

Switching of perpendicularly polarized nanomagnets with spin orbit torque without an external magnetic field by engineering a tilted anisotropy

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Spin orbit torque (SOT) provides an efficient way to significantly reduce the current required for switching nanomagnets. However, SOT generated by an in-plane current cannot deterministically switch a perpendicularly polarized magnet due to symmetry reasons. On the other hand, perpendicularly polarized magnets are preferred over in-plane magnets for high-density data storage applications due to their significantly larger thermal stability in ultrascaled dimensions. Here, we show that it is possible to switch a perpendicularly polarized magnet by SOT without needing an external magnetic field. This is accomplished by engineering an anisotropy in the magnets such that the magnetic easy axis slightly tilts away from the direction, normal to the film plane. Such a tilted anisotropy breaks the symmetry of the problem and makes it possible to switch the magnet deterministically. Using a simple Ta/CoFeB/MgO/Ta heterostructure, we demonstrate reversible switching of the magnetization by reversing the polarity of the applied current. This demonstration presents a previously unidentified approach for controlling nanomagnets with SOT.

spin orbit torque | perpendicular anisotropy | nanomagnets

Spin orbit coupling (SOC) and/or broken inversion symmetry in vertical heterostructures can generate accumulation of spins when a charge current is flowing through them. In doing so, it can exert a torque on an adjacent magnet (1–8). Indeed, high Z metals (Ta, Pt, W, etc.) with strong SOC have been used to inject spin currents into adjacent ferromagnetic layers and thereby to induce magnetic switching, oscillation, domain wall movement, etc. (1–5, 7–9). In a typical heterostructure such as Ta/CoFeB/MgO (from the bottom), an in-plane current flowing in the x direction (electrons flowing in the $-x$ direction) generates $\hat{\sigma} = \hat{y}$ polarized spins that accumulate at the Ta/CoFeB interface. Therefore, if the ferromagnet (CoFeB) is polarized in-plane, the spin accumulation can rotate it to the $+y$ direction by a Slonczewski-like torque $\vec{\tau}_{sl} = \vec{r}_{sl}^0 (\hat{m} \times \hat{\sigma} \times \hat{m})$ (10). The magnet can be switched to $-y$ direction by reversing the polarity of the current. Thus, a deterministic switching is possible by an in-plane current when the magnet is also polarized in-plane. However, if the magnet has a perpendicular magnetic anisotropy (PMA), an in-plane current and resultant in-plane spin accumulation cannot break the reversal symmetry. Consequently, no deterministic switching of the PMA magnet can be obtained. For this reason, an external magnetic field has to be applied in-plane in the same (or opposite) direction of the current flow that breaks the symmetry and makes it possible to switch a PMA magnet with an in-plane current (11). It is, however, desirable to switch the magnetization without needing an external magnetic field. At the same time, PMA is more suitable for scaling magnets to ultrascaled dimensions while retaining reasonable thermal stability (12–14). This means that alternate ways need to be found that can lead to symmetry breaking and make it possible to switch perpendicularly polarized magnets with in-plane current without having to apply an external magnetic field. In this paper, we report on such a scheme where we have fabricated magnetic nanodots from a

heterostructure stack of Ta/CoFeB/MgO such that their easy axes are slightly tilted from the film normal. This tilting breaks the symmetry with respect to an in-plane spin orbit torque (SOT). Indeed, pulsed currents flowing in-plane deterministically switch the magnetization up and down without any external magnetic fields. Notably, in a recent experimental work, Yu et al. (15) have shown that an in-plane current can switch perpendicularly polarized magnet by creating a thickness gradient in the film over the entire substrate. They showed that the structural symmetry breaking thus generated leads to a vertically oriented Rashba field that in turn switches the magnet. The physics of switching in our work, by contrast, is completely different. It is rather the non-conservative Slonczewski torque that is responsible for switching in our scheme and no distinguishable vertical Rashba field could be measured in our devices (as we discuss in details later).

The essential idea is described in Fig. 1. Imagine a magnet that has an easy axis on the x - z plane but slightly tilted away from the z axis to x axis as shown in Fig. 1A. Here, x - y is the film plane. The actual device is made of a heterostructure consisting of Ta/CoFeB/MgO as shown in Fig. 1B. For brevity, when the magnet is in the (x, z) quadrature, we shall call it the “up” position. Similarly when it is in the $(-x, -z)$ quadrature, we shall call it the “down” position. Now imagine a current is applied in the x direction such that a y polarized spin accumulation is generated at the Ta/CoFeB interface. Because y is orthogonal to the x - z plane, regardless of whether or not there is a tilt in the easy axis of the magnet in the x - z plane, no symmetry is broken and

Significance

When a spin orbit torque is applied to a magnet with perpendicular anisotropy, the induced spin accumulation is completely symmetric with respect to the magnetization. As a result, a deterministic switching cannot be achieved unless an external magnetic field is applied to break the symmetry. Here, we show that, by engineering a tilted anisotropy in a magnetic nanodot, the symmetry can effectively be broken and a deterministic switching of perpendicular magnetization can be achieved without needing the external magnetic field. These results are significant for the field of spintronics as the symmetry breaking provides new insight into the physics of spin orbit torque and the switching without a magnetic field could lead to significant impact in high-density storage applications.

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carefully controlling the thickness of the hard mask and Ar⁺ ion milling such that the thickness of the magnet gradually goes to zero on one side (details are described in *Methods* and *SI Appendix*, Fig. S2). Fig. 2A shows an SEM micrograph of a typical nanodot and the underneath Hall structure for transport measurement. Fig. 2B shows an atomic force microscopy image of the nanodot. A wedge-like feature is clearly seen. The depth profile shown in Fig. 2C indicates that MgO with lateral size of around 50 nm in the left part of nanomagnet was completely etched, and CoFeB wedge was created. Because MgO was etched out, the magnetization in the wedge region no longer has a perpendicular anisotropy and will follow the edge of the wedge to minimize magnetostatic energy, thereby providing an overall tilt to the perpendicular magnet, directed away from the edge (17–20). Micromagnetic simulations (*SI Appendix*, section 2) confirm that this indeed is the case, as we will discuss later.

We have performed anisotropic magnetoresistance (AMR) measurements to investigate potential tilting of the easy axis. Note that all of the experiments in this work were done at the ambient temperature. In the coordinate axes as shown in Fig. 3A and B, the dependence of resistivity on the rotation angle θ of the in-plane magnetic field (H_{imp}), as obtained from the AMR measurement, can be written as follows:

$$\rho(\theta) = \rho_{\perp} + (\rho_{\parallel} - \rho_{\perp}) \cos^2 \phi(\theta) = \rho_{\perp} + \Delta\rho \cos^2 \phi(\theta), \quad [1]$$

where ϕ is the angle between the magnetization M and the direction of current and (ρ_{\parallel}) is the resistivity when the magnetization is perpendicular (parallel) to the current direction. Let us consider the case when current is flowing in the $-x$ direction and a H_{imp} of a specific amplitude is rotated in-plane, starting from the $-x$ axis and ending on the $+x$ axis. Because the current is applied along the x axis, from Eq. 1, the resistance is the maximum when the H_{imp} is also along the $\pm x$ direction. Resistance is lower when H_{imp} is rotated away from x axis. Notably, if the magnetic easy axis is purely out-of-plane to begin with, applying the same magnitude of H_{imp} in the $+x$ or $-x$ direction should give the same resistance value. By contrast, if the magnet is tilted on the $+x$ axis, one would expect to see a larger resistance when the H_{imp} is in the $+x$ direction in comparison with when the field is in the $-x$ direction. This is exactly what we observe in our experiments. Starting from an initially up position, the resistance is larger when H_{imp} is in the $+x$ direction compared with when it is in the direction and the resistance goes down away from the x axis (Fig. 3C). On the other hand, starting from an initial down position, the resistance is larger when H_{imp} is in the $-x$ direction (Fig. 3D). This asymmetric resistance curve is also observed when magnetic fields of other magnitudes are applied such as 500, 800, 2,000 Oe, etc. (See *SI Appendix*, section 3, for details.) These data show that the easy axis of the magnetization is slightly tilted from the z axis and lies along (x, z) quadrature to $(-x, -z)$ quadrature. The observed behavior can be modeled well from Eq. 1 including the fact that the lowest value of resistance arises slightly away from the y axis and the sign of the shift depends on the starting polarization of the magnet, i.e., up or down. Finally, we have done AMR measurements with large H_{imp} ($\sim 3,000$ Oe). Fig. 3E and F show AMR data at this field starting from up and down positions, respectively. No asymmetry is observed between the AMR recorded for H_{imp} applied in the $+x$ and $-x$ direction. This is expected because now H_{imp} is strong enough to overcome the tilted anisotropy and the magnet follows the magnetic field. Thus, from the symmetric AMR along the x axis at high field and asymmetric AMR at low field, the existence of a tilted easy axis in the nanodots is confirmed. From the sign of the asymmetry, when the CoFeB wedge is in $-x$ as shown in Fig. 2C, one can also conclude that the easy axis lies on the $x-z$ plane, slightly tilted toward x from z , as shown in Fig. 1A. By AMR measurement, we can also estimate the angle of the tilt, which comes to be $\sim 2^\circ$ with an error of 1° due to the uncertainty in the angle of the in-plane field. At the same time, from comparing the slant in the R-H loop of an anomalous Hall resistance measurement for the wedged dot to a circular dot, another estimate for the tilt amplitude can be made. This measurement gives the tilt angle to be $\sim 5^\circ$ (*SI Appendix*, section 5). Based on these two measurements, we estimate the tilt angle to be between 2° and 5° .

With such a tilted easy axis of the nanodots, the switching scenario described in the preceding paragraphs, namely a deterministic switching from up to down with a current in the y direction and from down to up with a current in the $-y$ direction, should be possible. To demonstrate that we have performed anomalous Hall effect (AHE) measurements. The corresponding hysteresis loop is shown in Fig. 4A. Fig. 4B shows switching data at room temperature in the presence of zero external field. The magnet is saturated to down first by applying a strong vertical magnetic field in $-z$ direction. Next, the field is turned off and a current pulse (~ 1 -s duration) of 1.25 mA (2.5×10^7 A/cm²) is applied in the $-y$ direction. After the pulse is off, an AHE measurement is done by applying a small sense current. The AHE resistance, as shown in the bottom panel of Fig. 4B, shows that the magnet is now up. Next, a current pulse of the

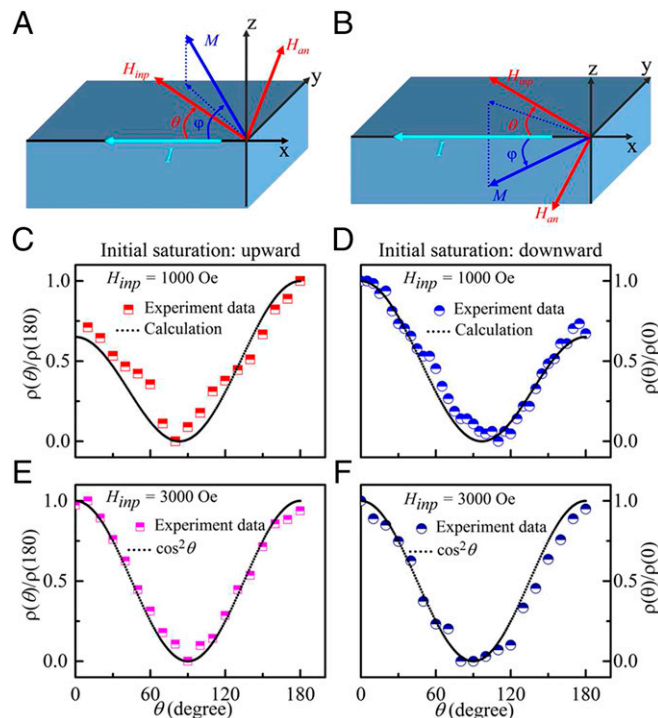


Fig. 3. Estimation of the tilt direction. The schematic configuration of the coordinate system used in the analysis of the AMR data. A and B show the configurations for the initial polarization of up and down, respectively. The sense current for AHE measurement flows along $-x$ direction pointing to the long axis of the nanomagnet. θ is the angle between the directions of the measurement current and the external in-plane field H_{imp} . $\theta = 0$ corresponds to the orientation where the applied field is aligned to the direction of current flow. θ was taken as positive for an anticlockwise rotation of the field from current. ϕ is the angle between M and I . Calculated and measured values of AMR resistivity are shown when $|H_{imp}| = 1,000$ Oe and the magnet was initially polarized (C) upward and (D) downward. Calculated and measured values of AMR resistivity are shown when $|H_{imp}| = 3,000$ Oe, and the magnet was initially polarized (E) upward and (F) downward.

of the nanodot. Second harmonic measurements in our devices show similar values for longitudinal and transverse field components to those that have been reported for similar structures in literature (6, 15, 21) (*SI Appendix, section 8*). To test the existence of an out-of-plane field in our devices, we have measured the threshold current needed to switch the magnetization in presence of an out-of-plane magnetic field. The basic idea is that if a current-induced Rashba field is present in the out-of-plane direction, the R_{AHE} vs. H_z hysteresis loop will be shifted left or right from its symmetric-around-zero position as shown in Fig. 4A. Fig. 4E shows the variation of threshold current flowing in the $+y$ direction (I_y^{th}) as a function of the absolute value of an applied field in the z direction. First, note that at zero current the H_c and $-H_c$ have the same absolute value, which is 175 Oe (note that this is a different device than that shown in Fig. 4A). If a Rashba field is present, a gap between H_z and $-H_z$ should open up as the amplitude of applied current goes up. Therefore, the most important region to look at will be where the current amplitude is close to the threshold amplitude needed for zero magnetic field switching. However, in Fig. 4E, we observe that at these high current levels there is hardly any difference between H_z and $-H_z$. A similar trend is seen when a current is applied in the $-y$ direction in Fig. 4F. This indicates that the contribution of an out-of-plane Rashba field, if any, is minimal in our devices.

In the intermediate range where neither the applied field nor the current alone is strong enough to switch the magnet, a gap between H_z and $-H_z$ can be observed. Note that, in this intermediate range, there is significant difficulty in ascertaining the exact value of the threshold current. Essentially, after setting a specific absolute value of the magnetic field, say 75 Oe, multiple experiments had to be done by varying the amplitude of the applied current pulse to find the threshold value. These repeated trials led to some uncertainty as to the exact value of the threshold current and also heating of the samples, which can also influence threshold current nontrivially. We note that, for a given absolute value of the magnetic field, the difference between the switching currents never exceeded 10% of the applied pulse amplitude (Fig. 4E and F). This is well within the experimental error in this intermediate region. As either the current or the field amplitude increases, the switching becomes easier and the uncertainty in the switching current goes down. In addition, if

one tries to associate a Rashba field for this intermediate region, for $+I_y$ it will be pointing in the $-z$ direction from Fig. 4E. However, for $-I_y$, it is still in the $-z$ direction (Fig. 4F), although reversing the current polarity should have reversed it. Because of the aforementioned two reasons, we believe that the gap observed between the magnetic field values in the intermediate region is an experimental artifact.

Detailed micromagnetic simulations were performed using the OOMMF package (22) to elucidate on the underlying physics (see *SI Appendix, section 2*, for details). A mesh was set up to capture the wedge in the magnet. With a magnet that closely follows the experimental dimensions (total 300 nm in length with 60-nm wedge, 90 nm wide, and 1.2 nm thick in the thickest area), the first set of tests were aimed at finding the equilibrium magnetization of the magnet. Starting from an initial condition of mostly up (with a slight component in the y direction to initiate the dynamics), when the magnet was allowed to relax, it went to an up position with a small component along the x direction, which is away from the wedge (Fig. 5A), exactly as it is observed in the experiments. The opposite happens when one starts from an initial condition of mostly down (Fig. 5B). This effect is understandable from magnetostatics. The magnetization in the wedge area has no perpendicular anisotropy because the MgO has been completely etched off. As a result, they follow the edge of the wedge such that the surface magnetic charge is minimized. Furthermore, they also try to align with the spins in the perpendicular region due to exchange coupling. As a result, they point in the $+x$ direction when bulk of the magnet is $+z$ and vice versa. What is reassuring is the fact that the overall tilt comes to be around 3° (*SI Appendix, section S.2.2*), which is very similar to what we have measured experimentally. Now starting from an initial magnetization along the x axis, as it will be when a current is applied along the y axis, depending on whether it started from $+x$ or $-x$, the magnetization goes to $+z$ or $-z$ direction after the current is turned off, exactly as we have seen in our experiments (Fig. 5C and D). Therefore, the simulations clearly show that the preferred direction of the tilt in our devices comes from minimizing the magnetostatic energy in the structure. Starting from that tilted position, the final state of the magnetization upon the application of current is also correctly predicted by the simulations. It also shows that the switching process is likely to be

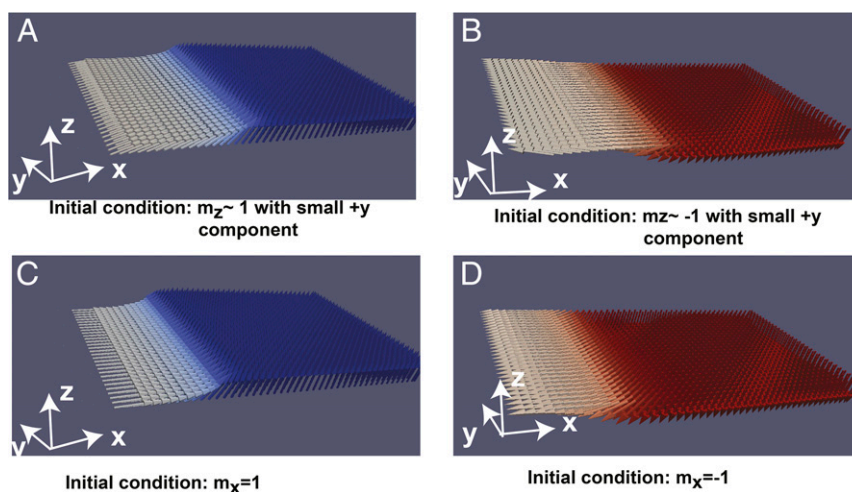


Fig. 5. Micromagnetic simulation. (A) Starting from the magnet saturated in $+z$ state with a small $+y$ component, the system is allowed to relax. The final equilibrium state of the magnet is such that the moments in the wedge region point toward $+x$. Thus, average magnetization tilts toward $+x$. (B) Similarly starting from the magnet saturated in $-z$ state, the magnetization tilts toward $-x$ at equilibrium. (C) Starting from the magnet in $+x$, caused by spin accumulation in $+x$ direction due to $-y$ directed current pulse, the magnet evolves to $+z$ state due to the tilt in the anisotropy axis. (D) Starting from the magnet in $-x$, caused by spin accumulation in $-x$ direction due to $+y$ directed current pulse, the magnet evolves to $-z$ state. The blue arrow indicates that the magnetic moment is in $+z$, whereas the red arrow indicates it is in $-z$ state.

