant collected at each centimeter were then correlated with  $MoS_2-PL$ and  $MoS<sub>2</sub>-Li$  concentrations using the same calibration curves mentioned previously. The remaining sand at each centimeter from the dissected column was then dried and the mass was recorded. Results were then plotted as the mass of nanoparticles retained in each centimeter of sand  $\lceil$ (mass of MoS<sub>2</sub>) retained/mass of  $MoS<sub>2</sub>$  injected)/mass of quartz sand dissected as a function of the depth of the column. These retention profiles were then developed to further understand the deposition behavior of  $MoS<sub>2</sub>$  throughout the column as a function of ionic strength.

#### Results and Discussion

#### Effects of pH and ionic strength on  $MoS<sub>2</sub>$  stability

To investigate the behavior of two leading  $MoS<sub>2</sub>$  nanomaterials to identify which one is least mobile in the



FIG. 1.  $MoS_2-PL$  evaluated as a function of pH for effective hydrodynamic diameter and electrophoretic mobility.  $MoS<sub>2</sub>-PL$  concentration was maintained at 10 mg/L and 10 mM KCl. Error bars indicate one standard deviation of triplicate measurements.

environment,  $MoS_2-PL$  and  $MoS_2-Li$  were extensively characterized. It has been reported in the past (Elimelech and Omelia, 1990; Wiesner and Lecoanet, 2004; Chowdhury *et al.*, 2013) that nanoparticle size (Chen and Elimelech, 2006), and surface charge (Keller *et al.*, 2010b) characterization can provide critical insights into predicting the ultimate fate of these materials in the environment. Therefore, a comprehensive study was performed to investigate the effects of solution chemistry on  $MoS_2$ -Li and  $MoS_2$ -PL across an environmentally relevant range of aquatic chemistries and the results are shown in Figures 1–4. The pH range studied was 4–10, and the IS spectrum was 0.1–100 mM KCl. This wide range of solution chemistry covers not only environmental conditions for most natural water (e.g., lakes, and streams), (Hemond, 2009; Crittenden *et al.*, 2012), but extreme conditions as well (i.e., pH 9– 10 and 31.6–100 mM KCl) (Hemond, 2009). The influence of pH on the diameter and EPM varied between the two types of  $MoS<sub>2</sub>$  tested as a function of their type. While it was expected that the size and surface charge of  $MoS<sub>2</sub>-PL$  would not significantly change across the environmental pH range (5–9) tested due to the surfactant PF-87 coating present, it was not known what the behavior would be at the lower and upper pH values (4, 10). In addition, it was anticipated that  $MoS<sub>2</sub>-Li$ would not be sensitive to pH (from 4 to 10) with respect to its diameter size and EPM as has been reported previously for other two-dimensional nanomaterials (e.g., graphene oxide) (Chowdhury *et al.*, 2013). Overall this hypothesis was correct



FIG. 3.  $MoS<sub>2</sub>-Li$  evaluated as a function of pH for effective hydrodynamic diameter and electrophoretic mobility. MoS<sub>2</sub>-Li concentration was maintained at  $10 \text{ mg/L}$  and 10 mM KCl. Error bars indicate one standard deviation of triplicate measurements.

and there was little effect on the effective diameter and EPM of  $MoS<sub>2</sub>$  as a function of pH within environmentally relevant parameters (pH 5–9). It was only when the pH was near or outside of this range when a significant difference in stability for  $MoS_2$ -Li and  $MoS_2$ -PL was observed (Figs. 1 and 3).

The effective diameter for  $MoS_2$ -Li (as measured by DLS) was much more sensitive to the change in pH than anticipated. Specifically, at pH 4, the effective diameter  $(401 \pm 106 \text{ nm})$ for  $MoS_2$ -Li was  $\sim$  25% larger than the effective diameter  $(298 \pm 12 \text{ nm})$  at pH 5 as seen in Figure 3. However, there was no significant effect ( $p > 0.05$ ) on the MoS<sub>2</sub>-Li EPM across the range of pH (4–10) tested (Fig. 3). Similar results have been reported for other planar nanomaterials (e.g., graphene oxide) (Chowdhury *et al.*, 2013; Lanphere *et al.*, 2013) where no significant effect on EPM or diameter was observed from pH 5–9. However, in recent studies, the zeta potential of spherical nanoparticles [ZnO, (Mohd Omar *et al.*, 2014), Ag (Yin *et al.*, 2014), and  $TiO<sub>2</sub>$  (Wang *et al.*, 2014)] were shown to be sensitive to changes in pH across an environmentally relevant range (5–9). For the  $MoS_2-PL$ , the opposite trend occurred when compared with the  $MoS_2-Li$  with regard to the influence on diameter and EPM. Specifically, there was no significant effect  $(p>0.05)$  in the diameter from pH 4–10, however, the EPM did significantly increase from  $-0.9 \pm 0.1$ to  $-1.4 \pm 0.1$  (10<sup>-8</sup> m<sup>2</sup>/[V·s]) at pH 8 and 9, respectively. These results suggests that the  $MoS<sub>2</sub>-PL$  will be the most stable across an environmental pH range (5–9) in natural



FIG. 2.  $MoS_2-PL$  evaluated as a function of ionic strength for effective hydrodynamic diameter and electrophoretic mobility.  $MoS_2-PL$  concentration was maintained at  $10 \text{ mg}/$ L and at an adjusted pH 5. Error bars indicate one standard deviation of triplicate measurements.



FIG. 4.  $MoS_2-Li$  evaluated as a function of ionic strength for effective hydrodynamic diameter and electrophoretic mobility. MoS<sub>2</sub>-Li concentration was maintained at  $10 \text{ mg/L}$ and at an adjusted pH 5. Error bars indicate one standard deviation of triplicate measurements.

waters and will not likely be removed during transport in sediment beds since the effective diameter will remain small as a result of the PF-87 coating. However, for the  $MoS_2-Li$ , the sensitivity to low  $pH \left( \langle 5 \rangle \right)$  with respect to the increase in effective diameter, may result in the  $MoS<sub>2</sub>-Li$  being settled down or sediments removed.

It was anticipated that the increase in IS would decrease the electrostatic double layer repulsion between the  $MoS<sub>2</sub> NPs$ and, therefore, result in the increased aggregation as has been reported in other nanoparticle studies (Keller *et al.*, 2010a; Petosa *et al.*, 2010, 2012; Chowdhury *et al.*, 2011) with TiO<sub>2</sub>,  $ZnO$ , and  $CeO<sub>2</sub>$ . However, it was unknown how these effects of IS would vary amongst the two types of  $MoS<sub>2</sub>$  characterized in this study or at what IS would cause a significant change.

In general, the  $MoS_2-PL$  and  $MoS_2-Li$  became less stable (as observed by an increase in effective diameter and a decrease in EPM) as a function of IS (Figs. 2 and 4, respectively). The average effective hydrodynamic diameters of the  $MoS<sub>2</sub>-Li$  and  $MoS<sub>2</sub>-PL$  were measured in DI water and were  $345 \pm 4$  and  $81 \pm 1$  nm, respectively. A significant effect  $(p<0.05)$  of IS on the effective diameter of MoS<sub>2</sub>-Li did not occur until 31.6 mM KCl was achieved when the diameter increased from  $340 \pm 13$  to  $1795 \pm 413$  nm, respectively. Furthermore, the  $MoS<sub>2</sub>-Li$  effective diameter increased almost by an order of magnitude from  $340 \pm 13$  nm to  $3057 \pm 402$  nm, across the range of 0.1–100 mM KCl.

This increase in  $MoS_2$ -Li aggregate diameter as a function of IS is likely a result of the decrease in electric double layer repulsion at higher IS and is in agreement with previous studies performed with carbon nanotubes (Chowdhury *et al.*, 2012b). However, the effective diameter of the  $MoS<sub>2</sub>-PL$ remained fairly constant  $(81 \pm 7 \text{ nm})$  across the IS range tested, and only increased by  $\sim$  17% from 0.10 to 100 mM KCl. This was expected as a result of the pluronic surfactant coating on the  $MoS_2$ , which limited its aggregation as a function of IS. Similar results have been reported for other engineered nanomaterials [e.g., fullerenes (Xie *et al.*, 2008), and  $TiO<sub>2</sub>$  (Chowdhury *et al.*, 2012a)] coated with natural organic matter where an increase in particle stability (as observed by a reduction in effective diameter) was reported as a function of NOM concentration due to steric repulsion.

Similar to the effects of IS on the effective diameter, it was expected that increasing the IS would decrease the overall EPM for both types of the  $MoS<sub>2</sub>$  as has been observed previously in the Chen and Elimelech (2009) study. As a baseline (lowest IS), the EPMs of  $MoS_2-Li$  and  $MoS_2-PL$  were measured in DI water and were  $-3.8 \pm 0.1$  and  $-1.8 \pm 0.0$  $(10^{-8} \text{ m}^2/[\text{V} \cdot \text{s}])$ , respectively. Similar IS conditions that influenced the effect on size for both types of  $MoS<sub>2</sub>$  were observed regarding the influence on EPM. For example, the  $MoS<sub>2</sub>-Li$  and  $MoS<sub>2</sub>-PL$  EPMs were not significantly affected until an IS of 31.6 and 10 mM KCl was achieved, respectively. For the MoS<sub>2</sub>-Li, the EPM decreased from  $-3.8 \pm 0.0$ to  $-2.5 \pm 0.4$  (10<sup>-8</sup> m<sup>2</sup>/[V·s]) across the range of 0.1–100 mM KCl. The EPM of the  $MoS_2-PL$  was also decreased from  $-1.7\pm0.1$  to  $-0.3\pm0.3$  (10<sup>-8</sup> m<sup>2</sup>/[V·s]) across the range tested. These behaviors for both types of  $MoS<sub>2</sub>$  with respect to their sensitivity to IS as observed by their decrease in EPM, are in agreement with traditional Derjaguin–Landau–Verwey– Overbeek (DLVO) theory regarding the stability for colloidal particles (Verwey, 1947).

**FIG. 5.** Packed bed column transport results for  $MoS_2-PL$ as a function of ionic strength. (A) 0–10 Pore volumes: breakthrough curves as a function of IS (1–31.6 mM KCl) are an average of triplicate measurements with error bars indicating one standard deviation. (B) 10–22 Pore volumes: breakthrough curve after DI was injected to determine the amount of  $MoS_2-PL$  reversibly bound.

## Packed bed column transport study

Column experiments were conducted with  $MoS<sub>2</sub>-Li$  and  $MoS<sub>2</sub>-PL$  to investigate the influence of ionic strength on the movement of  $MoS<sub>2</sub>$  in aquatic environments (e.g., sediment beds) across an environmentally relevant range of IS conditions (1–31.6 mM KCl at pH 5) and to determine which type of  $MoS<sub>2</sub>$  would be least mobile in aquatic environments. This specific range of solution chemistries was also chosen since it represents the regions where the most significant effects were observed during particle characterization (DLS and EPM) experiments. The transport results were plotted as breakthrough curves (BTCs) and are found in Figures 5 and 6. The BTCs represent the concentration of  $MoS<sub>2</sub>$  eluted, normalized by the injection concentration  $(C/C_0)$  as a function of PV that flowed through the column and provide a measure for the  $MoS<sub>2</sub>$  transport in porous media. The shapes of the BTCs also provide critical information to understand the potential deposition or retention phenomena (e.g., blocking, straining, ripening) that may be occurring during transport.

The normalized  $MoS_2$  concentration  $(C/C_0)$  was also examined and compared at each IS for  $MoS_2-PL$  and  $MoS_2-Li$ to identify which type of  $MoS<sub>2</sub>$  would be least mobile in sand

 $-1$  mM

в

 $-1 - 1$  mM



**FIG. 6.** Packed bed column transport results for  $MoS_2$ -Li as a function of ionic strength.  $(A)$  0–12 Pore volumes: breakthrough curves as a function of IS (1–31.6 mM KCl) are an average of triplicate measurements with error bars indicating one standard deviation. (B) 12–24 Pore volumes: breakthrough curve after DI was injected to determine the amount of  $MoS<sub>2</sub>-Li$  reversibly bound.



columns. In addition, the overall mass of  $MoS<sub>2</sub>-PL$  and  $MoS<sub>2</sub>-Li$  deposited in the column during transport was determined by integrating under the BTCs after each transport experiment and the results are shown in Table 1.

Overall, the deposition of  $MoS_2$ -Li and  $MoS_2$ -PL in the column increased for both nanomaterials as a function of IS, however, a significant effect was observed at different conditions dependent on the type of  $MoS<sub>2</sub>$ . Specifically, no significant effect ( $p > 0.05$ ) on MoS<sub>2</sub>-Li transport was observed until an IS of 31.6 mM KCl was achieved, and the mass deposited was greater compared to the  $MoS<sub>2</sub>-PL$ . For example, at 10 mM KCl the C/C<sub>0</sub> value was  $\sim$  0.95 for the MoS<sub>2</sub>-Li whereas at the same IS, the  $C/C_0$  value for the  $MoS_2-PL$  was  $\sim$  0.20. This difference was also observed between the two types of MoS<sub>2</sub> at 31.6 mM KCl where  $88.7\% \pm 2.5\%$  and  $96.9\% \pm 0.2\%$  of the particles were retained in the column for the  $MoS_2-Li$  and  $MoS_2-PL$ , respectively. This increase in deposition for the  $MoS<sub>2</sub>-PL$  is interesting since these particles are roughly 19 times smaller (as measured in hydrodynamic diameter) than the  $MoS<sub>2</sub>-Li$  counterpart. One might expect a greater retention in the column as the effective diameter of the aggregates increase due to the straining mechanism, which suggests that particles are retained and cannot pass through regions of the pore spaces due to their large size (Bradford *et al.*, 2002, 2003).

Straining can occur when the straining parameter (*dparticle/*  $d_{quartz}$ , which is the diameter of the particle ( $d_{particle}$ ) divided by the diameter of the quartz collector  $(d_{quartz})$ , reaches the theoretical threshold of 0.002 (Bradford *et al.*, 2003, 2005). Straining has been used in the past during similar column transport studies to help explain removal mechanisms observed for latex particles (Porubcan and Xu, 2011), nonspherical colloids (Xu *et al.*, 2006), and other twodimensional nanomaterials (e.g., graphene oxide) (Lanphere *et al.*, 2013). For  $MoS_2-PL$ , the effective diameter at the highest IS (31.6 mM KCl) correlated with a straining parameter  $(d_{M_0S_2}/d_{\text{quartz}})$  of 0.0003, which is almost an order of magnitude below the theoretical limit of 0.002 where straining can occur. This suggests that an alternate retention mechanism is likely occurring that is responsible for the  $MoS<sub>2</sub>-PL$  retention during transport in the porous media. The BTC observed for the  $MoS_2-PL$  at 10 mM KCl has been observed before in other transport studies and suggests that either blocking or straining is occurring. Blocking has been used in traditional colloid literature (Song and Elimelech, 1993; Liu *et al.*, 1995) to explain the deposition of colloids during transport in porous media. Blocking occurs when the particles entering the column undergo an initial rapid deposition on available collector sites, followed by a decrease in deposition rate as more deposition sites on the collectors become unavailable (Song and Elimelech, 1993; Liu *et al.*, 1995). Since the straining parameter is well below the theoretical limit, this suggests that blocking is occurring during  $MoS<sub>2</sub>-PL$  transport in porous media. As observed by the BTC shape in Figure 5, at 10 mM KCl the  $MoS<sub>2</sub>-PL$  undergo some massive initial deposition on the empty attachment sites located on the quartz collectors as they enter the column in the first centimeter, and then after these sites are blocked, a decrease in deposition occurs resulting in more  $MoS<sub>2</sub>$  being eluted from the column as a function of time. For the  $MoS<sub>2</sub>-Li$ transport experiments, the blocking shape was not observed from the BTCs suggesting that straining may be the dominant removal mechanism during transport. This is in agreement with the straining parameter  $(d_{MoS2}/d_{particle} = 0.006)$  calculated for  $MoS_2$ -Li at 31.6 mM KCl. From the transport experiments it is clear that each type of  $MoS<sub>2</sub>$  undergo unique interactions with the quartz collectors during transport with the  $MoS_2-PL$  being the least mobile as compared with  $MoS_2-$ Li under similar conditions.

## Remobilization of  $MoS<sub>2</sub>$

To further understand the possibility of remobilization or irreversible attachment, and to identify mechanisms by which particles are being retained (as observed from the transport experiment results described in the packed bed column transport study section) for  $MoS_2$ -Li and  $MoS_2$ -PL, DI water was injected for  $\sim$  12 PV following a column experiment at each IS. This results in the IS strength being reduced in the column and helps to shed light on the possible mechanisms responsible for attachment, deposition, and remobilization during transport. The results from this release portion of the experiments are shown in Figures 5B and 6B for  $MoS<sub>2</sub>-PL$ and  $MoS<sub>2</sub>-Li$ , respectively, and have been included in Table 1. The amount of mass remobilized (collected in the effluent) after the injection of DI water was divided by the amount that

TABLE 1. MASS BALANCE RESULTS FROM COLUMN EXPERIMENTS ARE ALSO SHOWN, WHERE  $M_E$ ,  $M_D$ ,  $M_R$ , and  $M_T$ , Represent the Relative Mass Percent of the Mass of the MoS<sub>2</sub> in the Effluent, Deposited Mass Recovered from Column Dissections, Remobilized Mass After DI Injection, and the Total Mass Collected

MoS <sub>2</sub>	<i>Ionic Strength (mM)</i>	$M_F~({\%})^{\rm a}$	$M_D$ $(\%)^{\rm b}$	$M_R$ $(\%)^c$	$M_T$ (%) <sup>d</sup>
$MoS_2-PL$	10 31.6	$96.7 \pm 1.0$ $13.2 \pm 0.2$ $3.1 \pm 0.2$	19.2 26.3 21.7	18.8 18.0 18.8	115.8 39.4 24.8
$MoS2-Li$	10 31.6	$84.0 \pm 7.6$ $81.5 \pm 6.5$ $11.3 \pm 2.5$	28.2 27.8 55.9	3.9 4.8 27.4	112.1 109.2 67.2

Transport experiments were conducted using a  $MoS<sub>2</sub>$  concentration of 10 mg/L and a flow rate of 2 mL/min.

 ${}^{\text{a}}_1\text{M}_E$ :% of mass of MoS<sub>2</sub> in effluent of mass injected.

 ${}^{\text{b}}\text{M}_{\text{D}}$ :% of mass of MoS<sub>2</sub> deposited in the column and collected from the column dissection.

 ${}^{\text{c}}\text{M}_{\text{R}}$ :% of MoS<sub>2</sub> remobilized after dispersed in (DI) injection of average retained mass from triplicate experiments.<br><sup>d</sup>M<sub>T</sub>:% of total mass (M<sub>E</sub>+M<sub>D</sub>) of MoS<sub>2</sub> in effluent (M<sub>E</sub>) plus mass collected from



FIG. 7. Retention profiles created from the column dissections after a transport experiment in the packed bed column at 1, 10, and 31.6 mM KCl. (A)  $MoS<sub>2</sub>-Li$  retention profile. (B)  $MoS_2-PL$  retention profile.

was retained in the column and multiplied by 100 to obtain the mass% remobilized  $(M_R)$ . For the MoS<sub>2</sub>-Li particles, the  $M_R$  increased from 4% to 27% across the IS range tested suggesting electrostatic interactions are occurring with the quartz collectors during transport and responsible for retention in the column. For the  $MoS_2-PL$ , however, the  $M_R$  remained constant  $(18.5\% \pm 0.5\%)$  across the IS range tested suggesting a maximum remobilization regardless of IS. This difference in remobilization phenomena between the two types of  $MoS<sub>2</sub> suggests that the PF-87 pluronic surfactant on$ the  $MoS<sub>2</sub>-PL$  is interacting with the quartz collectors, resulting in a strong bond and a limited  $M_R$ . The interactions between the PEO functional groups present on the  $MoS<sub>2</sub>-PL$ surface and the Si function groups on the quartz collector are likely forming strong intermolecular bonds and resulting in irreversible attachment. This irreversible attachment was not observed in other two-dimensional planar particles (i.e., graphene oxide) that exhibited  $\sim$  100% remobilization once DI water was injected during the release portion of a column experiment (Feriancikova and Xu, 2012; Lanphere *et al.*, 2013). The release experiments highlight the different behaviors for  $MoS<sub>2</sub>$  during transport in sand columns; with  $MoS<sub>2</sub>-PL$  being the least sensitive to change in IS and less mobile compared with  $MoS_2-Li$  when the remobilization  $(M_R)$  parameter was considered at similar conditions.

# Spatial distribution of  $MoS<sub>2</sub>$  in the column

To provide additional insight into the unique differences observed in the transport and release experimental results between the two types of  $MoS_2$ ; the distribution of  $MoS_2$ deposited at each centimeter during transport was quantified using previously described methods (Lanphere *et al.*, 2013) and is reported in Figure 7. Overall, the majority of  $MoS<sub>2</sub>-Li$ was equally distributed throughout each centimeter of the column for at the lower IS (1–10 M KCl) tested. However, once 31.6 mM KCl was achieved, a sharp increase (approximately seven times greater compared with 1 mM KCl condition) in  $MoS<sub>2</sub>-Li$  mass deposited in the top centimeter of the column was observed as seen in Figure 7B. This increase in deposition in the top centimeter of the column correlates well with the effective diameter ( $\sim$  1512 nm at 31.6 mM KCl) characterized in the DLS experiments, which results in a straining parameter  $(d_{M_0S2}/d_{quartz})$  of 0.006, suggesting that straining is the removal mechanism occurring at this condition. However, for the  $MoS<sub>2</sub>$ -PL there was no significant difference  $(p > 0.05)$  in spatial distribution as a function of IS. The retention profiles were very similar across the range of IS tested (1–31.6 mM KCl). The equal spatial distribution of  $MoS<sub>2</sub>-PL$  in the retention profiles suggest that they are being irreversibly attached to the quartz collectors and is in agreement to the trend observed during the release portion of the column experiments.

# DLVO theory

DLVO theory was used to identify whether the trends observed during the column transport experiments were due to  $MoS_2$ -quartz interactions (blocking) or  $MoS_2-MoS_2$  interactions (straining). Calculations were performed using traditional DLVO theory (Verwey, 1947; Derjaguin and Landau, 1993) to simulate  $MoS<sub>2</sub>$ -quartz interactions, assuming a constant potential and sphere-plate geometry (Hogg *et al.*, 1966) and can be seen in Tables 2 and 3. A Hamaker constant of  $6.26 \times 10^{-21}$  J was used for the MoS<sub>2</sub>-water– quartz systems (Feriancikova and Xu, 2012), and the EPM measurements for the  $MoS<sub>2</sub>$  nanomaterials were converted to zeta potential values using the Smoluchowski equation (Elimelech, 1995). The  $MoS<sub>2</sub>$ -Li interaction energy profile at 31.6 mM KCl suggest that a secondary minimum  $(-5.7 \text{ kT})$ exists at 10 nm. This secondary minimum at 31.6 mM KCl may help explain why there was twice as much deposition in the column compared with the 1 mM KCl condition. For the  $MoS<sub>2</sub>-PL DLVO$  profiles, there was only a small energy barrier (1.9 kT) at 10 mM KCl and no energy barrier at 31.6 mM KCl observed. This absence of an energy barrier for  $MoS<sub>2</sub>$ -quartz surfaces at 31.6 mM KCl likely contributed to the fraction (26%) of irreversibly bound particles retained in the column following the release portion of the transport experiments. This is also in agreement with what was observed in the retention profiles where 78% of the  $MoS<sub>2</sub>-PL$  particles remained on the quartz collectors even

Table 2. Results from Calculations Based Upon DLVO and the Sphere-Plate Assumption for MoS2-Li Nanoparticles and Quartz Collectors

$IS$ ( $mM$ KCl)	of $MoS_{2}$ (mV)	of quartz (mV)	barrier (kT)	Zeta potential Zeta potential $1^{\circ}$ energy $1^{\circ}$ energy barrier separation distance (nm)	$2^{\circ}$ min depth (kT)	$2^{\circ}$ <i>min</i> separation distance (nm)	Diameter <sup>a</sup> (nm)	$EPM^b$ $(10^{-8} \; m^2)$ $[V\cdot s]$
$\mathbf{1}$	$-52.4 \pm 3.2$	$-28.6 \pm 2.0$	260.4	2.5	$-0.018$	109.5	$338.5 \pm 9.6$	$-4.1 \pm 0.3$
10	$-50.7 \pm 1.7$	$-13.6 \pm 0.5$	52.8	2.5	$-0.32$	24	$342.2 \pm 4.9$	$-4.0 \pm 0.1$
31.6	$-41.7 \pm 1.0$	$-11.3 \pm 1.4$	97.6	2.0	$-5.7$	10	$1512.5 \pm 210.8$	$-3.3 \pm 0.1$

<sup>a</sup>Average hydrodynamic diameter values for MoS<sub>2</sub>-Li were experimentally recorded and were used in the Derjaguin-Landau-Verwey-Overbeek (DLVO) calculations. <sup>b</sup>

Electrophoretic mobility (EPM) values for MoS<sub>2</sub>-Li were used to calculate zeta potentials through the Smoluchowski equation (Elimelech, 1995).

IS (mM KCl)	of $MoS_{2}$ (mV)	of quartz (mV)	barrier (kT)	Zeta potential Zeta potential $1^{\circ}$ energy $1^{\circ}$ energy barrier $2^{\circ}$ min separation distance (nm)	depth (kT)	$2^{\circ}$ min separation distance (nm)	Diameter* (nm)	$EPM**$ $(10^{-8} \; m^2)$ $V^{-1}$ $s^{-1}$ )
$10^{-3}$	$-21.0 \pm 1.8$	$-28.6 \pm 2.0$	24.1	2.0	$-0.005$	95.5		$75.4 \pm 1.0$ - $1.6 \pm 0.1$
$10^{-2}$	$-11.4 \pm 1.0$	$-13.6 \pm 0.5$	1.9	2.5	$-0.13$	16		$83.4 \pm 2.8$ - 0.9 $\pm$ 0.1
$\frac{10}{10}$ – 1.5	$-5.3 \pm 1.1$	$-11.3 \pm 1.4$	None	2.0	None	N/A	$79.7 \pm 2.4$ - 0.4 $\pm$ 0.1	
			$(-2.6)$					

Table 3. Results from Calculations Based Upon DLVO and the Sphere-Plate ASSUMPTION FOR MOS<sub>2</sub>-PL NANOPARTICLES AND QUARTZ COLLECTORS

\*Average hydrodynamic diameter values were experimentally recorded and were used in the DLVO calculations.

\*\*EPM values for  $MoS<sub>2</sub>$ -Li were used to calculate zeta potentials via the Smoluchowski equation (Elimelech 1995).

after being chemically (reduction in IS) and physically (vortexed) perturbed.

To further explain the transport results observed in this study,  $MoS<sub>2</sub>-MoS<sub>2</sub>$  interactions were considered and additional DLVO interaction energy profiles were created. A secondary energy minimum  $(-1.8 \text{ kT at } 15 \text{ nm})$  was observed for the  $MoS_2$ -Li at 31.6 mM KCl, which correlates with the straining parameter  $(0.006)$  that was calculated and the seven-fold increase (compared to the 1 mM KCl condition) in mass deposited in the first centimeter as observed in the retention profiles (Fig. 7A). A similar trend was observed in previous reports for another two-dimensional nanomaterial (graphene oxide) (Feriancikova and Xu, 2012), and it was reported that the existence of a secondary minimum contributed to the removal of graphene oxide particles during transport through reversible chemical (electrostatic) interactions. For the  $MoS_2-PL$  a secondary minimum was not observed when  $MoS<sub>2</sub>MoS<sub>2</sub>$  interactions were considered.

Overall, DLVO interaction energy profiles helped provide additional information regarding the individual removal mechanisms for each type of  $MoS<sub>2</sub>$  during transport in sand columns. Specifically, the application of DLVO supports the claim that straining is the dominant mechanism for retention in porous media occurring at higher IS for  $MoS<sub>2</sub>-Li$  due to secondary energy minimums present, which results in aggregation of the  $MoS_2-MoS_2$  particles. It is also likely that the smaller  $MoS<sub>2</sub>-PL$  nanomaterials are experiencing deeper primary minimum interactions compared to the  $MoS<sub>2</sub>-Li$ since they are smaller in size and is in agreement with the trends observed in previous studies (Bradford and Torkzaban, 2012, 2013). This difference in primary interactions between  $MoS<sub>2</sub>-PL$  and  $MoS<sub>2</sub>-Li$  may help explain why greater amounts of  $MoS_2-PL$  (86.8%  $\pm$  0.2%) were retained in the column at similar IS (10 mM KCl) compared with  $MoS<sub>2</sub>-Li$  $(18.5\% \pm 6.5\%)$  even though the particles were smaller. Furthermore, the application of DLVO helps support the mechanistic explanation that  $MoS<sub>2</sub>-PL$  likely undergoes irreversible attachment to the quartz collectors due to the absence of (or minimal) energy barriers present during interactions with quartz collectors. These irreversible interactions contributed to the decreasing rate of  $MoS<sub>2</sub>-PL$  retention over time as observed by the blocking phenomena seen in the BTC shape during transport experiments. This is in agreement with the trends observed in the limited remobilization  $(M_R = 18.6\% \pm 0.5\%)$  during release experiments and the equal spatial distribution of  $MoS<sub>2</sub>-PL$ throughout the column in the retention profiles.

## SEM images of quartz sand and  $MoS<sub>2</sub>$

SEM images were taken to better understand the role of physical interactions between the  $MoS<sub>2</sub>$  and sand grains as a

FIG. 8. Scanning Electron Microscopy (SEM) images taken at different magnifications for quartz collectors and  $MoS_2$ -Li and  $MoS_2$ -PL particles used in the packed bed column experiments. (A, D) Quartz collector at 10 and 5  $\mu$ m. (**B**, **E**) MoS<sub>2</sub>-Li at 10 and 1  $\mu$ m. (**C**, **F**)  $MoS_2-PL$  at 10 and 1  $\mu$ ms.



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function of their different surface morphologies and to identify any potential physical removal mechanisms during transport. The SEM images shown in Figure 8 suggest that the physical geometry (two-dimensional planar flakes) of the  $MoS<sub>2</sub>$  likely plays an important role during transport through porous media. It is likely that the planar  $MoS<sub>2</sub>$  particles get wedged or removed during transport since their shape and size allow them to interact with the rough surface of the quartz sand. As seen in Figure 8A and D, the surface of the quartz collectors is very heterogeneous with craters and pits ranging from 1 to 2  $\mu$ m in diameter. In addition, the MoS<sub>2</sub>-Li (Fig. 8B, E) and  $MoS_2-PL$  (Fig. 8C, F) SEM images taken highlight the unique two dimension planar sheet-like structure of the  $MoS<sub>2</sub>$  particles. As a result of this geometry, van der Waals forces would have a greater influence since the large surface area of the  $MoS<sub>2</sub>$  is able to interact with the surface of the quartz surfaces. This is consistent with the DLVO calculations that revealed an absence in energy barrier for  $MoS<sub>2</sub>-PL$ and quartz surfaces, and may explain why there was greater  $MoS<sub>2</sub>-PL$  retention in the column during transport compared with  $MoS_2$ -Li even though the size (as measured in DLS experiments) was  $\sim$  19 times smaller at 31.6 mM KCl. Furthermore, it is likely that both types of  $MoS<sub>2</sub>$  are experiencing some wedging or straining as they transport through the column and interact with the rough surfaces on the quartz collectors. This assumption is consistent with others who have reported that wedging and straining are the mechanistic basis for colloid retention during transport in porous media (Bradford *et al.*, 2002, 2005; Johnson *et al.*, 2007).

### Summary and Conclusions

This study highlights the different fate and transport mechanisms for two common types of  $MoS<sub>2</sub>$  in aquatic environments. A comprehensive study was performed to identify the unique behaviors of  $MoS_2$ -Li and  $MoS_2$ -PL in aqueous environments to better understand which type of  $MoS<sub>2</sub>$  would be least mobile and likely be exposed to various aquatic species (e.g., benthic feeders) in sediment beds. Traditional methods used to synthesize  $MoS<sub>2</sub>-Li$  may result in increased transport in sediment beds and aquatic systems under environmentally relevant conditions (1–10 mM KCl). However, at higher IS ( $\geq$  31.6 mM KCl), MoS<sub>2</sub>-Li will become less stable and aggregate to sizes greater than  $1.5 \mu m$ . Under these conditions,  $MoS<sub>2</sub>-Li$  will likely be removed through straining and as a result of secondary minimum interactions in sediment beds may potentially bio-accumulate in benthic organisms. Conversely, the  $MoS<sub>2</sub>-PL$  will tend to transport farther in aqueous systems due to its stable characteristics at environmental conditions ( $pH 5-9$ ). If  $MoS<sub>2</sub>-PL$ comes into contact with quartz media in the subsurface, it will have a higher affinity to be deposited and attach as a result of the reduced energy barrier between  $MoS_2-PL + quartz$  surfaces as observed from transport experiments and column dissections in this study. Additionally, as observed in the SEM images, the two-dimensional planar geometry of both types of  $MoS<sub>2</sub>$  nanomaterials will likely contribute to the physical removal of individual flakes during transport in porous media as a result of being wedged into heterogeneous pits present on quartz surfaces. As a result of the chemical and physical mechanisms contributing to the removal during transport in porous media, this study confirms that  $MoS<sub>2</sub>-PL$ 

will be the least mobile type of  $MoS<sub>2</sub>$  in aquatic environments compared to  $MoS<sub>2</sub>-Li$ .

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## Author Disclosure Statement

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