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Author manuscript *Neural Comput.* Author manuscript; available in PMC 2016 April 01.

Published in final edited form as:

Neural Comput. 2015 April; 27(4): 898–924. doi:10.1162/NECO\_a\_00712.

## Neuronal calcium wave propagation varies with changes in endoplasmic reticulum parameters: a computer model

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## Abstract

Calcium ( $Ca^{2+}$ ) waves provide a complement to neuronal electrical signaling, forming a key part of a neuron's second messenger system. We developed a reaction-diffusion model of an apical dendrite with diffusible inositol triphosphate (IP<sub>3</sub>), diffusible Ca<sup>2+</sup>, IP<sub>3</sub> receptors (IP<sub>3</sub>Rs), endoplasmic reticulum (ER) Ca<sup>2+</sup> leak, and ER pump (SERCA) on ER. Ca<sup>2+</sup> is released from ER stores via IP<sub>3</sub>Rs upon binding of IP<sub>3</sub> and Ca<sup>2+</sup>. This results in Ca<sup>2+</sup>-induced-Ca<sup>2+</sup>-release (CICR) and increases Ca<sup>2+</sup> spread. At least two modes of Ca<sup>2+</sup> wave spread have been suggested: a continuous mode based on presumed relative homogeneity of ER within the cell; and a pseudosaltatory model where Ca<sup>2+</sup> regeneration occurs at discrete points with diffusion between them. We compared the effects of three patterns of hypothesized IP<sub>3</sub>R distribution: 1. continuous homogeneous ER, 2. hotspots with increased IP<sub>3</sub>R density (IP<sub>3</sub>R hotspots), 3. areas of increased ER density (ER *stacks*). All three modes produced  $Ca^{2+}$  waves with velocities similar to those measured in vitro (~50–90um /sec). Continuous ER showed high sensitivity to IP<sub>3</sub>R density increases, with time to onset reduced and speed increased. Increases in SERCA density resulted in opposite effects. The measures were sensitive to changes in density and spacing of IP<sub>3</sub>R hotspots and stacks. Increasing the apparent diffusion coefficient of Ca<sup>2+</sup> substantially increased wave speed. An extended electrochemical model, including voltage gated calcium channels and AMPA

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synapses, demonstrated that membrane priming via AMPA stimulation enhances subsequent  $Ca^{2+}$  wave amplitude and duration. Our modeling suggests that pharmacological targeting of IP<sub>3</sub>Rs and SERCA could allow modulation of  $Ca^{2+}$  wave propagation in diseases where  $Ca^{2+}$  dysregulation has been implicated.

## Keywords

Calcium waves; computer model; SERCA; IP<sub>3</sub> receptor; reaction-diffusion; calcium; electrochemical

## **1** Introduction

Calcium (Ca<sup>2+</sup>) is an important second messenger signal in many cell types, with diverse roles, from fertilization (Busa and Nuccitelli, 1985; Kretsinger, 1980) to regulating gene expression (West et al., 2001). Ca<sup>2+</sup> is involved in triggering destructive processes including apoptosis (Orrenius et al., 2003) and ischemia (Lipton, 1999). Cells, including neurons, therefore regulate cytosolic Ca<sup>2+</sup> concentration via buffers (Stern, 1992) and sequestration into mitochondria (Gunter et al., 2004) or endoplasmic reticulum (ER) (Berridge, 1998). In neurons, sequestration is modulated by neuronal activity (Pozzo-Miller et al., 1997), and elevated Ca<sup>2+</sup> opens certain ion channels. There is therefore a bidirectional interaction between chemical signaling and electrophysiology (Blackwell, 2013; De Schutter and Smolen, 1998; De Schutter, 2008).

 $Ca^{2+}$  is heavily buffered but travels long distances (>100 µm) to reach its targets. This poses a temporal problem if relying purely on diffusion.  $Ca^{2+}$  waves increase the rapidity of  $Ca^{2+}$ spread via  $Ca^{2+}$ -induced- $Ca^{2+}$ -release (CICR). CICR occurs in neurons and requires stores of  $Ca^{2+}$  held in the ER. ER is distributed throughout the cytosol in a connected way, through the dendrites and dendritic spines (Harris, 1994; Spacek and Harris, 1997). The ER SERCA pump pushes  $Ca^{2+}$  from the cytosol into the ER. When triggered by IP<sub>3</sub> and  $Ca^{2+}$ , the ER IP<sub>3</sub> receptors (IP<sub>3</sub>Rs) open and release some of the ER's  $Ca^{2+}$  into the cytosol. Regions of elevated  $Ca^{2+}$  then spread throughout portions of the dendritic tree (Ross et al., 2005).

Understanding  $Ca^{2+}$  wave modulation is difficult since cytosolic  $Ca^{2+}$  changes over time and passes between different intracellular compartments and extracellular space (Fall et al., 2004; Wagner et al., 2004). Adding to this complexity is the nonuniform distribution of IP<sub>3</sub>Rs, with clusters forming where local variations in IP<sub>3</sub>R or ER is heightened (Fitzpatrick et al., 2009).

At least two modes of  $Ca^{2+}$  wave spread have been identified: a continuous model that depends on continuous underlying substrate of regenerative potential, and a pseudosaltatory model where  $Ca^{2+}$  regeneration occurs at discrete points with diffusion between them. (We call it pseudo-saltatory here to distinguish from classical saltatory conduction in myelinated fibers, involving capacitative effects in addition to electrodiffusion.) It has been hypothesized that these two modes produce downstream functional differences in dendrites, where many mechanisms are responsive to the level of  $Ca^{2+}$ , for example  $I_h$  (Winograd et

al., 2008; Neymotin et al., 2013, 2014) and synaptic plasticity (Kotaleski and Blackwell, 2010).

To investigate  $Ca^{2+}$  waves in a spatiotemporal context relevant for neurons, we developed a model of  $Ca^{2+}$  waves which includes cytosol and ER. Baseline cytosolic  $Ca^{2+}$  and  $IP_3$  concentration are set to low values. The model includes ER SERCA pumps (pump  $Ca^{2+}$  from cytosol into ER), leak channels ( $Ca^{2+}$  leaks out of ER), and  $IP_3Rs$ . Our model generates  $Ca^{2+}$  waves with realistic physiological properties. We use our model to investigate how  $IP_3R$  density and clustering, alterations in SERCA density, alterations in ER stacking, and  $Ca^{2+}$  diffusibility alter waves. An additional complexity arises when we consider coupling to plasma membrane calcium channels which will contribute additional calcium flux triggered by rapidly-spreading regenerative voltage changes on the membrane. We have therefore also assessed our results in an extended electrochemical model, which included a variety of ion channels (leak, voltage-dependent calcium channels: VGCC, potassium, sodium) with synaptic activation. We used this model to demonstrate how priming due to AMPA-mediated membrane depolarization would enhance subsequent  $Ca^{2+}$  wave amplitude and duration.

## 2 Materials and methods

All simulations were run in the NEURON (version 7.3) simulation environment (Carnevale and Hines, 2006). NEURON has traditionally supported electrical modeling but has recently been extended to support reaction-diffusion (R×D) modeling as well (McDougal et al., 2013a,b). The full 1D R×D calcium wave model of the neuronal dendrite (depicted in Fig. 1) and the data analysis source-code is available on ModelDB (Peterson et al., 1996; Hines et al., 2004) (http://senselab.med.yale.edu/modeldb).

Our Ca<sup>2+</sup> dynamics are derived from Wagner et al. (2004), a spatial variant of Li and Rinzel (1994). Parameters are as in Table 1. We modeled a one-dimensional R×D system of intracellular neuronal Ca<sup>2+</sup> waves in an unbranched apical dendrite of a hippocampal pyramidal neuron (length of 1000 µm and diameter of 1 µm). Within the dendrite, we modeled cytosolic and endoplasmic reticulum (ER) compartments by using a fractional volume for each: suppose that for a given cell volume,  $f_{ER}$  denotes the fraction occupied by the ER (0.17), and  $f_{cyt}$  denotes the fraction occupied by the cytosol (0.83). Necessarily  $f_{cyt} + f_{ER}$  1. The inequality is strict if other structures are present, such as mitochondria.

ER-based calcium dynamics were previously modeled (De Young and Keizer, 1992). This previous work showed by a time-scale analysis that the qualitative dynamics of IP<sub>3</sub> receptors (IP<sub>3</sub>Rs) could be represented by only considering the slow Ca<sup>2+</sup> inactivation binding site state (Li and Rinzel, 1994). This work formed the basis of much subsequent work in intracellular Ca<sup>2+</sup> dynamics (Fall and Rinzel, 2006; Fall et al., 2004; Hartsfield, 2005; Peercy, 2008; Wagner et al., 2004). The model presented here is a variant that neglects dynamic IP<sub>3</sub> production (Wagner et al., 2004).

The ER  $Ca^{2+}$  model involves IP<sub>3</sub>Rs and SERCA pumps. As with the plasma membrane model, we combine the net effects of all other channels on the ER into a leak channel. We

denote by  $J_{IP3R}$ ,  $J_{SERCA}$ , and  $J_{leakER}$  the mass flux per unit volume due to the IP<sub>3</sub>R, SERCA pump, and leak channels, respectively. Dividing the mass flux by the volume fraction gives the change in concentration.

### State Variables

Cytosolic Ca<sup>2+</sup> concentration, ER Ca<sup>2+</sup> concentration, IP<sub>3</sub> concentration, and IP<sub>3</sub>R gating are denoted by  $Ca_{cyt}^{2+}$ ,  $Ca_{ER}^{2+}$ , IP<sub>3</sub>, and  $h_{IP_3R}$ , respectively.

$$\frac{\partial \mathrm{Ca}_{\mathrm{cyt}}^{2+}}{\partial t} = d_{\mathrm{Ca}_{\mathrm{cyt}}^{2+}} \cdot \Delta \mathrm{Ca}_{\mathrm{cyt}}^{2+} + \frac{J_{\mathrm{IP3R}} - J_{\mathrm{SERCA}} + J_{\mathrm{leakER}}}{f_{\mathrm{cyt}}} + c_{\mathrm{ionic}} \quad (1)$$

$$\frac{\partial \mathrm{Ca}_{\mathrm{ER}}^{2+}}{\partial t} = d_{\mathrm{Ca}_{\mathrm{ER}}^{2+}} \cdot \Delta \mathrm{Ca}_{\mathrm{ER}}^{2+} - \frac{J_{\mathrm{IP3R}} - J_{\mathrm{SERCA}} + J_{\mathrm{leakER}}}{f_{ER}} \quad (2)$$

$$\frac{\partial \mathrm{IP}_{3}}{\partial t}{=}d_{\mathrm{IP}_{3}}\cdot\Delta\mathrm{IP}_{3} \quad (3)$$

$$\frac{\partial h_{\mathrm{IP_{3}R}}}{\partial t} = \frac{h_{\infty\mathrm{IP_{3}R}} - h_{\mathrm{IP_{3}R}}}{\tau_{\mathrm{IP_{3}R}}} \quad (4)$$

where  $c_{\text{ionic}}$  denotes the net flux into the cytosol due to ion channels on the cell membrane and depends on both space and time. The first terms of the right hand sides is the contribution from diffusion. IP<sub>3</sub> diffusion within the cytosol is passive. Equation (1) allows for the coupling between Ca<sup>2+</sup> that enters the cytosol via the ion channels on the cell membrane and the intracellular Ca<sup>2+</sup> which is shuttled into ER via SERCA pumps.

#### Fluxes

Flux from the IP<sub>3</sub>R, SERCA pump, and leak channels are denoted by  $J_{IP3R}$ ,  $J_{SERCA}$ , and  $J_{leakER}$ , respectively.

$$\begin{split} J_{\rm IP3R} = & \overline{p}_{\rm IP_3R} \cdot m_{\rm IP_3R}^3 \cdot n_{\rm IP_3R}^3 \cdot h_{\rm IP_3R}^3 \cdot ({\rm Ca}_{\rm ER}^{2+} - {\rm Ca}_{\rm cyt}^{2+}) / \Xi \quad (5) \\ J_{\rm SERCA} = & -\frac{\overline{p}_{serca} \cdot {\rm Ca}_{\rm cyt}^{2+2}}{(k_{serca}^{2+} + {\rm Ca}_{\rm cyt}^{2+2}) \cdot \Xi} \quad (6) \\ J_{\rm leakER} = & \overline{p}_{leakER} \cdot ({\rm Ca}_{\rm ER}^{2+} - {\rm Ca}_{\rm cyt}^{2+}) / \Xi \quad (7) \end{split}$$

Here,  $\Xi = N_A/10^{18} \approx 602214.129$ , and is the number of molecules in a cubic micron at a concentration of 1 mM, where  $N_A$  is the Avogadro constant. The SERCA pump is a pump rather than a channel and so is modeled with Hill-type dynamics. The form of the fluxes  $J_{\text{IP3R}}$  and  $J_{\text{leakER}}$  parallels the forms for ion channels in the Hodgkin-Huxley equations.

Leak is ungated whereas  $J_{\text{IP3R}}$  is gated:  $m_{\text{IP3R}}$  and  $n_{\text{IP3R}}$  are fast gating variables depending on IP<sub>3</sub> and Ca<sup>2+</sup>, and  $h_{\text{IP3R}}$  is the slow Ca<sup>2+</sup> inactivation gating variable.

$$\begin{split} m_{\rm IP_{3R}} = & \frac{\rm IP_{3}}{\rm IP_{3} + k_{\rm IP_{3}}} \quad (8) \\ n_{\rm IP_{3R}} = & \frac{\rm Ca_{\rm cyt}^{2+}}{\rm Ca_{\rm cyt}^{2+} + k_{act}} \quad (9) \\ h_{\infty \rm IP_{3R}} = & \frac{k_{inh}}{k_{inh} + \rm Ca_{\rm cyt}^{2+}} \quad (10) \end{split}$$

Initial values of  $Ca_{cyt}^{2+}$ ,  $C_{AVG}$ , IP<sub>3</sub>, and  $h_{IP_{3R}}$  were set to 0.0001 mM, 0.0017 mM, 0.1 mM, and 0.8, respectively.  $Ca_{ER}^{2+}$  was then adjusted based on the cytosolic and ER fractional

volumes according to: 
$$\operatorname{Ca}_{\scriptscriptstyle{\operatorname{ER}}}^{2+} = \frac{C_{\scriptscriptstyle{\operatorname{AVG}}} - \operatorname{Ca}_{\scriptscriptstyle{\operatorname{cyt}}}^{2+} \cdot f_{\scriptscriptstyle{cyt}}}{f_{\scriptscriptstyle{ER}}}$$

This model supports both bistable and excitable waves. Both types of waves have been observed in the Xenopus oocyte (Fontanilla and Nuccitelli, 1998; Lechleiter et al., 1991), although not clearly defined in neurons.

## **Electrical dynamics**

Electrical dynamics were utilized for the subset of simulations described in the final figure. Electrical dynamics in the dendrite followed the standard parallel conductance model. Ion channels were based on a prior model of a hippocampal pyramidal cell (Safiulina et al., 2010) (http://senselab.med.yale.edu/modeldb/ShowModel.asp?model=126814) and included L-, N-, and T-type Ca<sup>2+</sup> channels, which allowed extracellular calcium to enter the dendrite at different voltage levels (McCormick and Huguenard, 1992; Kay and Wong, 1987). The equations for the ion channels follow, leak, Na<sup>+</sup> and K<sup>+</sup> currents were represented by the conductance approximation:  $I_{ion} = g_{ion} \cdot (V - E_{ion})$  (g conductance; E reversal potential) using  $E_{Na} = 50$  mV;  $E_k = -77$  mV;  $E_{leak} = -64$  mV; ( $g_{leak} = 39.4 \cdot 10^{-6}$  S/cm<sup>2</sup>), while the Goldman-Hodgkin-Katz (GHK) flux equation was used for Ca<sup>2+</sup> currents:  $I_{Ca} = p_{Ca} \cdot$  GHK<sub>Ca</sub> (p permeability). Channel dynamics were corrected for temperature by a Q<sub>10</sub> using the factor of  $qt = Q_{10}^{(T-25)/10}$  with T=37°Celsius and 25° was taken to be the temperature at which the experiment was done. Conductances and activation curves were not corrected for temperature (Iftinca et al., 2006). Voltage sensitive channels largely followed variants on the

Hodgkin-Huxley formalism, whereby  $\frac{dx}{dt} = \frac{x_{\infty} - x}{\tau_x}$  using steady-state value:  $x_{\infty} = \frac{\alpha_{\infty}}{\alpha_{\infty} + \beta_{\infty}}$ , the  $\tau_x$  forms either

$$\tau_{\text{EQN1}}: \tau_x = \frac{1}{qt \cdot (\alpha_\tau + \beta_\tau)} \text{ or } \tau_{\text{EQN2}}: \tau_x = \frac{\beta_\tau}{qt \cdot \mathbf{a}_0 \cdot (\mathbf{1} + \alpha_\tau)}$$

**L-type Ca<sup>2+</sup> channel**:  $p_L = \overline{p}_L \cdot m_L^2 \cdot h_L$  with  $p_L = 10^{-6} cm/s$ ;  $h_L$  (inactivation) was Ca<sup>2+</sup> dependent:  $h_L (\text{Ca}_{\text{cyt}}^{2+}) = \frac{k_i}{k_i + \text{Ca}_{\text{cyt}}^{2+}}$  with  $k_i = 0.001$  mM;  $m_L : \alpha_{\infty} = \frac{15.69 \cdot (-V + 81.5)}{\exp((-V + 81.5)/10.0) - 1.0}$ ;  $\beta_{\infty} = 0.29 \cdot \exp(-V/10.86)$ ;  $\tau_{\text{EQN2}}$  with  $\alpha_{\tau} = \exp(0.0378 \cdot 2 \cdot (V - 4))$ ,  $\beta_{\tau} = \exp(0.0378 \cdot 2 \cdot 0.1 \cdot (V - 4))$ ;  $a_0 = 0.1$ ;  $Q_{10} = 5$ .

**T-type Ca<sup>2+</sup> channel**:  $p_T = \overline{p}_T \cdot m_T^2 \cdot h_T$  with  $p_T = 10^{-6} cm/s$ ;

$$\begin{split} m_{T}:&\alpha_{\infty} = \frac{0.2 \cdot ((-V + 19.26))}{\exp(((-V + 19.26)/10.0) - 1.0)}; \beta_{\infty} = 0.009 \cdot \exp(-V/22.03); \tau_{\text{EQN2}} \text{ with } \alpha_{\tau} \\ = \exp(0.0378 \cdot 2 \cdot (V - (-28))); \beta_{\tau} = \exp(0.0378 \cdot 2 \cdot 0.1 \cdot (V - (-28))); h_{T}: \alpha_{\infty} = 10^{-6} \cdot 10^{$$

 $\exp(-V/16.26); \beta_{\infty} = \frac{1}{\exp((-V+29.79)/10)+1}; \tau_{EQN2} \text{ with } \alpha_{\tau} = \exp(0.0378 \cdot 3.5 \cdot (V-(-75))); \beta_{\tau} = \exp(0.0378 \cdot 3.5 \cdot 0.6 \cdot (V-(-75))); a_0 = 0.04; Q_{10} = 5.$ 

**N-type Ca<sup>2+</sup> channel**:  $p_N = \overline{p}_N \cdot m_N^2 \cdot h_N \cdot h2_N$  with  $p_N = 10^{-6} cm/s$ ;  $m_N$  used  $\alpha_{\infty} = 0.1967 \cdot (-V + 19.88)/(\exp((-V + 19.88)/10.0) - 1.0)$ ;  $\beta_{\infty} = 0.046 \cdot \exp(-V/20.73)$ ;  $\tau_{EQN2}$  with  $\alpha_{\tau} = \exp(0.0378 \cdot 2 \cdot (V - (-14)))$ ;  $\beta_{\tau} = \exp(0.0378 \cdot 2 \cdot 0.1 \cdot (V - (-14)))$ ;  $h_N$  following:  $\alpha_{\infty} = 1.6 \cdot 10^{-4} \cdot \exp(-V/48.4)$ ;  $\beta_{\infty} = 1/(\exp((-V + 39.0)/10) + 1)$ ; constant  $\tau_h = 80$ ; with  $h2_N = \{0.001\} / \{0.001 + Ca_{cyt}^{2+}\}$ ;  $a_0 = 0.03$ ,  $Q_{10} = 5$ .

The model also contained: a transient sodium channel  $I_{Na}$  and a delayed rectifier channel  $I_{K-DR}$  to allow for action potential generation; a calcium-dependent potassium channel which hyperpolarized the cell after calcium influx; and an A-type potassium channel for rapid inactivation. Equations for these channels follow.

Na channel:  $g_{Na} = \bar{g}_{Na} \cdot m_{Na}^3 \cdot h_{Na} \cdot s_{Na}$  with  $\bar{g}_{Na} = 0.11 \text{ S/cm}^2$ ;  $m_{Na}$  using  $\tau_{\text{EQN1}}$  with:  $\alpha = 0.4 \cdot (V - (-30 + 6))/(1 - \exp(-(V - (-30 + 6)/7.2))); \beta = 0.124 \cdot (-V - (30 - 6))/(1 - 6))$ 

$$\begin{split} &\exp(-(V-(30-6)))/7.2) \ h_{Na} \ \text{with special form} \ h_{\infty} = \frac{1}{1+\exp((V-(-50)-6)/4)} \ \text{and} \\ &\text{using } \tau_{\text{EQN1}} \ \text{with} \ \alpha_{\tau} = 0.03 \cdot (V-(-45+6))/(1-\exp(-(V-(-45+6))/1.5)) \ \text{and} \ \beta_{\tau} = 0.01 \cdot (V-(45-6))/(1-\exp(-(V-(45-6))/1.5)); \ Q_{10} = 2. \ s_{Na}: \ s_{\infty}(V) = 1; \ \tau_{\text{EQN2}} \ \text{with} \ \alpha_{\tau} = \exp(10^{-3} \cdot 12 \cdot (V+54) \cdot 9.648 \cdot 10^4/(8.315 \cdot (273.16+T))) \ \text{and} \ \beta_{\tau} = \exp(10^{-3} \cdot 12 \cdot 0.2 \cdot (V+54) \cdot 9.648 \cdot 10^4/(8.315 \cdot (273.16+T))) \ \text{and} \ \beta_{\tau} = \exp(10^{-3} \cdot 12 \cdot 0.2 \cdot (V+54) \cdot 9.648 \cdot 10^4/(8.315 \cdot (273.16+T))). \end{split}$$

**Delayed rectifier** *K* **channel**:  $g_{K-DR} = \bar{g}_{K-DR} \cdot n_{K-DR}$  where  $\bar{g}_{K-DR} = 0.01 \text{ S/cm}^2$ ;

 $n_{K-DR}$  following an atypical steady-state:  $n_{\infty} = \frac{1}{1+\alpha}$  with  $\alpha = \exp(10^{-3} \cdot -3 \cdot (V-19) \cdot 9.648 \cdot 10^4 / (8.315 \cdot (273.16+T)))$  using the same  $\alpha$  in  $\tau_{EQN1}$  with  $\beta = \exp(10^{-3} \cdot -3 \cdot 0.7 \cdot (V-19) \cdot 9.648 \cdot 10^4 / (8.315 \cdot (273.16+T))); a_0 = 0.02; Q_{10} = 1.$ 

**BK-type Ca<sup>2+</sup>dependent potassium channel** ( $K_{Ca}$ ):  $g_{KCa} = \bar{g}_{KCa} \cdot o_{KCa}$  with  $\bar{g}_{KCa} = 0.009 \text{ S/cm}^2$ ;  $o_{KCa}$ :

 $\alpha = \{0.28 \cdot \operatorname{Ca}_{\mathrm{cvt}}^{2+}\} / \{\operatorname{Ca}_{\mathrm{cvt}}^{2+} + k1 \cdot \exp(-2 \cdot 0.84 \cdot F \cdot V/R/(273.15+T))\};$ 

 $\beta = 0.48 / \{1 + \operatorname{Ca}_{\text{cyt}}^{2+} / k2 \cdot \exp(-2 \cdot F \cdot V / R / (273.15 + T))\} \text{ with } k1 = 0.48 \cdot 10^{-3}; k2 = 0.13 \cdot 10^{-6}; \text{ using } \tau_{\text{EON1}}.$ 

A-type potassium channel,  $K_A$ :  $g_A = \bar{g}_A \cdot n_A \cdot l_A$  where  $\bar{g}_A = 0.02 \ S/cm^2$ ;  $n_A$  with  $\tau_{EQN2}$ and an atypical steady-state value:  $n_\infty = \frac{1}{1+\alpha}$  with  $\alpha = \exp(10^{-3} \cdot \zeta (V) \cdot (V-17) \cdot 9.648 \cdot 10^4/(8.315 \cdot (273.16 + T)))$ ;  $\beta_n(V) = \exp(10^{-3} \cdot \zeta (V) \cdot 0.55 \cdot (V-17) \cdot 9.648 \cdot 10^4/(8.315 \cdot (273.16 + T)))$ ;  $\zeta (V) = -1.5 + (-1/(1 + \exp((V + 34)/5)))$ ;  $a_0 = 0.05$ ;  $Q_{10} = 5. \ l_A$ : with steady-state:  $l_\infty = \frac{1}{1+\alpha}$ ;  $\alpha = \exp(10^{-3} \cdot 3 \cdot (V + 50) \cdot 9.648 \cdot 10^4/(8.315 \cdot (273.16 + T)))$ ; time-constant:  $\tau = 0.26 \cdot (V + 44) \cdot qt$ ;  $Q_{10} = 1$ .

#### Synapses

The AMPA receptor was modeled with a double-exponential mechanism: rise 0.05 ms, decay 5.3 ms,  $E_{AMPA}$  0 mV. A single AMPA synapse was placed at 500 µm – midway on the dendrite; synaptic weight = 0.5 µS. The metabotropic synapse was not modeled in detail but was given as a local 12.5× increase in IP<sub>3</sub> (1.25 µM) in a 4 micron segment at the same location.

#### Simulation variations

In one set of simulations, we changed the density of IP<sub>3</sub>R and SERCA, by scaling  $p_{IP_3R}$  and  $p_{serca}$ , respectively. IP<sub>3</sub>R hotspots were modeled by scaling  $p_{IP_3R}$  at discrete locations, while keeping IP<sub>3</sub>R density in between the hotspots at 0.8× baseline IP<sub>3</sub>R density. ER stacks were simulated by scaling  $p_{IP_3R}$ ,  $p_{serca}$ , and  $p_{leakER}$  simultaneously, while keeping ER density in between the stacks at 0.8× baseline density. Changes in Ca<sup>2+</sup> buffering were simulated by

scaling the Ca<sup>2+</sup> diffusion coefficients,  $d_{Ca^{2+}_{cyt}}$  and  $d_{Ca^{2+}_{ER}}$ , in equal amounts. We also modulated  $d_{IP3}$  in a set of simulations. The range of parameter values examined were constrained to generate Ca<sup>2+</sup> waves with wave features (onset, speed, duration, amplitude) that were around ranges for experimentally-observed limits in neurons (Nakamura et al., 1999; Ross et al., 2005; Fitzpatrick et al., 2009).

#### **Electrochemical simulation variations**

One set of simulations (Fig. 11) was run with the electrical dynamics (see **Electrical dynamics** above) to determine how AMPA-ergic stimulation's alteration of membrane potential prior to an IP<sub>3</sub> stimulus impacts the subsequent  $Ca^{2+}$  waves. Simulations were run for 12 s. Background IP<sub>3</sub> concentration in these simulations was set between 0–0.01 mM to prevent spontaneous oscillations. AMPA activations (150 spikes with interspike interval of 25 ms) finished 250 ms prior to IP<sub>3</sub> stimulus (amplitude: 2.5 mM).

#### Data analysis

 $Ca^{2+}$  and  $IP_3$  concentrations were recorded in a 3D array using temporal and spatial resolution of 5 ms and 1  $\mu$ m. Cytosolic  $Ca^{2+}$  wave features were extracted by thresholding the array using an amplitude threshold of 2× baseline cytosolic  $Ca^{2+}$  concentration (2.5× for the electrical simulations). Wave onset was defined as the delay after IP<sub>3</sub> stimulus until the

 $Ca^{2+}$  passed threshold. Amplitude was maximum  $Ca^{2+}$  concentration. Speed was calculated using onset time at stimulation location and final onset time at location of wave termination (note that use of final onset time leads to artifactual lowering of wave speed in pathological condition of multiple  $Ca^{2+}$  waves). Duration was calculated as median at each position from wave onset to offset (going below threshold).

## 3 Results

This study involved over 4800 15–30 second simulations, testing how variations in levels of  $IP_3R$  density, SERCA density,  $IP_3$  hotspots, and ER stacks, altered wave initiation, amplitude, speed, and duration. An additional set of 32 twelve-second simulations were run, testing how the number of AMPA inputs prior to an  $IP_3$  stimulus impacted the calcium waves. Simulations were run using the NEURON simulator on Linux on a 2.27 GHz quad-core Intel XEON CPU. Thirty seconds of simulation ran in 1.5–2 minutes, depending on simulation type.

We simulated the arrival of IP<sub>3</sub> from a metabotropic receptor activation cascade as an instantaneous augmentation of IP<sub>3</sub> concentration to 1.25  $\mu$ M IP<sub>3</sub> (12.5× background) in a 4 micron segment in mid-dendrite (Fig. 2). IP<sub>3</sub> then spread gradually along the dendrite, providing a permissive effect for Ca<sup>2+</sup> activation of the IP<sub>3</sub> receptors (IP<sub>3</sub>Rs). Activation of IP<sub>3</sub>Rs permitted release of Ca<sup>2+</sup> from the endoplasmic reticulum (ER) stores into the cytosol (Fig. 2a). The elevation in cytosolic Ca<sup>2+</sup> was mirrored by a depression in ER Ca<sup>2+</sup> as the wave passed (Fig. 2b). The cytosolic Ca<sup>2+</sup> wave started once adequate Ca<sup>2+</sup> had built up to spillover and activate neighboring IP<sub>3</sub>Rs and initiate sustained positive feedback of Ca<sup>2+</sup> induced-Ca<sup>2+</sup>-release (CICR). The wave then spread bidirectionally, as Ca<sup>2+</sup> and IP<sub>3</sub> simultaneously diffused laterally. Ca<sup>2+</sup> waves were not able to reverse direction and propagate back towards their source due to an IP<sub>3</sub>R refractory period provided by inactivation gating. Parameters including *d*<sub>IP<sub>3</sub></sub> (1.415  $\mu$ m<sup>2</sup>/ms),

 $d_{_{Ca_{cyt}^{2+}}}(0.08\mu m^2/ms), d_{_{Ca_{ER}^{2+}}}(0.08\mu m^2/ms), p_{IP_{3R}}$  (120400.0 molecules/mM/ms),  $p_{serca}$  (1.9565 molecules/ms), and  $p_{leakER}$  (18.06 molecules/mM/ms) were adjusted from modeling (Wagner et al., 2004) and experiments (Allbritton et al., 1992), to produce a wave speed of 77 µm/s, which is comparable to neuronal Ca<sup>2+</sup> wave speeds measured experimentally (68±22 µm/s in Nakamura et al. (1999)). These parameters provided our baseline simulation for further parameter explorations. Median Ca<sup>2+</sup> elevation lasted ~ 1 s and peaked at 1.6 µM, both also comparable with experimental results (Ross et al., 2005; Fitzpatrick et al., 2009).

#### 3.1 IP<sub>3</sub>R density modulates excitability

IP<sub>3</sub>R was necessary for wave generation (Fig. 3).  $Ca^{2+}$  waves did not begin below IP<sub>3</sub>R density of 92.2% of baseline, when the first wave was obtained (Fig. 3a and left red dots on c,d,e,f). This wave had initiation delays ~100 ms longer compared to the baseline wave (Fig. 2a and middle red dot on Fig. 3c), with only a slightly lowered amplitude, lower speed, and lower duration. The major effect, time to onset, was strongly dependent on IP<sub>3</sub>R due to the need to gradually source a sufficient amount of  $Ca^{2+}$  to produce enough lateral spillover to initiate a positive feedback cycle, as the lower IP<sub>3</sub>R density produced a lower total  $Ca^{2+}$  flux from ER to cytosol. ER  $Ca^{2+}$  must also be sourced at a rate that exceeds the reuptake

governed by SERCA pumping. With increased density, time to onset decreased rapidly to a low of 25 ms.

Speed increased proportionately with increased IP<sub>3</sub>R (Fig. 3d). At the lowest value, wave speed was 72  $\mu$ m/s. Speed depended on both the rate of ER sourcing at each point of wave regeneration, and the rate of diffusion required to reach the next set of IP<sub>3</sub>Rs laterally. Proportionate increase in speed with density was due to the reduced requirement for additional Ca<sup>2+</sup> for positive feedback, since less Ca<sup>2+</sup> is needed to trigger the release of Ca<sup>2+</sup>. At each location, IP<sub>3</sub>R must source adequate Ca<sup>2+</sup> not only for positive feedback and spillover, but also to exceed the local sink back into ER provided by the SERCA pump. Ca<sup>2+</sup> wave duration and amplitude both had a small positive association with increases in IP<sub>3</sub>R (Fig. 3e,f), due to the higher density of IP<sub>3</sub>Rs liberating more Ca<sup>2+</sup> for the wave, which then remained elevated longer.

## 3.2 SERCA opposes IP<sub>3</sub>R effects

The SERCA pump returns  $Ca^{2+}$  back into the ER, opposing the release from  $IP_3R$  conductance. Therefore, the parameter dependence of measures was generally opposite to that seen with  $IP_3R$  density (Fig. 4). The least excitable wave was produced at the highest SERCA density tested (Fig. 4b). The most excitable wave shown here, at  $0.66 \times$  baseline SERCA (Fig. 4a), was caused by diminished  $Ca^{2+}$  reuptake after release. Further reduction in SERCA resulted first in spontaneous  $Ca^{2+}$  waves and then in simultaneous release from all ER sites at once due to regenerative responses starting from baseline levels (not shown).

Patterns of wave measures were similar to, but reversed from, those seen with IP<sub>3</sub>R increase. Onset to the first wave varied from 30 to 330 ms. Shorter onset times were produced by faster Ca<sup>2+</sup> accumulation with less reuptake. Similarly, the most excitable wave had a faster propagation speed (84  $\mu$ m/s) and longer duration due to longer-lasting Ca<sup>2+</sup> supporting faster spread. Amplitude showed a slight inverse relationship with increasing SERCA density, due to the higher SERCA reuptake of Ca<sup>2+</sup> into ER stores diminishing Ca<sup>2+</sup> availability for the wave (Fig. 4f).

At the maximal SERCA consistent with wave propagation,  $1.07 \times$  baseline, onset to wave initiation was increased dramatically (330 ms), wave speed decreased to ~75 µm/s, and duration decreased to 0.82 s (Fig. 4b). Ca<sup>2+</sup> waves were no longer able to initiate at higher levels, as Ca<sup>2+</sup> was sucked back into the ER before a Ca<sup>2+</sup> wave could ignite.

## 3.3 IP<sub>3</sub>R / SERCA balance regulates wave induction

The balance between IP<sub>3</sub>R and SERCA density determined whether  $Ca^{2+}$  waves could be elicited via IP<sub>3</sub> stimulation (Fig. 5). Excessive  $Ca^{2+}$  or insufficient IP<sub>3</sub>R would not allow a wave (regions 1,4). Maximum  $Ca^{2+}$  level in region 1 was low compared to regions 2 and 3, since ER  $Ca^{2+}$  had already partially leaked out of the ER prior to the IP<sub>3</sub> stimulus, and then diffused within the dendrite, preventing the large, localized rise in cytosolic  $Ca^{2+}$  which formed a  $Ca^{2+}$  wave.

Moving from region 1 towards region 2,  $IP_3R$  density decreased and SERCA density increased, both of which reduced cytosolic  $Ca^{2+}$ , which was no longer sustained above

threshold throughout the simulation. This allowed the IP<sub>3</sub> stimulus to elicit distinct waves. However, the high levels of IP<sub>3</sub>R produced multiple waves without return to a subthreshold state. These repetitive waves caused an artifactual lowering of measured wave speed (see Methods for details of calculation).  $Ca^{2+}$  amplitude was higher than in region 1 due to repeated higher elevations of cytosolic  $Ca^{2+}$  compared to balanced fluxes in region 1.

In region 3, there was a match across  $IP_3R$  and SERCA density, allowing for a single wave to be elicited from the  $IP_3$  stimulus. Within this region, for any given SERCA density, an increase in  $IP_3$  density increased  $Ca^{2+}$  flux from ER to cytosol, reducing time to onset, and increasing speed, duration, and amplitude. In region 4, SERCA dominated the dynamics, shuttling  $Ca^{2+}$  back into the ER too quickly to allow  $IP_3$  stimulation to initiate a  $Ca^{2+}$  wave.

## 3.4 IP<sub>3</sub> diffusibility regulates wave onset and speed

Altering the IP<sub>3</sub> diffusion constant ( $d_{IP_3}$ ) had a pronounced effect on Ca<sup>2+</sup> wave initiation (Fig. 6). With low  $d_{IP_3}$  (0.1415 µm<sup>2</sup>/ms), the IP<sub>3</sub> stimulus remained localized in the central stimulus region (~500 µm), producing a high prolonged local elevation with more rapid activation of local IP<sub>3</sub>Rs and shorter wave onset time (40 ms). Ca<sup>2+</sup> near the stimulus then remained elevated longer than at other dendritic locations. Since the IP<sub>3</sub> stimulus spread to neighboring dendritic locations more slowly, Ca<sup>2+</sup> wave speed was slightly reduced (73.5 µm/s).

High  $d_{IP_3}$  (1.981 µm<sup>2</sup>/ms) significantly delayed the onset of the Ca<sup>2+</sup> wave (230 ms; Fig. 6b), due to IP<sub>3</sub> diffusing quickly away, producing less local IP<sub>3</sub>R stimulation and increasing the time required to trigger the wave. However, once the wave was initiated, it spread slightly faster (77.6 µm/s), the pre-arriving IP<sub>3</sub> giving a slight boost at subsequent locations along the dendrite. Amplitude had a minimal inverse relationship with increasing  $d_{IP_3}$  (Fig. 6f) because local IP<sub>3</sub> spread out more quickly, reducing local IP<sub>3</sub>R activation. At  $d_{IP_3}$  values larger than 1.981 µm<sup>2</sup>/ms, IP<sub>3</sub> was so diffuse, that it could not elicit the Ca<sup>2+</sup> wave.

Overall, increasing  $d_{\text{IP}3}$  caused delayed onet (Fig. 6c), a minor increase in speed (Fig. 6d), and no appreciable change in duration or amplitude (Fig. 6ef).

#### 3.5 Pseudo-saltatory waves via IP<sub>3</sub>R hotspots

The prior simulations were performed using uniform density of ER mechanisms at all locations. However, there is evidence of inhomogeneities, for example at dendritic branchpoints, where elevations in local IP<sub>3</sub>R density (IP<sub>3</sub>R hotspots) might contribute to assisted propagation of Ca<sup>2+</sup> waves at these sites of potential failure (Fitzpatrick et al., 2009). In neurons, hotspots have average center-to-center spacing of approximately 20  $\mu$ m (edge-toedge spacing of 10  $\mu$ m) but show considerable variability in the spacing (Fitzpatrick et al., 2009). In the following simulations, we varied hotspot IP<sub>3</sub>R density while keeping the interhotspot IP<sub>3</sub>R density at 80% of baseline.

Both IP<sub>3</sub>R hotspot strength and spacing altered Ca<sup>2+</sup> wave speeds (Fig. 7). We started with a 20% reduction in density between hotspots (Fig. 7a), with the hotspots having 93% of the IP<sub>3</sub>R density used in Fig. 2. This alteration reduced wave speed from 77 to 68  $\mu$ m/s. However, the propagation pattern differed qualitatively, with its saltatory nature readily seen

as spots of high  $Ca^{2+}$  concentration which occur at the locations of IP<sub>3</sub>R hotspots. Local velocity parallels amplitude in showing an increase at the hotspots, where activation produces a local highly varying gradient (peaked high second spatial derivative) leading to a higher immediate diffusion speed. Slower wave progression occurs between hotspots where lower IP<sub>3</sub>R concentration only slightly boosts the wave progression from diffusion. Increased IP<sub>3</sub>R density at hotspots increased wave speed further to 90 µm/s (2.0× baseline; Fig. 7b). Further augmentation provided an approximately linear increase in speed due to faster and larger release of  $Ca^{2+}$  from the ER stores at the hotspots (Fig. 7c). Below ~93% of hotspot IP<sub>3</sub>R density, the waves did not initiate, since there was insufficient  $Ca^{2+}$  release to trigger waves (0 speeds at left of Fig. 7c). Wave propagation also did not occur in the absence of IP<sub>3</sub>R between hotspots, with the threshold for  $Ca^{2+}$  wave initiation having interhotspot IP<sub>3</sub>R densities of ~0.66 and ~0.93× baseline, respectively.

Hotspot spacing also modulated Ca<sup>2+</sup> wave speeds, with larger spacing producing slowing as Ca<sup>2+</sup> and IP<sub>3</sub> had to travel further via mildly-boosted diffusion between hotspots before being fully reboosted. At 1.87× baseline IP<sub>3</sub>R density with a spacing of 15  $\mu$ m, the wave had a speed of 100  $\mu$ m/s (Fig. 7d). Increasing center-to-center spacing between hotspots to 100  $\mu$ m resulted in reduction of wave propagation speed to 66  $\mu$ m/s. At the densities shown, hotspot spacing had a larger impact on wave speed (~66–100  $\mu$ m/s) than IP<sub>3</sub>R density (~68–90  $\mu$ m/s).

Intracellular Ca<sup>2+</sup> concentration is heavily regulated via buffering mechanisms, in part presumed to provide careful regulation of Ca<sup>2+</sup>-triggered signaling cascades. Ca<sup>2+</sup> buffering also modulates Ca<sup>2+</sup> diffusion efficacy and apparent Ca<sup>2+</sup> diffusion coefficient ( $\mathbf{D}_{Ca(App)}$ ). Because we were not modeling buffering directly in these simulations, we altered  $\mathbf{D}_{Ca(App)}$ instead. Changing  $\mathbf{D}_{Ca(App)}$  dramatically altered propagation speed over a wide range (Fig. 8). With diminished  $\mathbf{D}_{Ca(App)}$ , wave speed was substantially reduced to 25 µm/s (Fig. 8a). The converse of this was that Ca<sup>2+</sup> was relatively immobile, with duration at one location slightly increased (Fig. 8e). This heightened local Ca<sup>2+</sup> elevation also allowed for a shorter onset to wave initiation (20 ms; Fig. 8c). Increasing  $\mathbf{D}_{Ca(App)}$  above baseline level had opposite effects: wave speed was augmented (288 µm/s) but time to onset was delayed (40 ms), both due to faster spread of Ca<sup>2+</sup>. We also note that while onset varied nearly linearly, speed showed much greater alteration.

#### 3.6 Pseudo-saltatory waves via ER stacks

There are at least two ways that hotspots could occur in dendrites: Type 1. increased density of  $IP_3R$  at particular locations on homogeneous distribution of ER; Type 2. increased "density" (accumulation) of ER at particular locations. The former case was explored in the prior section. The latter case has been identified as locations of ER lamellar specialization that are sometimes described as ER *stacks*. We next explored such stacks as an alternative type of hotspot, noting that this increase in local ER provides increased density of SERCA and leak as well as  $IP_3R$  at the hotspot locations. In these simulations, the ER density between stacks was at  $0.8 \times$  baseline, consistent with continuous ER throughout the dendrite (Martone et al., 1993; Terasaki et al., 1994).

Increasing ER stack density from ~0.8× (Fig. 9a) to ~2× baseline (Fig. 9b), increased wave speed from 68 µm/s to 86 µm/s (Fig. 9c). This is due to more release of  $Ca^{2+}$  from the leak and IP<sub>3</sub>R channels, which was opposed by heightened SERCA activity. The duration of  $Ca^{2+}$  amplitude elevation was reduced at the ER stacks. In addition, as the ER stack density increased, the onset to wave initiation shortened (220 to 30 ms), due to heightened leak and IP<sub>3</sub>R extrusion of  $Ca^{2+}$  into cytosol. However,  $Ca^{2+}$  elevation duration decreased as the ER stack density increased (965 to 795 ms), due to heightened SERCA pumping.

Increasing the center-to-center spacing of ER stacks, while maintaining a fixed ER density of  $1.86 \times$  baseline (Fig. 9d–f), tended to decrease wave speed (93 to 71 µm/s). This was due to lower overall availability of ER for releasing Ca<sup>2+</sup> and spreading the wave. Interestingly, with heightened center-to-center spacing (15 and 100 µm) of ER stacks, the duration of heightened Ca<sup>2+</sup> elevation was increased (755 to 960 ms), due to lower overall presence of SERCA pumps.

Waves could not initiate below a minimum of  $0.8 \times$  ER stack density due to insufficient Ca<sup>2+</sup> sourcing. Effects of density and spacing of stacks were generally similar to those seen with IP<sub>3</sub>R hotspots. Onset time depended primarily on ER stack density. Variation of **D**<sub>Ca(*App*)</sub> with ER stacks (Fig. 10), produced results very similar to those seen with Type 1 hotspots (Fig. 8), except that the effects on time to onset were more pronounced. With ER stacks, the onsets were significantly larger (35–70 ms) compared to those for IP<sub>3</sub>R hotspots (20–40 ms). These heightened onset times with ER stacks were due to the higher SERCA activity, which reduces cytosolic Ca<sup>2+</sup> availability.

## 3.7 Electrical priming enhances Ca<sup>2+</sup> waves

Adding electrical dynamics (voltage-gated calcium channels – VGCCs, Na<sup>+</sup> and K<sup>+</sup> channels, AMPA synapses) to the dendrite allowed for a more complex set of interactions (Fig. 11). VGCCs admitted extracellular calcium into the cytosol, which was then taken up into the ER via SERCA pumps. Compared to the prior simulations, this increased ER steady-state calcium concentration (0.03 mM), allowing for larger contributions to  $Ca^{2+}$  waves. Cytosolic calcium equilibrated to 65 nM, due to a balance between SERCA uptake and VGCC calcium influx. In these simulations, the background IP<sub>3</sub> level was set to 0 mM to prevent spontaneous oscillations due to ongoing activation of IP<sub>3</sub>R. Calcium waves were more localized due to this absence of background IP<sub>3</sub>.

Activation by a 2.5 mM IP<sub>3</sub> stimulus at 7 s combined with the higher background cytosolic calcium levels to trigger IP<sub>3</sub>Rs to release calcium from the ER. ER calcium efflux then triggered a localized calcium wave (Fig. 11a middle,bottom), which spread approximately 150  $\mu$ m, consistent with experimental data (Fitzpatrick et al., 2009; Hong and Ross, 2007; Ross et al., 2005). The spreading calcium wave produced electrical effects, by contributing to the opening of calcium-dependent potassium channels, which hyperpolarized the membrane to -77 mV. Once the calcium wave had passed, the voltage recovered to baseline.

Prior AMPAergic stimulation (150 pulses; 25 ms interspike interval; 3000–6750 ms) produced spikes and further supplemented the ER calcium stores via a sequence of electrochemical interactions (Fig. 11b,c top):  $Ca^{2+}$  entry through VGCCs augmented

cytosolic Ca<sup>2+</sup> (early rise of 65–120 nM at center), augmenting ER Ca<sup>2+</sup> via uptake (Fig. 11b center). This gradually increased Ca<sup>2+</sup> throughout the dendrite (0.03–0.045 mM at center), *priming* the ER. Subsequent IP<sub>3</sub> activation was then able to liberate more ER calcium, inducing a higher-amplitude (7.5  $\mu$ M) calcium wave (Fig. 11b,c bottom). The larger amount of calcium liberated allowed the wave to spread 50  $\mu$ m further, and produced a longer duration (2 s), compared to the wave produced without synaptic priming. The heightened availability of Ca<sup>2+</sup> also enhanced the electrical effects, producing greater voltage suppression due to activation of the hyperpolarizing Ca<sup>2+</sup>-dependent potassium channels.

Increasing the number of AMPA stimuli between 0–150 (interspike interval of 25 ms; AMPA activation at 3 s; IP<sub>3</sub> activation at 7 s) produced nearly linear increases in the IP<sub>3</sub>-induced calcium wave amplitude (5.5–7.5 $\mu$ M) and duration (1.6–2.0 s) (not shown). Wave speeds tended to decrease slightly (57–63  $\mu$ m/s). There was no significant impact of AMPA stimulation on the wave onset, which was relatively fast for all simulations, at 5 ms past the IP<sub>3</sub> stimulation. Using increased background IP<sub>3</sub> (0.01 mM) and increasing the number of AMPA stimuli similarly produced higher amplitude, faster Ca<sup>2+</sup> waves with IP<sub>3</sub> stimulation (not shown).

## 4 Discussion

We have developed a reaction-diffusion (R×D) model of neuronal Ca<sup>2+</sup> waves with multiple compartments (ER and cytosol), multiple diffusing species (Ca<sup>2+</sup> and IP<sub>3</sub>), and multiple ER membrane mechanisms (IP<sub>3</sub> receptors, SERCA pumps, and ER leak). IP<sub>3</sub>Rs opened in response to a cytosolic Ca<sup>2+</sup> and IP<sub>3</sub> stimulus, admitting more Ca<sup>2+</sup> into the cytosol. Ca<sup>2+</sup> and IP<sub>3</sub> then diffused and bound to neighboring IP<sub>3</sub>Rs, triggering further release of Ca<sup>2+</sup> and initiating the process of Ca<sup>2+</sup>-induced-Ca<sup>2+</sup>-release (CICR). The Ca<sup>2+</sup> waves had properties matching those observed *in vitro* (Ross et al., 2005), including wave speed (~50–90 µm/s), amplitude (2 µM), and duration (~1 s; Fig. 2). We used our model to compare wave propagation in 3 scenarios: continuous IP<sub>3</sub>R distribution, IP<sub>3</sub>R hotspots, and ER stacks.

Because the cytosolic  $Ca^{2+}$  required for CICR depended on a supply from ER stores via the action of IP<sub>3</sub>Rs, the Ca<sup>2+</sup> waves demonstrated primary sensitivity to IP<sub>3</sub>R density, such that wave initiation required a minimal level of IP<sub>3</sub>R. IP<sub>3</sub>R density correlated negatively with onset, and positively with amplitude, speed, and duration of Ca<sup>2+</sup> waves (Fig. 3). SERCA pumps were responsible for setting the pace of reuptake of cytosolic Ca<sup>2+</sup> into the ER and had a large impact on the properties of Ca<sup>2+</sup> waves (Fig. 4). Low SERCA triggered the generation of hyper-excitable and spontaneous waves. High SERCA diminished wave spread and speed, while increasing onset to wave initiation.

We compared two hypothesized mechanisms for saltatory waves: 1. heightened density of  $IP_3R$  ( $IP_3R$  *hotspots*), and 2. heightened density of ER (ER *stacks*), where ER stacks heighten local leak and SERCA, as well as  $IP_3R$ .  $IP_3R$  hotspots were more effective in setting the pace of wave propagation and onset, whereas ER stack augmentation of  $Ca^{2+}$  waves was less pronounced due to the opposing effects of SERCA (leak and  $IP_3R$  increased

excitability, while SERCA decreased excitability). We hypothesize that  $IP_3R$  hotspots will predominate since they are the more effective mechanism for boosting  $Ca^{2+}$  waves.

Intracellular Ca<sup>2+</sup> concentration is heavily buffered. Although our model did not contain an explicit representation of buffering mechanisms, we used our model to test the effects of alterations in buffering capacity by modulating the apparent Ca<sup>2+</sup> diffusion constant ( $\mathbf{D}_{Ca(App)}$ ). We found that  $\mathbf{D}_{Ca(App)}$  had potent effects on Ca<sup>2+</sup> wave speed and propagation efficacy. Effects were similar for both IP<sub>3</sub>R hotspots (Fig. 8) and ER stacks (Fig. 10). Interestingly, at high Ca<sup>2+</sup> diffusibility, there were competing effects on excitability: higher speed but longer duration to wave initiation (Fig. 8c & d and Fig. 10c & d).

Finally, we extended our chemical signaling model with electrical components, including a set of calcium and voltage-dependent ion channels and synapses. We used this extended model to demonstrate that electrical activation via AMPA synapses and voltage-gated calcium channels primes ER calcium stores, and contributes to enhanced  $Ca^{2+}$  waves on subsequent IP<sub>3</sub> stimulation (Fig. 11).

#### 4.1 Predictions

Our model allows us to make the following experimentally-testable predictions.

- 1. Lowering apparent IP<sub>3</sub>R densities will lead to reduced spread of Ca<sup>2+</sup> waves. This is testable by using IP<sub>3</sub>R antagonists (Taylor and Tovey, 2010).
- 2. Reductions in SERCA with antagonist (*e.g.*, thapsigargin) will lead to hyperexcitability with multiple or spontaneous  $Ca^{2+}$  waves.
- **3.** Dendritic branch-point hotspots will be IP<sub>3</sub>R hotspots rather than ER stacks, allowing optimal boosting at locations where the wave might otherwise fail. Testable by using immunohistochemistry to look for IP<sub>3</sub>Rs and correlating with electron microscopy evaluation of stacking/ER concentration.
- 4. Increased buffering with added diffusible buffers provided via whole cell patch will reduce  $Ca^{2+}$  wave velocity. This is testable by modulating buffering properties and effectiveness via supply of  $Mg^{2++}$  (for parvalbumin modulation), BAPTA (for calbindin modulation), or EGTA (ethylene glycol tetraacetic acid) (Storm, 1987). Additionally, the use of dye indicators, used to follow wave progression, will alter velocity and other wave properties.

## 4.2 Roles of ER and of Ca<sup>2+</sup> waves

Since neuronal dendritic trees are very large (hundreds of microns), diffusion alone would make it impossible for dendritic  $Ca^{2+}$  to reach presumed targets in the soma and other dendritic locations for use in modulation of physiological processes including: synaptic plasticity, transcription regulation, and membrane current regulation (*e.g.*, *I*<sub>h</sub>). Our modeling emphasizes that  $Ca^{2+}$  waves can have large variability in speed and distance of propagation, which will have major effects on how  $Ca^{2+}$  is distributed and which targets are hit at what concentrations. However, we note that an entirely different role for ER has been proposed by Shemer et al. (2008). Instead of a role for ER in enhancing  $Ca^{2+}$  waves, they suggested

that the ER forms a "cable-within-a-cable" which is electrically active and would provide more rapid distal to proximal communication via electrical signaling comparable to that of the plasma membrane.

Pathologically,  $Ca^{2+}$  dysregulation may occur via multiple pathways, including heightened IP<sub>3</sub>R density, lowered SERCA density, or altered buffering. Dysregulation of  $Ca^{2+}$  homeostasis has been implicated in Alzheimer's disease (Lytton et al., 2014; Rowan and Neymotin, 2013; Rowan et al., 2014) and in ischemia, where  $Ca^{2+}$  signaling is an important element in the triggering of apoptosis (Green and LaFerla, 2008; LaFerla, 2002; Stutzmann, 2005; Thibault et al., 1998; Zündorf and Reiser, 2011; Taxin et al., 2014).

## Acknowledgments

Research supported by NIH grant R01 MH086638 and NIH grant T15 LM007056. The authors would like to thank Larry Eberle (SUNY Downstate) for Neurosim lab support, Tom Morse (Yale) for ModelDB support, and Herman Moreno (SUNY Downstate) & the Shepherd lab (Yale) for helpful comments. The authors declare no competing financial interests.

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## Figure 1.

Schematic of the dendrite model (1000 µm length; 1 µm diameter) showing the extracellular space, cytosol and endoplasmic reticulum (ER). Thick black lines represent the plasma and ER membranes. The extracellular space provides a source and sink for multiple molecules.  $Ca^{2+}$ ,  $Na^+$ ,  $K^+$ , and IP<sub>3</sub> are depicted as black circles. These molecules enter and exit the dendrite and its sub-compartments via the channels/receptors (gray pores) in directions indicated by arrows. SERCA pump moves cytosolic  $Ca^{2+}$  into the ER. Lightning symbols represent stimulus locations (AMPA-mediated depolarization; mGluR-mediated IP<sub>3</sub> augmentation). Dotted lines between molecules and receptors indicate the receptor is modulated by the molecule (*e.g.*,  $Ca^{2+}$  activates  $K_{Ca}$  and IP<sub>3</sub>R).

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## Figure 2.

 $Ca^{2+}$  wave propagation with baseline parameters. Elevated IP<sub>3</sub> stimulus placed at middendrite (500 µm on y-axis) after 2 s past start of simulation. (a) cytosolic [Ca<sup>2+</sup>] shows a wave of increased concentration (b) ER [Ca<sup>2+</sup>] shows a mirror image wave of decreased concentration as Ca<sup>2+</sup> is released to cytosol.

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## Figure 3.

 $Ca^{2+}$  wave propagation is sensitive to the density of IP<sub>3</sub>R. From left to right, the three red dots in c–f correspond to features of the following  $Ca^{2+}$  waves: (a) low IP<sub>3</sub>R density (0.9× baseline); baseline IP<sub>3</sub>R density (Fig. 2); (b) high IP<sub>3</sub>R density (1.8× baseline). (c) time to wave onset (0 indicates no wave), (d) wave propagation speed, (e) duration (2× Ca<sup>2+</sup> elevation; median duration across dendritic locations) (f) peak amplitude. (c–f each show 91 different equally spaced parameter values) (Colorscale in (a) and (b) same as in Fig. 2.)

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## Figure 4.

Ca<sup>2+</sup> wave propagation is sensitive to the density of SERCA. Tested by varying SERCA density relative to baseline (n = 45). (**a**) hyper-excitable Ca<sup>2+</sup> wave produced with 0.66× baseline SERCA, (**b**) diminished Ca<sup>2+</sup> wave produced with 1.03× baseline SERCA, (**c**) time to Ca<sup>2+</sup> wave onset (0 indicates no wave at the SERCA level), (**d**) speed, (**e**) duration (median duration of Ca<sup>2+</sup> elevation across dendritic locations), (**f**) peak amplitude. Left (right) red dots in c–f corresponds to activity shown in (**a**) ((**b**)). (Colorscale in (**a**) and (**b**) same as in Fig. 2.)

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## Figure 5.

Ca<sup>2+</sup> wave propagation is sensitive to density of IP<sub>3</sub>Rand SERCA. 2838 simulations varying IP<sub>3</sub>R (86 levels) and SERCA density (33 levels) relative to baseline shows different dynamics: region 1: sustained Ca<sup>2+</sup> elevation precluded IP<sub>3</sub>-evoked wave; region 2: multiple Ca<sup>2+</sup> waves; region 3: single Ca<sup>2+</sup> wave; region 4: insufficient IP<sub>3</sub>R to source Ca<sup>2+</sup> wave. (a) time to Ca<sup>2+</sup> wave onset where present, (b) speed where wave present, (c) duration where wave present (median duration of Ca<sup>2+</sup> elevation across dendritic locations), (d) maximum calcium level.

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#### Figure 6.

 $m Ca^{2+}$  wave propagation is sensitive to IP<sub>3</sub> diffusion coefficient ( $d_{IP_3}$ ). From left to right, red dots in c-f correspond to features of Ca<sup>2+</sup> waves with (**a**) low (0.1415 µm<sup>2</sup>/ms); baseline (1.415 µm<sup>2</sup>/ms; Fig. 2); (**b**) high (1.981 µm<sup>2</sup>/ms)  $d_{IP_3}$ . (**c**) time to wave onset, (**d**) wave propagation speed, (**e**) duration (2× Ca<sup>2+</sup> elevation; median duration across dendritic locations) (**f**) peak amplitude. (Colorscale in (**a**) and (**b**) same as in Fig. 2.)

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## Figure 7.

Comparison of waves (and wave speeds) generated from varying IP<sub>3</sub>R hotspot density (a–c) and separation (d–f). (a) 20  $\mu$ m spacing; 0.93× density. (b) 20  $\mu$ m spacing; 2× density. (c) wave speed as a function of IP<sub>3</sub>R hotspot density, (d) 15  $\mu$ m spacing; 1.87× density. (e) 100  $\mu$ m spacing; 1.87× density. (f) wave speed as a function of IP<sub>3</sub>R hotspot spacing. (Colorscale in a,b,d,e same as in Fig. 2.) Red points in (c) ((f)) are from waves in (a),(b) ((d),(e)).

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## Figure 8.

Comparison of waves generated from varying  $\mathbf{D}_{Ca(App)}$  in the presence of IP<sub>3</sub>R hotspots. Ca<sup>2+</sup> waves generated when using a Ca<sup>2+</sup> diffusion coefficient of (**a**) 0.008 µm<sup>2</sup>/ms and (**b**) 0.8 µm<sup>2</sup>/ms. (**c**) time to Ca<sup>2+</sup> wave onset (0 indicates no wave at the Ca<sup>2+</sup> diffusion level), (**d**) speed, (**e**) duration (median duration of Ca<sup>2+</sup> elevation across dendritic locations), (**f**) peak amplitude. Left (right) red dots in c–f corresponds to activity shown in (**a**) ((**b**)). (Colorscale in (**a**) and (**b**) same as in Fig. 2.)

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## Figure 9.

Comparison of Ca<sup>2+</sup> waves generated from varying ER stack properties. Simulations were run by varying density and spacing of ER stacks. ER density in between stacks was set at  $0.8 \times$  baseline. (a) Ca<sup>2+</sup> wave from  $0.8 \times$  ER stack density with 20 µm spacing, (b) Ca<sup>2+</sup> wave from  $2.0 \times$  ER stack density with 20 µm spacing, (c) Ca<sup>2+</sup> wave speed as a function of ER stack density (fixed 20 µm spacing), (d) Ca<sup>2+</sup> wave from 15 µm ER stack spacing with ~1.87× density, (e) Ca<sup>2+</sup> wave from 100 µm ER stack spacing with ~1.87× density, and (f) wave speed as a function of ER stack spacing. Note the bumps in Ca<sup>2+</sup> release around ER stacks. (Colorscale in a,b,d,e same as in Fig. 2.) Red points in (c) ((f)) are from waves in (a), (b) ((d),(e)).

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## Figure 10.

Comparison of waves generated from varying  $D_{Ca(App)}$  in the presence of ER stacks. (a) Local Ca<sup>2+</sup> event generated with Ca<sup>2+</sup> diffusion coefficient of 0.0 µm<sup>2</sup>/ms. (b) Ca<sup>2+</sup> waves generated when using a Ca<sup>2+</sup> diffusion coefficient of 0.8 µm<sup>2</sup>/ms. (c) time to Ca<sup>2+</sup> wave onset, (d) speed, (e) duration (median duration of Ca<sup>2+</sup> elevation across dendritic locations), (f) peak amplitude. Left (right) red dots in c–f corresponds to activity shown in (a) ((b)). (Colorscale in (a) and (b) same as in Fig. 2.)

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## Figure 11.

Electrical stimulation with increased number of AMPA activations enhances  $Ca^{2+}$  waves induced by IP<sub>3</sub> (2.5 mM at 7 s). (a) Control simulation:  $Ca^{2+}$  wave with no AMPA inputs prior to the IP<sub>3</sub> stimulus. (b)  $Ca^{2+}$  wave with train of 150 AMPA inputs (onset: 3 s; interspike interval: 25 ms) prior to the IP<sub>3</sub> stimulus. (c) Comparison of voltage (top), ER  $Ca^{2+}$  (middle), and cytosolic  $Ca^{2+}$  (bottom) in control (black) and simulation with 150 AMPA inputs (red).

## Table 1

Baseline parameters for the  $Ca^{2+}$  wave model.

 $f_{cyt} = 0.83$ 

 $f_{ER} = 0.17$ 

 $\bar{p_{\mathrm{IP_{3R}}}} = 120400.0$  molecules/mM/ms

 $k_{\rm IP3} = 0.00013 \text{ mM}$ 

 $k_{act} = 0.0004 \text{ mM}$ 

 $k_{inh} = 0.0019 \text{ mM}$ 

 $p_{leakER} = 18.06$  molecules/mM/ms

 $\bar{p_{serca}} = 1.9565$  molecules/ms

 $k_{serca} = 0.0001 \text{ mM}$ 

 $\tau_{IP_{3}R} = 400 \text{ ms}$ 

$$d_{_{\rm Ca^{2+}_{\rm cvt}}}{=}0.08\mu{\rm m^2/ms}$$

 $d_{_{\rm Ca}^{2+}_{\rm ER}}{=}0.08\mu{\rm m}^2/{\rm ms}$ 

 $d_{\rm IP3} = 1.415 \ \mu {\rm m}^2/{\rm ms}$