

# **HHS Public Access**

Author manuscript *Magn Reson Med*. Author manuscript; available in PMC 2017 February 01.

Published in final edited form as:

*Magn Reson Med*. 2016 February ; 75(2): 606–615. doi:10.1002/mrm.25610.

# **Separation of Cellular and BOLD Contributions to T2\* Signal Relaxation**

# **Xialing Ulrich**1 and **Dmitriy A. Yablonskiy**1,\*

<sup>1</sup>Department of Radiology, Washington University, St. Louis MO, 63130, USA

# **Abstract**

**Purpose—**The development of a reliable clinical technique for quantitative measurements of the parameters defining the BOLD effect, i.e. oxygen extraction fraction, OEF, and deoxygenated cerebral blood volume, dCBV, is needed to study brain function in health and disease. Herein we propose such a technique that is based on a widely available gradient recalled echo (GRE) MRI.

**Methods—**Our method is based on GRE with multiple echoes and a model of signal decay (Yablonskiy, MRM 1998) that takes into account microscopic cellular (R2), mesoscopic (BOLD), and macroscopic (background field gradients) contributions to the GRE signal decay with additional accounting for physiologic fluctuations.

**Results—**Using 3T MRI, we generate high resolution quantitative maps of R2\*, R2, R2', and tissue concentration of deoxyhemoglobin, the latter providing a quantitative version of SWI. Our results for OEF and dCBV in gray matter are in a reasonable agreement with the literature data.

**Conclusion—**The proposed approach allows generating high resolution maps of hemodynamic parameters using clinical MRI. The technique can be applied to study such tissues as gray matter, tumors, etc.; however, it requires further development for use in tissues where extra- and intracellular compartments possess substantially different frequencies and relaxation properties (e.g. white matter).

# **Keywords**

BOLD; tumor; stroke; neurological diseases; MRI

# **Introduction**

The presence of deoxygenated blood changes the local magnetic field around and within the blood vessels creating inhomogeneous distribution of local frequency shifts of water molecules. As a result, a faster transverse decay rate (R2\*) of MRI signal can be observed, i.e. BOLD (blood-oxygen-level dependent) effect (1). In a gradient recalled echo (GRE) experiment the contribution of the BOLD effect to the tissue R2\* relaxation rate constant is usually called R2'. The other part of R2\* is usually called R2 which describes the part of the

<sup>\*</sup>Corresponding Author: Dmitriy A. Yablonskiy, PhD, Mallinckrodt Institute of Radiology, Washington University, 4525 Scott Ave. Room 3216, St. Louis MO, 63110, yablonskiyd@mir.wustl.edu, Tel.: +1(314)362-1815, Fax: +1(314)362-0526.

GRE signal decay resulting from water interaction with cellular components of biological tissue. Hence it is commonly assumed that

$$
R2^* = R2 + R2'
$$
 [1]

In other words, the parameter R2 represents the GRE signal decay that would exist in the absence of the BOLD effect. One should keep in mind that the parameter R2 in Eq. [1], that we call "cellular contribution" to the R2\*, generally does not represent irreversible (with respect to the spin echo) part of the MRI signal transverse relaxation. This would only be the case for a single compartment tissue in the static dephasing regime (2). For a multicompartment tissue, Eq. [1] should be written in the form of  $R2^* = R2^* + R2'$ , however in this manuscript we are using traditional notations, as defined in Eq. [1] (see further comments in the Discussion section). Separating the MRI signal decay due to the BOLD effect, that is R2', from R2\*, could provide important information on tissue metabolic properties which is important for understanding normal human brain operation (3) as well as the pathophysiology of neurological disorders such as stroke (4), Alzheimer's disease (5–7), Huntington's disease (8), Parkinson's disease (9,10) and others (11–14). It can also be of great importance for the evaluation of hypoxia within brain tumors and other organs (15,16).

Quite a few techniques have been proposed in the past to use BOLD effect for studying brain hemodynamic properties (see (17) and references therein). One of such techniques, quantitative BOLD (qBOLD) (18), is based on the theory of MRI signal formation in the presence of blood vessel network (2) and experimental method – Gradient Echo Sampling of Spin Echo (GESSE) proposed and tested on phantoms in (19). Similar method was later explored in human brain in (20,21). While results in phantom accurately reflect phantom parameters, the human data (20,21), being encouraging, reflect blood volume and blood oxygenated level only in an "apparent" manner. This was corrected in (18) by introducing a realistic model of human brain tissue that accounted for multi-compartment tissue structure (intra- and extra-cellular water and intravascular blood). qBOLD technique was also validated in animal studies by correlating qBOLD measurements of blood oxygenation level with direct measurements (22).

To fully explore the advantages of the qBOLD technique, one would require very high SNR in the data because of the high demand in the multi-parameter fitting of the qBOLD model (18,19,23–25). The robustness of qBOLD quantification could potentially be improved by independent measurements of some model parameters (26–28). Because of the complexity of the qBOLD mathematical model, many researchers are trying to evaluate the BOLD effect by estimating R2' part of tissue transverse relaxation rate constant  $R2^*$ , by separately measuring R2\* from a gradient recalled echo (GRE) experiment and R2 from a spin echo (SE) experiment (29–33).

However the R2' obtained in such a manner is very sensitive to the measurement protocol (29). Indeed, the signal we measure in the GRE experiment is a combination of signals from multiple tissue compartments having different  $R2^*$  and different frequency shifts (18). For the same reason, the "irreversible" transverse decay rate constant, R2, which represents the

properties mostly related to the tissue cellular content, is also an average parameter depending on details of the SE MRI pulse sequence. Importantly, the "weighted" contributions of the compartments in GRE and SE experiment are different. For example, if we measure R2\* using multi-gradient echo sequence with short TR and TE, then the signal is highly weighted by short T1 and T2 compartments. If one applies multiple spin echo scans to measure R2, which usually requires long TR and SE times, then the signal is highly weighted by long T2 compartments. Thus combining these R2\* and R2 measurements would not allow calculating the true BOLD-related R2'. To limit the problem, one should try using consistent TR and TE for both GRE and SE experiments and sample the signal within a time frame where signal variations between compartments are minimal.

In this paper, we present a simple and robust technique for separating the BOLD-related R2' contribution to R2\* without using SE sequences. Our method relies on the GRE sequence with multiple gradient echoes and theoretical BOLD model (2,19). We also use recentlydeveloped techniques allowing correction of GRE signal for compounding effects of magnetic field inhomogeneities (34) and physiological fluctuations (35). This new method provides high resolution maps of R2\*, R2, R2' and hemodynamic parameters, i.e. tissue concentration of deoxyhemoglobin, oxygen extraction fraction (OEF), and deoxygenated cerebral blood volume (dCBV).

#### **Theory**

The MR signal decay from an imaging voxel in a gradient echo experiment can be represented as a product of several factors (19):

$$
S(TE)=S_0 \cdot \exp(-R2 \cdot TE + i \cdot 2\pi \cdot \Delta f \cdot TE) \cdot F_{BOLD}(TE) \cdot F(TE) \quad [2]
$$

where TE is the gradient echo time,  $R2=1/T2$  is the tissue transverse relaxation rate constant (describing GRE signal decay in the absence of BOLD effect), Δ*f* is the frequency shift (dependent on tissue structure and also macroscopic magnetic field created mostly by tissue/air interfaces), function *FBOLD (TE)* describes GRE signal decay due to the presence of blood vessel network with deoxygenated blood (veins and adjacent to them part of capillaries), and function F(TE) describes effect of macroscopic magnetic field inhomogeneities. In this paper we use voxel spread function (VSF) method (34) for calculating F(TE). Details are provided below in the Methods section.

In our previous publications, we used a statistical approach (2) assuming that a statistically significant amount of randomly positioned and randomly oriented blood vessels are present in each imaging voxel. In this case, *FBOLD (TE)* is (2):

$$
F_{BOLD}(TE) = \exp[-\zeta \cdot f_s(\delta \omega \cdot TE)] \quad [3]
$$

where ζ is the deoxygenated cerebral blood volume fraction (dCBV) and δω is the characteristic frequency determined by the susceptibility difference between deoxygenated blood and surrounding tissue (2):

$$
\delta\omega{=}\frac{4}{3}\pi\cdot\gamma\cdot B_0\cdot Hct\cdot\Delta\chi_0\cdot(1-Y) \quad \text{[4]}
$$

In this equation,  $\chi_0 = 0.27$  *ppm* (36) is the susceptibility difference between fully oxygenated and fully deoxygenated blood,  $Y$  is the blood oxygenated level (with  $Y = 0$  being fully deoxygenated), Hct is the blood hematocrit, and  $\gamma$  is the gyromagnetic ratio. Function  $f_s$ describes the signal decay due to the presence of blood vessel network which was defined in (2). Herein we use a mathematical expression for the function  $f_s$  in terms of a generalized hypergeometric function  $_1F_2$  (17):

$$
f_s(\delta \omega \cdot TE) = {}_1F_2\left(\left[-\frac{1}{2}\right]; \left[\frac{3}{4}, \frac{5}{4}\right]; -\frac{9}{16} (\delta \omega \cdot TE)^2\right) - 1 \quad \text{[5]}
$$

This non-linear function is presented in Fig. 1a. The non-linearity of  $f_s$  is very important for our method because it allows us to separate BOLD (R2') and cellular (R2) contributions to the GRE signal decay. In fact, because of this non-linearity, the R2\* parameter can only be introduced for characterizing the GRE signal decay if  $TE \gg \delta \omega^{-1}$ . In this case the function *fs* becomes linear (2) and we can calculate BOLD-related R2' as:

$$
R2' = \zeta \cdot \delta\omega \quad \text{[6]}
$$

The residual of the function *f<sup>s</sup>* (δω·*t*)−δω·*t* is shown in Figure 1b demonstrating substantially non-linear quasi-oscillating behavior at short times.

In the current paper, we still use an assumption that blood vessels are randomly oriented in each imaging voxels (which, of course, can be violated in voxels containing large vessels) but we are not using statistical requirement of a large number of randomly positioned blood vessels. In this case, *FBOLD (TE)* can be obtained by averaging Eq. [35] in (2) for a single blood vessel with respect to random orientations of the vessel's axis. The result is:

$$
F_{BOLD}(TE) = 1 - \frac{\zeta}{1 - \zeta} \cdot f_s(\delta \omega \cdot TE) + \frac{1}{1 - \zeta} \cdot f_s(\zeta \cdot \delta \omega \cdot TE) \tag{7}
$$

The dependence of the function *FBOLD (TE)* in Eq. [7] on dCBV (ζ) is shown in Fig. 2. One can see that this function's general behavior is similar to the original exponential function in Eq. [3]: it describes a non-linear behavior for a short period of time and quasi-linear after that (as seen in Figure 1). One should keep in mind that Eq. [7] does not represent an exact solution for the extravascular BOLD signal but is rather an empirical formula that provides a reasonable approximation for both limiting cases – small volume fractions of blood vessels (see Fig. 2) and a case of a few large vessels (as it was originally derived in (2)).

By making use of Eqs. [4] and [6] we can also evaluate the concentration of deoxyhemoglobin per unit tissue volume (18):

$$
C_{deoxy} = \frac{3}{4} \cdot \frac{\zeta \cdot \delta \omega \cdot n_{Hb}}{\gamma \cdot \pi \cdot \Delta \chi_0 \cdot B_0}
$$
 [8]

Where  $n_{Hb}$  is the total intracellular Hb concentration equal to  $5.5 \times 10^{-6}$  *mol/mL* (18).

By fitting equation [2] to the complex signal using nonlinear regression algorithm, we are able to find the five parameters:  $S_0$ , R2,  $f$ ,  $\zeta$  and  $\delta\omega$  for each voxel in the brain. In addition, we also calculate a standard relaxation rate constant R2\* using a simplified model without separating BOLD effect but accounting for macroscopic field inhomogeneities:

$$
S(t) {=} S_0 \cdot \exp(-R2^* \cdot TE {+} i \cdot 2\pi \cdot \Delta f \cdot TE) \cdot F(TE) \quad \text{[9]}
$$

#### **Methods**

#### **Image acquisition**

All studies were approved by the Institutional Review Board of Washington University in St. Louis. MRI data were obtained from eight healthy volunteers (4 Female and 4 male, age 23–30) using Siemens 3T Trio MRI Scanner (Siemens, Erlangen, Germany) and a 12 channel head coil. A 3D multi gradient echo sequence was used to obtain the data. Sequence parameters were: resolution  $1\times1\times2$  mm<sup>3</sup> (read, phase, slab), FOV 256 mm $\times192$  mm, repetition time TR = 50ms, 10 gradient echoes with first echo time  $TE_1 = 4$  ms, echo spacing  $TE = 4$ ms. Additional phase stabilization echo (the navigator data) was collected for each line in k-space to correct for image artifacts due to the physiological fluctuations (35). Field inhomogeneity effects were removed by using an advanced version (37) of the voxel spread function (VSF) approach (34). Standard clinical MP RAGE (38) images with high resolution (voxel size:  $0.9 \times 0.9 \times 1.5$  mm<sup>3</sup>) were also collected for segmentation purposes. After data acquisition, raw k-space data were read into MATLAB (The MathWorks, Inc.) for processing.

#### **Image processing**

**Correction for physiologic motion artifacts—**The correction of data for physiological artifacts (i.e. the background frequency fluctuations due to the physiological motions such as breathing, etc.) is accomplished using procedures developed in (35). In brief, the phase stabilization data (navigators) are collected as the last gradient echo for each line in k-space by rewinding phase encoding gradients prior to acquisition. Then each echo train, including the navigator, is Fourier transformed in the read-out direction only and we subtract a corresponding phase drift from the k-space data for each gradient echo time TE in the same echo train before doing the Fourier transform in the phase encoding directions.

$$
S(x, k_y, k_z, TE, ch) \Rightarrow S(x, k_y, k_z, TE, ch) \cdot \exp\left[-i \cdot d\varphi_{nav}(x, k_y, k_z, ch)\frac{TE}{TE_{nav}}\right]
$$
 [10]

Here,  $d\phi_{nav}$  is the phase difference of a current navigator from the 1<sup>st</sup> navigator recorded and *TEnav* is the navigator echo time.

**RF channel combination—**After we correct the k-space data for physiological artifacts, we apply FFT in the phase-encoding directions to get images. 3D spatial Hanning filter is then applied to the data in the image domain. To achieve an optimal signal-to-noise ratio, we use the following equation to combine the data of all channels (39):

$$
S_n(TE) = \sum_{ch=1}^{M} \lambda_{ch} \cdot \overline{S}_n^{ch}(TE_1) \cdot S_n^{ch}(TE); \lambda_{ch} = \frac{1}{M \cdot \varepsilon_{ch}^2} \sum_{ch'=1}^{M} \varepsilon_{ch}^2
$$
 [11]

where the sum is taken over all *M* channels (ch), *S* denotes complex conjugate of *S*,  $\lambda_{ch}$  are weighting parameters and ε*ch* are noise amplitudes (r.m.s.). Index *n* corresponds to the voxel position (*n*=x,y,z). This algorithm allows for the optimal estimation of quantitative parameters, and also removes the initial phase incoherence among the channels (39,40).

**Correction for B0 field inhomogeneities—**Background magnetic field inhomogeneities are known to adversely affect measurement of R2\* and parameters of qBOLD model (19). Herein we use an advanced approach (37) of the VSF (voxel spread function) method (34) that provides a solid platform for correcting images. VSF method relies on the solution of the VSF equation

$$
S_n(TE) = \sum_m \Psi_{nm}(TE) \cdot \sigma_m(TE) \quad [12]
$$

where  $S_n(TE)$  is the measured MRI signal in the voxel *n* at the gradient echo time TE, σ*m*(*TE*) is the ideal "non-contaminated" signal from the voxel *m* that would exist in the absence of Gibbs artifacts and magnetic field inhomogeneities, and the matrix Ψ*nm*(*TE*) defines the signal leakage to the voxel *n* from the neighboring voxels *m* (see definition in (37)). In (37) the VSF equation was solved using similarity approximation – an assumption that the ideal signal  $\sigma_m(TE)$  from the voxel *m* can be approximated as

$$
\sigma_m(TE) = \sigma_n(TE) \cdot |S_m(TE_1)/S_n(TE_1)| \quad \text{[13]}
$$

where  $TE<sub>1</sub>$  is the first gradient echo time. This approximation is justified if neighboring voxels belong to "similar" tissues but it could lead to image artifacts at the boundaries between tissues with different MR properties where the MR signal can decay at neighboring voxels with different transverse relaxation rates *R*2\*. To address this issue we proposed (37) to improve similarity approximation by a substitution

$$
\sigma_m(TE) = \sigma_n(TE) \cdot |S_m(TE_1) \cdot \exp[-R2_m^*(TE - TE_1)] / S_n(TE_1) \cdot \exp[-R2_n^*(TE - TE_1)] | [14]
$$

Since information on  $R2^*$  is not initially available, we use an iterative procedure - the original similarity approximation, Eq. [13], is used to calculate initial values of R2\* that are used subsequently in Eq. [14]. This allows solution of VSF equation in the form of

$$
S_n(TE) = \sigma_n(TE) \cdot F_n(TE) \quad [15]
$$

with F-function calculated based on Eq. [14] in the following form:

$$
F_n(TE) = \sum_m \left| S_m(TE_1) \right| \cdot \exp\left(-R2_m^* \cdot TE\right) / \left[ \left| S_n(TE_1) \right| \cdot \exp\left(-R2_n^* \cdot TE\right) \right]
$$
 [16]

After calculation, Eq. [16] is normalized to 1 at TE = 0.

**Data fitting procedure—**The combined MR signal, in Eq. [11], is analyzed on a voxelby-voxel basis by fitting Eq. [2] with F-function calculated according to the algorithm in Eq. [16]. While the hypergeometric function in Eq. [5] is in the MatLab library, its calculation is very slow. To speed up calculation, we use the Taylor expansion of this up to the  $30<sup>th</sup>$  order (note that only 15 terms are present in this expansion as function is even). The square root of the fitting amplitude produces image of  $S_0$  in Eq. [2] which, for our sequence parameters, is T1-weighted image. Non-linear least square curve fitting algorithm is used to find the optimum parameters. To avoid the singularities in Eq. [2] and save computing time, we set the boundaries for fitting parameters to be within the following values: dCBV is from 0.001 to 0.99, and Y is from 0.1 to 0.9 (10% to 90%). The fitting results are not sensitive to the initial values of the parameters within the boundaries. All codes for image processing and fitting are written in MATLAB (R2011a, the MathWorks Inc, USA). To accelerate the computation, we also apply a brain mask that excludes the skull and the outside regions.

**Image segmentation—Our** method has an advantage of producing inherently coregistered T1W images with good WM/GM contrast and quantitative R2\* maps. R2\* maps have little contrast between GM and WM, so we use our T1W images to improve coregistration with MPRAGE using web-based tools, FSL (41–43). Image segmentation is done using FreeSurfer (Martinos Ctr for Biomed Imaging).

# **Results**

Example of the images and quantitative maps are shown in Figure 3. In this figure, S0 is the calculated signal intensity that would correspond to  $TE = 0$ , which for our sequence parameters represents T1-weighted image. While R2\* map shows some WM/GM contrast, this contrast is substantially enhanced on the R2 map as a result of removing the BOLD contribution. Our results in cortical gray matter for all the subjects are summarized in Table 1.

Examples of the signal decay and fitting curves are shown in Figure 4 for two representative voxels – in white and gray matter. While in GM fitting residuals are small and random, in WM our fitting returns the lower boundaries for Y and dCBV, and residuals show systematic (not related to noise) deviation from zero. This means that the BOLD effect in WM cannot be separated by the model described in Eq. [2]. This will be further discussed in the following section. Hence, we show only results for GM.

Our data for deoxyhemoglobin concentration are quite reasonable. It is important to note that,  $C_{\text{deoxy}},$  Eq. [8], and R2', Eq. [3] (3), are different only by a numerical factor, but  $C_{\text{deoxy}}$ has clear physical meaning. As shown in Figure 3,  $C_{\text{deoxy}}$  is mainly seen in the cortical gray matter that has greater concentration of blood vessels as compared to WM. The concentration  $C_{\text{deoxy}}$  has the highest values in the areas of large veins. To further validate

our results, we compare our results for C<sub>deoxy</sub> with the results obtained by SWI calculated from GEPCI data using an algorithm described in (39). Example of the data is shown in Figure 5. The brighter area of Figure 5(a) means higher concentration of deoxygenated hemoglobin, which is consistent with the location of veins seen in SWI.

# **Discussion**

The development of a reliable technique for quantitative measurements of the vascular characteristics affecting BOLD effect, such as oxygen extraction fraction, OEF, and deoxygenated cerebral blood volume, dCBV, can benefit diagnosis of brain tumors, stroke, and other neurological diseases (27,44–47). In this paper we have demonstrated successful separation of cellular (R2) and vascular  $(R2')$  contributions to the transverse relaxation rate of GRE signal. Our method is based on the MRI signal acquisition using sequence with multiple gradient echoes and analytical model of signal decay, Eq. [2], that takes into account cellular (R2), extravascular (*FBOLD*), and macroscopic (*F*) contributions to the GRE signal decay – Ref. (19). The extravascular BOLD effect is described in the framework of the theory that is an extension of the previously developed theory of BOLD effect (2) to tissue voxels that might have elevated volume fraction of blood vessel network with deoxygenated blood – dCBV. The macroscopic field inhomogeneity effects are accounted for in the framework of the VSF approach (34) with enhanced image contrast (37). Also, artifacts due to the physiologic fluctuations are minimized using technique developed in (35).

The restrictions of applying our method are mostly related to two factors. First, our method of separating cellular from vascular contributions to the tissue transverse relaxation relies on the specific behavior of the vascular (BOLD) contribution to the signal decay as described in Eqs. [5], [7], and illustrated in Figure 1 and Figure 2. As first proposed in (2) and demonstrated on phantom studies in (19), this behavior is not mono-exponential which allows separating it from the mono-exponential R2 decay. However, as is seen in Figure 1, we are sampling only very few initial points of the decay curve where this non-linear behavior is most pronounced. As demonstrated in (25) this might lead to rather large evaluation errors in parameters estimations. However, the use of the improved BOLD function, Eq. [7], slightly alleviates this problem because Eq. [7] shows nonmonoexponential behavior in a broader range of TEs – see Figure 2. We should also mention that the BOLD model (2) that we are using was developed using the static dephasing regime when diffusion phenomena are not important. Computer Monte-Carlo simulations by Dickson et al (48) demonstrate that the diffusion effects do not significantly affect evaluation of BOLD parameters for a GRE experiment with relatively short TEs even for small blood vessels if additional δ*R2* correction is used (see Figure 7 in (48)).

Second, the model in Eq. [2] does not account for the multi-compartment structure of the brain tissue as was accounted for in the qBOLD approach (18,22). However, qBOLD, relying on the model with multiple parameters (compartments volume fractions, relaxation rate constants and frequencies) requires very high SNR and a very big number of sampling points which can be achieved only for images with rather big voxels (typical resolution  $4\times4\times4$  mm<sup>3</sup>). In this paper we use images with much higher resolution -  $1\times1\times2$  mm<sup>3</sup> and

found parameters in GM rather similar to the qBOLD technique – see Table 2. This is because in GM relaxation rate constants in the intracellular and extracellular compartments are rather similar (18) and the frequency shift between compartments (about 5 Hz at 3T) creates only very small additional signal decay in the time scale of 40 ms where our data are acquired. Indeed, as we can see from Figure 10 in (17), the multi-compartment structure of MR signal becomes pronounced only at gradient echo times around and beyond 40 ms. However, this is not the case for the GRE experiment in WM where contribution from the myelin water to the GRE signal is important as it has very strong signal decay and very high frequency shift from intra- and extra-cellular water (49–52). These two effects cause substantial non-linear behavior of the signal decay curve that is not accounted for in our model. Especially the large R2 of the myelin water compartment causes the signal to curve in the opposite direction to the curvature caused by the BOLD effect shown in Figure 1. Thus, the existence of myelin "interferes" with the BOLD effect. And indeed, our model fails in most parts of WM as it is demonstrated in Figure 4. This was not the case for qBOLD approach in (18) where data were acquired using GESSE (gradient echo sampling of spin echo) sequence (19) with the spin echo time of 60 ms where contribution of myelin water signal is practically negligible due to its short T2.

In our model, Eq. [2], we also neglected contribution of the intravascular signal. While contribution of this signal is proportional to dCBV (as the contribution of extravascular  $BOLD - Eq. (7)$ , the vascular signal decays much faster due to the short blood  $T2* (30 \text{ ms})$ for Y=0.6 at 3T) and additional signal decay due to the presence of randomly oriented blood vessels with the orientation-dependent frequencies (53). The latter leads to an additional signal decay with the time scale of about  $3\pi/(2\delta\omega)$  (53) which is about 25 ms for Y=0.6 at 3T. These two mechanisms lead to substantially faster intravascular signal decay rate (apparent T2\* less than 15 ms) as compared to T2 of GM that defines extravascular BOLD contribution to the GRE signal decay (about 60 ms per Tables 1 and 2).

Despite these potential sources of errors, our results for the model parameters in GM are quite reasonable. Mean value of dCBV is 4.6% which includes contribution from blood vessels of different sizes, including large veins. The peak value of the distribution is much smaller – it is about 2.7% as seen in Table 1 and the example shown in Figure 6. This value is more consistent with the literature data for the total blood volume (arterial + venous) of 4.5%, 3%, and 2.7% reported by Carpenter et al. (54), Diringer et al. (55), and Derdeyn et al. (4), correspondingly. The mean value of blood oxygenation level in our subjects is 53% and the peak value is 73%. This values are in the range of MRI estimates of Y about 65% based on the measurements in the superior sagittal sinus by Lu and Ge (56), Jain et al (57) and Barhoum et al (58). Assuming that oxygenated level of the arterial blood is about 100%, we can estimate oxygen extraction fraction, OEF, from our data to be between 47% of mean value and 25% of peak value which is also in the range of the literature data. Indeed, Carpenter et al. (54), Yamauchi et al. (59), Diringer et al. (55), and Raichle et al. (60) reported OEF values of 35%, 42.6%, 41%, and 40%, correspondingly. An example of the 3D histogram of the voxels as a function of blood oxygenation level Y, and dCBV is shown in Figure 6. The histogram has a peak with Y around 0.55 and dCBV around 0.03, however, bigger dCBV, presumably originating from the voxels with large blood vessels, also present in this asymmetric histogram contributing to higher mean dCBV values.

The mean value of  $C_{\rm deoxy}$  over the whole cortex that we found is 21.1 µM with the peak value of 13  $\mu$ M and the average value in large veins about 200  $\mu$ M (red color in Figure 5). For pure blood with Hct = 40%, Y = 0.6, we can estimate  $C_{deoxy} = Hct \cdot (1-Y) \cdot n_{Hb}$  to be about 880 µM. This number is higher than our measurements in large veins which is expected – our static dephasing regime model of BOLD signal can only describe signal dephasing in the blood in an apparent manner. More sophisticated models for the blood signal (see detail discussion in (17) can be implemented in our approach for obtaining quantitative measurements of blood oxygenation level in large veins. For voxels with dCBV between 2% and 3% we can expect  $C_{\text{decay}}$  to be between 18  $\mu$ M and 26  $\mu$ M which is consistent with our measurements. It is also consistent with  $C_{\text{decay}}$  of 14.3  $\mu$ M measured in the neonatal brain surface by optical method  $(61)$ .

Several papers have reported on different methods of separating R2 and R2′ contributions to the R2\* signal decay. Their results are summarized in Table 2 together with our data. It is often assumed that the transverse relaxation of MR signal can be characterized by simple exponential functions:

$$
S(TE)_{_{GRE}} = S_0 \cdot \exp(-R2^* \cdot TE + i \cdot 2\pi \cdot \Delta f \cdot TE) \ S(TE)_{_{SE}} = S_0 \cdot \exp(-R2 \cdot TE); R2^* = R2 + R2' \quad [17]
$$

where R2 is assumed to represent an irreversible (with respect to  $180^{\circ}$  refocusing RF pulse) part of the signal decay and R2′ is assumed to represent a reversible part of the signal decay. However, as we already noticed in the Introduction, this assumption is not strictly correct because signal decay in biological tissue in either SE or GRE experiment is never monoexponential due to a number of reasons. First, biological tissue structure is always multi-compartment and not accounting for this effect can lead to substantial discrepancies in R2 and R2′ estimates even if considering them as apparent parameters. For example, Paling et al. (31) report R2′ values of 9.2 s<sup>-1</sup> and Sedlacik et al. (33) report R2′ values of 7.1 s<sup>-1</sup> obtained with 3T MRI. These values are too big to be attributed to vascular contribution to the R2\* relaxation rate constant. The issue here is that in these papers the authors independently measured R2\* and R2 assuming a monoexponential decay for both experiments ignoring the fact that the relative contributions of multi-compartments to the GRE and SE experiments are different (the compartments have different relaxation rate constants and different frequencies). Secondly, the extravascular BOLD contribution to the GRE signal decay is not monoexponential, as described by our Eqs. [7], [5] and illustrated in Figure 1 and Figure 2. Also, SE measurements can be affected by the SE sequence timing (62). Ignoring this effect can also lead to errors in estimating BOLD contribution – compare different ways of estimating R2′ presented in Table 1 and Table 2. Third, accounting for macroscopic field inhomogeneities is also essential for accurate estimates of tissue R2′. In our paper we use a VSF approach (34) with enhanced image contrast (37) as described by Eqs. [12]–[16]. This is an essential step in accurate estimate of BOLD contribution because macroscopic magnetic field inhomogeneities affect GRE signal decay in a non-monoexponential manner (34) which can affect accuracy of all parameters estimates (19,34). Importantly, as was demonstrated in (34), an often-used sinc-function approach to correct for in-plane field inhomogeneities in 2D acquisition and for any direction in 3D acquisition, is not appropriate and only leads to additional errors in parameters estimates.

# **Conclusion**

In this paper we have developed a technique for separating cellular and BOLD-related contributions to the R2\* MRI signal decay. Our method is based on MRI signal acquisition using a single sequence with multiple gradient echoes (no spin echo sequences are involved) and analytical model of the signal decay that takes into account microscopic (cellular), mesoscopic (extravascular BOLD), and macroscopic (background field gradients) contributions to the GRE signal decay. The method provides reliable and realistic results for GM and is expected to work for any tissue where multiple compartments have similar magnetic properties (i.e. relaxation times and frequencies) which is the case in GM. We can also hypothesize that tissue tumors might satisfy such conditions. However, the method does not produce reliable BOLD results in WM due to the presence of myelin water compartment with substantially different magnetic properties. An approach that would account for this multi-compartment structure is necessary to provide accurate estimates in WM. Such an approach can also improve our results in GM.

### **Acknowledgement**

The authors are grateful to Drs. Alexander Sukstanskii and Anne Cross for helpful discussions and Dr. Jie Wen for help with data acquisition. This work was supported by a grant 5R01NS055963.

### **References**

- 1. Ogawa S, Lee TM, Kay AR, Tank DW. Brain magnetic resonance imaging with contrast dependent on blood oxygenation. Proceedings of the National Academy of Sciences of the United States of America. 1990; 87(24):9868–9872. [PubMed: 2124706]
- 2. Yablonskiy DA, Haacke EM. Theory of NMR signal behavior in magnetically inhomogeneous tissues: the static dephasing regime. Magnetic resonance in medicine : official journal of the Society of Magnetic Resonance in Medicine / Society of Magnetic Resonance in Medicine. 1994; 32(6): 749–763.
- 3. Raichle ME, Mintun MA. Brain work and brain imaging. Annu Rev Neurosci. 2006; 29:449–476. [PubMed: 16776593]
- 4. Derdeyn CP, Videen TO, Yundt KD, Fritsch SM, Carpenter DA, Grubb RL, Powers WJ. Variability of cerebral blood volume and oxygen extraction: stages of cerebral haemodynamic impairment revisited. Brain. 2002; 125(Pt 3):595–607. [PubMed: 11872616]
- 5. Iadecola C. Neurovascular regulation in the normal brain and in Alzheimer's disease. Nat Rev Neurosci. 2004; 5(5):347–360. [PubMed: 15100718]
- 6. Iadecola C. Rescuing troubled vessels in Alzheimer disease. Nat Med. 2005; 11(9):923–924. [PubMed: 16145570]
- 7. Vlassenko AG, Vaishnavi SN, Couture L, Sacco D, Shannon BJ, Mach RH, Morris JC, Raichle ME, Mintun MA. Spatial correlation between brain aerobic glycolysis and amyloid-beta (Abeta) deposition. Proc Natl Acad Sci U S A. 2010; 107(41):17763–17767. [PubMed: 20837517]
- 8. Powers WJ, Videen TO, Markham J, McGee-Minnich L, Antenor-Dorsey JV, Hershey T, Perlmutter JS. Selective defect of in vivo glycolysis in early Huntington's disease striatum. Proc Natl Acad Sci U S A. 2007; 104(8):2945–2949. [PubMed: 17299049]
- 9. Beal MF. Mitochondrial dysfunction in neurodegenerative diseases. Biochim Biophys Acta. 1998; 1366(1–2):211–223. [PubMed: 9714810]
- 10. Schapira AH. Mitochondrial dysfunction in neurodegenerative disorders. Biochim Biophys Acta. 1998; 1366(1–2):225–233. [PubMed: 9714816]

- 11. Leenders KL, Frackowiak RS, Quinn N, Marsden CD. Brain energy metabolism and dopaminergic function in Huntington's disease measured in vivo using positron emission tomography. Mov Disord. 1986; 1(1):69–77. [PubMed: 2973559]
- 12. Santens P, De Reuck J, Crevits L, Decoo D, Lemahieu I, Strijckmans K, Goethals P. Cerebral oxygen metabolism in patients with progressive supranuclear palsy: a positron emission tomography study. Eur Neurol. 1997; 37(1):18–22. [PubMed: 9018027]
- 13. Shishido F, Uemura K, Inugami A, Tomura N, Higano S, Fujita H, Sasaki H, Kanno I, Murakami M, Watahiki Y, Nagata K. Cerebral oxygen and glucose metabolism and blood flow in mitochondrial encephalomyopathy: a PET study. Neuroradiology. 1996; 38(2):102–107. [PubMed: 8692416]
- 14. Tanaka M, Kondo S, Okamoto K, Hirai S. [Cerebral perfusion and oxygen metabolism in Parkinson's disease: positron emission tomographic study using oxygen-15-labeled CO2 and O2]. Nihon Rinsho. 1997; 55(1):218–221. [PubMed: 9014453]
- 15. Tatum JL, Kelloff GJ, Gillies RJ, Arbeit JM, Brown JM, Chao KS, Chapman JD, Eckelman WC, Fyles AW, Giaccia AJ, Hill RP, Koch CJ, Krishna MC, Krohn KA, Lewis JS, Mason RP, Melillo G, Padhani AR, Powis G, Rajendran JG, Reba R, Robinson SP, Semenza GL, Swartz HM, Vaupel P, Yang D, Croft B, Hoffman J, Liu G, Stone H, Sullivan D. Hypoxia: importance in tumor biology, noninvasive measurement by imaging, and value of its measurement in the management of cancer therapy. Int J Radiat Biol. 2006; 82(10):699–757. [PubMed: 17118889]
- 16. Davda S, Bezabeh T. Advances in methods for assessing tumor hypoxia in vivo: implications for treatment planning. Cancer Metastasis Rev. 2006; 25(3):469–480. [PubMed: 17029029]
- 17. Yablonskiy DA, Sukstanskii AL, He X. Blood oxygenation level-dependent (BOLD)-based techniques for the quantification of brain hemodynamic and metabolic properties - theoretical models and experimental approaches. NMR in biomedicine. 2013; 26(8):963–986. [PubMed: 22927123]
- 18. He X, Yablonskiy DA. Quantitative BOLD: mapping of human cerebral deoxygenated blood volume and oxygen extraction fraction: default state. Magnetic resonance in medicine : official journal of the Society of Magnetic Resonance in Medicine / Society of Magnetic Resonance in Medicine. 2007; 57(1):115–126.
- 19. Yablonskiy DA. Quantitation of intrinsic magnetic susceptibility-related effects in a tissue matrix. Phantom study. Magnetic resonance in medicine : official journal of the Society of Magnetic Resonance in Medicine / Society of Magnetic Resonance in Medicine. 1998; 39(3):417–428.
- 20. An H, Lin W. Quantitative measurements of cerebral blood oxygen saturation using magnetic resonance imaging. Journal of cerebral blood flow and metabolism : official journal of the International Society of Cerebral Blood Flow and Metabolism. 2000; 20(8):1225–1236.
- 21. An H, Lin W, Celik A, Lee YZ. Quantitative measurements of cerebral metabolic rate of oxygen utilization using MRI: a volunteer study. NMR in biomedicine. 2001; 14(7–8):441–447. [PubMed: 11746936]
- 22. He X, Zhu M, Yablonskiy DA. Validation of oxygen extraction fraction measurement by qBOLD technique. Magnetic resonance in medicine : official journal of the Society of Magnetic Resonance in Medicine / Society of Magnetic Resonance in Medicine. 2008; 60(4):882–888.
- 23. Sohlin, M.; Schad, LR. Theoretical prediction of parameter stability in quantitative BOLD MRI: dependence on SNR and sequence parameters. Proceedings of 18th Annual Meeting of ISMRM; Honolulu, Hawaii. 2009. p. 1623
- 24. Sedlacik J, Reichenbach JR. Validation of Quantitative Estimation of Tissue Oxygen Extraction Fraction and Deoxygenated Blood Volume Fraction in Phantom and In Vivo Experiments by Using MRI. Magnetic resonance in medicine : official journal of the Society of Magnetic Resonance in Medicine / Society of Magnetic Resonance in Medicine. 2010; 63(4):910–921.
- 25. Wang X, Sukstanskii AL, Yablonskiy DA. Optimization strategies for evaluation of brain hemodynamic parameters with qBOLD technique. Magn Reson Med. 2013; 69(4):1034–1043. [PubMed: 22623013]
- 26. Christen T, Lemasson B, Pannetier N, Farion R, Segebarth C, Remy C, Barbier EL. Evaluation of a quantitative blood oxygenation level-dependent (qBOLD) approach to map local blood oxygen saturation. NMR in biomedicine. 2011; 24(4):393–403. [PubMed: 20960585]

- 27. Christen T, Lemasson B, Pannetier N, Farion R, Remy C, Zaharchuk G, Barbier EL. Is T2\* enough to assess oxygenation? Quantitative blood oxygen level-dependent analysis in brain tumor. Radiology. 2012; 262(2):495–502. [PubMed: 22156990]
- 28. Christen T, Zaharchuk G, Pannetier N, Serduc R, Joudiou N, Vial JC, Remy C, Barbier EL. Quantitative MR estimates of blood oxygenation based on T2\*: a numerical study of the impact of model assumptions. Magn Reson Med. 2012; 67(5):1458–1468. [PubMed: 22183768]
- 29. Ni W, Christen T, Zun Z, Zaharchuk G. Comparison of R2' measurement methods in the normal brain at 3 tesla. Magnetic resonance in medicine : official journal of the Society of Magnetic Resonance in Medicine / Society of Magnetic Resonance in Medicine. 2014
- 30. Gelman N, Gorell JM, Barker PB, Savage RM, Spickler EM, Windham JP, Knight RA. MR imaging of human brain at 3.0 T: preliminary report on transverse relaxation rates and relation to estimated iron content. Radiology. 1999; 210(3):759–767. [PubMed: 10207479]
- 31. Paling D, Tozer D, Wheeler-Kingshott C, Kapoor R, Miller DH, Golay X. Reduced R2' in multiple sclerosis normal appearing white matter and lesions may reflect decreased myelin and iron content. Journal of neurology, neurosurgery, and psychiatry. 2012; 83(8):785–792.
- 32. West J, Blystad I, Engstrom M, Warntjes JB, Lundberg P. Application of quantitative MRI for brain tissue segmentation at 1.5 T and 3.0 T field strengths. PloS one. 2013; 8(9):e74795. [PubMed: 24066153]
- 33. Sedlacik J, Boelmans K, Lobel U, Holst B, Siemonsen S, Fiehler J. Reversible, irreversible and effective transverse relaxation rates in normal aging brain at 3T. NeuroImage. 2014; 84:1032– 1041. [PubMed: 24004692]
- 34. Yablonskiy DA, Sukstanskii AL, Luo J, Wang X. Voxel spread function method for correction of magnetic field inhomogeneity effects in quantitative gradient-echo-based MRI. Magn Reson Med. 2013; 70(5):1283–1292. [PubMed: 23233445]
- 35. Wen J, Cross AH, Yablonskiy DA. On the role of physiological fluctuations in quantitative gradient echo MRI: implications for GEPCI, QSM, SWI. Magn Reson Med. 2014
- 36. Spees WM, Yablonskiy DA, Oswood MC, Ackerman JJ. Water proton MR properties of human blood at 1.5 Tesla: magnetic susceptibility,  $T(1)$ ,  $T(2)$ ,  $T^*(2)$ , and non-Lorentzian signal behavior. Magn Reson Med. 2001; 45(4):533–542. [PubMed: 11283978]
- 37. Ulrrich, X.; Yablonskiy, DA. Enhancing Image Contrast in Human Brain by Voxel Spread Function Method. ISMRM Annual Meeting; Milan, Italy. 2014. p. 3197
- 38. Mugler JP, Brookeman JR. Three-dimensional magnetization-prepared rapid gradient-echo imaging (3D MP RAGE). Magnetic Resonance in Medicine. 1990; 15(1):152–157. [PubMed: 2374495]
- 39. Luo J, Jagadeesan BD, Cross AH, Yablonskiy DA. Gradient Echo Plural Contrast Imaging Signal model and derived contrasts: T2\*, T1, Phase, SWI, T1f, FST2\*and T2\*-SWI. NeuroImage. 2012; 60(2):1073–1082. [PubMed: 22305993]
- 40. Quirk JD, Sukstanskii AL, Bretthorst GL, Yablonskiy DA. Optimal decay rate constant estimates from phased array data utilizing joint Bayesian analysis. J Magn Reson. 2009; 198(1):49–56. [PubMed: 19181549]
- 41. Jenkinson M, Beckmann CF, Behrens TE, Woolrich MW, Smith SM. Fsl. NeuroImage. 2012; 62(2):782–790. [PubMed: 21979382]
- 42. Woolrich MW, Jbabdi S, Patenaude B, Chappell M, Makni S, Behrens T, Beckmann C, Jenkinson M, Smith SM. Bayesian analysis of neuroimaging data in FSL. NeuroImage. 2009; 45(1 Suppl):S173–S186. [PubMed: 19059349]
- 43. Smith SM, Jenkinson M, Woolrich MW, Beckmann CF, Behrens TEJ, Johansen-Berg H, Bannister PR, De Luca M, Drobnjak I, Flitney DE, Niazy RK, Saunders J, Vickers J, Zhang YY, De Stefano N, Brady JM, Matthews PM. Advances in functional and structural MR image analysis and implementation as FSL. NeuroImage. 2004; 23:S208–S219. [PubMed: 15501092]
- 44. Christen T, Bolar DS, Zaharchuk G. Imaging brain oxygenation with MRI using blood oxygenation approaches: methods, validation, and clinical applications. AJNR American journal of neuroradiology. 2013; 34(6):1113–1123. [PubMed: 22859287]

- 45. Blockley NP, Griffeth VE, Simon AB, Buxton RB. A review of calibrated blood oxygenation level-dependent (BOLD) methods for the measurement of task-induced changes in brain oxygen metabolism. NMR in biomedicine. 2013; 26(8):987–1003. [PubMed: 22945365]
- 46. Brown JM, Wilson WR. Exploiting tumour hypoxia in cancer treatment. Nature reviews Cancer. 2004; 4(6):437–447. [PubMed: 15170446]
- 47. Siemonsen S, Fitting T, Thomalla G, Horn P, Finsterbusch J, Summers P, Saager C, Kucinski T, Fiehler J. T2' imaging predicts infarct growth beyond the acute diffusion-weighted imaging lesion in acute stroke. Radiology. 2008; 248(3):979–986. [PubMed: 18647849]
- 48. Salomir R, De Senneville BD, Moonen CTW. A fast calculation method for magnetic field inhomogeneity due to an arbitrary distribution of bulk susceptibility. Concepts in Magnetic Resonance Part B-Magnetic Resonance Engineering. 2003; 19B(1):26–34.
- 49. Sati P, van Gelderen P, Silva AC, Reich DS, Merkle H, de Zwart JA, Duyn JH. Microcompartment specific T2\* relaxation in the brain. NeuroImage. 2013; 77(0):268–278. [PubMed: 23528924]
- 50. Wharton S, Bowtell R. Fiber orientation-dependent white matter contrast in gradient echo MRI. Proceedings of the National Academy of Sciences of the United States of America. 2012; 109(45): 18559–18564. [PubMed: 23091011]
- 51. Sukstanskii AL, Yablonskiy DA. On the role of neuronal magnetic susceptibility and structure symmetry on gradient echo MR signal formation. Magn Reson Med. 2014; 71(1):345–353. [PubMed: 23382087]
- 52. Yablonskiy DA, Sukstanskii AL. Biophysical mechanisms of myelin-induced water frequency shifts. Magn Reson Med. 2014; 71(6):1956–1958. [PubMed: 24700617]
- 53. Schmainda KM, Rand SD, Joseph AM, Lund R, Ward BD, Pathak AP, Ulmer JL, Baddrudoja MA, Krouwer HGJ. Characterization of a first-pass gradient-echo spin-echo method to predict brain tumor grade and angiogenesis. American Journal of Neuroradiology. 2004; 25(9):1524–1532. [PubMed: 15502131]
- 54. Carpenter DA, Grubb RL Jr, Tempel LW, Powers WJ. Cerebral oxygen metabolism after aneurysmal subarachnoid hemorrhage. Journal of cerebral blood flow and metabolism : official journal of the International Society of Cerebral Blood Flow and Metabolism. 1991; 11(5):837– 844.
- 55. Diringer MN, Yundt K, Videen TO, Adams RE, Zazulia AR, Deibert E, Aiyagari V, Dacey RG Jr, Grubb RL Jr, Powers WJ. No reduction in cerebral metabolism as a result of early moderate hyperventilation following severe traumatic brain injury. J Neurosurg. 2000; 92(1):7–13. [PubMed: 10616076]
- 56. Lu H, Ge Y. Quantitative evaluation of oxygenation in venous vessels using T2-Relaxation-Under-Spin-Tagging MRI. Magnetic resonance in medicine : official journal of the Society of Magnetic Resonance in Medicine / Society of Magnetic Resonance in Medicine. 2008; 60(2):357–363.
- 57. Jain V, Magland J, Langham M, Wehrli FW. High temporal resolution in vivo blood oximetry via projection-based T2 measurement. Magnetic resonance in medicine : official journal of the Society of Magnetic Resonance in Medicine / Society of Magnetic Resonance in Medicine. 2013; 70(3): 785–790.
- 58. Barhoum S, Rodgers ZB, Langham M, Magland JF, Li C, Wehrli FW. Comparison of MRI methods for measuring whole-brain venous oxygen saturation. Magn Reson Med. 2014
- 59. Yamauchi H, Fukuyama H, Nagahama Y, Nabatame H, Ueno M, Nishizawa S, Konishi J, Shio H. Significance of increased oxygen extraction fraction in five-year prognosis of major cerebral arterial occlusive diseases. J Nucl Med. 1999; 40(12):1992–1998. [PubMed: 10616876]
- 60. Raichle ME, MacLeod AM, Snyder AZ, Powers WJ, Gusnard DA, Shulman GL. A default mode of brain function. Proceedings of the National Academy of Sciences of the United States of America. 2001; 98(2):676–682. [PubMed: 11209064]
- 61. Cooper CE, Elwell CE, Meek JH, Matcher SJ, Wyatt JS, Cope M, Delpy DT. The noninvasive measurement of absolute cerebral deoxyhemoglobin concentration and mean optical path length in the neonatal brain by second derivative near infrared spectroscopy. Pediatric research. 1996; 39(1): 32–38. [PubMed: 8825383]

62. Kiselev VG. Transverse relaxation effect of MRI contrast agents: a crucial issue for quantitative measurements of cerebral perfusion. Journal of magnetic resonance imaging : JMRI. 2005; 22(6): 693–696. [PubMed: 16261568]





(a) Function  $f_s$ , Eq. [5], versus ( $\delta \omega \cdot t$ ). The dots are values at our echoes ( $t=T E_n$ ) assuming Y= 60%. (b) *f<sup>s</sup>* (δω·*t*)−(δω·*t*) versus (δω·*t*), where *t* is the time after RF excitation pulse.

Ulrich and Yablonskiy Page 17



**Figure 2.** 

The dependence of the function *FBOLD*, describing the BOLD-related contribution to the extra-vascular signal decay, Eq. [7], on the deoxygenated blood volume fraction ζ. The function *f*<sup>s</sup> defining BOLD signal in a standard exponential form, Eq. [3], is also shown. The plots demonstrate that Eq. [7] is more sensitive to  $\zeta$  as compared to function  $f_s$  in Eq. [3].

Ulrich and Yablonskiy Page 18



#### **Figure 3.**

Fitting results and parameter maps from one subject.  $S_0$  is a T1-weighted image which is a square root of the fitting amplitude; R2\* is the apparent transverse decay rate obtained by fitting to data a mono-exponential function, Eq. [9], without accounting for the BOLD effect; R2 is the transverse decay rate after removal the BOLD effect, Eq. [2]; also shown are the R2' map (defined by the differences of R2\* and R2) and  $C_{\text{deoxy}}$  map, Eq. [8]. Both these maps have higher values at the location of large veins. Residual map shows values mostly less than 1% of the total signal.

![](_page_18_Figure_2.jpeg)

#### **Figure 4.**

Example of the fitting curves and residuals. The fitting results for gray matter voxel (short arrow) are R2=13.0 s<sup>-1</sup>, f=0.41 Hz,  $\delta \omega$ =201 s<sup>-1</sup> (which implies 1-Y=38.28% if Hct=0.4), and  $\zeta$  = 4.55%. For the white matter voxel (long arrow), the fitting results are R2=21.7 s<sup>-1</sup>, f=−0.002 Hz, δω=52 s<sup>-1</sup> (1-Y=10% if Hct=0.4), and  $\zeta$ = 0.1% (δω and  $\zeta$  here hit the lower boundaries of ranges of the parameters).

![](_page_19_Figure_2.jpeg)

#### **Figure 5.**

A comparison between the maximum intensity projection of Cdeoxy maps over five slices (left) and the minimum intensity projection of SWIs over the same slices (middle). The brighter regions in Cdeoxy maps correspond to the darker ones in SWIs, which correspond to the voxels occupied by large veins. MIP of Cdeoxy maps also provides the information on capillaries which are not seen in the MIP image of SWI. The histogram (right) is an example of the distribution of Cdeoxy in cortical gray matter. All data are from the subject 1.

![](_page_20_Figure_2.jpeg)

#### **Figure 6.**

The plot represents an example of the 3D histogram of voxels as a function of blood oxygenation level Y, and dCBV (ζ) with bright white color corresponding to higher values. One can see that the peak value corresponds to Y around 0.55 and dCBV around 0.03.

Author Manuscript

**Author Manuscript** 

# **Table 1**

Values of the parameters for the cortical gray matter for eight subjects. The last row represents average values for all 8 subjects. Numbers are Mean (or Values of the parameters for the cortical gray matter for eight subjects. The last row represents average values for all 8 subjects. Numbers are Mean (or peak) +/- STD. Y values are calculated assuming Hct = 0.4. peak) +/– STD. Y values are calculated assuming Hct =  $0.4$ .

![](_page_21_Picture_405.jpeg)

# **Table 2**

Summary of mean values of R2\*, R2 and R2' for cortical gray matter obtained at 3T MRI by different methods. There are two ways of calculating R2' in Summary of mean values of R2\*, R2 and R2′ for cortical gray matter obtained at 3T MRI by different methods. There are two ways of calculating R2' in this paper: [a] R2'=R2\*-R2, [b] R2'= $\zeta$ -80 this paper: [a]  $R2' = R2^* - R2$ , [b]  $R2' = \zeta \cdot \delta \omega$ 

![](_page_22_Picture_190.jpeg)