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# Flexible pulse-controlled fiber laser

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## SUBJECT AREAS:

FIBRE LASERS

ULTRAFAST LASERS

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Controlled flexible pulses have widespread applications in the fields of fiber telecommunication, optical sensing, metrology, and microscopy. Here, we report a compact pulse-controlled all-fiber laser by exploiting an intracavity fiber Bragg grating (FBG) system as a flexible filter. The width and wavelength of pulses can be tuned independently by vertically and horizontally translating a cantilever beam, respectively. The pulse width of the laser can be tuned flexibly and accurately from  $\sim 7$  to  $\sim 150$  ps by controlling the bandwidth of FBG. The wavelength of pulse can be tuned precisely with the range of  $>20$  nm. The flexible laser is precisely controlled and insensitive to environmental perturbations. This fiber-based laser is a simple, stable, and low-cost source for various applications where the width-tunable and/or wavelength-tunable pulses are necessary.

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Fiber-based lasers attract extensive attention due to their various advantages such as compactness, reliability, and high stability<sup>1–12</sup>. A variety of mode-locking techniques have been developed to make pulse lasers as versatile tools for many applications in fiber telecommunication, optical frequency comb generation, metrology, and microscopy<sup>13–23</sup>. So far, a number of saturable absorbers (e.g. the nonlinear polarization rotation<sup>24–30</sup>, nonlinear optical loop mirror<sup>31–33</sup>, semiconductor saturable absorber mirror<sup>34–38</sup>, graphene<sup>39–43</sup>, and carbon nanotube (CNT)<sup>44–47</sup>) have been proposed to implement the mode-locking operation. Among them, CNTs are particularly interesting for pulse generation because of highly environmental stability and being insensitive to the polarization of pulses evolving in the laser cavity<sup>45–51</sup>.

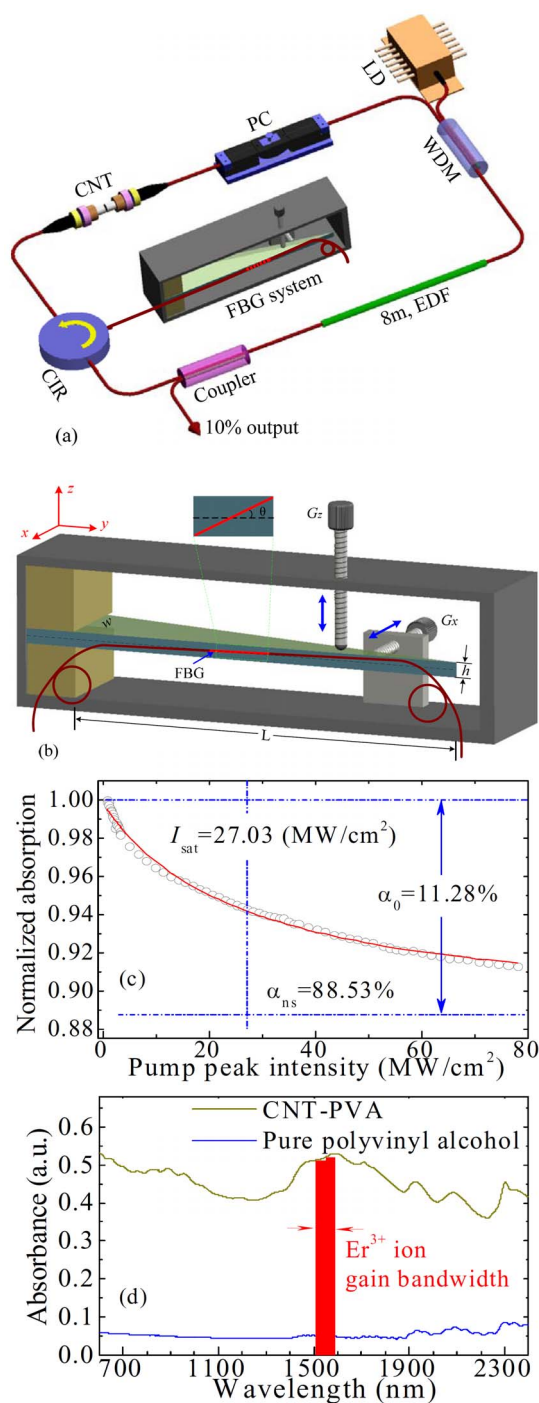
To control the laser property, a filter is widely employed into the laser cavity. For instance, a birefringent-plate-based filter can stabilize high-energy pulses in the all-normal-dispersion fiber lasers<sup>52</sup>. But its operation wavelength is fixed, as well as its spectral bandwidth is inflexible because it attributes to the thickness of birefringent plate<sup>26,52</sup>. Wavelength tuning can be realized by using an intracavity bandpass filter<sup>46,53</sup>, whereas the spectral bandwidth and the pulse width are still inflexible. Fortunately, fiber Bragg grating (FBG) is an ideal component for fiber lasers because it can provide the changeable dispersion and the tunable transmittance wavelength together with negligible nonlinearity<sup>54,55</sup>. Then, FBG offers the great flexibility for controlling the wavelength of the generated pulses. However, the pulse width of lasers is usually fixed although it can be tuned slightly by changing the pump power or adjusting the components of cavity<sup>45,46</sup>. To address this issue, we have proposed a flexible technique by means of controlling FBG.

In this article, we report a compact pulse-controlled fiber laser for the first time to our best knowledge, in which the pulse width and the pulse wavelength can be controlled precisely by adjusting FBG. The controlled scalable range of pulse width in the proposed fiber laser is accurately tunable from  $\sim 7$  to  $\sim 150$  ps. The wavelength of pulse is precisely tunable with the range of  $>20$  nm. This laser is insensitive to environmental perturbations and thus is viable for various practical applications.

## Results

**Experimental setup of controlled flexible laser.** The schematic diagram of controlled flexible fiber laser is shown in Fig. 1(a). The proposed laser consists of a FBG system, a circulator (CIR), a CNT saturable absorber (SA), a polarization controller (PC), a wavelength-division multiplexer (WDM), an 8-m-long erbium-doped fiber (EDF) with 6 dB/m absorption at 980 nm, a fused coupler with 10% output ratio, and a piece of single-mode fiber (SMF). The dispersion parameters of EDF and SMF are about 11.6 and  $-22$  ps<sup>2</sup>/km at 1550 nm, respectively. The total cavity length is  $\sim 21.5$  m.

The key component of this laser, shown in Fig. 1(b), is a FBG system that introduces the accurately width-tunable and wavelength-tunable operations for pulse generation. A uniform FBG is glued in a slanted direction onto the lateral side of a right-angled triangle cantilever beam with length  $L = 18$  cm, width  $w = 3$  cm, and thickness  $h = 1.2$  cm. The flexible cantilever beam is made of polyurethane with the high resistance against fatigue. The FBG is carefully attached by using the UV curable epoxy. The center of FBG lies on the neutral plane



**Figure 1** | (a) Laser setup. EDF, erbium-doped fiber; WDM, wavelength-division multiplexer; PC, polarization controller; LD, laser diode; CIR, circulator; FBG, fiber Bragg grating; CNT, carbon nanotube. (b) FBG system. A uniform FBG is glued in a slanted direction onto the lateral side of a right-angled triangle cantilever beam. The flexible cantilever beam is made of polyurethane. (c) Nonlinear absorption characterization of the CNT-SA. The solid curve is fitted from the experimental data (circle symbols). (d) Absorption spectra of the pure polyvinyl alcohol (PVA) and the CNT-PVA composite. The red stripe illustrates the spectral gain region of the  $\text{Er}^{3+}$ -doped fiber.

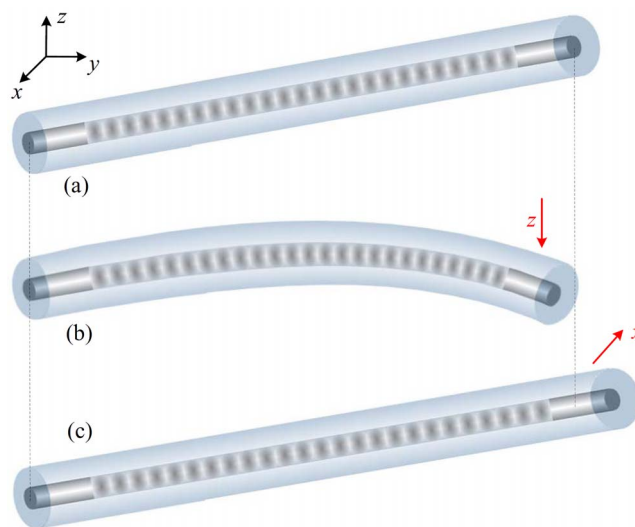
of the beam, as shown in inset of Fig. 1(b). The angle,  $\theta$ , between the axis of the FBG and the neutral layer of the beam is about  $15^\circ$ .

The integrated CNT-based SA is made by sandwiching a  $\sim 2$  mm<sup>2</sup> sample between two fiber connectors. The fabrication procedure is

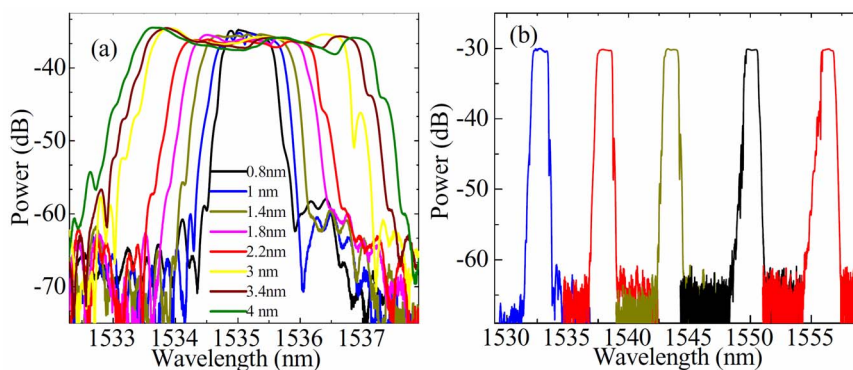
shown in our previous report<sup>45</sup>. Figure 1(c) shows the normalized nonlinear absorption of CNT-SA, which is experimentally measured with a homemade ultrafast laser at 1550 nm. The experimental data are fitted as the solid curve of Fig. 1(c) on the basis of a simplified two-level SA model<sup>45</sup>. Figure 1(c) illustrates that the linear limit of saturable absorption ( $\alpha_0$ ), the nonsaturable absorption ( $\alpha_{\text{ns}}$ ), and the saturation intensity ( $I_{\text{sat}}$ ) are about 11.28%, 88.53%, and 27.03 MW/cm<sup>2</sup>, respectively. Figure 1(d) demonstrates the absorption spectra of the pure polyvinyl alcohol (PVA) and the CNT-PVA composite, which are measured by a spectrometer (JASCO V-570 UV-vis-NIR). It is worth noting that the tube diameter of CNTs is ranging from 0.8 to 1.3 nm here, whereas it is less than 2 nm in Ref. 45.

**Bandwidth-tunable and wavelength-tunable operations.** The beam is bent when the screw  $G_z$  is translated along  $z$ -axis, as shown in Fig. 1(b). The half of the grating is under varying tension, whereas the other half is under varying compression. If the center of the grating is located exactly at the neutral layer of the beam, there will be no strain at the center of the grating<sup>57,58</sup>. In this case, the vertical distance shift at the FBG center is equal to zero and then the central wavelength shift  $\Delta\lambda_c$  is also equal to zero according to Eqs. (1) and (2) (see METHODS). Figure 2(b) shows the induction principle of tension and compression strain along the FBG based on the symmetrical bending. As a result, the bandwidth  $\Delta\lambda_B$  of FBG can be flexibly controlled by the tension and compression strain at each side of the FBG without the shift of the central wavelength. Figure 3(a) shows some typically experimental results of reflection spectra of the FBG with the variation of translation. The bandwidth of FBG is changed in the range from about 0.8 to 4 nm. Obviously, the central wavelength approximately remains fixed although the bandwidth-tunable operation is realized.

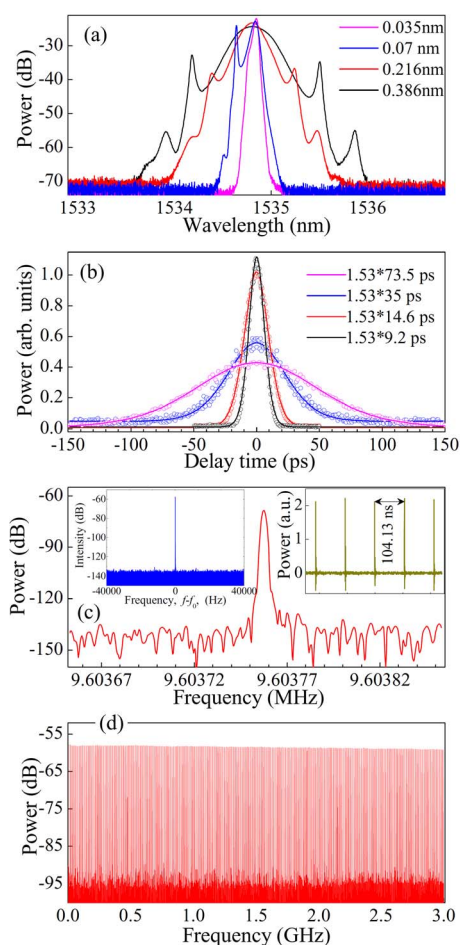
When the screw  $G_x$  in Fig. 1(b) is translated horizontally (i.e., along the direction of  $x$ -axis), the variation of FBG bandwidth  $\Delta\lambda_B$  is slight whereas the central wavelength shifts distinctly. Figure 2(c) demonstrates the induction principle of strain along the uniform FBG. The experimental results are shown in Fig. 3(b), where the central wavelength of FBG is tuned with the range of  $>20$  nm but the spectral profile changes slightly. The theoretical explanation can be achieved from Eqs. (3) and (5) (see METHODS). Obviously, the bandwidth- and wavelength-tunable operations can be realized independently by the vertical and horizontal translations, respectively.



**Figure 2** | Schematic diagram of bandwidth-tunable and wavelength-tunable operations by flexibly controlling FBG. (a) Without the translation in the free state, (b) vertically translating the screw  $G_z$  along the direction of  $z$ -axis, and (c) horizontally translating the screw  $G_x$  along the direction of  $x$ -axis.



**Figure 3** | Typically reflection spectra of FBG. (a) Vertically translating the screw  $G_z$  along  $z$ -axis. The FBG bandwidth  $\Delta\lambda_B$  is changed whereas the central wavelength approximately is fixed.  $\Delta\lambda_B$  is about 0.8, 1, 1.4, 1.8, 2.2, 3, 3.4, and 4 nm from inner to outer, respectively. (b) Horizontally translating the screw  $G_x$  along  $x$ -axis. The central wavelength of FBG is tuned with the range of  $>20$  nm, whereas the spectral profile changes slightly.

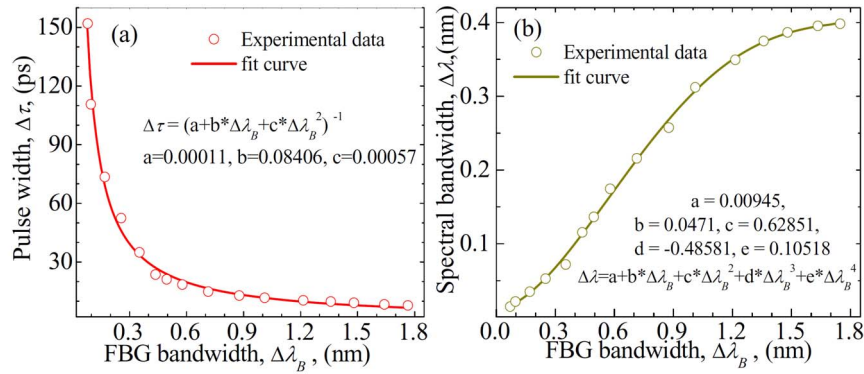


**Figure 4** | Typical laser characteristics. (a) Optical spectra at the FBG bandwidths ( $\Delta\lambda_B$ ) of 0.17, 0.35, 0.71, and 1.48 nm (from inner to outer) by vertically translating the screw  $G_z$  in Fig. 1(b). (b) Autocorrelation traces of the experimental data (circle symbols) and  $\text{sech}^2$ -shaped fit (solid curves). The FWHM spectral bandwidths ( $\Delta\lambda$ ) and the corresponding pulse widths ( $\Delta\tau$ ) are about 0.035 nm and 73.5 ps, 0.07 nm and 35 ps, 0.216 nm and 14.6 ps, and 0.386 nm and 9.2 ps, respectively. (c) Fundamental RF spectrum with the resolution of 2 Hz and the span of 200 Hz. Inset in the top left corner: RF spectrum with the resolution of 10 Hz and the span of 80 kHz. Inset in the top right corner: oscilloscope traces with the separation of  $\sim 104.13$  ns, corresponding to 9.60376 MHz of the fundamental harmonic frequency that is independent of the pump power. (d) Wideband RF spectrum up to 3 GHz.

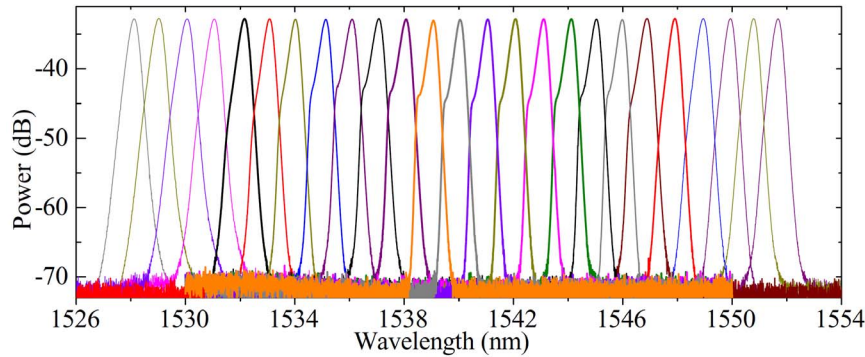
**Experimental results of controlled flexible laser.** Self-starting mode-locking operation starts at the pump power of  $\sim 19$  mW. With the appropriate setting of the polarization controller and the pump power, the proposed laser delivers the pulses with different widths by vertically adjusting the FBG (i.e., translating the screw  $G_z$  along  $z$ -axis, as shown in Fig. 1(b)). The typical output spectra are shown in Fig. 4(a) with the central wavelength of  $\sim 1535$  nm. The corresponding autocorrelation traces of the experimental data and the  $\text{sech}^2$ -shaped fit are shown in Fig. 4(b). The full width at half maximum (FWHM) spectral bandwidths are about 0.035, 0.07, 0.216, and 0.386 nm at the FBG bandwidths of 0.17, 0.35, 0.71, and 1.48 nm, respectively. The corresponding pulse widths ( $\Delta\tau$ ) are about 73.5, 35, 14.6, and 9.2 ps, respectively. Then, the calculated time-bandwidth products (TBPs) are about 0.33, 0.31, 0.40, and 0.45, respectively, which are close to the value of the transform-limited  $\text{sech}^2$ -shaped pulses. The fluctuation of TBPs originates from the variations of the FBG bandwidth, the optical spectrum, and the total dispersion of laser cavity. The radio frequency (RF) spectra in Fig. 4(c) give a signal-to-noise ratio of  $>60$  dB ( $>10^6$  contrast), showing low-amplitude fluctuations and good mode-locking stability<sup>59</sup>. No spectrum modulation is observed over 3 GHz in Fig. 4(d), indicating no Q-switching instability.

Figures 5(a) and 5(b) show the relationships of the pulse width  $\Delta\tau$  and the spectral bandwidth  $\Delta\lambda$  with respect to the FBG bandwidth  $\Delta\lambda_B$ , respectively. The corresponding TBPs are ranging from  $\sim 0.31$  to  $\sim 0.47$  for the experimental data. The symbols and solid curves denote the typically experimental data and the fit curves, respectively. In the experiments, the maximum of pulse width  $\Delta\tau$  can be up to  $\sim 152$  ps and the minimum of  $\Delta\tau$  can be down to  $\sim 7$  ps. Note that the tuning range of  $\Delta\tau$  is limited by the FBGs. We can observe from Fig. 5(a) that the pulse width  $\Delta\tau$  approximately decreases with a rational function according to the FBG bandwidth  $\Delta\lambda_B$ . The rational equation that produces the best fit is as  $\Delta\tau = (0.00011 + 0.08406 \cdot \Delta\lambda_B + 0.00057 \cdot \Delta\lambda_B^2)^{-1}$ . It is seen from Fig. 5(b) that the spectral bandwidth  $\Delta\lambda$  approximately increases with a polynomial function of degree 5 with respect to  $\Delta\lambda_B$ . The polynomial equation that produces the best fit is as  $\Delta\lambda = 0.00945 + 0.0471 \cdot \Delta\lambda_B + 0.62851 \cdot \Delta\lambda_B^2 - 0.48581 \cdot \Delta\lambda_B^3 + 0.10518 \cdot \Delta\lambda_B^4$ . Figure 5(b) illustrates that the spectral bandwidth  $\Delta\lambda$  of pulses approximately linearly increases along with  $\Delta\lambda_B$  for  $\Delta\lambda_B < 1$  nm, whereas it changes slightly for  $\Delta\lambda_B > 1.5$  nm. The origin of such behavior can be explained as follows. The different parts of FBG reflect the different wavelengths so that the round-trip distances for different frequencies of a pulse are different. When the spectral bandwidth of FBG is large enough, only a part of FBG spectra is employed in the mode-locking operation of fiber lasers and, then, the FBG with larger reflection bandwidth has a slight influence on the laser bandwidth.





**Figure 5** | (a) Pulse width  $\Delta\tau$  and (b) spectral bandwidth  $\Delta\lambda$  with respect to the FBG bandwidth  $\Delta\lambda_B$ . The beam is translated vertically, i.e., translating the screw  $G_z$  along the direction of  $z$ -axis.



**Figure 6** | Output spectra of laser by horizontally translating the screw  $G_x$  along the direction of  $x$ -axis.

Figure 6 demonstrates the typical output spectra by horizontally translating the screw  $G_x$  along the direction of  $x$ -axis. The experimental results show that the spectral bandwidth  $\Delta\lambda$  and pulse width  $\Delta\tau$  change slightly although the central wavelength of pulses is tuned evidently, indicating the stability of our output pulses. It is seen from Fig. 6 that the tuning range of wavelength is  $>20$  nm with the FWHM spectral bandwidth of  $\sim 0.3$  nm. The fluctuation of spectral profile originates from the dispersion variation of FBG when the central wavelength of FBG is tuned. The tuning range of wavelength is determined by the FBG. The experimental observations show that our laser can long-termed stably work for both width-tunable and wavelength-tunable operations. It attributes to the intrinsic merit of the CNT-based SA, which is insensitive to the environmental perturbations and the polarization of pulses.

## Discussion

In the experiments, the bandwidth of FBGs limits the tuning range of pulse width, as shown in Fig. 5(a). Theoretically, the proposed laser can deliver the pulses with the width of  $<1$  ps if FBG in Fig. 1 is optimized. It is seen from Fig. 5(a) that the pulse width  $\Delta\tau$  of laser can be less than 1 ps when the FBG bandwidth  $\Delta\lambda_B$  is more than 11 nm. Although it is hard to fabricate a uniform FBG with the bandwidth of  $>2$  nm, the bandwidth of the chirped FBGs can be up to 30 nm easily<sup>54</sup>. We can observe from Fig. 3(a) that the bandwidth  $\Delta\lambda_B$  of the uniform FBG can be extended up to 4 nm by vertically translating the screw  $G_z$ , as shown in Fig. 1(b). However, the reflectivity of FBG decreases along with the increase of  $\Delta\lambda_B$ . For instance, the reflectivity of FBG is less than 30% for  $\Delta\lambda_B \approx 2.5$  nm. As a result, this laser will not operate when  $\Delta\lambda_B$  is extended to be more than 2.5 nm.

## Methods

**Principle of controlled operation.** When the beam is bent or elongated (Fig. 2), linearly varying strain along its thickness or length is achieved. The strain response of

FBG originates from both the physical elongation of the grating and the change in the refractive index due to photoelastic effects. The central wavelength variation of FBG,  $\Delta\lambda_c$ , induced by the strain  $\varepsilon$  along its axial direction is given by<sup>56,57</sup>

$$\Delta\lambda_c = \lambda_0(1 - P_e)\varepsilon. \quad (1)$$

Here  $\lambda_0$  is the initial Bragg wavelength of the grating.  $P_e$  is the effective photoelastic coefficient (approximately equal to 0.22), which is relative to the fiber Poisson ratio and the effective refractive index of the fiber core.  $1 - P_e$  is the strain tuning coefficient. By vertically translating the screw  $G_z$  along  $z$ -axis in Fig. 1(b), the strain introduced to the beam during bending is transferred to the FBG and induces an axial strain gradient along the grating, i.e.,

$$\varepsilon = C\kappa\Delta z \cos(\theta), \quad (2)$$

where  $\kappa$  is the curvature of the neutral layer of beam,  $\Delta z$  is the vertical distance measured from the neutral layer, and  $C$  ( $0 < C < 1$ ) is a constant introduced to describe the efficiency of strain transfer from the beam to the grating. The variation of Bragg wavelength is proportional to the local axial strain along the grating and the chirp of grating can be achieved when the beam is bent. The strain can be approximately proportional to  $\kappa$  and the grating length<sup>57,58</sup>. Then, the reflection bandwidth variation of FBG can be expressed as<sup>57</sup>

$$\Delta\lambda_{BV} = C\lambda_0\kappa l(1 - P_e) \sin(2\theta)/2, \quad (3)$$

where  $l$  is the length of the grating.

By horizontally translating the screw  $G_x$  along  $x$ -axis in Fig. 1(b), the strain introduced to the grating can be approximated by

$$\varepsilon = C \frac{wf}{L^2} \cos(\theta), \quad (4)$$

where  $f$  is the horizontal displacement at the end of beam. Substituting Eq. (4) into Eq. (1), we can achieve the central wavelength change of FBG when the translation is along the direction of  $x$ -axis, i.e.,

$$\Delta\lambda_c = C\lambda_0(1 - P_e) \frac{wf}{L^2} \cos(\theta). \quad (5)$$

In this case, the reflection bandwidth variation of FBG can be approximated to zero, i.e.,  $\Delta\lambda_{BV} \approx 0$ . In fact, when the screw  $G_x$  in Fig. 1(b) is translated horizontally, the angle between the axis of FBG and the vertical plane of translation is zero, i.e.,  $\theta = 0$ . Thus,  $\Delta\lambda_{BV}$  is equal to zero according to Eq. (3).



**Measurement method.** An optical spectrum analyzer (Yokogawa AQ-6370), an autocorrelator (APE PulseCheck SM1600), a 6-GHz oscilloscope (LeCroy WaveMaster 8600A), a radio-frequency (RF) analyzer (Agilent E4447A), and a 10-GHz photodetector are used to measure the laser output performances.

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## Author contributions

X.L. proposed the laser system and wrote the manuscript text. Y.C. performed the main experimental results. All authors discussed the results and substantially contributed to the manuscript.

## Additional information

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