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Biophysics of Mitosis

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Abstract

Mitosis is the process by which eukaryotic cells organize and segregate their chromosomes in preparation for cell division. It is accomplished by a cellular machine composed largely of microtubules and their associated proteins. This article reviews literature on mitosis from a biophysical point of view, drawing attention to the assembly and motility processes required to do this complex job with precision. Work from both the recent and the older literature is integrated into a description of relevant biological events and the experiments that probe their mechanisms. Theoretical work on specific subprocesses is also reviewed. Our goal is to provide a document that will expose biophysicists to the fascination of this quite amazing process and provide them with a good background from which they can pursue their own research interests in the subject.

A. INTRODUCTION

A1. A Mitosis Primer

Mitosis is the process by which all eukaryotic cells segregate their already duplicated chromosomes in preparation for cell division, or “cytokinesis”. The cell’s preparations for mitosis include the duplication of its DNA. The resulting “sister” DNA double helices are tied together by multiple copies of a protein complex laid down as DNA replicates. Cells then synthesize many additional macromolecules, so at the end of “interphase” (the period between cell divisions), they contain all the materials needed to form two viable cells. Mitosis and cytokinesis then separate this biochemically doubled cell into two essentially identical objects, each equipped to grow and divide again. Thus, mitosis is one part of a cell’s growth and division cycle (Figure 1A).

The first physical problem that a soon-to-divide cell must solve is the restructuring of its chromosomes, so each is sufficiently compact to be separable from its sister within a space no bigger than the cell. In a human cell the DNA molecules range in length from ~1.9 – 8.5 cm, whereas the nucleus that contains these 46 duplicated objects is an approximately spherical compartment with diameter usually $< 8 \mu\text{m}$. Thus, each DNA duplex must be

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reduced in length by >1,000-fold. Such “condensation” is achieved in multiple steps, initially by wrapping the DNA around nucleosome “core particles” (octamers of the histone proteins); this makes the material called “chromatin.” Fibers of chromatin are then coiled and looped in ways that are not yet well understood until each chromosome is only a few micrometers long and usually <1 μm thick. For a review of chromosome condensation, see (Belmont, 2006).

The majority of chromatin condensation occurs during the first stage of mitosis, called “prophase”. As each piece of chromatin becomes shorter and thicker, it becomes visible in the light microscopy as a thread, hence the term mitosis (mitos = thread in Greek). With continued condensation chromosomes commonly display their dual nature: the sister DNA duplexes formed at replication become distinct. These are called “chromatids”. At this stage sister chromatids are still linked by the “cohesins” protein complexes laid down during DNA replication. These attachments are of great importance for the logic of mitosis because accurate segregation of sister chromatids depends on their being attached until the moment when all chromatids will begin segregation at essentially the same time.

Commonly there is one place along the length of each chromosome where sister chromatids are particularly tightly coupled; this is the “primary constriction” or “centromere”. As the chromosomes condense, cohesins dissociate from much of the chromosome arms, but near the centromere the chromatids remain bound until their segregation begins. The centromere is also the place where each chromosome develops specializations for its attachment to the machine that will effect its segregation (Figure 1B).

Prophase ends when the already condensed chromosomes begin to interact with the segregation machine, the “mitotic spindle”; this name came from on the fact that in many cell types the structure resembles the spindle that was used years ago to twist wool into yarn. The spindle is an assembly of microtubules; it organizes the chromosomes into a disk-shaped array, then pulls the two identical parts of each chromosome to opposite ends of the cell (Figure 2). In many cells, including mammals, chromosome-spindle interaction is permitted by a disassembly of the envelope that separated nucleus from cytoplasm throughout interphase. The spindle, which forms from cytoplasmic components, can now influence chromosome position and behavior. In many micro-organisms, however, the spindle forms within the nuclear envelope, which never breaks down (a “closed mitosis”). Thus, there is nothing fundamental for mitosis about the mixing of nucleus and cytoplasm.

Spindles are composed largely of microtubules (MTs), which are polymers of the GTP-binding protein, tubulin in complex with many associated proteins. Spindle formation initiates the process of chromosome organization in the mitotic stage called “prometaphase.” At this time, the essential event is the attachment of all chromosomes to spindle MTs in such a way that each chromatid of every chromosome is associated with MTs that are associated with one and only one end of the mitotic spindle. A second event of prometaphase is the migration of all chromosomes to the spindle mid-plane, or “equator”, a process called “congression”. Once these motions have been accomplished, the cell is said to be in “metaphase” (Figure 2A).

Normal cells include quality control processes that determine whether each chromosome is properly attached to the spindle before segregation is allowed to begin; this is the “spindle assembly checkpoint.” Shortly after this checkpoint has been satisfied, the cohesins that have been holding sister chromatids together are cleaved by a protease. MT-generated forces acting on the now-independent sister chromatids move them to opposite ends of the cell in a process called “anaphase.” If the nuclear envelope dispersed during spindle formation, it now reforms on the still-condensed chromosomes by the application of vesicles derived largely, if not entirely, from the previously dissociated nuclear envelope. As these membranes are fusing to define the two nuclear compartments, the cell initiates cytokinesis, the process that divides the cytoplasm into two approximately equal parts, each of which contains its own nucleus. At the same time, the chromosomes de-condense, and the daughter cells return to interphase (Figures 1,2).

Cells employ a special mitotic machine for chromosome segregation because each chromosome is present in only the two copies made by DNA replication, and each daughter cell needs one copy of every chromosome to be able to grow and divide again. At the same time, cells contain vast numbers of structures, like mitochondria, ribosomes, cytoplasmic membranes, and enzymes, etc. that must also be segregated to daughter cells for the progeny to be healthy. No special device is used for these segregation processes, presumably because the objects are sufficiently numerous that the likelihood all copies will be in one half of the cell at cytokinesis is very small. In some cells, structures like mitochondria are present in only a few copies; in such cases, the structure in question will commonly fragment into many pieces for cell division, so the laws of large numbers will help to assure their equipartition. Some specialized cells, e.g., certain algae, contain only one copy of an essential organelle, like a chloroplast or a flagellum; its segregation to daughter cells is then coupled to the position of the mitotic spindle and thus to the process of cytokinesis. Thus, cell division is both parsimonious and clever.

A2: The problems considered here

The above account shows that mitosis includes many processes. This article will focus on the organization and segregation of already-condensed chromosomes. These processes depend on the mitotic spindle, so our presentation is organized in terms of spindle formation and action. For treatments of membrane dynamics during mitosis, see (Altan-Bonnet et al., 2004; Tang et al., 2008); for cytokinesis, see (Pollard, 2010). Our review of the remaining mitotic phenomena will rely on some knowledge of MT dynamics, reviewed in (Desai & Mitchison, 1997) and of MT-dependent motor enzymes (Gatlin & Bloom, 2010; Vale, 2003). We will not, however, deal with MT-related structures, such as centrioles or flagella; nor will we consider the important differences in chromosome structure between mitosis and meiosis, which allow the latter process to reduce the chromosome number from a diploid to a haploid complement in preparation for sexual reproduction.

Mitosis itself is of sufficient interest that it has been reviewed with almost overwhelming frequency; a complete list of even the most relevant reviews would try the reader’s patience. We draw attention, though, to a few recent papers that will be helpful for the interested

reader: (Cheeseman & Desai, 2008; Dumont & Mitchison, 2011; McAinsh et al., 2003; Walczak & Heald, 2008; Welburn & Cheeseman, 2008)

A3. Structure of a typical mitotic spindle

Spindle organization is most easily understood at metaphase, after the forming spindle has already imposed order on chromosome position. Below we will see how this structure forms and how it changes as it does its job. Once fully formed, most spindles have two identical ends, called “poles”. An axis of rough rotational symmetry runs between these poles, and there is a 2-fold axis of approximate symmetry perpendicular to the cylindrical axis. The duplicated but not-yet-segregated chromosomes lie at or near the spindle midplane, or “equator”, forming the “metaphase plate” (Figure 2A). Each of the two chromatids in every metaphase chromosome is attached to one spindle pole, while its sister is attached to the other; meanwhile, the sister chromatids are still mechanically coupled by cohesins. This arrangement anticipates the functional symmetry of anaphase, in which each chromatid will move toward the pole to which it is attached. The resulting segregation of sister chromatids to sister poles is the essence of mitosis.

The details of spindle structure are more complex than the overview. Each MT is a polar polymer in the sense that all the α - β -tubulin dimers from which it is made point in one direction; every β -tubulin is distal to its associated α -tubulin relative to the spindle pole with which the MT is associated. The MT end with β -tubulin exposed is usually more dynamic than the α -tubulin end, faster in both growth and shortening; it is called the “plus-end”. MT polarity is important for spindle tubulin dynamics and for the action of the many motor enzymes that use ATP to power their motility over the MT surface.

In vivo tubulin polymerization is commonly initiated by a lock-washer shaped, multi-protein complex called the γ -tubulin ring complex (γ -TuRC); each of these contains about 13 copies of the γ isoform of tubulin and several associated proteins (Wiese & Zheng, 2006). This complex defines the position of MT nucleation, the polar orientation of the resulting polymer, and the lattice into which tubulin assembles. Most copies of the γ -TuRC are concentrated near the spindle poles, thanks to tethers made from long, α -helical, coiled-coil proteins. The presence of two foci of MT initiation results naturally in a bipolar structure with the initiating sites at its ends.

Spindle MT lengths are heterogeneous, so some plus ends lie quite distant from their initiating site, while others are close (Figure 2a). Moreover, some MTs are initiated elsewhere in the spindle. The distributions of MTs lengths and positions are complex (Mastrorade et al., 1993; Yang et al., 2007a), but all these complexities maintain the two-fold symmetry of the spindle, which is fundamental to its ability to segregate chromosomes.

The plus ends of some spindle MTs become associated with chromosomes in ways that are essential for mitotic success. The centromere of each chromatid binds many proteins specific to this locus, forming a specialization called the “kinetochore.” This is the most important part of the chromosome for spindle MT attachment; it is also responsible for chromosome segregation at anaphase. Kinetochores are central to our understanding of the biophysics of mitosis, so they are considered in detail below. Suffice it to say here that each kinetochore

shows an affinity for MTs with a stronger binding to the MT plus end (Huitorel and Kirschner, 1988); the resulting kinetochore-MT bonds attach that chromatid to the spindle, whereupon it is pulled toward the pole with which it has become associated. Conversely, kinetochore binding at an MT end makes that tubulin polymer more stable (Mitchison, 1989b). The kinetochores on sister chromatids are commonly arranged back-to-back, so they naturally tend to form connections with sister poles through the MTs whose plus ends they bind. In reality, however, this “proper” attachment does not always form, and we will see that cells have developed error correction processes to bring errant chromosomes into a proper, two-fold symmetric attachment to the spindle.

In many cells, including those of most vertebrates, the spindle poles initiate some MTs that grow away from the sister pole and do not contribute to the spindle per se. These form star-like arrays called “asters” (Figure 2). Astral MTs can be long enough to reach the edge of the cell, where they interact with proteins in the cell “cortex”, the gelatinous, actin-rich layer just beneath the inner surface of the plasma membrane. Not all cells have asters, so they cannot be essential for mitosis, but they do help to position the spindle near the cell’s middle, and they can contribute both to spindle mechanics and the position of cytokinesis.

A4. Mechanical requirements for mitosis

Clearly, spindles can generate whatever forces are necessary to organize the chromosomes during prometaphase and to segregate them in anaphase. A look at chromosome velocities suggests that these forces might be very small. Though chromosomes can be big on a cellular scale, their movements are slow, usually ranging from 0.01 – 1 $\mu\text{m}/\text{sec}$, where the high end of this range is seen only during the early stages of mitosis and is usually brief. Based on both reasonable estimates and measurements of cytoplasmic viscosity, the forces necessary to overcome viscous drag on chromosomes is at most a few tens of piconewtons (pN). Such forces could be generated by only a few motor enzymes of the kinesin or dynein families. Thus, force development seems not to be a major task for the spindle. When mitosis goes wrong, however, chromosomes can become entangled, so their segregation required more force. Indeed, direct measurement has shown that spindles can actually produce $\sim 10,000\times$ the force necessary to overcome the viscous drag that would act on a chromosome (Nicklas, 1983). Thus, both the generation and regulation of spindle forces are topics of great interest.

The principal job of the spindle is to segregate the chromosomes without mistakes. A daughter cell that lacks an entire chromosome is not likely to thrive, and even a cell with one extra chromosome may be compromised, as is seen in the pathologies faced by children born with Down or Patau Syndromes (three chromosomes 21 or 13). The accuracy of chromosome segregation has been measured in budding yeast, and the chance of chromosome loss is $\sim 10^{-5}/\text{cell division}$ (with 16 chromosomes per haploid cell) (Hartwell & Smith, 1985). This is an impressive accomplishment for a machine that functions at the micrometer scale, using engines that are truly nano-machines. An important aspect of our review will be an effort to understand how spindles can be so accurate.

A5. The mechanics to be understood

The processes considered below are: 1) How are spindle MTs formed and organized so a metaphase structure is obtained? 2) How do spindle MTs interact with chromosomes so sister kinetochores become attached to sister poles? 3) How do chromosomes get moved to the spindle midplane and form the metaphase plate? 4) How does a cell decide when to start anaphase and where to place the daughter cells and will form at cytokinesis? 5) How does a cell exert forces on a chromosome to pull sister chromatids to sister poles? and 6) How does a spindle elongate to push the separated chromosomes to opposite ends of the cell?

B. BUILDING AND MAINTAINING A BIPOLAR SPINDLE

B1. Introduction

Here we consider spindles of two kinds: those with and those without “centrosomes” at their poles. Centrosomes are structured initiators of MT polymerization formed by clustering multiple copies of the γ -TuRC into one region of the cell. They duplicate during interphase, usually about the time of DNA replication, so a cell enters mitosis with two of them. The active γ -TuRCs associated with each centrosome define the position, number, and polarity of the MTs that form. However, some cells form spindles without centrosomes; this includes all higher plants and oocytes from many animal species. Acentrosomal spindle formation can follow either of two pathways: one emphasizes MT initiation in the vicinity of chromosomes (as seen in frog oocytes), the other takes advantage of both a cylindrical array of cytoplasmic MTs and MT initiation from polar regions, even though centrosomes are absent (common in higher plants). Kinetochores too can contribute to spindle MT formation, though the extent of this pathway varies among organisms. The striking thing about this diversity is that all these assembly pathways lead to essentially the same metaphase structure, suggesting that this arrangement is a favorable place on the energy landscape. For an alternative treatment of this interesting problem, see (Duncan & Wakefield, 2011).

B2. Centrosome-mediated spindle formation

Cells that contain two identical MT initiation sites have an obvious route by which to form a two-fold symmetric MT array; simultaneous activation of these sites will form just the structure needed.

Centrosome morphology is diverse, but all centrosomes include a structured core to which multiple copies of the γ -TuRC are connected. This assembly defines the splays of MTs that grow as the spindle forms, including the asters seen in many animal cells (Figures 2A and 3A). A collection of long, α -helical coiled-coil proteins links the structured material at the center of the centrosome with multiple copies of the γ -TuRC. In vertebrate cells one such linker is the calmodulin-binding protein, pericentrin (Dicthenberg et al., 1998), a.k.a. kendrin. These tethers probably also function to hold the minus ends of MTs in place, giving the whole assembly mechanical coherence. However, centrosome function is more complicated than a simple tethering of MT initiators; it requires the action of multiple protein kinases, e.g., Aurora A (Cowley et al., 2009; Glover et al., 1995), whose localization to the centrosome requires the MT-associated protein TPX2 (Bird & Hyman, 2008; Garrett et al., 2002). The kinase, Plk1, is also pole-associated and important for spindle formation (Barr et

al., 2004), as are the protein phosphatases, PP6 (a negative regulator of Aurora kinase) and PP2A (which has many substrates) (Bollen et al., 2009). It seems likely, however, that cells contain some additional MT-initiating capacity that is not dependent on the γ -TuRC, because when the gene for γ -tubulin is knocked down, many centrosome-associated MTs still form, though their organization is generally incorrect (Mahoney et al., 2006). It is not yet certain whether such polymer initiation is due to incomplete depletion of γ -tubulin, spontaneous MT formation as occurs *in vitro*, or some alternative initiation complex.

Centrosomes accumulate several MT-dependent motor enzymes, including those with either plus- or minus end-directed activities (Gaglio et al., 1997). Kinesin-13s, which are ATPases that work not as motors but as MT disassembly engines, are also common (Desai, 1999). Some centrosomes also concentrate up to three classes of AAA ATPases like katanin that can sever MTs (Zhang et al., 2007). This may explain how some MTs lose their minus end caps, as seen in the blastomeres of a nematode (O'Toole et al., 2003). Tethering such MTs to the centrosome may depend on "patronin", which stabilizes MT minus ends in *Drosophila* S2 cells against enzymatically driven shortening (Goodwin & Vale, 2010). Thus, the centrosome is a site cell regulation, MT initiation, anchorage, and severing.

One of the best-studied centrosomes is that found in budding yeast. It is a plaque built right into the nuclear envelope (which remains intact throughout mitosis, Figure 3B). Its nuclear surface binds multiple copies of an α -helical coiled-coil protein, SPC110, which in turn links a small γ -tubulin complex to the nuclear-facing surface of the centrosome (Kilmartin & Goh, 1996). Although SPC110 is not homologous with pericentrin, the idea of clustering γ -tubulin into a defined geometry is the same. Moreover, while the spindle is forming, the cytoplasmic face of the centrosome uses a different fibrous protein to bind the γ -TuRCs that initiates MTs that serve as asters. This example shows how cells can use quite different proteins to accomplish essentially the same function. It also shows the use an interphase structure (the nuclear envelope) to define the placement of key MT organizing structures (the centrosomes), so they will be able to make MTs that influence events in both the cytoplasm and the nucleus.

The MT-initiating activity of a centrosome is under cell cycle control. In mammals 5–10 times more MTs grow from a mitotic centrosome than from one in interphase (Snyder & McIntosh, 1975). Indeed, the amount of pole-localized γ -tubulin increases as the cell goes into mitosis (Khodjakov & Rieder, 1999). These changes are accompanied by changes in the phosphorylation of many centrosome-localized proteins, but how these modifications alter the ability of the structural material in the centrosome to bind γ -TuRCs is not yet known. The salient point is that all centrosome-based spindles form from a pair of MT initiating structures, making a bipolar framework of MTs that can then attach the chromosomes.

B3. Mechanisms for separating duplicate centrosomes at the onset of spindle formation

In cells that are growing and dividing, centrosomes usually lie close to the nucleus through most of interphase. They commonly duplicate around the onset of S-phase but remain close to one another, functioning as a single MT organizing center until the spindle begins to form. Centrosome separation commonly occurs during prophase, while the nuclear envelope is still intact. This process has been proposed to result from pushing forces generated by

interactions between MTs that grow from sister centrosomes. Such a model is supported by the observation that inhibition of Kinesin5, a homo-tetrameric, plus-end directed motor enzyme, blocks centrosome separation, leading to monopolar spindles, (Brust-Mascher & Scholey, 2007). A mechanism by which this plus end-directed MT motor might contribute to centrosome separation is diagrammed in Figures 4A and B. In some cells, e.g. those cultured from newt heart, the spindle forms as the centrosome separates (Taylor, 1959). In this case motor-driven processes may still pertain, but spindle MT growth, including the elongation of kinetochore-associated MTs, is also a part of the separation process (Toso et al., 2009).

An alternative view emerges from the observation that inhibition of dynein, a minus-end directed motor, can block centrosome separation in at least some cells. Microinjection of function-blocking anti-dynein causes the pair of adjacent centrosomes to function as a single spindle pole and form a monopolar spindle (Vaisberg et al., 1993). Insights into how this minus end-directed motor might function in centrosome separation have come from several studies on dynein's action in vivo (Waters et al., 1993). Dynein can couple to actin, probably through its cellular partner, the dynactin complex. This binding tethers some dynein to the cell cortex (Dujardin & Vallee, 2002). Here, it can interact with the ends of MTs initiated by centrosomes and pull each centrosome toward the cell surface. When two vicinal centrosomes are both affected in this way, the balance of forces will help to pull the centrosomes apart (Figure 4C). A dynein-mediated mechanism of this kind has been realized in vitro, acting on a single centrosome in a shallow, cylindrical well (Laan & Dogterom, 2010). The resulting forces center a single centrosome, but such pulling forces could obviously contribute to the separation of duplicate centrosomes in some systems.

B4. Spindle formation without centrosomes

The best-studied system of acentrosomal spindle formation is found in frog oocytes. Complex biological activities persist in cell-free extracts of egg cytoplasm (Hyman & Karsenti, 1996). Sperm, sperm heads, isolated somatic nuclei, and even phage DNA can be introduced into this medium, whereupon mitotic spindles will form. When centrioles are present, e.g., from the basal bodies of sperm flagella, they organize centrosomes and form MT asters that interact with chromosomes and produce a functional spindle. When centrosomes are not present, spindles still form following a wide range of treatments. The most interesting for our purposes is the addition of chromosomes, or even naked DNA, which the egg extract turns into chromatin by adding histones and other chromosomal proteins (Heald et al., 1996).

In this extract chromatin binds a GTP exchange factor for the small monomeric G-protein, Ran (Carazo-Salas et al., 1999; Kalab & Heald, 2008). On the other hand, proteins that activate RanGTPase (Ran-GAPs), thereby stimulating Ran's self-terminating hydrolytic activity, are uniformly dispersed. The result is a gradient in RanGTP, with its highest concentrations immediately around the chromosomes (Athale et al., 2008). At least two protein factors respond to this gradient in ways that influence MT behavior. TPX2, mentioned previously as a centrosomal protein with RanGTP-stimulated MT-binding and initiating activity, can promote MT formation without γ -tubulin (Gruss & Vernos, 2004). CDK-11 is a RanGTP-stimulated, cyclin-dependent kinase that confers stability on MTs

(Yokoyama et al., 2008). Together with several MT-associated proteins (MAPs), these activities lead to a centrosome-independent growth of MTs in the vicinity of chromatin (Athale et al., 2008).

Current data show Ran-GEF distributed uniformly over the chromosomes, but the MT-initiating activity of chromosomes is stronger near kinetochores than elsewhere for reasons that are not yet well understood. This behavior was first observed by Pease, using high hydrostatic pressure to depolymerize spindle fibers, then a release of pressure to allow them to reform (Pease, 1946). Similar results have been obtained in many ways, e.g., in cultured animal cells by the washout of a drug that poisons tubulin polymerization (Geuens et al., 1989; Witt et al., 1980). Earlier experiments on chromosomes in lysed cells showed that MTs formed near kinetochores, although the organization of these MTs was poor (Snyder & McIntosh, 1975). How these MTs are nucleated, how they become organized into bundles with the correct polar orientation, and how they contribute to spindle formation under normal circumstances are still unanswered questions.

A second set of activities in frog egg extract can reorganize the chromatin-initiated MTs into a bipolar array (Walczak et al., 1998). Motor-mediated MT sliding will cluster MT minus ends into one or a few foci (Gaglio et al., 1996). The long, fibrous protein, NuMA, first identified as a structural component of both nuclei and spindles, contributes to the focusing and organization of the MT minus ends as they are clustered by dynein. When both components are present in mitotic extracts, nicely ordered asters are formed with high efficiency (Merdes et al., 2000). These structures have sufficient mechanical integrity to function as spindle poles (Charlebois et al., 2011). When chromosomes or chromatin-binding structures are present, they often become arranged quite symmetrically near the midplane of a bipolar, spindle-shaped MT array (Hyman & Karsenti, 1996) (Figure 5A).

This self-organizing behavior is so strong in frog egg extracts that the chromosomes can be replaced by metal micro-beads. These nonphysiological structures become coated with chromatin, and the system forms a bipolar spindle with the beads near its midplane (Heald et al., 1996). The organization of the resulting MT array depends on the size of the chromatin-coated objects (Dinarina et al., 2009), but the noteworthy point is that a cohort of cytoplasmic activities can lead to a metaphase structure without centrosome-mediated initiation of MTs and without any specialized MTs connecting to kinetochores. Apparently the materials present in frog oocyte extracts can form a bipolar spindle simply by MT polymerization, cross-linking, and fiber rearrangement.

This acentrosomal pathway is not specific to egg cells. Most of the players in oocyte spindle formation are found in somatic mammalian cells, and extracts from cultured human cells can reproduce many of the activities described above (Chakravarty et al., 2004; Gaglio et al., 1997). Work with extracts from mitotic mammalian cells has also identified contributions from motor enzymes that are either minus-end directed (a kinesin-14 called HSET) or plus-end directed (a kinesin 5 called Eg5) (Gaglio et al., 1996). It follows that centrosome-independent pathways may contribute to spindle formation, even when centrosomes are present. This possibility is supported by the observation that centrioles can be eliminated from *Drosophila* by a loss-of-function mutation, yet a fly will form, albeit without cilia or

flagella and with somewhat disorganized mitoses (Basto et al., 2006). Centrosomes can also be ablated from cultured cells by laser irradiation, and again, spindles will form (Khodjakov et al., 2000). It is intriguing, though, that gradients of Ran-GTP are not essential for all kinds of acentrosomal spindle formation. The S-2 cell line cultured from *Drosophila* can form spindles when centrosome formation is blocked, as expected from the above. In this system, however, experiments with RNAi have shown that components of the Ran-GTP pathway can also be knocked down, and spindles will still form and function (personal communication, Sara Pereira, Institute for Molecular and Cell Biology, University of Porto, Portugal).

Thus, current evidence supports the view that cells can drive towards metaphase by multiple routes. The presence of convergent or overlapping pathways will reappear as we treat additional aspects of spindle function; apparently an important process, like chromosome segregation, is not left to a single mechanism. In the context of spindle formation this redundancy suggests that the two-fold symmetric organization of spindle components lies at a minimum on a multidimensional energy landscape. Theoretical treatments of this problem are presented and discussed below.

Acentrosomal spindles are also found in higher plants. The literature contains many careful descriptions of spindle formation in these cells, but they do not yet provide experimental detail like that mentioned above. Among the best studied plant spindles are those formed in endosperm tissue of the African Blood Lilly, *Scadoxus multiflorus*, *spp katherinae*. Form-birefringence and fluorescent tubulin have been used to localize MTs in living cells, and labeled tubulin antibodies have localized MTs in fixed material (De Mey et al., 1982). All methods reveal a sheath of MTs that assembles during prophase around the still-patent nuclear envelope. As prophase continues, the MT cylinder disappears and is replaced by MTs running roughly parallel to the axis of the former cylinder. The latter MTs emerge from areas on opposite sides of the nucleus. Simultaneously, the nuclear envelope breaks down, and the spindle proper forms (Inoue & Bajer, 1961; Smirnova & Bajer, 1994) (Figure 5B). The prophase cylinder anticipates the spindle axis, but it seems to have little to do with spindle formation proper; the nucleus-invading MTs are the principal structural intermediate in the formation of these barrel-shaped spindles.

Plants contain γ -tubulin (Liu et al., 1993; Murata et al., 2005), though accurate protein localization in these species has been plagued by cross-reacting material, putting the specificity of immunolocalization into question; γ -tubulin seems to be distributed through much of an already-formed plant spindle, nor is not obviously concentrated at the poles before spindle formation begins. Plants also contain a gene that appears to be orthologous to TPX2 (Vos et al., 2008), so the Ran-GTP pathway may play a role in the regulation of plant spindle formation. Moreover, at least some plant chromosomes can initiate MTs, as cited above in the paper by Pease, where spindle fibers grew from kinetochores after release from treatment by high pressure.

As in animal cells, plant spindles need a kinesin-5 to form and maintain a bipolar spindle, and a minus end-directed kinesin-14 helps to focus MT minus ends at the spindle poles (Bannigan et al., 2008). Given this information, it seems that acentrosomal spindle

formation in plants is more like centrosomal spindle formation in animal cells than is the acentrosomal spindle formation in oocytes.

B5. Additional Factors in Spindle MT Initiation and Organization

An additional mechanism for the initiation of spindle MTs has recently been discovered. augmin is a protein complex that binds a γ -TuRC to the wall of a preexisting MT, facilitating the initiation of new MTs along-side old ones (Goshima et al., 2008). This activity was first found in fruit flies but has now been detected in a wide range of cells. One might imagine that augmin would produce a branched MT organization, analogous to that described for actin microfilaments in the cell cortex, where the Arp2/3 complex provides a mechanically analogous initiation for these protein polymers. Branched MTs are rarely seen by electron microscopy, so there is probably a release step that follows the initiation of a new MT by augmin. Note, however, that some plant spindles contain regions where MT organization resembles a “fir tree” (Bajer & Mole-Bajer, 1986), which may be a result of augmin or augmin-like activity. Such activity is clearly important for spindle action, because augmin knock-downs produce small or ineffective spindles (Uehara & Goshima, 2010). Just why the augmin-initiated MTs make a more effective spindle is not yet clear. It may be that the accurate segregation of many chromosomes requires more MTs than can form from each centrosome, or that the MT density achieved by centrosome initiation alone is not great enough to drive chromosome motion to the metaphase plate (see below).

B6. Other fibrous components of the mitotic spindle

Studies on spindle structure and chemistry have focused on MTs since their discovery in the early 1960's. However, other fibrous materials have also been reported as spindle components and may come to be seen as important in spindle function. Spindle actin was identified by several investigators, but it seems likely that these results were misleading; a study that used phalloidin and DNAase 1 to localize both fibrous and soluble actin, respectively, found that spindles contained only the latter (Barak et al., 1981). Thus, on current evidence, actin is unlikely to play a significant role in bulk spindle function. Note, however, that some drugs thought to block actin function have an impact on chromosome-spindle interaction in some systems (Snyder & Cohen, 1995). This case is still open, but there is not yet sufficient evidence to build a model for actin's role in spindle function.

The “Nuclear-Mitotic Apparatus” (NuMA) protein is a 236 kDa coiled-coil protein discovered in the 1980's. It redistributes from the nucleoplasm during interphase to the polar regions of the mitotic spindle during mitosis (Compton & Cleveland, 1994). As described above, it contributes to the ability of dynein to organize MTs into aster-like arrays in vitro. Its contribution to spindle function depends upon extensive ADP ribosylation (Chang et al., 2005), but its precise mechanical role is not yet defined.

Six other fibrous proteins have been found in spindles: chromator, skeletor, megator, EAST, titin, and lamin-B. In spindles of *Drosophila* embryos the first four of these assume the shape of the spindle and persist in that arrangement after spindle MTs have been dissolved experimentally (Qi et al., 2004). Titin, the micrometer-long muscle protein, has been found in the spindles of insect spermatocytes (Fabian et al., 2007). Lamin-B, a component of the

nuclear lamina, helps spindles to form in frog egg extracts, and it requires RanGTP, to do so (Tsai et al., 2006). Moreover, some proteins that spend interphase in association with nuclear pore complexes (and may therefore interact with lamin-B) are important for aspects of kinetochore function (Rasala et al., 2006). Special regions of the spindle, like the “midbody” that forms at the spindle equator in late anaphase, have also been thought to include a matrix component that contributes to their form and function (Kapoor & Mitchison, 2001). Thus, there is a repertoire of protein components that may be important for spindle formation, at least in some cells.

At this stage in the history of understanding spindle function, the roles of all these non-MT proteins are hard to pin down. If spindles were to contain a fabric other than MTs, be it fibers or clusters of membrane components, these could be of great importance to spindle mechanics, e.g., by serving as a framework that binds the tails of MT-dependent motors. These enzymes could now interact mechanically with MTs and push them in directions defined by MT polarity. Such forces would be proportional to MT length, a property that would be of great value in getting prometaphase chromosomes to the spindle equator, as discussed below. It is therefore important for students of this subject to keep an open mind about the mitotic roles of structural proteins other than those directly associated with MTs. Given the lack of detail now available, however, these materials will not be considered further here.

B7. Factors controlling spindle size

Cells range widely in size, and spindles are commonly scaled appropriately. A metaphase spindle in a yeast cell is only slightly longer than 1 μm , and it contains only a few dozen MTs; the spindle in a frog zygote is $\sim 40 \mu\text{m}$ long and contains uncounted thousands of MTs. How does a cell know the appropriate size for its spindle? Some aspects of this problem are doubtless rooted in the amounts of various spindle proteins synthesized before mitosis, but more subtle and interesting aspects of the problem are now being uncovered. For example, spindles of related frog species can differ ~ 3 -fold in size. When extracts are made from these different eggs, the spindles formed *in vitro* are similar in size to the ones formed *in vivo*. When such extracts are mixed, spindles of intermediate sizes are formed, suggesting the presence of control factors in the isolated ooplasm (Brown et al., 2007). Certainly, MAPs can alter the mean length of MT populations, both positively and negatively, and spindle organizing proteins, like TPX, are also important (Bird & Hyman, 2008; Greenan et al., 2010). How cells define spindle length, when more is involved than simply the mean MT length, is an interesting unsolved problem; for reviews, see (Cai et al., 2009b; Dumont & Mitchison, 2009a; Goshima & Scholey, 2010).

B8, In vitro systems and mathematical models for mitotic spindle formation

It is well established that MTs can be organized into astral arrays simply by adding a single type of motor enzyme to a sample of Taxol-stabilized MTs. Brain kinesin-1 induces large asters in which MT plus ends are gathered at the center of the array (Urrutia et al., 1991). Such rearrangements obviously require that kinesin binds at least two MTs, but in that experiment the mechanism for such coupling was unknown. More recently, asters have been

formed by mixing MTs with artificial motors containing several kinesin heads, bound together by streptavidin-biotin linkages (Nedelec et al., 1997).

Networks of interconnected asters can form from MTs mixed with plus and minus end-directed motor enzymes (Surrey et al., 2001). These structures showed an important difference from simple asters in that many MTs ran in approximately parallel bundles from one “spindle pole” to another. In the MT bundles, polymers associated with opposite poles ran antiparallel to one another.

Although many processes are involved in spindle formation, these studies suggested that its essence can be realized by relatively simple mechanisms based on the action of one or a few motor protein. Such work has defined a question that has driven most of the theoretical research in this area: what components are necessary and sufficient to form spindle-like structure from MTs? Several mathematical models have been based simply on MTs and plus- or minus-end directed motors. In these approaches MTs are usually represented as linear, infinitely thin, polar objects that behave mechanically like inextensible elastic rods with known flexural rigidity (Nedelec, 2002). A plus- or a minus end-directed motor is modeled as an extensible molecule with motor heads on either one or both its ends. It is free to diffuse then bind and move along the MTs. Binding and unbinding are stochastic events governed by probabilities. Many thousands of motors are usually allowed to work simultaneously. Motor complexes with both motor heads bound exert forces on the MTs according to Hooke’s law, i.e. the amount of force they generate is proportional to the extent of molecule stretching. When a head is bound, it moves along the MT at a velocity that depends inversely on the force that opposes its movement, and the relationship is usually considered to be linear, dropping to zero at the “stall force”.

In Table 1 we summarize two important models that exemplify the formation of spindle-like structures by two distinct processes: from two MT asters (Nedelec, 2002) or from randomly distributed MTs interacting with an array of beads covered with plus end-directed motors (Schaffner & Jose, 2006).

The first model presents a simple and elegant mechanism for spindle formation, but it requires a “heteromotor” – a protein with plus end-directed motor activity at one end and minus end-directed activity at the other. Such an enzyme has not yet been found in nature. Many cells contain kinesin-5s that crosslink MTs and slide them relative to one another, but both motor heads are plus end-directed. Interestingly, however, recent research suggests that the directionality of Cin8, a yeast kinesin-5, will change when the motor is working at low density (Roostalu et al., 2011). This capability has not yet been incorporated into a model for spindle formation.

Civelekoglu-Scholey and colleagues (Civelekoglu-Scholey et al., 2011) have found that a steady-state spindle length can be established and maintained with known motor enzymes by opposing the activities of kinesin-5s and kinesin-14s. This result expands on Nedelec 2002, although stable spindle lengths were found at steady-state within only a narrow, albeit physiological, range of motor concentrations, and the MTs had to be very dynamic. Such a regime might have been overlooked in Nedelec’s work, since they also considered forces

due to MT bending that are not considered in (Civelekoglu-Scholey et al., 2011). The range of motor concentrations that is successful in defining a stable spindle can be extended by adding an elastic matrix that surrounds the spindle, e.g., one that might come from one of the spindle matrix components mentioned above.

The two papers in Table 1 are good examples of models in which a single mechanism of spindle formation is extensively characterized by computer simulations, based on reasonable premises and initial conditions. They describe plausible mechanisms for spindle formation *in vivo*, but as mentioned earlier, a cell is likely to use several mechanisms at one time. This complexity may provide accuracy, reliability and robustness for such an important process. A model that incorporates several mechanisms has not yet been created, but obviously models based on single mechanisms are a necessary first step: one needs to know how the separate parts of a system will function before modeling the complexity of the whole machine.

More models for spindle formation have recently been developed (Burbank et al., 2007; Channels et al., 2008; Loughlin et al., 2010). These improve our understanding of details in the mechanisms outlined above, but in our opinion, they do not add significant new insight to the mechanisms for spindle formation. For a review of these papers, see (Mogilner & Craig, 2011).

Available models do not, however, address the roles of other components in spindle assembly. An important question for future work is the extent to which chromosomes and kinetochore MTs contribute to the organization of a bipolar structure. Perhaps the stabilization of MTs by kinetochores would make the presence of the speculative “heteromotor” unnecessary. This kind of question can be addressed mathematically; it is a good example of the ways in which further theoretical work could make a significant contribution to spindle research. An important next step for the theoretical approach to spindle formation should involve the development of more complex models that helped us to evaluate the contributions of each component to success in the overall process.

C. ATTACHING CHROMOSOMES TO SPINDLE MTs

C1. Introduction

There are diverse pathways by which the MTs of a forming spindle can get access to condensing chromosomes. When centrosomes are present, they commonly lie near the nucleus, so as the envelope disperses, the MTs naturally encounter the chromosomes; attachment can ensue. In “closed” spindles, which form within a nuclear envelope, the passage of soluble spindle components through the nuclear pores and their subsequent assembly within the nucleus makes spindle-chromosome interaction inevitable. In cells without centrosomes the chromosomes participate in MT initiation, so interaction of some kind is immediate. The key to understanding chromosome-MT interaction, then, is not to see how interaction happens but to learn how sister kinetochores bind to MT associated with sister poles. We start on this question with spindles that employ centrosomes, then tackle other systems.

C2. Random encounters

Centrosome-initiated MTs radiated from two nearby foci, so some of their plus ends inevitably enter the chromosome mass. Some dynamically unstable MTs will therefore encounter kinetochores by chance, allowing an attachment that can develop into a stable, kinetochore-MT connection without any significant chromosome motion. There is considerable cell biological evidence that MTs ending on kinetochores are stabilized by this interaction. For example, kinetochore-associated MTs (KMTs) are more stable than other spindle MTs to subunit dilution or to treatment with cold. Moreover, the $t_{1/2}$ for spindle MTs in general is ~ 30 s (Saxton et al., 1984), while that for KMTs is >200 s (Gorbsky et al., 1988). Thus, the latter will naturally persist while other MTs continue rapid turnover, leading to a gradual selection of MTs that are kinetochore-bound. Experiments with isolated chromosomes have shown that the affinity of mammalian kinetochores is greatest for MT plus ends (Huitorel & Kirschner, 1988), so this attachment is in some sense a minimum energy configuration. In short, chromosomes “select” the MTs that happen to encounter their kinetochores by chance (Kirschner & Mitchison, 1986).

This insightful idea is widely accepted in principle, but a close look at its predictions has identified problems. Quantitative analysis predicts that the time required for dynamically unstable MTs to associate correctly with all kinetochores (92 in a human cell, more in some species, less in others) exceeds the duration of prometaphase, suggesting that other mechanisms are at work (Holy & Leibler, 1994). These issues are discussed more thoroughly below, but one uncertainty in such work is that one does not know the functional size of a kinetochore. Electron microscopy reveals that kinetochores on chromosomes not yet attached to the spindle include a “corona” of fibers, composition not yet identified, that extends from the obvious kinetochore (McEwen et al., 1998). This specialization probably increases the functional size of a kinetochore and increases the probability of kinetochore-MT interaction. Another ameliorating factor is the previously mentioned initiation of MTs by kinetochores, a phenomenon that has now been seen in yeast (Tanaka et al., 2005), mammals (Maiato et al., 2004), and elsewhere. These MTs may further extend the function of the corona. There are, however, unanswered questions about how the polarities of kinetochore-initiated MTs conform to the pattern of MT polarity seen in fully-formed spindles: plus MT ends at kinetochores (Euteneuer & McIntosh, 1981). In vitro the MTs that grow at kinetochores are poorly organized (Witt et al., 1980). In vivo the polarity of kinetochore-initiated MTs in fruit flies are normal (Maiato et al., 2004), but in budding yeast they are opposite (Kitamura et al., 2010). How these MTs interact with centrosome-initiated MTs has not yet been determined, nor is it known whether they become part of the metaphase kinetochore-centrosome connection or are simply transient structures. Nonetheless, one can imagine that either polarity of MT may accelerate the association of kinetochores with centrosome-initiated MTs.

C3. Chromosome repositioning to increase the probability of MT encounters

The kinematics of unusual chromosomes during their attachment to the spindle is revealing about additional pathways for the chromosome attachment process. While most chromosomes in both vertebrates and plants become attached to the spindle at or near the places where they lay at the time of spindle MT initiation, a few chromosomes move

poleward at speeds that exceed the rate of anaphase by >10-fold. This rapid movement appears to be the result of one kinetochore interacting with the walls of spindle MTs, rather than with their ends (Alexander & Rieder, 1991). In animal cells the motion is probably driven by the motor enzyme dynein (Yang et al., 2007b), while in budding yeast it is a result of a Kinesin-14 (Tanaka et al., 2005). Such motility is most commonly seen with chromosomes that failed to form a bipolar or “amphitelic” attachment at the onset of spindle formation, or that lay by chance at a significant distance from the spindle as it formed. It appears to be a process by which cells can bring these errant chromosome to a place where there are many short MTs growing and shortening by dynamic instability, an environment that makes proper, end-on attachments to kinetochores more likely.

When kinetochore-MT attachment is initially with the MT wall, the chromosome will commonly change to the geometry that is more common in true metaphase: MTs ending at the kinetochore. This transition has not yet been studied well, but it probably involves the replacement of at least some lateral connections by end-on ones, given all the dynamic MT plus ends in the spindle and the greater strength of kinetochore bonds to MT plus ends. It may also involve the plus end-directed motor activity that is characteristic of most kinetochores, an action that should facilitate the formation of kinetochore – MT plus-end connections. One can imagine that MT severing activities are also involved, but none of the relevant AAA ATPases has yet been localized to kinetochores. Thus, there is an important transition in the formation of proper chromosome attachments that remains to be fully understood. While some chromosomes move normally in anaphase when their kinetochores appear still to be associated with the walls of at least some MTs, end-on attachment seems to be preferred, and recent descriptive work is elucidating the transition by showing that individual chromosomes are not stably attached to the spindle until late prometaphase (Magidson et al., 2011).

C4. Attaching sister kinetochores to sister spindle poles

An important problem in the formation of proper chromosome-spindle connections is the establishment of “amphitelic” attachment, i.e., the association of sister kinetochores with sister poles (compare Figures 2A and 6A). Careful observation has shown that attachment errors are common in early prometaphase, but most such mistakes are corrected by the time of anaphase onset. Three kinds of incorrect attachments have been identified: 1) monotelic, in which only one kinetochore has become associated with MTs, and those MTs all come from a single spindle pole; 2) syntelic, in which both kinetochores are attached to MTs associated with a single pole; and 3) merotelic, in which one kinetochore is associated with MTs from both spindle poles. Since incorrect chromosome attachment can lead to chromosome loss, a serious medical issue, the mechanisms of error correction are a subject of intense current study; they are now being reviewed on their own quite frequently (Gregan et al., 2011; Matos & Maiato, 2011), so they will not be treated here in detail. Nonetheless, we can say that proper amphitelic attachments are favored by several redundant mechanisms; all may participate in achieving this essential state, perhaps with different importance in different types of cells.

- a. Kinetochores on mitotic chromosomes commonly face in opposite directions, thanks to the processes of chromosome replication and condensation; this places

sister kinetochores back-to-back (Figures 2a, 3a). If one kinetochore is associated with MTs from the east pole, its sister is likely to associate with MTs from the west.

- b.** MT-kinetochore attachments are initially unstable, but when they come under tension, their stability increases. This was first demonstrated experimentally by micromanipulation of large meiotic chromosomes in grasshopper spermatocytes (Nicklas et al., 1982) and has been confirmed by a variety of less direct experiments in other systems. For example, when metaphase spindle MTs are stabilized and elongated by the addition of Taxol, some of the chromosomes detach from the spindle (De Brabander et al., 1986).

Tension on sister kinetochores is a result of each MT's kinetochore attachment site experiencing a pole-directed force as soon as MT binding has occurred. Evidence for this force is seen both in the afore-mentioned migration of a monotelic chromosome toward the pole to which it is attached and in the stretching apart of sister kinetochores of amphitelic chromosomes (Nicklas, 1997). A dramatic example of stabilization associated with stretching is seen in the attachment of chromosomes to pole-initiated MTs in diatoms; here chromosomes are commonly monotelically attached in early prometaphase, whereupon they move irregularly toward and away from the pole to which they are attached. As soon as amphitele is achieved, each chromosome becomes so stretched that its kinetochores reach the spindle poles while the chromosome arms are still at the metaphase plate (Pickett-Heaps et al., 1980). Experimental evidence for pole-directed forces is also seen in the response of a metaphase chromosome to micro-beam ablation of one of its kinetochores; the chromosome immediately starts moving toward the pole attached to the unirradiated kinetochore (McNeill & Berns, 1981). Thus, an amphitelic chromosome, which is properly attached to the spindle, is under tension; monotelic or syntelic attachments should experience less of this kinetochore-associated force. The proper arrangement of chromosomes is selected for, not directed, because improper connections are unstable while the right ones persist (Kirschner & Mitchison, 1986).

The correction of monotelic errors seems to be simply a matter of time. So long as the cell does not start anaphase too soon, the unattached kinetochore is likely to encounter dynamic MTs growing from the opposite pole and make the connections that will establish amphitele. Correction of synteles, on the other hand, requires the release of MTs associated with one or the other of the two kinetochores. The lack of tension at the kinetochores offers a plausible explanation for this instability, but now we must answer two questions: how is tension generated, and how does it promote stability? Tension generation is discussed in our section on anaphase, but one part of the mechanism has a simple, plausible solution: kinetochores bind minus end-directed motors, like dynein in vertebrates and a kinesin 14 in fungi. These enzymes probably interact with MT walls and pull the attached kinetochore poleward. If a chromosome has made an amphitelic attachment, it is pulled toward both poles at once, generating tension at the kinetochore; in synteles there is no obvious antagonistic force.

How does such tension increase the stability of MT binding? A priori, one might expect the opposite effect, since tension should make it easier for thermal motions to break the MT-kinetochore bond. The true mechanism(s) for tension-induced stability is not yet known, but

several plausible ideas have been proposed. For example, if a minus end-directed motor on one kinetochore is trying to walk toward one pole, but the sister kinetochore is similarly engaged, the motors will stall. It is easy to imagine that this will leave the enzymes in a transition state from which their MT dissociation rate is low. Thus, strain-dependent off-rate constants could contribute stability to the kinetochore-KMT connection (McIntosh et al., 2002).

A second plausible mechanism for tension-induced stability is based on the assumption that KMTs are attached to the kinetochore by their plus ends, so chromosome motion is coupled to tubulin polymerization and depolymerization. In vitro a catastrophe, which initiates rapid MT shortening, is accompanied by outward flaring of the strands of tubulin from which the MT wall is composed, the “protofilaments” (Mandelkow et al., 1991). The bent configuration of tubulin is associated with its having GDP bound, while GTP promotes a straighter tubulin dimer and hence straighter protofilaments (Wang & Nogales, 2005). If GDP tubulin must bend into its lower energy structure to manifest its high rate of dissociation from the MT, then anything that keeps protofilaments from bending will promote MT stability. A kinetochore-MT attachment to protofilaments would allow tension to inhibit protofilament bending and thus inhibit tubulin depolymerization (Grishchuk, 2009).

Another mechanism that contributes to the instability of syntelic attachment is based on the protein kinase, aurora-B, which is localized in the vicinity of the centromere. Its substrates include proteins that help to bind KMTs to kinetochores, such as the hetero-tetrameric complex, Ndc80. Protein phosphorylation by aurora-B reduces the affinity of this (DeLuca et al., 2006) and related proteins to the negatively-charged surface of a MT, loosening the kinetochore-KMT connection, thereby giving the chromosome another chance to find MTs that are associated with the right pole. Indeed cells deficient in this kinase make a higher-than normal number of errors in chromosome attachment, while the stability of kinetochore-KMT attachments is enhanced by at least one protein phosphatase, PP1 (Francisco et al., 1994). How, then, are these activities regulated so only inappropriate connections are released? A clever model has suggested that the tension of a bipolar attachment pulls the MT-binding substrates away from aurora kinase, allowing the phosphatase to increase connection stability (Lampson & Cheeseman, 2010; Tanaka et al., 2002) (Figure 6B). Direct evidence for this proposal has been obtained by fluorescence resonance energy transfer (FRET) (Fuller et al., 2008). It is noteworthy, however, that careful microscopy has shown that the strain important for a chromosome being ready for segregation is not an increased distance between sister kinetochores, which is comparatively easy to measure; it is strain within a kinetochore that tells the cell that this chromosome is properly attached (Khodjakov & Pines, 2010; Maresca & Salmon, 2010) (Figure 6C).

Regardless of mechanism, all these corrections and retrials take time. A few species get extra time by starting to establish bipolar spindle attachments rather during S-phase. Careful microscopy of budding yeast nuclei has shown that during G1, the centromeres are all close to the cell's single, nuclear envelope-associated centrosome. As DNA replication begins, they are released into the nucleoplasm, probably as the centromeric DNA of that chromosome is replicated (Kitamura et al., 2007). At about this time, the centrosomes also

duplicate, providing the two MT initiating structures that will become the spindle poles. Attachment of the sister kinetochores that form on the duplicated centromeres can now occur with MTs growing from the two centrosomes, and they have the rest of the cell cycle to get it right.

In budding yeast, however, as well as in cells that don't start chromosome attachment until prophase (usually hours after DNA is replicated) there is a control mechanism called the "spindle assembly checkpoint" (SAC), which detects unattached kinetochores and signals to delay anaphase onset until proper attachment has been achieved (Hoyt et al., 1991; Li & Murray, 1991). This function is discussed below.

Correction of merotelic attachments is more problematic because tension between multiple MT attachments is within a kinetochore, not across the centromere (from one kinetochore to its sister). These errors lead to one chromatid being pulled in two directions at anaphase, which can leave that chromosome near the spindle equator as other chromosomes segregate, so these errors too must be corrected. Again aurora B kinase is involved (Cimini et al., 2006), but just how these corrections are achieved is still uncertain. One possibility is that tension within the plane of a kinetochore activates or inactivates factors different from tension perpendicular to the kinetochore's plane (McIntosh & Hering, 1991), allowing the necessary activities to loosen the KMT-kinetochore connection, offering another chance for correct attachments.

C5. Molecular constituents of kinetochore/spindle interaction

Our knowledge of the molecules involved in establishing a proper kinetochore-MT connection has grown tremendously in the last several years. Many kinetochore components have been identified by a fruitful combination of genetics, immunology, and proteomics with molecular and cellular biology. This field is now so big that it requires a review of its own, and indeed there are many fine papers that catalogue the protein components of kinetochores in different organisms. An exhaustive list of even the reviews would take a lot of space. Here we draw attention to readable compendia on budding yeast (McAinsh et al., 2003) and vertebrates (Walczak & Heald, 2008; Wan et al., 2009; Welburn & Cheeseman, 2008).

In short, a kinetochore is built upon special DNA, the centromere, which is defined in part by an AT-rich nucleotide composition (albeit, in budding yeast there is a specific centromere sequence, and in other organisms there are DNA sequences that turn up frequently). Some organisms show clear evidence that the DNA in this region is modified, e.g., by methylation, leading to a mark upon the DNA that is inheritable (an "epigenetic" trait). This genetic locus becomes wound around nucleosomes that contain a centromere-specific isoform of histone H3, commonly called centromere protein A (CENP-A). The loading of these nucleosomes occurs during S-phase and requires a special set of protein factors (Cleveland et al., 2003). Chromatin that contains CENP-A can in turn bind additional proteins that form the "inner centromere", a platform that binds additional protein complexes that form microtubule binding sites. The latter include motor enzymes: kinesins that walk toward the MT plus end (the kinesin 7 called CENP-E in mammals or kinesin 8s in both mammals and yeasts), one or more kinesin that promotes MT depolymerization (kinesin 13s in mammals, a kinesin 14

in yeasts, as well as kinesin 8s in both yeasts and mammals), and motors that walk toward the MT minus end (cytoplasmic dynein 1 in vertebrates and kinesin 14s in yeasts). When dynein is present, it is accompanied by dynactin and bound to the kinetochore through a chain of additional cofactors.

There are also at least three kinetochore protein complexes that bind MTs but are not motors: the highly conserved Ndc80 complex, which is sometimes joined by KNL1 and other proteins to form the “KMN network” (Cheeseman et al., 2008). Mammalian kinetochores contain the SCA complex of three MT-binding proteins (Welburn et al., 2009) as well as the long fibrous protein CENP-F (a.k.a. mitosin) (Feng et al., 2006). In yeasts there is a hetero-decameric protein complex, called Dam1 or DASH, which can oligomerize into rings or partial rings on the MT surface (Miranda et al., 2005; Westermann et al., 2006).

Kinetochores also accumulate most of the proteins that associate with the plus ends of growing MTs, the Tip Associated Proteins or “TIPs” (EB1, Klp170, CLASP, XMAP215, a.k.a. Tog, etc.). These are kinetochore-associated only when MTs are present, so they are probably best regarded as a special class. However, some EB1-associated proteins may be of particular importance for the formation stable kinetochore-MT connections, e.g., the formin mDia3, which binds MTs in a phosphorylation-sensitive manner (Cheng et al., 2011). The presence of this actin “polymerase” at kinetochores is particularly intriguing, given that XMAP215 has been shown to serve as a MT polymerase with an activity on tubulin that is analogous to formin’s activity on actin (Brouhard et al., 2008; Widlund et al., 2011). Finally, kinetochores bind several enzymes that regulate the phosphorylation of other proteins: the Aurora B kinase and its co-factors, several kinases that participate in the SAC (Mad1, Bub1, BubR1/Mad3, and MPS1), and some critically important phosphatases (PP1, PP2A).

C6. Chromosome attachment to spindles without centrosomes

As described above, acentrosomal spindle formation in higher plants resembles the same process in centrosomal spindles, so the processes of chromosome attachment are probably quite similar too. However, there is not enough information about these events to treat them independently here. Spindle formation in vertebrate oocytes, on the other hand, is clearly distinct, and issues of kinetochore-MT interaction must be addressed.

The RanGTP exchange factor, RCC1 appears to be associated with the entirety of a chromosome (Carazo-Salas et al., 1999), suggesting that Ran-mediated MT formation will occur all over every chromosome. How, then, do MTs become associated with kinetochores to make functional spindle attachments? This question has not yet been answered by students of this pathway. One way to think about the problem, though, is to assume that additional properties of the kinetochore bias the location of MT initiation or subsequent chromosome binding. For example, the localization of aurora B kinase to the centromere may activate special processes there, and indeed there is evidence that a gradient in aurora B activity spreads out over the chromosome for surprising distances (Welburn et al., 2010). Another possibility can be developed from the fact that kinetochores have a higher affinity for MT plus ends than other MT parts (Huitorel & Kirschner, 1988), so short MTs that grow from initiation all over the chromosomes may be captured at the kinetochore as they diffuse locally. There is ample evidence that MTs bound to kinetochores by their plus ends can

elongate by the addition of tubulin at the kinetochore, accompanied by a motion of the MT away from the kinetochore, both in vivo (Mitchison & Salmon, 1992) and in vitro (Mitchison & Kirschner, 1985). If this process is even slightly faster than normal MT growth, it will make KMTs grow in preference to other MTs. The kinetochore-distal ends of these polymers may then be clustered by the same minus end-directed motors known to be important for the bunching of other MTs to form asters.

Two problems in evaluating this and related proposals are the size of the spindles that form in frog ooplasm, both in vivo and in vitro (which makes detailed microscopy difficult) and the lack of an obvious kinetochore structure on chromosomes condensed in these extracts. Nonetheless, there are many students of this system, so we can hope that the necessary information will soon be forthcoming.

C7. Models for chromosome attachment

There are two important aspects of MT-kinetochore attachment that can be analyzed mathematically: 1) how to attach all kinetochores to MT within a reasonable time, and 2) how to make these attachments amphitelic. As sketched above, a successful idea for how the tips of MTs growing from centrosomes might find randomly distributed kinetochores was introduced by Kirschner and Mitchison in 1986; they called it “search-and-capture”, reviewed in (O’Connell & Khodjakov, 2007). This idea has been analyzed quantitatively, as summarized in Table 2. Dynamically unstable MTs are indeed effective in searching for kinetochores, compared to polymers with a single state dynamic; rapid transitions between periods of growth and shortening reduce search-and-capture times 1000-fold (Holy & Leibler, 1994). However, this work left an important question open: can this mechanism provide realistic times of attachment for tens of kinetochores distributed through 3D space? This question has now been addressed (Paul et al., 2009; Wollman et al., 2005). The more recent of these models presents an advanced version of the earlier one, and it is summarized in Table 2. With some added model features search-and-capture works for a realistic number of chromosomes within an appropriate amount of time, so this result might be considered well established.

It is worth noting, however, that both these analyses have focused on kinetochore interactions with MT ends. Electron microscopy of prometaphase spindles in organisms with obvious kinetochores has shown that initial kinetochore-MT interactions are more commonly with MT walls (Roos, 1973). Moreover, the most thorough study of a single kinetochore’s attachment to spindle MTs in vivo found lateral attachment as the first stage in the process (Alexander & Rieder, 1991). If the average length of dynamically unstable, centrosome-initiated MTs is greater than the average kinetochore-centrosome distance, then many MTs will present their sides to the randomly oriented kinetochores, greatly increasing the interaction surfaces that a kinetochore and its corona might encounter. This consideration will reduce the predicted time for attachment of all kinetochores to randomly growing, dynamically unstable MTs and may reduce the apparent need for additional mechanisms, such as the features added in Paul et al. (2009).

Paul et al. address the important issue of correcting errors in chromosome-spindle attachment. They assumed that syntelic attachments dissolve within seconds, following

initial capture by sister kinetochores of MTs from the same pole. Amphitelic attachments, on the other hand were postulated to be secured and able to block further captures. This mechanism improved the number of correct attachments, but the number of merotelic attachments remained great. Two possibilities for further correction were considered. First, merotelic attachments were dissolved over time, just like syntelic attachments. That brought the number of incorrect attachments down significantly, but it increased the requisite time. A second suggestion was that following capture of one MT, the chromosome rotates so the kinetochore with one MT attached now faces the pole from which that MT grew. Now its sister faces away from that pole. Moreover, chromosome arms and kinetochores were presumed to be impermeable to MTs, which impeded future attachments of the rotated kinetochore to MTs associated with the pole away from which it had turned. This mechanism proved efficient in eliminating merotelic attachments. Together with the rapid dissolution of syntelic attachments, this hypothesis brought correct capture of all kinetochores to a realistic time.

As described above, the correction of syntelic attachments depends on aurora B and is sensitive to tension at the kinetochores (Lampson & Cheeseman, 2010). This tension in turn depends of the number of the MTs attached per kinetochores. A fair estimate of the time required to correct all syntelic attachments will require developing a mathematical model that includes an aurora B-dependent mechanism. Whether this mechanism can explain the dissolution of syntelic attachments within seconds, as suggested in Paul et al., is an important question that should be addressed. Correcting merotelic attachments is even more complex. Chromosomes do not appear to be impermeable to MTs, as suggested in Paul et al., 2009; MTs can clearly penetrate chromosome arms (McIntosh & Landis, 1971), though the frequency of this event is unknown and penetration of kinetochores occurs only under non-physiological conditions (Euteneuer & McIntosh, 1981). This issue requires further experimental and theoretical investigation.

D. GETTING PROMETAPHASE CHROMOSOMES TO THE SPINDLE EQUATOR

D1. Introduction

Most mitotic cells take the time (and expend the energy) to cluster their soon-to-be-segregated chromosome on the spindle equator before starting anaphase, the process called “congression”. Given that precision is all important for successful mitosis, putting all the moving objects on the same starting line makes sense. Budding yeast has been said to lack these motions (Straight et al., 1997), but here chromosome stretch and motility are sufficient that the initial observation may not be right (Goshima & Yanagida, 2001). Congression is certainly common, even if not universal. Indeed, a metaphase condition will evolve over time, even with unusual meiotic chromosomes that contain an imbalance of kinetochores (Nicklas & Arana, 1992), suggesting that the gathering of chromosomes on the spindle midplane has adaptive value. Here, we consider both descriptive and experimental evidence about how congression works, then we describe models that have been proposed to explain it.

D2. Descriptions of chromosome congression

Many chromosomes are situated about midway between the poles as the spindle begins to form, so little or no congression is necessary. Chromosomes that are not near the midplane when they become bioriented, undergo a series of motions that result ultimately in their displacement to the spindle equator, producing the metaphase configuration. This is the essence of congression. These motions are along, not across nearby spindle MTs, so a bioriented chromosome is always moving one kinetochore toward the pole it faces while the sister kinetochore is moving away. Since there are pole-directed forces acting on each kinetochore, one could imagine that congression results from an imbalanced tug-of-war in which the kinetochore more distant from its pole can pull harder than its sister, leading to the metaphase condition. Early models for congression were built on just this hypothesis (Ostergren, 1945; Ostergren, 1951; Ostergren, 1961), but careful observations in a variety of natural and experimental conditions have shown that this appealing idea is not sufficient to account for the available data.

When an amphitelic chromosome is moving, it must be experiencing an imbalance in non-viscous forces; this could result from inequality in either pulling or pushing, making it difficult to interpret observations *in vivo*. However, mono-oriented chromosomes on either a monopolar or a bipolar spindle also display motions toward and away from the pole with which they are associated, demonstrating that a single pole can both pull and push. Presumably the poleward (P) or antipoleward (AP) motions of chromosomes are the result of imbalance in these forces. A hint about these forces comes from a very simple observation; when near one pole or on a monopolar spindle, chromosomes acquire a distinctive shape; while moving P, the centromere leads and the chromosome arms lag behind (Ostergren & Bajer, 1961), a structure consistent with the previously-described pull at the centromere. While moving AP, however, this chromosome shape is largely maintained; the chromosome arms lead, suggesting that they are responding to a push (Rieder & Salmon, 1994).

The complexity of prometaphase chromosome motions demonstrates, however, that a simple balance of pushing and pulling is not sufficient to explain congression. In many cell types, prometaphase chromosomes oscillate, a process called “directional instability” (Skibbens et al., 1993): both kinetochores switch within a very few minutes between P and AP motions. The rates of these movements are approximately constant ($\sim 1\mu\text{m}/\text{min}$) and depend neither on the direction of movement nor position in the spindle. Net displacement to the spindle midplane is achieved by changes in the durations of P and AP movements that bias chromosome position toward the equator (Ke et al., 2009; Waters et al., 1996). A recent, more quantitative analysis of prometaphase chromosome oscillations in human cells shows that the speeds of oscillation and of fluctuations in sister kinetochore separation during late prometaphase and metaphase are set by microtubule depolymerases, e.g., a kinesin 13, whereas the periods of these motions depend on the stiffness of the mechanical linkage between sisters. Metaphase plates become thinner as cells progress toward anaphase as a result of reduced oscillation speed at a relatively constant oscillation period (Jaqaman et al., 2010).

The cohesin links between sister chromatids are stiff, but chromatin itself can be stretched; its elastic modulus in grasshopper spermatocytes is 4.3×10^2 Pa (Nicklas, 1983). Chromatin deformation can therefore be used as a measure of tension at kinetochores and an internal standard for mitotic forces (Nicklas, 1988). Subsequent work has defined the mechanical properties of chromatin even better (Poirier et al., 2000). The distance between sister kinetochores during congression is greater than that seen on a resting chromosome (Waters et al., 1996), indicating that oscillating sister kinetochores are under tension most of the time. The distance between sister kinetochores on mono-oriented chromosomes is again greater than that seen on unattached chromosomes, albeit to a lesser extent than with bi-oriented chromosomes. The tension on a monotelic or syntelic chromosome cannot be generated by forces on opposing kinetochores, confirming that there are other forces acting on chromosomes.

Chromosome motion during congression is coupled to the dynamics of the KMTs. In mammalian cells there are >6 MTs per kinetochore and in some plants the number of KMTs exceeds 100 (Jensen & Bajaj, 1973). The switch in direction of chromosome direction implies that all these KMTs reverse their polymerization states essentially simultaneously. Moreover, when a chromosome shows directional instability, the KMTs attached to sister kinetochores must also be coordinated, so the switch from polymerization to depolymerization on one kinetochore is matched by the opposite change on its sister (Skibbens et al., 1993). The regulation of these switches of MT dynamics is an important problem that remains poorly understood.

D3. Perturbations of congression reveal the forces that drive the process

Ostergren suggested that the forces pulling on chromosomes were proportional to spindle fiber length (MTs were unknown at the time) (Ostergren, 1945). With the discovery of spindle MTs, this idea was reformulated in terms of the length and numbers of MTs attached to each kinetochore. Experiments with grasshopper spermatocytes in meiosis undertook to change these parameters, either by making multivalent chromosomes (Hays et al., 1982) or by irradiating one side of a chromosome with a focused laser beam (Hays & Salmon, 1990). The results showed that when the number of KMTs changed, the chromosome shifted to a new equilibrium position with more short MTs balancing a few longer ones - direct support for Ostergren's hypothesis.

Data now available show, however, that this hypothesis is not sufficient to account for congression. A model based only on poleward pulling forces predicts that chromosomes on monopolar spindles should move in close to the pole. However, chromosomes on monopolar spindles behave much like those on a bipolar spindle: after a chromosome has moved close to the single pole, it switches direction and starts moving AP (Cassimeris et al., 1994), confirming that a single pole can exert AP forces. Moreover, chromosomes on a monopolar spindle oscillate about a position a few micrometers from the pole, with the chromosome arms projecting away from the pole (Rieder et al., 1986). To characterize these pushing forces Rieder et al. cut the arms of mono-oriented chromosomes with a laser micro-beam. The arms (chromosome fragments that lacked centromeres) moved AP to the periphery of the cell, while the centromere-containing fragment moved closer to the pole. These AP

movements required non-KMTs (Ault et al., 1991) and are said to result from “polar ejection forces” (PEFs). The strength of PEFs, and thus the steady-state position of a monotelic chromosome relative to the pole, seems to depend on the density of the non-KMTs (Cassimeris et al., 1994). Electron microscopy of spindle cross-sections indicates that MT density is largest near the poles and decreases towards the spindle midplane (McIntosh & Landis, 1971), so AP forces probably depend on position along the spindle axis. This provides an attractive mechanism for chromosome positioning at the spindle equator.

Two mechanisms for PEFs are currently considered. Growing MTs can generate pushing forces of up to 5pN per MT by rectifying Brownian fluctuations (Laan & Dogterom, 2010). Dynamically unstable, pole-initiated MTs could therefore push on chromosomes, simply by growing into the chromatin. Another possible mechanism is based on the action of kinesins that interact with chromatin. There are two of these “chromokinesin”: a kinesin-4 and “Kid”, which is a kinesin-10. These plus end-directed, MT-dependent motor proteins associate with chromosome arms (Antonio et al., 2000; Wang & Adler, 1995) and could obviously engage with spindle MTs to generate forces pushing away from a nearby pole.

One PEF induced by a chromokinesin has been better characterized (Brouhard & Hunt, 2005). In these in vitro experiments Kid contributed to MT binding with chromosome arms, resulting in ATP-dependent movement of MTs relative to chromosomes. The forces generated by Kid were 2 – 3 pN per MT, and the movements showed a peak velocity of $3.6 \pm 2.4 \mu\text{m}/\text{min}$. These movements included frequent stalls and backward slippages, behaviors consistent with experiments in which Kid-MT interactions were characterized directly (no chromosomes present) by optical trapping (Yajima et al., 2003).

The sufficiency of chromokinesins for generating PEFs and the ability of these forces to explain congression have been examined by several kinds of experiments. When Kid was depleted from extracts of frog eggs and spindles were induced, the chromosome arms became disorganized and stretched along the length of the entire spindle, suggesting that PEFs were significantly reduced (Antonio et al., 2000; Funabiki & Murray, 2000). Nonetheless, the chromosomes congressed. Similar results were obtained in cultured human cells; when Kid was depleted by RNAi (Tokai-Nishizumi et al., 2005) or inactivated by the microinjection of antibodies (Levesque & Compton, 2001); chromosome arms became misaligned, but congression and cell cycle progression were unaffected. Interestingly, however, chromosome oscillations were suppressed.

To find out whether pushing forces can be generated at kinetochores distinct from chromosome arms, congressing chromosomes have been severed between sister kinetochores (Khodjakov & Rieder, 1996). If the cutting was done when the chromosome was close to the spindle equator, the P moving kinetochore always continued its motion, while its AP moving sister always stopped, suggesting that it was not experiencing a push. However, when a similar experiment was carried out on a prometaphase chromosome situated $<3 \mu\text{m}$ from the pole, the AP moving kinetochore could either stop and switch to moving P, or it could continue moving AP. These results suggest that the behavior of kinetochores depends in significant ways on the density of MTs in their environment.

The observations summarized above have led to an hypothesis in which congression occurs as a result of both P and AP forces applied to chromosomes from both poles (Rieder & Salmon, 1998). Each kinetochore is presumed to be connected to MTs from only one pole, and these are postulated to generate position-independent pulling forces. Sister kinetochores are attached to sister poles, so the pulling forces normally balance. AP forces (the PEFs) are exerted on the whole chromosome, and since these depend on the number of interactions with MT from each pole, they are position-dependent and drive congression. This mechanism can explain congression in many organisms and under many conditions, but the situation in vivo appears to be even more complex, so there may be additional pathways by which chromosomes can congress.

As an example of this complexity, a kinetochore-localized kinesin-7, CENP-E, has been identified as another factor important for congression. Immunodepletion of CENP-E arrested cells in mitosis, and eventually they died by apoptosis (Schaar et al., 1997). These cells contained some chromosomes that lay near the poles, having failed either to congress or to establish bipolar connections. Chromosomes that did biorient in these cells oscillated, but some of these failed to align on the metaphase plate, suggesting that CENP-E plays a role in congression (Putkey et al., 2002; Wood et al., 1997). The majority of chromosomes did, however, congress properly, so CENP-E cannot be essential for the process.

It has recently been shown that even bi-orientation is not essential for congression. Chromosomes with a mature attachment to only one pole can attach laterally to the KMTs of other chromosomes and move to the spindle equator by sliding along these fibers (Kapoor et al., 2006). Notably, when CENP-E was depleted by RNAi only 15% of such chromosomes reached the equator, versus 73% in control cells; apparently CENP-E is the driver for this novel congression pathway.

D4. Models for chromosome congression

Aspects of congression would be easy to understand if KMTs generated pole-directed forces proportional to their length and number, and several models have been suggested for how this might be accomplished. For all such models, one must postulate that KMTs can lose subunits at or near the spindle poles, but since the minus ends of at least some KMTs are uncapped by metaphase (O'Toole et al., 2003), this proposal is plausible.

In addition, a motor must interact with KMTs along their lengths and push on some sort of support, e.g., spindle membranes or a matrix component. By this hypothesis a plus end-directed motor whose head binds any spindle MT and whose tail binds the secondary structural component would exert a net pole-directed force on each kinetochore that is proportional to the length and number of its associated KMTs. Now the spindle equator would be a minimum energy position for the chromosomes. An alternative to this idea can be based on hypothetical interactions between KMTs and nonKMTs. Since these two MT classes in a half-spindle are parallel, generating forces through their interaction might seem to be impossible by symmetry (McIntosh et al., 1969). However, spindle MTs almost certainly resemble cytoplasmic MTs in being built with a lattice that includes a singularity (Mandelkow et al., 1986). The "seam" that runs along the MT wall in vivo (McIntosh et al., 2009) gives them a sidedness that would allow force generation between parallel MTs, just

as happens between adjacent doublet MTs in a cilium or flagellum. In this hypothesis, a minus end-directed motor specialized for this task would bind by its tail to the seam on KMTs, treating them as cargo; nonKMTs would serve as the substrate upon which the motor walks. Such a force-generating mechanism would again pull each KMT poleward with a force proportional to its length. The preference of this motor's tail for KMTs could be explained simply if the on-rate for its binding were slow: KMTs are more stable than non-KMTs, so only the former would persist long enough to bind a meaningful population of motor tails.

The difficulty of finding a single hypothesis that explains all evidence about congression has led to the suggestion that kinetochores are somehow “smart”, i.e. can sense chromosome position and adjust KMT state accordingly (Mitchison, 1989a). Smart kinetochores might sense spatial gradients of either MT dynamics (Gardner et al., 2008) or of some unknown factor. Mathematical approaches have been used extensively to try to understand either such “intelligence” or the mechanisms that drive chromosomes to the metaphase plate.

First principles suggest that chromosome motion will be determined by a balance of the applied forces, implying that changes in the direction of movement and/or an alignment of chromosomes at the spindle equator will require either a spatially dependent force or a spatially dependent regulator of force generating activity. Table 3 presents three different approaches to this problem, summarizing our current understanding of the problem. Each of these models presents a simple mechanism that can explain aspects of chromosome oscillations and positioning during prometaphase. Their biggest value is that they give us a quantitative understanding of how each mechanism can contribute to chromosome congression. For a recent review of the first two models see (Vladimirou et al., 2011).

In cells, however, it is likely that these and other mechanisms work together (Kops et al., 2010). The evidence cited above shows the complexity of the phenomenology to be explained. Thus, models based on a single mechanistic premise are unlikely to survive. As with models for spindle formation, they represent a good place to start, but further work will be required to make them realistic. An advantage of the experiments *in silico* is that they allow one to switch pathways on and off or to alter spindle parameters, so one can see exactly how such changes affect the result. Further mathematical modeling should allow distinctions to be made between these and other possibilities, although it will require building more complex models that include more players and mechanisms.

E. BIOPHYSICAL PROPERTIES OF THE METAPHASE MITOTIC SPINDLE

E1. Introduction

As chromosomes attach and congress, the spindle approaches the 2-fold symmetric structure seen in textbooks. The onset of true metaphase is not well-defined, given chromosome oscillations throughout prometaphase, but since chromosome alignment is pretty good in late prometaphase, and cells look ready for anaphase, such stages are generally called “metaphase”. This period is sufficiently constant in structure that it has provided the material for many experimental studies.

E.2 Stability of the spindle to physical perturbation

As early as the 1950s, spindles were known to be labile, because cell lysis led to their rapid dissolution. Efforts to stabilize spindles for study in a cell-free environment have gone through several phases, beginning with Mazia's pioneering work with sulfhydryl reagents, passing through the use of glycols and other water-miscible solutes (Rebhun et al., 1975), then on to the use of conditions that promote the polymerization of purified tubulin in vitro (Weisenberg & Rosenfeld, 1975). All these stabilized spindles were biochemically complex and physiologically quite inert, incapable of most of the functions that would have provided assays for the significance of their individual components. Nonetheless, several potentially interesting MAPS were found, for example (Murphy, 1980), and the discovery of Taxol as a stabilizer of MTs at high dilution led to more (Kuriyama, 1986). However, the complexity of these spindles was too great for the then-available methods for protein separation and characterization.

Spindle properties in vivo were studied to more effect. With polarization optics, the number and/or density of spindle MTs could be measured, as reviewed in (Inoue & Salmon, 1995), so the effects of environmental changes could be assessed quantitatively. Reductions in temperature or increases in hydrostatic pressure reduced spindle birefringence, supporting the idea that the fibrous components of spindles were stabilized by hydrophobic bonds. The data on spindle birefringence as a function of temperature were interpreted with two models for spindle assembly: one very simple but hard to reconcile with the polymerization reactions that probably accounted for spindle fiber formation (Inoue et al., 1975), and one based on condensation-polymerization, the then-accepted mechanism for tubulin polymerization (Salmon, 1975). These innovative studies provided a flavor of biophysical analysis and helped the community to pay attention to the energies and reaction parameters that are certainly important for spindle behavior in vivo, but they were frustrated by the complexity of the system they were trying to analyze with comparatively simple tools. For example, either pressure or lowered temperature could have been affecting the ionic composition of cytoplasm, inducing MT dissolution by an indirect pathway.

Early studies of spindle reformation after dissolution by high hydrostatic pressure (hundreds of bar) did provide convincing evidence that spindle fiber growth in plant cells was controlled at least in part by kinetochores (Pease, 1946). Work with hydrostatic pressure also demonstrated that different parts of a given spindle (and different stages during spindle action) were differentially resistant to dissolution by such treatments: the MTs of the "midbody" that formed in the spindle midzone during late anaphase were highly resistant to dissolution by pressure (Salmon et al., 1976). This work showed that spindle MT dynamics were a moving target, varying from place to place in the spindle and changing as mitosis proceeded.

The lability of spindle MTs in vivo was analyzed in an informative way by the use of colchicine, a tubulin-binding drug that blocks MT assembly. Rapid application of colchicine to a metaphase spindle with subsequent analysis of spindle birefringence gave the first direct evidence that spindle MT turnover is exceptionally fast (Salmon et al., 1984). Birefringence decreased within seconds of drug addition, suggesting that metaphase MTs were turning over at a remarkably high rate. This observation was confirmed and extended by studies that

used fluorescent tubulin injected into living cells, followed by fluorescence photobleaching (NB: this method is commonly called fluorescence recovery after photobleaching (FRAP), but this is a misnomer; the fluorescence does not recover, it redistributes). With this assay, spindle MT turnover half-times were ~30 sec in mitosis and >120 sec in interphase (Saxton et al., 1984). These rates were so much faster than could be explained by the then-accepted ideas of MT assembly by condensation polymerization that something was clearly wrong in the way spindles were being considered. The discovery of MT dynamic instability resolved this conundrum and provided a sensible way to think about spindle MT dynamics in vivo (Mitchison & Kirschner, 1984).

E3. Spindle dynamics as seen by local perturbations

Many investigators have recognized the value of probes that would disturb or report on the behavior of a local region within the spindle, rather than the spindle as a whole. Initial work used microbeams of ultraviolet light to reduce spindle birefringence in a small fraction of the spindle's volume (Forer, 1965). The ensuing processes of birefringence recovery were complex and led to diverse and conflicting interpretations. A brief summary is that areas of reduced birefringence were generally repaired in a polar fashion with fibers growing into the perturbed area from the side near the kinetochores, more evidence for growth of MTs at or near the kinetochores, as seen after perturbation by high pressure.

UV microbeaming was re-instituted for a study of the mechanics of metaphase spindles in diatoms. Here clear evidence was obtained that the non-kinetochore MT (nonKMTs) that form a bundle running between the spindle poles provided mechanical support that held the spindle poles apart: irradiation of one side of the interpolar spindle led to spindle collapse, with the structure folding around its weakened side (Leslie & Pickett-Heaps, 1983). Thus, in at least some cells the pole-directed forces that act on chromosomes have a reaction counterpart in forces that act on the poles to pull them toward the metaphase plate. However, renewed efforts with this method to understand the dynamics of kinetochore MTs in bigger spindles has continued to yield complex and often contradictory results (Spurck et al., 1997). It seems likely that the distributions of resistive forces in large and small spindles are different, a problem that deserves further work.

Local spindle perturbations by micro-irradiation of non-MT spindle parts has produced more consistent results and provided insight into several aspects of spindle behavior. As mentioned above, the tension acting on each metaphase kinetochore was demonstrated by ultraviolet irradiation: the chromosome moved toward the pole connected to the unirradiated kinetochore. This result was confirmed by laser irradiation of kinetochores (Skibbens et al., 1995), and tension throughout the centromere was demonstrated by laser ablation of the chromatin between sister kinetochores; this treatment allowed both kinetochores to move poleward in metaphase, as if anaphase had already begun.

Local perturbations of spindles containing fluorescent tubulin have also revealed interesting aspects of spindle MT behavior. The MTs in a kinetochore-associated fiber turn over faster during metaphase than after anaphase has started (Gorbsky & Borisy, 1989; Zhai et al., 1995). Moreover, a zone of reduced fluorescence, generated by microirradiation of metaphase KMTs, stayed in place as the chromosomes moved in anaphase, suggesting that

MT dynamics at kinetochores are important for chromosome motion (Gorbsky et al., 1987). This result was questionable at the time, because tubulin conjugated directly to a fluorophore, like rhodamine or fluorescein, makes MTs photolabile (Vigers et al., 1988). Several alternative approaches have, however, confirmed the finding. For example, tubulin labeled with a photoactivatable fluorophore will form spindles *in vivo* and can then be activated locally, so the position of the resulting zone of fluorescence can be followed over time (Mitchison, 1989b). These results confirmed the rapid turnover of non-KMTs and showed that tubulin adds to metaphase KMTs at the kinetochores. Newly added subunits migrate toward the spindle poles at about 1 $\mu\text{m}/\text{min}$, a motion called “flux”. Moreover, a spindle can be cut with a microneedle, and MT dynamics at the kinetochore continue unperturbed (Nicklas, 1989), confirming the importance of kinetochores for coupling KMT dynamics to chromosome motion.

E4. Spindle dynamics as seen by global studies with specific reagents or tools

The special character of KMTs seems to result simply from their interaction with kinetochores, posing the question, what aspect of kinetochore structure/function can have such a marked influence on MT behavior? The formation of bundles of KMTs, known as K-fibers, depends on Ndc80, a heterotetrameric kinetochore protein complex. Knockdown of any subunit in the Ndc80 complex inhibits end-on MTs attachments and the formation of well-defined K-fibers, which results in severe defects in chromosome congression (DeLuca et al., 2006). For example, in cells whose Nuf2 subunit was knocked down, the chromosomes congressed properly in only 25% of the cells (Cai et al., 2009a). But the complexity of the situation *in vivo* is demonstrated by the fact that when both Nuf2 and HSET (a spindle-localized kinesin-14) were knocked down at the same time, the number of cells with bi-oriented and congressed chromosomes increased to 75%. The movement of these chromosomes to the spindle equator seemed to occur by kinetochores sliding along bundles of MTs. Recent work also implicates the long, α -helical coiled-coil kinetochore protein, CENP-F (mitosin), in the formation of a stable kinetochore-MT association; knock-down by RNAi again leads to instability in chromosome attachment to the spindle (Yang et al., 2005).

Information about the dynamics of all spindle MTs has recently been extended by using highly sensitive cameras with sufficient signal-to-noise ratio to observe spots containing only a few fluorophores per resolution element. With such a device one can follow the behavior of MTs that are sparsely labeled and appear speckled (Waterman-Storer et al., 1998). This method allows the motions of many MTs to be followed simultaneously, demonstrating that metaphase MTs are continuously in motion, even when the spindle itself appears stable (Vallotton et al., 2003; Yang et al., 2007a). Thus, tubulin is added at the pole-distal end of each MT and is removed at or near the pole; probably every spindle MT is in flux toward the pole nearest its minus end, albeit non-KMTs in most animal spindles turn over so fast at their dynamically unstable plus ends that their flux is hard to see.

Speckle imaging of fruit fly spindles in which various motor enzymes have been inactivated by mutation or antibody injections has demonstrated that flux depends on the activity of kinesin-13s operating at or near the spindle poles to promote MT depolymerization (Rogers

et al., 2004). It seems, therefore, that flux is driven at least in part by ATP-dependent, enzymatically facilitated MT depolymerization. Other studies suggest that kinesin-5 acting in the spindle mid-region contributes to spindle MT flux in at least some systems (Brust-Mascher et al., 2004). Thus, kinematics show spindle MTs to be dynamic in a polar manner, and experiments suggest that this is not simply MT “treadmilling”, i.e., the polymerization of tubulin at one MT end and its depolymerization at the other, as occurs in vitro (Margolis & Wilson, 1981). It is a more “active” process in which ATP is consumed to make the metaphase spindle into a bipolar treadmill. It must be added, however, that flux is not universal: in yeasts where the minus ends of all spindle MTs appear to be capped by a γ -tubulin complex, flux has not been observed. Thus, it is not necessary for mitosis. Moreover, the details of flux can be surprisingly complex; in spindles formed in frog extracts without chromosomes the speed of MT flux is a function of position along the spindle axis (Yang et al., 2008). This observation demands that such spindles are made from MTs that are short relative to the distance from pole-to-pole. Hypotheses for the role of flux in spindle function are considered below.

E5. Mechanical properties of spindles

A growing spindle can elongate in spite of a force applied to impede pole separation (Hays, 1985). The magnitude of this force has been hard to measure, given the technical difficulties of the observations, but the existence of an elongation force demonstrates that a spindle is a mechanical entity, even though it is composed of so many individual parts. However, even a metaphase spindle, which has achieved a steady state length, is dynamic, as seen in its responses to mechanical perturbations with a micro-needle, e.g., (Nicklas et al., 1982).

More recent work has extended the earlier observations. When a spindle is flattened, its structure rearranges to accommodate the perturbation and retain 2-fold symmetry (Dumont & Mitchison, 2009a). Spindles are visco-elastic for small perturbations but respond with plastic deformation to bigger ones. They are mechanically anisotropic, as suggested by their structure: axial compliance is about 10-fold less than that in orthogonal directions. The application of 4 nN will induce spindle shortening of about 1 μm . Apparently, spindle MTs will shorten in response to compressive forces, but we must also conclude that these MTs are interconnected to form a mechanical unit.

Both the chemistry and structure of MT-MT connections in spindles are still poorly understood. There are certainly bridges between spindle MTs (Hepler et al., 1970), and in later mitosis the MT cross-linking protein Ase1 plays an essential role in maintaining connections between MTs growing from the two spindle poles (Zhu & Jiang, 2005). MT neighbor density analyses show that the relative positions of MT cross-sections are not random, even in metaphase (Mastrorarde et al., 1993), but the molecules that connect spindle MTs are still elusive. They are likely to be important for both the structure and mechanics of spindles, because they will help to define the ease with which one spindle MT can move relative to its neighbors, a property of great importance in the motion of MT-attached chromosomes. For a thoughtful discussion of spindle fiber “friction” and its role in the spindle “force map”, see (Dumont & Mitchison, 2011).

Micromanipulation of chromosomes in living cells has made certain properties of spindle linkages clear: KMTs are well-attached to kinetochore, albeit these connection can be broken with a firm microneedle (Nicklas & Gordon, 1985). The force required to pull a chromosome off of its KMTs increases throughout prometaphase (Nicklas & Ward, 1994). KMTs are also tethered at or near the spindle pole, so a chromosome can easily be moved across the spindle, with its fiber pivoting at the pole. It is easy to push a mono-oriented chromosome closer to the pole to which it is connected, presumably by deforming its KMTs, but pulling it farther away from the pole is hard (Begg & Ellis, 1979). Increasing the pole-to-kinetochore distance probably requires the addition of new tubulin subunits, a comparatively slow plastic deformation.

E6. Summary and Conclusions

The metaphase spindle is a dynamic, multi-order steady state. Dynamically unstable MTs are turning over remarkably fast at their plus ends, allowing pole-associated MTs to explore the volume surrounding each centrosome. MTs that associate with kinetochores, or anything else that slows their rate of turnover (e.g., interactions with MTs from the opposite pole) are in flux toward the spindle pole, with net tubulin addition at the pole-distal ends and net loss at or near the poles. Tubulin depolymerization is probably facilitated, if not driven, by ATP-dependent motors that modulate MT dynamics in their vicinity. Thus, metaphase spindles behave as mechanically coherent, two-fold symmetric treadmills in which MTs are constantly migrating toward both spindle poles.

F. BIOPHYSICAL ASPECTS OF SPINDLE CONTROL

F1. Introduction

The principal job of the spindle is clear: accurate segregation of the duplicated chromosomes. The timing and position/orientation of mitosis are, however, important issues that are subject to control. These kinds of regulation bring to our attention the biophysical problems of a cell's sensing its progression through the growth and division cycle and of its perceiving how its spindle is oriented relative to organism-based coordinates.

When the spindle should act is a two-part question: 1) when should the whole process of cell division begin and 2) when is it right to start anaphase. Mitosis must not begin until after the chromosomes have been duplicated and the cell has accumulated enough materials to provide the dowry of macromolecules that will allow subsequent cell generations to thrive. These issues of cell cycle control are well studied; for a brief, insightful review, see (Nurse, 2002). However, the act of chromosome segregation is like a Rubicon for cell cycle progression; once this boundary has been crossed, there is no turning back; this decision too must be regulated. The mechanism controlling anaphase onset is the "spindle assembly checkpoint" (SAC); this mechanism too has been extensively studied and reviewed, although its details are not yet well understood. Below, we sketch only the biophysical issues of anaphase control. **Placement and orientation** of spindles also present some biophysical issues, and these are described next.

F2. Timing of anaphase onset

The processes of chromosome attachment and bi-orientation seem to work by selection: the correct (biologically advantageous) structures are stable while others are not. Thus, getting the last chromosome properly attached may take time. It therefore makes sense for improperly attached chromosomes to delay anaphase onset. The event of properly attaching the last chromosome could, in principle, remove this inhibitory signal and initiate the mitotic step that is irrevocable. With such a mechanism there would be no need for a sophisticated process to count the chromosomes already attached (McIntosh, 1991).

Careful, descriptive work has provided support for this idea: while the duration of prometaphase in cultured mammalian cells is variable, the time between bi-orientation of the last chromosome and the onset of anaphase is comparatively precise (Rieder et al., 1994). Genetic screens have identified proteins that can delay anaphase onset when a proper chromosome attachment has not yet occurred (Hoyt et al., 1991; Li & Murray, 1991). Much subsequent work has shown that many of these proteins are kinetochore associated, making it comparatively easy for them to monitor the state of chromosome attachment. Thus, a SAC that assesses chromosome attachment and terminates a wait-anaphase signal as soon as that chromosome is “on board” is now the widely accepted view of this kind of spindle regulation.

The pathway for SAC action is understood in principle: kinetochores on unattached chromosomes activate a protein complex that exchanges rapidly between kinetochore and cytoplasm (Howell et al., 2000); it provides a short-lived, diffusible signal that inhibits the trigger for chromosome segregation. So long as one chromosome remains unattached, this signal is continuously renewed, and anaphase is delayed. Interestingly, however, the diffusion of this signal appears to be confined to the spindle, because when two spindles inhabit the same cell, their SACs are independent (Rieder et al., 1997).

The anaphase trigger is a protease, commonly called “separase”; it degrades the cohesion links between sister chromatids, allowing them to initiate anaphase motions (Wirth et al., 2006). A diffusible, kinetochore-activated inhibitor of this protease keeps the cell in metaphase until the last chromosome is attached, whereupon activation of the inhibitor stops, residual inhibitor degrades, proteolytic activity turns on, and anaphase can start, as reviewed in (Ciliberto & Shah, 2009; Nezi & Musacchio, 2009; Varetto & Musacchio, 2008). The two key unanswered questions here are: how does a single chromosome emit a signal strong enough to influence the entire spindle, and how does proper chromosome attachment attenuate this signal, i.e., keep the separase inhibitor active?

The issue of signal strength cannot be approached very well until we know more about the chemistry by which the signal works, so it will not be discussed here. Two proposals for the relationship between chromosome attachment and signal attenuation have captured the attention of the field and merit description: 1) tension at a bi-oriented kinetochore turns off the ability of that kinetochore to activate the inhibitor (McIntosh, 1991) and 2) the presence of KMTs bound at kinetochores is all that is needed to turn the activation off (Khodjakov & Pines, 2010; Yang et al., 2009). Early evidence favored the tension hypothesis; the experimental application of tension to a mono-oriented chromosome allowed anaphase to

begin (Li & Nicklas, 1995; Maresca & Salmon, 2010), but subsequent work has provided support for the MT occupation hypothesis (Varetti & Musacchio, 2008). Since tension promotes the stability of KMT attachment, however, it tends to increase the number of KMTs, so distinguishing between these hypotheses is not easy, and the issue remains unsolved. Indeed, there are indications that the SAC includes two pathways, one that is silenced by tension and the other by MT occupancy (Skoufias et al., 2001). In either case the mechanism that silences the “wait-anaphase” signal involves strain at the kinetochore, be it due to the binding of additional MT ends or to the stresses that they apply.

Recent evidence suggests that the device that responds to this strain is the outer region of the kinetochore itself, not some part of the connection between sister chromatids (Maresca & Salmon, 2010), so a fascinating problem in biological mechanics remains to be solved. A clue about mechanism may lie in the fact that mammalian cells show a distinct sharpening of the metaphase plate just before anaphase starts. This might reflect a change in spindle action as part of the anaphase trigger. Work that will sort this problem out is underway in several labs.

F3. Control of spindle position

Spindle placement within the cell defines where the cell will divide during cytokinesis and thus the relative sizes of its daughters. In higher plants a cell will live where it is born (although its size and shape can change), so the position of cleavage at cytokinesis is critical for organism morphogenesis. In animals cell placement is less important because daughter cells can move.

In both plants and animals the most common position of the spindle is at the cell’s center, so daughter cells are of very similar sizes; the variation in volume is usually only a few percent and probably reflects the limited precision of the spindle positioning system. In some cells, however, the spindle moves toward one side of the cell before anaphase begins and cytokinesis produces cells of markedly different size. For example, the division of some embryonic cells is asymmetric, and the daughter cells differ in both size and developmental fate. A possible correlation between size and fate has fascinated biologists for years and led to ideas about how physical asymmetry might be coupled to biological asymmetry. In some cases, factors that reside in the cytoplasm of this “mother” cell have been shown to move to one end of the cell before division, so asymmetry in size is accompanied by asymmetry in the cytoplasmic dowry that the two daughter cells inherit (Ciosk et al., 2006). These cases are of biological importance to our understanding of embryonic differentiation.

The molecular mechanisms that establish an asymmetric cell division are just beginning to be understood. There are proteins resident in the cell cortex (the gelatinous layer just inside the plasma membrane) whose positions correlate with the asymmetry of the subsequent division and are probably causal for it. These include the products of the “Par” genes, which were discovered through the phenotypes of their loss of function alleles in embryonic cell divisions of worms and flies (Kemphues et al., 1988). They affect the distribution of actin and related proteins in the cell cortex, and this in turn leads to mitotic asters of different sizes, causing the spindle to shift from its customary position at the cell’s center (Munro et al., 2004).

In some plant cells, like the ones that divide to produce the cells that form the pores in a leaf (the “stomata”), the cell’s nucleus migrates many micrometers from a central position to establish the site of future cell division (Pickett-Heaps, 1969). Thus, the mechanisms that link physical and biological asymmetry at cell division are probably numerous. All such mechanisms seem to involve interactions between the spindle and the cell cortex with ensuing spindle motions. Exactly how these fascinating processes work are important problems for future work.

F4. Control of spindle orientation

The orientation of the spindle can also be significant, both in animals and plants. For example, the spindles that form in many stem cells show consistent orientations relative to nearby markers, like a basement membrane. When the spindle axis is perpendicular to this surface, the division tends to be asymmetric and “differentiative”; such behavior is characteristic of cells whose divisions gives rise to one daughter that will grow and divide again while the second daughter will specialize to perform essential functions for the organisms. When the division axis is parallel to the nearby surface, mitosis is generally symmetric or “multiplicative”, meaning that it simply increases the number of cells present (Siller & Doe, 2009).

In plant cells the orientation of the spindle axis is particularly important because cell walls and cell-cell adhesion block subsequent cell motion; cell division is therefore a major mechanism in plant morphogenesis (Rasmussen et al., 2011). The mechanisms that allow a cell’s neighborhood to convey information through the plasma membrane to the cytoskeleton and turn the spindle appropriately are still mysterious; they comprise additional fascinating subjects for future work.

G. CHROMOSOME SEGREGATION AT ANAPHASE

G1. Introduction

Anaphase is the most dramatic stage of mitosis because all the chromosomes segregate in approximate synchrony. While their speeds are slow, a time-lapse recording reveals a biological phenomenon that is truly breathtaking (see, e.g., the videos available at http://www.cellimagelibrary.org/images?basic_video=true&k=mitosis&simple_search=Search). Ironically, in spite of the attention that anaphase has received over ~120 years, we still do not know how it works. Even the descriptions of this mitotic stage are complex, now that we can look not just at the chromosomes but at the spindle fibers themselves. This section will therefore review the relevant processes in stages, starting with kinematics and moving on to the experiments and theories that tell us something about spindle mechanics.

G2. The kinematics of anaphase spindle components

Chromosome segregation in most organisms begins with a rather abrupt transition that is initiated by the proteolytic severing of cohesion links between sister chromatids. The ensuing separation of duplicate chromatids is sufficiently synchronous that most of the chromosomes begin their anaphase motions within seconds of one another. This precision is

important for the accuracy of chromosome segregation, since a lack of synchrony would produce lagging chromosomes that might then be excluded from the forming nuclei.

Within a given cell the speed of each anaphase chromosome is rather constant. Occasionally, anaphase chromosome movements falter, so a detailed analysis shows some similarity to the oscillations of prometaphase, suggesting that the complex behaviors of earlier mitosis are still possible. However, continuous motion at a constant speed is more common. Big chromosomes move at about the same speed as small ones, so viscous drag cannot be rate-limiting for speed (Nicklas, 1965). It is common, however, for the chromosomes on the spindle's periphery to move a little faster than those near the spindle axis, particularly in cells with focused spindle poles. The leading edge of each chromosome set therefore becomes cup-like as it approaches the pole.

As chromosomes near the poles, they often slow, but at about this time the poles commonly begin separating, so chromosome motion relative to cell-fixed coordinates continues at an essentially constant speed. The movement of chromosomes to poles and the increasing separation of the poles are called anaphases A and B, respectively; both contribute to chromosome segregation. Their timing and extent can differ considerably from one cell type to another. In many vertebrate cells the anaphases are of approximately equal extents, and in these cases, A commonly precedes B. An image of anaphase chromosomes in a typical mammalian cell is shown in Figure 7. In many fungi with elongate cells, metaphase spindles are short, so anaphase A is brief in both time and distance. Anaphase B, on the other hand, is many times longer, leading to a final separation of the nascent nuclei that can be even greater than 15 times the length of the metaphase spindle. Since daughter nuclei often reform during this motion, its later stages have been called "anaphase C". We think this additional term is unnecessary, since the mechanism for late-stage nuclear motion appears to be the same as earlier spindle elongation.

The kinematics of anaphase MTs is more complex than that of the chromosomes and poles, so it must be considered in parts, depending on the connections formed by the MTs under consideration. KMTs shorten to allow anaphase A. Initial work on vertebrate cells, based on a variety of methods including the labeled tubulins described above, showed that anaphase KMTs lose subunits at their kinetochores; the term, "Pacman" anaphase, was introduced to connote KMTs being consumed at their kinetochores (Mitchison et al., 1986).

The advent of speckle imaging, based on MTs unevenly labeled with fluorescent tubulin, has allowed a much more thorough assessment of anaphase MT movements. The site of anaphase KMT depolymerization is variable, both with time in any given cell and from one cell type to another. In mammals, initial KMT shortening is $\sim 3/4$ at the kinetochores, but in late anaphase A it becomes more rapid at the poles than at the kinetochores (Yang et al., 2008). In blastomeres of the nematode, *C. elegans*, there is essentially no KMT subunit loss at the kinetochores, and almost all of anaphase is achieved by spindle elongation (Labbe et al., 2004). In spermatocytes of the crane fly, *Nephrotoma*, the KMTs are actually adding tubulin at the kinetochores as they shorten during anaphase A, thanks to comparatively rapid depolymerization at the poles (LaFountain et al., 2004). Thus, an understanding of anaphase must recognize these different strategies and accept that biological variation is significant.

The anaphase kinematics of spindle MTs that do not end on kinetochores (non-KMTs) is very different. In cells with centrosomes (e.g. most animal cells and protozoans) some non-KMTs form asters. These retain the rapid dynamic instability of metaphase, but during anaphase they commonly grow, becoming so long they extend to the cell cortex. In animal cells these MTs somehow instruct the cytokinetic machinery where to form the furrow that will divide the cell during cytokinesis.

Other non-KMTs extend into the body of the spindle. The subset of these MTs that is shorter than $\sim 1/2$ the metaphase spindle length retains rapid dynamics. The longer non-KMTs pass the chromosomes (or occasionally penetrate them) and interact with their counterparts from the opposite pole; these “interpolar” MTs (ipMTs) become stable at or shortly after the onset of anaphase (Saxton & McIntosh, 1987) and form a mechanical framework that connects the two spindle poles (McIntosh & Landis, 1971). In some cells there is also an anaphase initiation of new MTs immediately behind the now-separated chromosomes (Uehara & Goshima, 2010). These contribute to the number and distribution of MT in the spindle “interzone”, i.e., the region between the separating chromosomes. In plant cells these interzonal MTs are of particular importance because they form the “phragmoplast”, an array of interdigitating, antiparallel MTs that gathers Golgi-derived vesicles at the cell’s midplane. These vesicles fuse to build a large, flat, membranous compartment that grows in the plane perpendicular to the spindle axis, forming the “cell plate”, which ultimately fuses with the plasma membrane, effecting plant cell cytokinesis (Segui-Simarro et al., 2004).

As anaphase proceeds, ipMT bundles in most cells become obvious by either light or electron microscopy. In some small spindles, however, one or more bundles of ipMTs is evident even before anaphase onset; in diatoms, for example, the interpolar bundle is very prominent, even in prometaphase (McDonald et al., 1977). In a few cells (e.g., nematode blastomeres) such bundles are not apparent, even in late anaphase; the distance between spindle poles increases with no detectable connection between them.

As the spindle lengthens in anaphase B, the ipMTs in most cells also elongate. Micro-injection and photobleaching experiments indicate that tubulin adds to the pole-distal ends of these MTs (Masuda et al., 1990; Saxton & McIntosh, 1987). One might think that pole-distal polymerization would simply increase the extent of ipMT overlap, but these elongating polymers slide relative to one another, so their increasing length simply enables them to slide farther without separating into unconnected sets. Mechanisms for this sliding elongation are discussed below.

G3. Responses of anaphase A to experimental perturbations and their implications for mechanism

Many organisms, e.g., plants and insects, will grow at a wide range of temperatures, permitting a study of temperature effects on the speed of chromosome motion: the rate of anaphase A increases by a factor of 2 – 3 with each 10°C increase in temperature (Schaap & Forer, 1979). While chromosome speeds at a given temperature differ among organisms, the extent of temperature dependence is quite consistent. These data have been used to estimate the activation energy of the rate-limiting step for anaphase (assuming that it is the same at

all temperatures). The average value is 17 kcal/mol, consistent with the not surprising statement that chromosome speed is regulated by an enzyme.

Using an innovative form of local perturbation, Nicklas studied the effects of local temperatures on the kinetics of spindle-dependent process. Grasshopper spermatocytes were maintained on a microscope in a cool room (~17° C) but exposed to the tip of a tiny Nichrome wire that allowed local heating in chosen parts of the spindle (Nicklas, 1979). The shape and magnitude of the resulting temperature gradients were assessed by several criteria, allowing confident statements about temperatures in different parts of the same spindle. Spindle birefringence responded to these gradients as expected, yielding more signal in the warmer half spindle during both metaphase and anaphase. Surprisingly, however, the rates of anaphase in the two half spindles were essentially the same and were defined by the temperature in the spindle interzone. If MT dynamics in these spindles resembles those in crane flies, where flux toward the poles is the major mechanism for anaphase A, then these results are compatible with the hypothesis that flux is driven by sliding motors in the zone of MT overlap (Margolis & Wilson, 1981; McIntosh et al., 1969), and anaphase A is a result of a frictional drag between these interpolar MTs and the KMTs, coupled with MT depolymerization at the spindle poles.

Treatments with drugs that inhibit tubulin polymerization have not revealed any change in the rate of anaphase A, but the addition of Taxol, which stabilizes MTs and drives tubulin polymerization, can cause anaphase chromosomes to back up, suggesting that KMT depolymerization is mechanically important for anaphase A (Bajer et al., 1982).

There have been several waves of efforts to make experimental models for anaphase, based on isolated spindles or lysed mitotic cells. Conditions were then sought that would induce anaphase-like motions. Using buffers that supported tubulin polymerization in vitro, several labs have described quite robust spindle elongation in vitro. Lysis of comparatively stable spindles allowed the demonstration that ATP was needed for ipMT sliding (McDonald et al., 1986), but informative work on anaphase A was more difficult. Some studies of lysed cells found that ATP supported chromosome-to-pole motion (Hogan et al., 1993), but another study obtained similar motions without ATP, simply by the addition of Ca^{2+} , which is known to depolymerize MTs in vitro. This result was interpreted to mean that the role of ATP was simply to keep KMTs labile, perhaps by phosphorylating a MT-associated protein (Spurck & Pickett-Heaps, 1987).

More extensive and reliable chromosome motion was achieved with spindles assembled in cytoplasmic extracts of frog eggs (Murray et al., 1996), but even with this system, the only conclusion reached was that the speed of anaphase was similar to the speed of metaphase MT flux toward the spindle poles (Desai et al., 1998). The complexity of all these “models” has made more penetrating experiments difficult.

Genetic perturbations of spindles in vivo have provided a more informative experimental approach to mechanism, though it has taken the marriage of genetics with cell and molecular biology to get information about anaphase. Early work with mutant fungi was informative about the proteins necessary for spindle formation, as sketched above, but it helped less with

anaphase, perhaps because no one has yet reported really tight, temperature-sensitive (TS) alleles of relevant mitotic components. Thus, if spindles formed, they tended to work; if they didn't form, the effect of a mutation on anaphase was not apparent.

Analogous work with fruit flies has been more informative, thanks to the better cytology of these larger spindles, the availability of some TS alleles, and the possibility of antibody micro-injection to work in parallel with conditional mutants. Even here, though, the data from different labs are not in good agreement. For example, one lab found that the perturbation of dynein after spindles had formed slowed anaphase chromosome motion, suggesting that this minus end-directed, kinetochore-associated motor plays a role in anaphase A (Sharp et al., 2000). Analogous effects on anaphase were obtained by RNAi of proteins like Rod, which are important for the binding of dynein to kinetochores (Yang et al., 2007b). On the other hand, some labs have found that while dynein is localized to kinetochores in early mitosis, it leaves those sites during prometaphase and is almost undetectable by late metaphase or anaphase (King et al., 2000). Further, the disruption of dynein by other methods in other labs has failed to show an effect on anaphase (Goshima & Vale, 2003; Vorozhko et al., 2008). Moreover, dynein is important for many cellular processes, so indirect effects from its perturbation cannot be excluded. This problem is compounded when the antigen probed by injected antibodies or the component knocked down was not dynein itself, but a protein like Rod that is important for the binding of dynein to kinetochores and other cargos as well. Thus, the case for dynein as a motor for anaphase A is still equivocal.

Certain fungi have, however, supported some quite definitive studies. The genetics of budding yeast are so well-developed that students of this cell have built strains in which only one chromosome of 16 is marked for viewing by fluorescence microscopy, and the centromere of this chromosome is modified to be dysfunctional until external conditions are altered. After all the other chromosomes have progressed to a drug-induced metaphase arrest, activation of the centromere on the marked chromosome has allowed the study of a single kinetochore as it attaches to some hyper-elongated spindle MTs (Tanaka et al., 2005). This work has provided clear evidence that a kinetochore-associated, minus end-directed kinesin-14 (Kar3) contributes directly to chromosome-to-pole motion in this cell type. Many investigators now believe that kinetochore-localized, minus end-directed motors drive anaphase A, at least in yeast.

In keeping with the adage that biology is never simple, however, there is convincing evidence in both fission and budding yeasts that kinetochore-localized, pole-directed motors are dispensable for normal rates of chromosome-to-pole motion (Grishchuk & McIntosh, 2006; Tanaka et al., 2007). These results have led to a revival of the idea that MT dynamics might provide the machinery for mitotic movements. Consistent with this view, work in vitro has demonstrated that MT depolymerization can generate forces that are strong enough to drive anaphase A (Coue et al., 1991; Grishchuk et al., 2005; Lombillo et al., 1995a). Recent work has identified and characterized several kinetochore-associated proteins that can couple MT shortening in vitro to the motion of a cargo against a load (Asbury et al., 2006; Efremov et al., 2007; Grishchuk et al., 2008a).

These coupling proteins come in three classes: 1) those that oligomerize on a MT wall. This includes the Ska complex in mammals (Welburn et al., 2009) and the Dam1 complex (Grishchuk et al., 2008b) in fungi, which can even form rings around MTs (Miranda et al., 2005; Westermann et al., 2006); 2) fibrous protein complexes, like Ndc80 (McIntosh et al., 2008; Powers et al., 2009), which include one or more sites that can bind the MT lattice while another protein moiety binds components deeper in the kinetochore; and 3) motor enzymes, like a kinesin 7 or 8, which display an ATP-dependent, plus end-directed motility on MT walls but can also follow the shortening end of a depolymerizing MT in an ATP-independent manner (Grissom et al., 2009; Lombillo et al., 1995b). Thus, students of mitosis must take seriously the possibility that energy stored in tubulin polymers is an important factor in the development of forces for anaphase A.

G4. Putting kinematics and experiments together to understand biological mechanism

Understanding mitotic motions is made even more complex by the fact that all kinetochores so far analyzed include an activity that can use ATP to enhance MT depolymerization. The most direct examples are the kinesin 13s, which localize to the inner kinetochore of mammalian chromosomes (Walczak et al., 1996; Wordeman, 2005). Inactivation of these enzymes by mutation or RNAi causes chromosomes to slow (Rogers et al., 2004) or to lag behind their neighbors in anaphase (Maney et al., 1998; Stout et al., 2006).

Both mammalian and fungal kinetochores harbor a kinesin 8, a plus end-directed motor that shortens MTs in vivo (Cottingham et al., 1999; West et al., 2002). The budding yeast kinesin 8 will also induce MT shortening in vitro (Gupta et al., 2006), and some evidence suggests that the strength of this activity is proportional to MT length (Varga et al., 2006). The kinesin 8 from fission yeast does not induce MT depolymerization in vitro but it increases the probability of MT catastrophe in vivo (Tischer et al., 2009). Meanwhile, the evidence about mammalian kinesin 8 is conflicting (Mayr et al., 2007; Wordeman & Stumpff, 2009).

Thus, the almost universal shortening of KMTs during anaphase A may not be due simply to solution conditions that favor tubulin depolymerization. ATP-dependent motor enzymes, either instead of or in addition to solution conditions, may control anaphase tubulin depolymerization. This result is consistent with the temperature sensitivity of chromosome speeds in anaphase, which shows the high activation energy characteristic of enzyme action. Note also that the growth of non-KMTs during anaphase shows that mitotic cytoplasm at this time does not promote depolymerization of all MTs; kinetochores (and poles) must provide special depolymerizing environments, whether they depends on motor enzymes or not.

A recent advance holds promise of increasing our understanding of kinetochore-MT interaction: the study of isolated centromeres interacting with MTs in vitro (Akiyoshi et al., 2010). Mini-chromosomes that contain little more DNA than the centromere itself have been isolated from budding yeast and attached to micro-beads that can be grasped and manipulated with a laser trap. Ongoing studies are probing the roles of specific kinetochore proteins in the properties of this interaction, including the ability of forces applied to the kinetochore to manipulate MT dynamics. This approach is likely to be fruitful in revealing the contributions of individual kinetochore components to mitotic chromosome behavior.

MT kinematics in anaphase A demonstrates that we must also consider tubulin depolymerization at the spindle poles. This phenomenon has not yet received the attention given to kinetochore-mediated MT shortening, but in the long run it may be just as important. Such motion probably does not occur in yeasts, where all spindle MTs are capped throughout mitosis (O'Toole et al., 1999), suggesting that they do not depolymerize at the poles. Polar MT depolymerization definitely does occur in spindles from fruit flies, crane flies, mammals (see above), and probably in many other cells as well. Spindle poles contain both MT depolymerases, like the kinesin 13s, and plus end-directed motor activities, as well as MT severing activities that could cut MTs loose from their γ -tubulin caps (Zhang et al., 2007). The evidence for such motions in anaphase is clear in some cells, but how the poles hold on to shortening MTs and help to pull the chromosomes poleward remains to be discovered. This is likely to be a fruitful area for future biophysical work.

Given that some part of anaphase A is associated with MT shortening by depolymerization at the poles, there are two additional mechanical systems that must be considered: interactions among spindle MTs and between MTs and spindle matrix component(s), as mentioned in our discussion of congression. Considering first MT-MT interactions, anaphase A could result from forces generated between antiparallel ipMTs in the interzone, as suggested above: the motion of ipMTs could push on KMTs through passive MT-MT links. If coupled with MT depolymerization at the poles, anaphase A would ensue. Alternatively, anaphase A could result from force-producing interactions between KMTs and nonKMTs in the region between chromosomes and poles. The obvious question about such a mechanism is how would a motor, e.g., one that was minus end-directed, know that its cargo-binding end should associate with KMTs only? One plausible mechanism is based on the fact that KMTs are less dynamic than non-MTs. If the on-rate for tail-binding by this motor were slow, it would be less likely to associate with the highly dynamic non-KMTs. The net mechanical effect would then be pole-directed forces on KMTs.

A second possibility, as with congression, is a plus end-directed motor whose cargo-binding end interacts with spindle matrix. Now the tethered motor could push any spindle MT poleward by walking toward its plus end. Until spindle matrix is better established and understood, however, this proposal must be viewed as highly speculative.

In summary, the anaphase motion of kinetochores along MTs may be driven in part by motor activity, in part by MT depolymerization, and in part by motors that bias MTs towards disassembly. All these factors must be considered in coming to grips not only with the biological variation of anaphase A but also with the time-dependent variation in the kinematics of chromosomes relative to KMTs.

G5. Experimental perturbations of anaphase B and their implications for mechanism

IpMTs can contribute to anaphase B in either of two ways: 1) in diatoms, yeast, and fruit flies the available evidence shows convincingly that the interpolar spindle forms a pushing engine; sliding motors, like kinesin 5, cross-bridge antiparallel MTs and push them poleward by walking toward the MT plus ends (Figure 4). The region where ipMTs overlap also contains Ase1, a protein of the PRC1 family that cross-links anti-parallel MTs without generating sliding forces (Janson et al., 2007). This activity helps to maintain the mechanical

stability of ipMT bundles, because its removal by mutation leads to a high frequency of broken spindles (Pellman et al., 1995). Such MT cross-linking may also provide a resistance to sliding analogous to friction (Bormuth et al., 2009), albeit the properties of this kind of protein bonding do not conform to some of the laws that describe macroscopic friction.

There is also evidence for minus end-directed motor activity in the region of ipMT overlap (Hoyt et al., 1993; Sharp et al., 1999); these enzymes provide a force antagonistic to the plus end-directed sliding, pulling the sliding MTs back in the opposite direction. MT motion in anaphase B is therefore a result of a force balance (Brust-Mascher et al., 2004; Cottingham et al., 1999). Additional evidence shows that the extent of anaphase B is usually greater than the initial extent of MT overlap, so MT polymerization by plus-end addition of tubulin is also involved in spindle elongation. This polymerization is likely to be permissive for motions of greater extent, not a cause for spindle elongation, because the interdigitating ipMTs have nothing for their plus-end growth to push against.

2) A different function for ipMTs is seen in anaphase B of filamentous fungi. Here, experiments with UV micro-irradiation have shown that the interzone spindle works as a brake on the rate of spindle elongation (Aist et al., 1991); the dominant forces for the elongation of these spindles are pulls that come from dynein associated with the cell cortex that interacts with astral MTs extending from the poles (Fink et al., 2006). An analogous scenario has been found in blastomeres of the nematode, *C. elegans*, where genetic perturbations of the interpolar spindle accelerate anaphase B. Here kinesin-5 works as a brake on spindle elongation (Saunders et al., 2007), and anaphase B is again driven by outward pulling on astral MTs. One might imagine that such forces would again depend on dynein motors anchored at the cell cortex, but genetic evidence suggests that the pulling in these cells comes from a different mechanism (Schmidt et al., 2005).

In cells where astral MTs and cortical pulling provide the forces for anaphase B, does the interzone contribute anything important to the process? The slowing of spindle elongation by a resistance to sliding in the interzonal MTs might enhance mitotic precision, but ipMTs may also play a more important role. A variety of studies have shown that prior to cytokinesis the entire cell cortex is actin-rich and can presumably bind dynein. The ipMTs may serve to steer the motions of the two half-spindles, so they move away from one another and cannot end up in the same daughter cell.

Cortex-based pulling is not a universal anaphase mechanism, since anaphase B is part of chromosome segregation by spindles formed in plant endosperm and egg extracts, where mechanically robust cortices do not exist. Once again, we find variation in just how a mitotic process is effected, but the result is the same: the two sets of separating chromosomes are moved far enough apart that cytokinesis will assure one complete set of sister chromosomes in each of the two sister cells.

G6. The magnitude of anaphase forces on chromosomes

The force a spindle can exert on a single anaphase kinetochore has been measured by using calibrated micro-needles to deliver counter-forces that affected the rate of pole-directed chromosome motion. The speed of anaphase A in a grasshopper spermatocyte spindle was

unaffected by forces <100 pN, but above this value, chromosome speed dropped linearly to zero at a force of around 700 pN (Nicklas, 1983). This is $\sim 10^4$ greater than would be necessary to overcome viscous drag on a chromosome moving at anaphase speeds through a medium with the viscosity of the spindle.

Structural studies on grasshopper spermatocytes have found about 15 KMTs per chromosome, so the measured anaphase force corresponds to ~ 47 pN/MT. A force of this magnitude could be generated by 8 – 10 ATP-dependent motor enzymes operating in parallel on each KMT, but it could also, in principle, be produced simply by MT depolymerization. Experimental work on forces generated by tubulin depolymerization in vitro has been interpreted to mean that a single MT can generate as much as ~ 50 pN (Grishchuk et al., 2005), the same order as the theoretical maximum, based on the idea that much of the energy from hydrolysis of tubulin-bound GTP is stored in the strain required to straighten a GDP-tubulin dimer so it will fit into the MT wall (Molodtsov et al., 2005; Wang & Nogales, 2005). Not all workers agree with this interpretation, however, as discussed in the theory section below.

An alternative way of thinking about the generation of this high force can be built from the details of insect spindle function. KMTs in the spermatocytes of crane flies (which are distantly related to grasshoppers) are actually polymerizing at the kinetochore during anaphase A (LaFountain et al., 2004). If the same should be true for grasshoppers, the motile force measured above was produced somewhere other than at the kinetochore (Dumont & Mitchison, 2009b). The measured force would have been due to drag on KMTs as they move through the system that binds KMTs to kinetochores, not active force-production at the kinetochores themselves. Where, then, was this force generated? The effects of local heating on the rates of insect spermatocyte chromosome motion suggest that the motive force for anaphase A in these cells comes from the interzone (Nicklas, 1979), consistent with the model mentioned above that forces for chromosome-to-pole motion may be generated by interdigitating ipMTs and conveyed to the KMTs through the “friction” of passive cross-links.

Unfortunately, there are as yet no comparable measurements of anaphase forces in other systems, though it is generally assumed that all mitotic machines are capable of similar force production. Given that the rapid flux of anaphase KMTs in crane flies is unusual, there is a significant need for comparable measurements of stall forces on chromosomes in other systems to make sure that other spindles are as strong as the one that has been described. Regardless of mechanism, however, the presence of such a powerful force-generating system implies the existence of a governor that either limits force generation until chromosome speed is slowed or defines the rate of chromosome movement, even when the load is low. When MT-dependent motor enzymes were first proposed as an origin of mitotic forces, MT dynamics were suggested as the governor for movements (McIntosh et al., 1969). Now that tubulin depolymerization is being considered as a source for mitotic forces, this suggestion becomes more complex. Something other than viscous drag is clearly limiting the rate of chromosome motion, but both the chemistry and the physics of the regulatory device remain to be determined.

G7. Mathematical models for anaphase

To initiate anaphase, sister chromatids are cut from one another by separate. The ensuing chromosome-to-pole motion is likely to be driven by a subset of the mechanisms that drove chromosomes during congression. In fact most of the models discussed in section D will produce anaphase A if sister chromatids are artificially cut in metaphase, a fortunate situation, since experimental severing of a metaphase chromosomes does induce anaphase-like behavior in each chromatid. For mathematical formulations of these mechanisms, see Table 3 from section D. Note, however, that the anaphase kinematics presented above makes it clear that one must consider more than simply kinetochore-mediated force generation; pole-localized mechanical processes and perhaps action in the interzone are also relevant, so all existing models will need expansion to deal with the realities of biological diversity.

Spindle elongation (anaphase B) is driven by mechanisms different from those of anaphase A. Several groups are taking a mathematical approach to analyze the balance of forces described above (Civelekoglu-Scholey & Scholey, 2011; Wollman et al., 2008). Key components for such work include asters that interact with each other and the cell cortex via molecular motors. Such motors exert pulling and pushing forces on spindle MTs, resulting in displacements of the spindle poles. One of the best studied examples of these phenomena is found in *Drosophila* embryos. When these large, multinucleate cells enter mitosis, their peripheral region contains an almost cylindrical shell that harbors several hundred mitotic spindles, each ~10 μm long and located just under the cell surface. The cell cortex or other external factor cannot be limiting to the spindle elongation here, so the length of the spindle must be controlled by intrinsic, spindle-associated mechanisms, reviewed in (Goshima & Scholey, 2010). Here, as elsewhere, key roles are played by kinesins-5 and -14, molecules that crosslink and slide anti-parallel ipMTs at the spindle midzone, acting outward or inward, respectively (Peterman & Scholey, 2009), although other factors, like MT flux, must contribute. For mathematical formulations of these models, see Table 4, which summarizes two models describing spindle elongation via two distinct mechanisms: one in which pole separation is driven mainly by the action of kinesins that crosslink ipMTs, and another where the key role is played by cortex associated dyneins.

H. PERSPECTIVE

This account of mitotic phenomenology is far from complete, but the examples chosen should serve to whet the appetite of any curious biophysicist. Mitosis contains many fascinating problems concerning the roles of nanomachines in biomechanics. The current accounts of spindle parts and the properties of these components provide many factual constraints on rigorous modeling efforts, and the data on spindle behavior under various experimental conditions provide a sound initial test-bed for many model predictions. At the same time, the mechanical properties of spindle components, like kinetochores and centrosomes, are beautiful subjects for advanced experimental work. We hope that this account will help to draw new talent to the field and stimulate exciting work.

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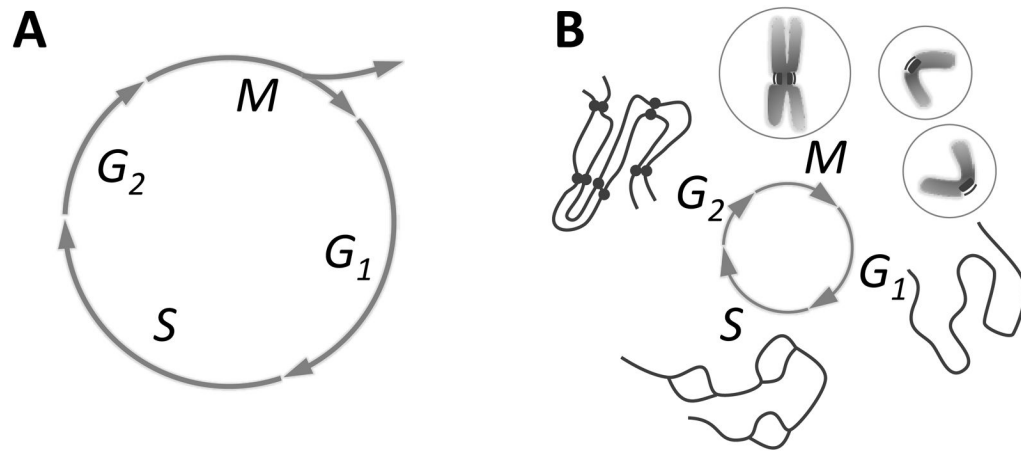


Figure 1. The cell growth and division cycle

Fig. 1A: Mitosis (M) is seen in the context of the whole cell cycle, represented as a circle. “Interphase”, the time between divisions, includes S phase, when the cell’s DNA is replicated, and gaps (G₁ and G₂) before and after S phase. M is followed by cell division, or “cytokinesis”, which divides the already duplicated cell into two essentially identical pieces. The timing of cytokinesis relative to mitosis and interphase varies among different cells, so it is shown only as the forking arrows that imply the emergence of two cells from one. Fig. 1B: Chromosome structure as a function of time in the cell cycle: in G₁ one de-condensed chromosome is shown, representing all chromosomes in the cell. These are replicated during S (“bubbles” are places where replication has already occurred, but replication is not yet complete). In G₂ there are two copies of each chromosome, still held together by cohesions (bars between “chromatids” in the diagram). In early M (prophase) the chromosomes condense, whereupon one can see each chromosome as a double object with two chromatids and a primary constriction (centromere) where cohesions still hold the chromatids together. This is the site where kinetochores form. During M the chromatids separate; in cytokinesis a cell is formed around each chromosome set. Circles on either side of M represent first the one, then the two cells formed by division.

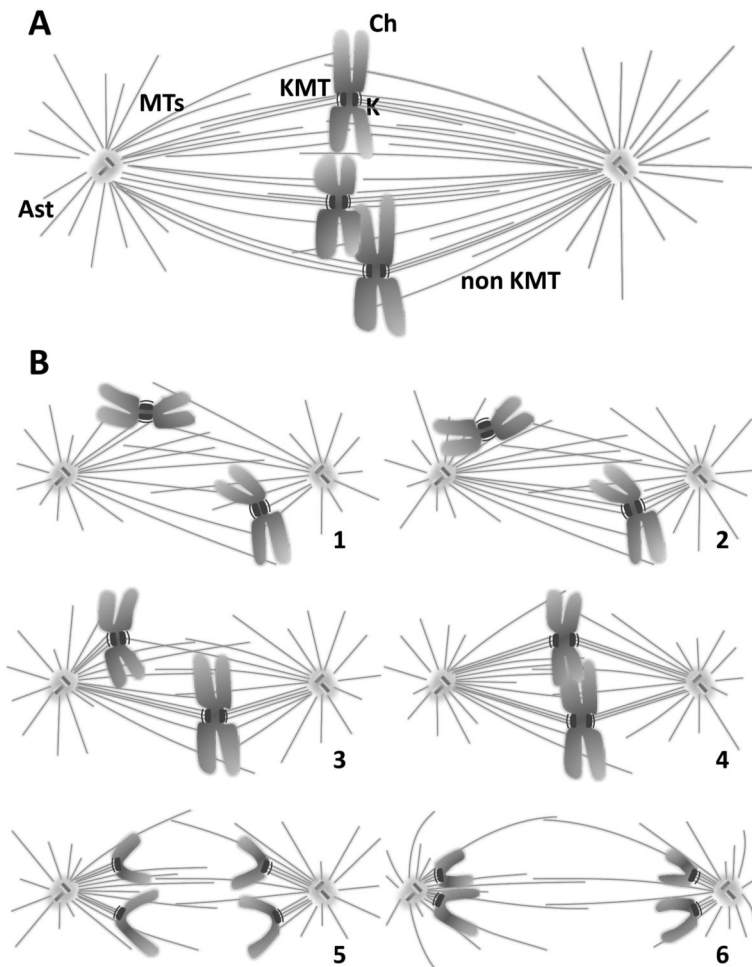


Figure 2. Formation of a typical mitotic spindle

Fig. 2A: A typical spindle at metaphase, showing three chromosomes (Ch) situated near the spindle equator. The spindle poles (SP) and the microtubules (MTs) form the body of the spindle. All MTs are polar polymers, vectorial nature not diagrammed, which are oriented with their plus ends distal to the pole with which they are associated. MTs that associate with kinetochores (Ks), called KMTs, share this polar orientation. Some KMTs reach all the way to the pole but some do not. Many MTs do not end on Ks (nonKMTs), either because they miss that target or because they pass right through the chromatin. Some MTs project away from the spindle and form the asters (Ast). Fig. 2B: Diagrams showing spindle formation (1–4), including chromosome-spindle attachment (1–3) and chromosome congression to the metaphase plate (3,4). Chromosome segregation in anaphase A (5) is followed by spindle elongation, or anaphase B (6). Cytokinesis ensues (not shown).

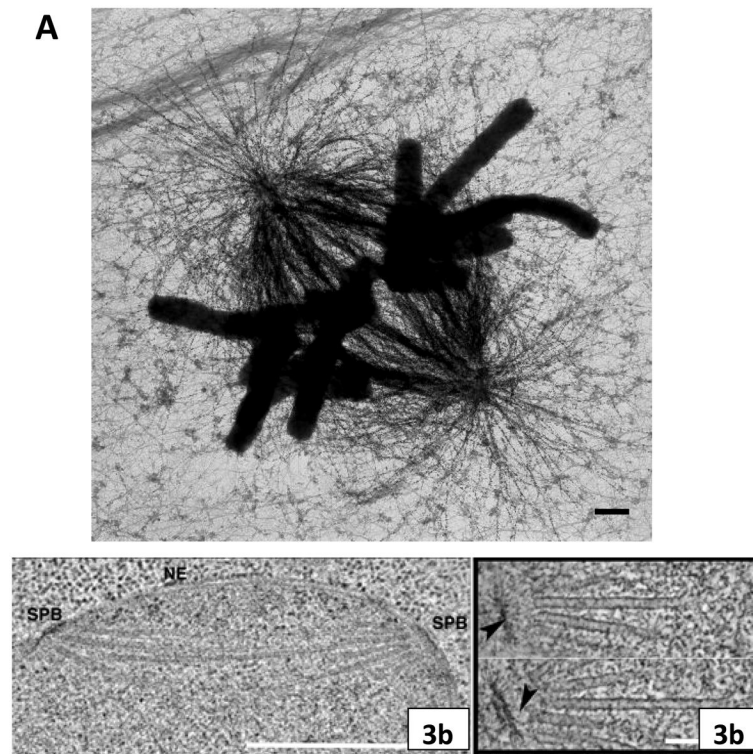


Figure 3. Spindle structure as seen in the electron microscope

Fig. 3A: High voltage electron micrograph of a mitotic mammalian cell, strain PtK₁, in metaphase. Chromosomes are stained with uranyl acetate, MTs by colloidal gold attached to an antibody specific for tubulin. Bundles of MTs converge on each kinetochore; at their other ends they focus at the poles. Astral MTs are clear. Background fibers are mostly intermediate filaments that are not part of the spindle. Bar = 1 μ m. (Micrograph courtesy of Mary Morphew, Univ. Colorado, reprinted from McIntosh, *Mol. Biol. Cell*, Nov., 2011).

Fig. 3B: Electron micrographs of spindles from the yeast, *Saccharomyces cerevisiae*. Left shows the spindle pole bodies (SPB) of a metaphase cell as part of the nuclear envelope (NE); the MTs course through the nucleoplasm, but the chromosomes are not sufficiently condensed to be visible. Bar = 0.5 μ m. Right shows two slices from a tomographic reconstruction of one spindle pole and MTs emanating from it. Their pole-proximal “minus” ends are capped and connected to the pole by slender strands (arrowheads). MT plus ends flare out as the protofilaments bend near the MT end. Bar = 75 nm. (Micrographs courtesy of Eileen O’Toole, Univ. Colorado.)

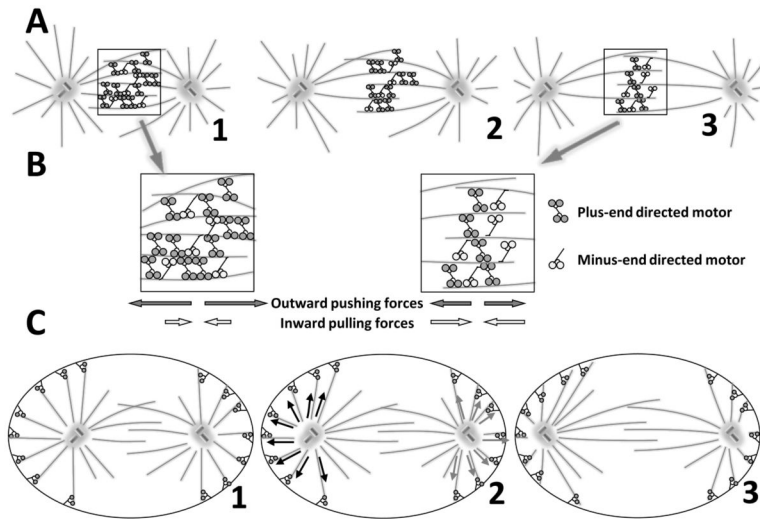


Figure 4. Mechanisms of centrosome separation

Fig. 4A: Spindle pole separation driven by motors that cross-bridge antiparallel MTs. The dominant effect is from plus end-directed motors, which walk in that direction on the pairs of antiparallel MTs to which they are connected. This action pulls the MT ends nearer to one another (Fig. 2A: 2,3). Increased pole separation is permitted by MT elongation. Fig 4B: Diagram showing zoomed in views of the 1st and 3rd states of spindle elongation from Fig. 4A. MT sliding is actually driven by a balance of forces generated by both plus and minus end directed motors. The spindle reaches its steady state length when the number of motors pulling in each direction becomes equal. Fig. 4C: Spindle pole separation driven by pulls from outside the spindle acting on centrosome-associated astral MTs. Dark dots represent sites of cortex-attached minus end-directed motor activity that pulls on the MTs and thus on the spindle poles, forcing them apart. On panel 2 black arrows show forces generated by motors that act on a left pole and grey arrows for forces that act on the right pole.

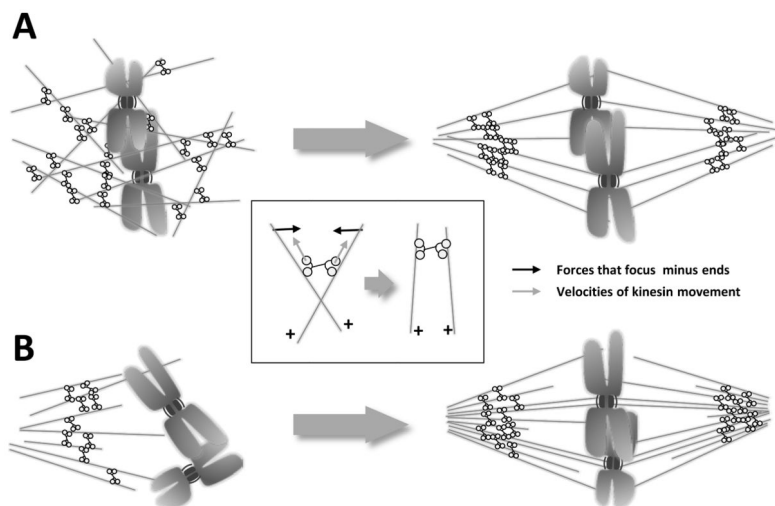


Figure 5. Formation of acentrosomal spindles

Fig. 5A: Spindle formation without centrosomes in extracts from amphibian oocytes. MTs form near chromosomes, due to protein activities regulated by RanGTP. MT-dependent motor activities then rearrange these polymers, clustering their minus ends, so a bipolar spindle is formed. The insert shows how minus-end directed motors that crosslink MTs can align and focuses them into pole-like structure. Fig. 5B: Spindle formation without centrosomes in higher plant cells. There are regions near the nucleus where MT initiation is probable, as with centrosomes, even though no such structures are seen. In this diagram, only one such pole is shown, but in reality, MTs grow into the nuclear region from both sides. As the nuclear envelope disperses, these MTs grow into the nucleoplasm, forming a bipolar MT array in which some polymers interact directly with kinetochores.

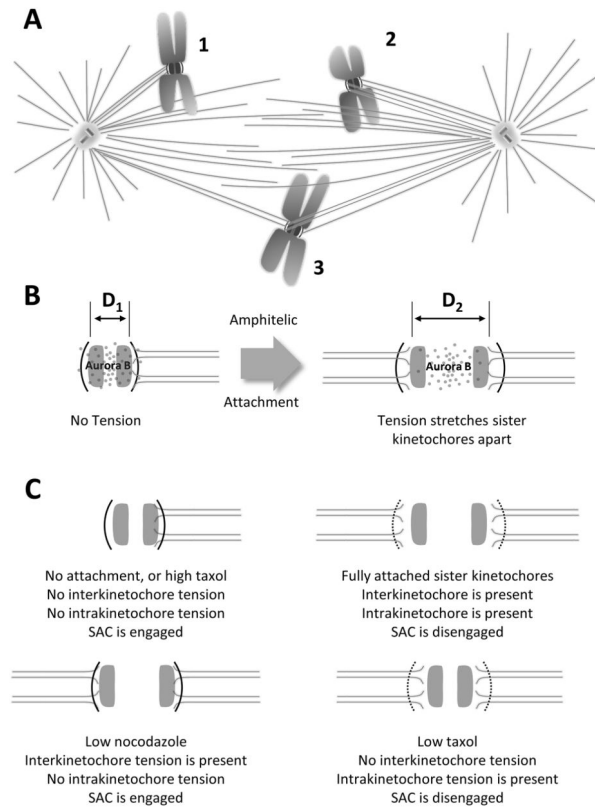


Figure 6. Forming correct chromosome spindle interactions

Fig. 6A: Diagram of chromosomes that are incorrectly connected to the spindle by three different arrangements: 1) monotele, 2) syntele, and 3) merotele. Fig. 6B: Model for a mechanism by which tension at the centromere might reduce the access of Aurora B kinase to substrates important for kinetochore-MT attachment (after Tanaka et al., 2002). Chromosome strain in response to tension moves MT attachment sites out of the region affected by Aurora kinase activity (stippling). Each kinetochore is shown as a layered structure: the dark arc represents the outer kinetochore and shaded bar represents the inner. Fig 6C: Situations where inter- and intra- kinetochore tension can be independently controlled *in vivo* indicated how the SAC might sense tension within a single kinetochore and not across sister kinetochores. Kinetochore are shown as in 6B.

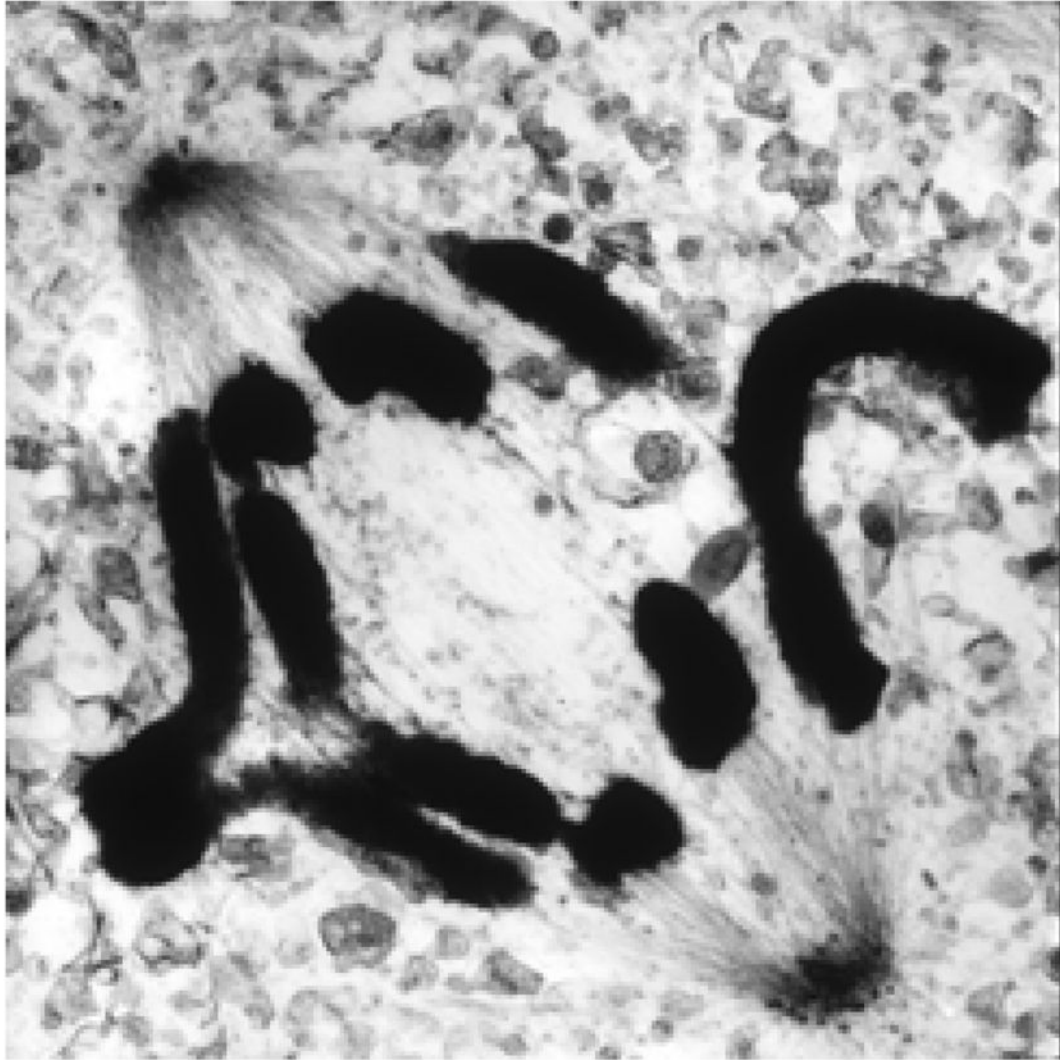


Figure 7. The morphology of anaphase

Electron micrograph of a mammalian cell in anaphase, showing spindle MTs during chromosome separation. This cell has also already started anaphase B. A 250 nm thick slice from a plastic-embedded PtK₁ cell that was lysed in the presence of an equilibrium mixture of tubulin and MTs, then fixed with aldehydes and OsO₄, followed by staining with uranyl acetate and lead citrate and imaging in a high voltage electron microscope. Courtesy of Mary Morpew, Univ. of Colorado.

Table 1

Model	Basic assumptions of the model	Main result of the model
Nedelec 2002. JCB	This model studies the formation of spindles from two MT asters that contain different numbers of dynamic MTs (20–80) and 500 – 15000 motors that can freely diffuse, bind and unbind the aster MTs. Several kinds of motors are considered. One kind has one end with either a plus- or a minus-end directed activity, while the other end can bind MTs but cannot move. Another kind of motor is composed of dimers with motors at both ends. These motors can either be both minus-end directed, both plus-end directed or “heterocomplexes” with one end minus-end directed and the other plus-end directed. A screen was performed to evaluate the relative position of two asters at steady state when the MTs were exposed to various motor complexes.	“Heteromotors”, with plus-end activity on one end and minus- on the other, can yield a steady-state that stabilizes antiparallel MTs. This leads to spindle-like asters interaction. In all other scenarios considered, the steady state asters either fuse or separate.
Schaffner and José, 2006. PNAS	This model explores the formation of spindle like structures directed by the action of chromatin-covered beads in the absence of MT organizing centers, as in experiments of Heald et al., 1996. The model considered a 5µm long linear array of beads covered with chromokinesins, which are plus end-directed kinesins. Around 400 MTs are assumed to be randomly distributed and oriented in the space around beads. A minus-end directed motor that has properties similar to dynein diffuses in the surrounding solution and can bind and unbind MTs. It is assumed that this motor has two heads and it can crosslink MTs. Simulations are results of the activities of these two motors. Kinesin pushes the minus ends of MTs away from the beads. Dynein molecules crosslink MTs and focus their minus ends together into a pole.	The model shows that when MTs are long, the action of these two motors is enough to form a bipolar spindle. Processivity of dynein, and its on and off rates, are key regulators that allow for true bipolar spindle formation, in contrast to multipolar or monopolar structures. For short MTs, an additional mechanism is required.

Table 2

Model	Basic assumptions of the model	Main result of the model
Holy Leibler 1994. PNAS	The search is thought to be composed of two events: a “success” when a growing polymer finds a kinetochore and a “failure” when the polymer grows and shrinks without finding the kinetochore. There is a probability and average duration of both successful and unsuccessful events that depends on the characteristics of polymer dynamics, distance to a kinetochore, and effective interaction area. Two types of polymers are considered and compared: single-state polymers, whose evolution results from a competition between monomers adding and dissociating at well-defined rates, and a two-state polymer that can be either growing or shrinkage with random switches between the two states. Distance to the kinetochore was 2 – 50 μ m. The effective interaction area between the kinetochore and the MT tip was supposed to be determined by rapid fluctuations of the MT tip due to Brownian motion. This area depended on the MT length, and for a 10 μ m long MT, tip fluctuation was somewhat less than 1 μ m.	Single state and two state polymers are compared in their efficiency of capturing kinetochores. Average search time is shown to decrease 103 fold by going from a single state to a two state polymer.
Paul et al., 2009. PNAS	The model includes two spindle poles in 3D, with ~250 MTs emanating from each. MTs are considered as rods with no thickness that undergo dynamic instability. If an MT does not attach to a kinetochore it is likely to shrink all the way to the pole (no rescues), whereupon it starts growing in a different direction. Chromosomes are represented as cylinders randomly distributed and oriented with the nuclear volume. KT's are modeled as cylindrical objects perpendicular to the chromosomes with diameters of 0.88 μ m. They are placed in pairs on opposite sides of the chromosomes and assumed to bind 10 MTs. Chromosomes can move, though their positions do not result from generated forces. Instead, chromosomes jump to new random positions at a rate characteristic of chromosome movement during cell division. Kinetochores are said to nucleate MTs, which increases the effective length of the kinetochore cylinder. For additional assumptions about correction of erroneous attachments mechanisms, see text.	The model supports previous results (Wollman et al., 2005) and suggests that dynamic instability is not sufficient to make “search and capture” fast enough. Random chromosome movements and larger kinetochores bring the time down to realistic values. The model also shows that if there are no specific corrections for erroneous attachments, they prevail over the correct ones.

Table 3

Model	Basic model assumptions and description	Main result of the model
(Joglekar & Hunt, 2002)	Two types of forces that act on a chromosome: PEFs and poleward pulling forces generated at the kinetochores. PEFs are inversely proportional to the square of the distance from the spindle pole. Sister kinetochores are connected by a spring. Interaction between each kinetochore and disassembling KMTs is described as in (Hill, 1985). Kinetochore position is obtained at each time from the force balance equation: $\mathbf{F}_{PEF} + \mathbf{F}_{kin} + \mathbf{F}_{spring} = 0$, where \mathbf{F}_{PEF} = PEF; \mathbf{F}_{kin} = the force generated by KMTs; and \mathbf{F}_{spring} = force that connects sister kinetochores. The chromosome moves toward the pole that faces the kinetochore with more depolymerizing KMTs attached. Depolymerizing KMTs on the opposite pole slip off Hill's sleeve, and the resulting tension between kinetochores is low. As the chromosome moves closer to a pole, the ejection force opposing the leading kinetochore increases, becoming large enough that the additional force generated by even one or two shortening KMTs on the trailing kinetochore is sufficient to cause all shortening MTs on the leading kinetochore to detach. As a result the chromosome abruptly reverses its direction of motion. Thus PEFs bias the direction of chromosome motion toward the spindle equator.	The model accounts for the oscillations of both bi-oriented and mono-oriented chromosomes. It predicts that the amplitude of these oscillations should be bigger for smaller chromosomes. This was tested and found to be true in Ke, et al. 2009.
(Civelekoglu-Scholey et al., 2006)	The attachment of MTs to chromosomes is mediated by molecular motors. The direction and velocity of a chromosome's movement are determined by a balance of forces acting on its kinetochore: $\mathbf{F}_1 + \mathbf{F}_2 + \mathbf{F}_3 + \mathbf{F}_4 + \mathbf{F}_5 = 0$, where 1 = antagonistic effect of plus and minus end directed motors bound to the kinetochore and moving along their KMT tracks; 2 = forces generated by polymerizing KMT plus ends; 3 = elastic tension due to flexible cohesin bonds between sister kinetochores during metaphase, which pull the kinetochores toward one another; 4 = PEFs; and 5 = viscous drag on the chromosome. All KMTs are considered to be independent; they bind and unbind from a kinetochore frequently making it change its direction of movement. Partial synchronization is achieved, however, by using rescues, whose frequency is promoted by tension on an individual KMT attachment. Tension between sister kinetochores tends to prevent a KMT tip from depolymerizing out of the kinetochore, and makes this MT to stay attached at the kinetochore longer.	The model provides a quantitative description of the experimentally observed behavior and rates of metaphase/anaphase kinetochore and KMT dynamics in <i>Drosophila</i> embryos and partially explains phenotypes observed by deletion of dynein. The model makes predictions for chromosome behavior in other organisms.
(Gardner et al., 2005; Sprague et al., 2003)	There is only one KMT for each kinetochore in the budding yeast. Each kinetochore is assumed to be at the KMT tip. Since this cell contains 16 chromosomes, 32 KMTs were simulated (16 from each pole). Each MT is assumed to behave according to the two state dynamic instability model, which is described by 4 parameters: velocities of growth and shrinkage and frequencies of catastrophe and rescue. 3 basic models are considered. First - all 4 parameters are constants; second - rescue frequency is assumed to increase exponentially with the increase in stretch between sister kinetochores; third - rescue or catastrophe frequency is assumed to depend linearly on the concentration of an activated but unknown promoter. Based on the theoretical gradient of this catastrophe promoting molecule, KMT catastrophe frequency was low at the spindle poles and maximal at the spindle equator. Similarly, rescue frequency was high at the spindle poles but decreased rapidly as the KMT plus ends approached the spindle equator.	Only combinations of models that included both tension-dependent rescue or catastrophe frequencies and a gradient in MT dynamics successfully reproduced the experimental results on localization of kinetochores in yeast.

Table 4

Model	Basic model assumptions and description	Main result of the model
(Wollman et al., 2008)	There are four types of the MTs considered based on their plus end interaction scenario. Plus ends of astral MTs are connected to the cellular cortex via dynein molecules that pull on the poles via MTs; plus ends of the interpolar MTs are free and interpolar MTs from opposite poles interact with each other via kinesin-5 and kinesin-14 molecules that antagonize each other pushing poles apart and pulling them back together respectively; plus end of the kinetochore MTs terminate at the kinetochore, where three motors generate forces: dynein, CENP-E and a depolymerizing Kinesin-13; plus end of the chromosome associated MTs terminate at the chromosome arms bound via chromokinesin molecules. The model assumes that each motor can be either in 'on' or 'off' state. This state is assumed to change once during spindle elongation at random time that is used as a model fitting parameter. In total there are 8 mitotic motors considered, and each is assumed to have linear force-velocity relationship. Stall forces and unloaded velocities for each motor are considered to be fitting parameters. In total there are 39 parameters resulting in an astronomical number of possible ways in which multiple mitotic motors of various characteristics can be integrated to build different mitotic spindles. Stochastic optimization process is used to determine sets of parameters that yield plausible dynamics of pole separation. Number of model variants producing realistic behavior is significantly reduced by using experimental results from mutants as constraints. Six distinct molecular scenarios potentially underlying spindle elongation have been identified, each scenario comprising several slightly different model variants.	Changes in scenarios and dynamics of the spindle are regulated by motors being turned on and off at different times. The most general scenario is that metaphase + is maintained by balance of forces between kinesin-5, kinesin-14 and depolymerization activity of kinesin-13 that effectively converts MT sliding to poleward flux. Anaphase B is the result of the switching off of MT depolymerization at the poles by kinesin-13.
(Hara & Kimura, 2009)	The model considers two asters of MTs inside the fixed cellular cortex. Astral MTs reach cortex where they interact with dynein molecules evenly distributed along cortex. Two models are considered. In the first constant-pulling model dynein generates constant force equal for each MT. In the second force-generator-limited model, the amount of force generated at a single MT is proportional to the surface area covered by MT. This assumption is simplified by assuming that the force acting on a single microtubule is proportional to the squared length of the microtubule.	Both models give quantitative agreement with experimental data on the spindle size and spindle elongation rate versus cell length. The force-generator-limited model describes spindle elongation for wt cells; constant-pulling model – in cells lacking G α subunit of the G protein.