

Atomic resolution probe for allostery in the regulatory thin filament

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Edited by James A. Spudich, Stanford University School of Medicine, Stanford, CA, and approved February 5, 2016 (received for review October 1, 2015)

Calcium binding and dissociation within the cardiac thin filament (CTF) is a fundamental regulator of normal contraction and relaxation. Although the disruption of this complex, allosterically mediated process has long been implicated in human disease, the precise atomiclevel mechanisms remain opaque, greatly hampering the development of novel targeted therapies. To address this question, we used a fully atomistic CTF model to test both Ca²⁺ binding strength and the energy required to remove Ca²⁺ from the N-lobe binding site in WT and mutant troponin complexes that have been linked to genetic cardiomyopathies. This computational approach is combined with measurements of in vitro Ca²⁺ dissociation rates in fully reconstituted WT and cardiac troponin T R92L and R92W thin filaments. These human disease mutations represent known substitutions at the same residue, reside at a significant distance from the calcium binding site in cardiac troponin C, and do not affect either the binding pocket affinity or EF-hand structure of the binding domain. Both have been shown to have significantly different effects on cardiac function in vivo. We now show that these mutations independently alter the interaction between the Ca²⁺ ion and cardiac troponin I subunit. This interaction is a previously unidentified mechanism, in which mutations in one protein of a complex indirectly affect a third via structural and dynamic changes in a second to yield a pathogenic change in thin filament function that results in mutation-specific disease states. We can now provide atom-level insight that is potentially highly actionable in drug design.

cardiac thin filament | hypertrophic cardiomyopathy | calcium homeostasis | molecular modeling | steered molecular dynamics

ardiac contraction is regulated by the binding of Ca²⁺ to cardiac troponin C (cTnC) (1, 2) (Fig. 1). The Ca²⁺ binds in an EF-hand motif, consisting of two α -helices separated by a loop region containing seven coordinating oxygens that interact with the Ca^{2+} , as seen in Fig. 2. Once Ca^{2+} is bound, interactions between cTnC and the switch peptide domain of cardiac troponin I (cTnI) cause conformational changes that reduce the ability of cTnI to bind to actin (3). When bound to actin, cTnI shifts the equilibrium location of cardiac tropomyosin (cTM) to a position that prevents the interaction of actin with myosin, thereby inhibiting the power stroke that drives contraction of cardiac muscle (4, 5). Thus, Ca²⁺ binding can be viewed as the initial step in a process that results in the eventual hydrolysis of ATP, with the sliding of the thin filament over the thick filament and the generation of mechanical work (5, 6). In disease states, changes in the ability of the cardiac muscle to be properly regulated by and to regulate Ca^{2+} are often observed and play a central role in pathogenic remodeling and sudden cardiac death (SCD) (7). What is not well understood is the proximal biophysical cause by which mutation affects function at the molecular level.

Due to both the size (over 300,000 atoms with a length greater than 80 nm) and the inherent flexibility of some portions of the cardiac thin filament (CTF), there is not a single experimental tool that is able to reveal the full structure. X-ray crystallography and NMR spectroscopy have been unable to resolve the more mobile regions on the atomistic scales (8–10). Fig. 3 illustrates our full atomistic CTF model that was constructed and implemented to allow us, for the first time to our knowledge, to follow these atomistic-level changes. Using computational methods, we will describe how these changes can be followed even when they occur on a time scale far longer than standard computational times.

Hypertrophic cardiomyopathy (HCM) is caused by mutations in sarcomere proteins and is the most common cardiac genetic disorder, affecting at least 1 in 500 individuals (11). It is a leading cause of SCD in young people (12). Many of the known mutations in proteins that comprise the regulatory thin filament disrupt intracellular Ca²⁺ homeostasis, having an impact on both cardiac function and SCD risk, forming the basis for current therapeutic interventions (13, 14). Independent gene mutations, even at a great distance from the Ca²⁺ binding domain in cTnC, can result in dramatically different phenotypes. These differences extend to the single-residue level; for example, substitutions at the same amino acid residue of cardiac troponin T (cTnT), Arg92Leu (R92L) and Arg92Trp (R92W), are phenotypically distinct in both patients and animal models (15). Despite extensive study, the biophysical mechanisms that underlie genotype-specific, clinically relevant distinction remain unclear, and limit the development of targeted therapies for this most common genetic cardiomyopathy.

Here, we report the application of our computational model to mutations that are predicted to affect Ca^{2+} dissociation. The N-terminus of cTnI can be found near the N-lobe binding site of cTnC. In particular, the cTnI Glu32 (E32) side chain is above the Ca^{2+} binding pocket. cTnI E32 contributes another coordinating oxygen site because the Glu has a negatively charged carboxylate that is able to interact with the positively charged Ca^{2+} . As Ca^{2+} exits the binding pocket, the E32 can help to stabilize the ion in the

Significance

The cardiac thin filament (CTF) is a complex and highly dynamic multiprotein complex that regulates heart function at the molecular level. Mutations in proteins that comprise the thin filament cause cardiomyopathic remodeling that is due, in part, to dysregulation of Ca²⁺ binding. Many of the known mutations in multiple thin filament proteins occur at a significant physical distance from the Ca²⁺ binding domain; thus, the elucidation of molecular mechanisms has been challenging. We have developed and now apply an atomistic model of the CTF to this question and have obtained molecular insight that both provides disease mechanisms and facilitates future rational drug design. Finding the molecular cause of disease is the first step in designing molecular interventions.

Author contributions: J.C.T. and S.D.S. designed research; M.R.W. and S.J.L. contributed new reagents/analytic tools; M.R.W., S.J.L., J.C.T., and S.D.S. analyzed data; M.R.W., S.J.L., J.C.T., and S.D.S. wrote the paper; M.R.W. performed computations; and S.J.L. conducted experiments.

The authors declare no conflict of interest.

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This article is a PNAS Direct Submission.

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This article contains supporting information online at www.pnas.org/lookup/suppl/doi:10. 1073/pnas.1519541113/-/DCSupplemental.



Fig. 1. Troponin complex. The cTnT is represented in yellow, the cTnI is represented in blue, and the cTnC is represented in red. The Ca²⁺ is represented as the cyan spheres.

pocket, therefore influencing the rate at which the Ca^{2+} is released. Our computations show that the R92L mutation causes an increased frequency of interaction between the Ca²⁺ and the cTnI E32, whereas the R92W mutation leads to limited interactions between the cTnI E32 and the Ca²⁺. Furthermore, there are almost no changes in the geometry of the EF-hand binding pocket, and so the differential binding strength of the Ca²⁺ to cTnC is not responsible for the changes observed experimentally in WT cTnT, R92L, and R92W complexes. The current results demonstrate the complexity of the "action at a distance" effect of mutations in a large, highly regulated multiprotein complex. Here, mutation in the cTnT protein alters the function of the cTnC protein via the indirect mediation of cTnI. This alteration occurs at a distance over 100 Å from the site of mutation and could not be discerned without an allatom model. Such high-resolution mechanistic insight can be harnessed in the design of novel small-molecule therapeutics, an active subject of current efforts to move the treatment of this disorder from the merely palliative to cure (16).

Results

Construction of the Atomistic CTF Model. The motivation for development of the model has been the work from our group and others



Fig. 2. EF-hand Ca^{2+} binding pocket. (A) Ca^{2+} N-lobe binding pocket showing the EF-hand structure. The cTnC is represented in red. The Ca^{2+} is represented as a silver sphere. Coordinating residues are represented in a ball and stick design: Oxygens are represented as red spheres, nitrogens are represented as blue spheres, carbons are represented as cyan spheres, and hydrogens are represented as white spheres. (*B*) Rotation of previous 90°. (C) View of the binding pocket from directly above.

Fig. 3. Full atomistic CTF model. Actin is represented in gray and silver. Tropomyosin is represented in green and orange. The cTnT is represented in yellow, the cTnI is represented in blue, and the cTnC is represented in red. The Ca^{2+} is represented as the cyan spheres. The cTnT R92 mutation site and Ca^{2+} binding pocket are identified for location comparison.

showing that mutations cause complex effects, often at a significant distance from the site of amino acid replacement (17-20). The function of the thin filament as a regulatory agent relies on the ability of the individual protein components to interact dynamically in response to physiological demands. We thus created a full atomistic CTF model (Fig. 3) to study both structural and dynamic changes as a result of mutation. As has been shown in previous studies, mutations cause changes to both local and long-range interactions (17, 18, 20, 21). The long-range effects are difficult to discern when such complex systems are not studied in atomistic detail. This difficulty is especially true when fragments of the complex are used. Alterations in the Ca²⁺ dissociation rate from the cTnC binding pocket have been previously linked to HCM mutations (22). As such, we measured this parameter for the WT, R92L, and R92W substitutions in the fully reconstituted systems. Details of the model building are given in SI Materials and Methods.

Mutations Alter the Ca²⁺ Dissociation Rates. In vitro Ca²⁺ dissociation rates were measured for WT and mutant reconstituted CTFs using stopped-flow kinetics. The R92L and R92W mutations were introduced into human cTnT (hcTnT), and the E32A mutation was introduced into human cTnI (hcTnI), after which all were bacterially expressed and purified. Human cTnC (hcTnC) was mutated at site T53C, as shown in Fig. 1A; it was then bacterially expressed and purified, and then labeled with a thiol-reactive, environmentally sensitive fluorescent probe, 2-(4'(iodoacetamido)anilino)naphthalene-6-sulfonic acid (IAANS), at site T53C (hcTnC T53C_{IAANS}). This probe on hcTnC allows for monitoring of conformational changes within the site II Ca^{2+} binding pocket of hcTnC in the presence and absence of Ca^{2+} . Bacterially expressed human Ala-Ser cTM and hcTnI WT were purified, and skeletal muscle actin was harvested from rabbit muscle and further purified. The full thin filament was reconstituted (0.8 hcTn complex_{IAANS}/1 cTM/7.5 filamentous actin) and dialyzed against working buffer containing 200 µM Ca2+. The Ca2+-bound thin filaments were rapidly mixed with a buffer containing EGTA, and the resulting change in fluorescence was detected. The data were fit with a single-exponential decay, and a dissociation rate (k_{off}) was calculated.

In agreement with literature values (22), WT thin filaments (Tn WT_{T53C IAANS}) resulted in a Ca²⁺ dissociation rate of 116 \pm 3.3 s⁻¹ (Fig. 4). Thin filaments containing the R92L mutation (Tn R92L_{T53C IAANS}) resulted in a decrease in the Ca²⁺ dissociation rate, 84.8 \pm 1.8 s⁻¹ (Fig. 4), suggesting that Ca²⁺ release from the thin filament is more inhibited compared with WT thin filaments. In contrast, thin filaments containing the R92W mutation (Tn R92W_{T53C IAANS}) resulted in an increase in the Ca²⁺ dissociation rate, 150.8 \pm 4.7 s⁻¹ (Fig. 4), compared with both WT and R92L. This increase in dissociation rate suggests that there is a decrease in inhibition of Ca²⁺ release at the level of the thin filament.

 Ca^{2+} Coordination Strength Is Not Affected by Mutations. Given the variation in the WT and two mutations in the Ca^{2+} off rate, we used our atomistic model to illuminate the molecular-level mechanism



Fig. 4. In vitro Ca²⁺ dissociation kinetics for reconstituted CTFs. The decrease in IAANS fluorescence was measured upon removal of Ca²⁺ from the thin filament via EGTA. Dissociation rates were calculated by fitting the decay curves with a single exponential decay. Decay curves were an average of three to five individual traces collected at least five times each. ****P* < 0.001; *****P* < 0.0001.

influencing Ca²⁺ dissociation. An initial reasonable hypothesis was that the mutation in cTnT creates the transmission of a structural effect to the N-lobe binding pocket of cTnC. Such effects at a distance have been noted before in simpler studies from our group and others (17, 18, 20, 21). To investigate this hypothesis, ensembles of 1-ns molecular dynamics (MD) simulations were produced, starting from the individual equilibrated structures. The seven coordinating oxygen-to-Ca²⁺ mean distances were calculated from 5-ps intervals. The ensemble mean distances were compared between the mutants and the WT. In the WT, a water molecule is intercalated into the space between the cTnC S69 and the Ca²⁺ during the equilibrium process. The water forms a hydrogen bond with the coordinating oxygen side chain, allowing the oxygen of the water to become the new coordination atom. In previous computational work (23, 24), similar effects were found, and there is a small change in coordinating distance; however, overall, none of these effects can be the cause of the variation in measured dissociation rates. Inspection of the trajectories showed that the N-lobe of the cTnC did have some overall movement, whereas the structure of the Ca²⁺ binding pocket remained stable. One MD simulation of each ensemble was allowed to continue to evolve in time for 5 ns to observe how the distances changed over a longer period.

When comparing the ensemble averaged 1-ns trajectories with the 5-ns simulation of the individual models, the mean coordinating oxygen distances were almost identical. The largest change in the mean distance was for the cTnT 71 carboxyl oxygen of the R92W, which changed 0.063 Å (Table S1). When comparing the ensemble averaged coordinating oxygen distances of the WT and mutants, 75% of the mean values changed by less than 0.01 Å between WT and all of the mutants, as can be seen in Fig. 5.

When comparing the individual mean coordinating oxygen distances results for the WT and R92 mutations, the ensemble mean distances varied by 0.2 Å at most (Fig. 5). The oxygen from the water molecule in the WT binding pocket was less than 0.2 Å closer to the Ca^{2+} than the cTnC S69 side chains of the R92 mutations, which can be explained by the hydrogen bond created by the bisecting water molecule decreasing the distance. All of the coordination oxygen distances were within the oxygen distances expected of such a complex, that is, a mean Ca^{2+} -to-oxygen distance of 2.44 Å (25).

As shown in both Fig. 5 and Fig. S1, the coordinating oxygento- Ca^{2+} distances did not change significantly upon mutation; therefore, it is reasonable to assume that the effect on the rate of Ca²⁺ release is not due to the coordinating oxygens of the cTnC binding pocket and the strength to which they bind Ca²⁺. This result is complementary to similar work comparing the WT with the cTnI S23D/S24D double mutation (23) and the WT with the cTnI R145G mutation (24). However, with the full CTF model, we are able to investigate the effects of mutations that are not within the core region of cTn. The R92L and R92W cTnT mutations are over 100 Å away from the binding pocket, yet they are still able to cause differential Ca²⁺ release rates. Given the similarity of these structures, the physiological alterations caused by the R92L and R92W mutations would have to be caused by another physical/chemical feature of the CTF.

Key Role of cTnI E32 in Ca²⁺ Dissociation. To elucidate the mechanism for changes in the Ca²⁺ dissociation rate, we studied energetic barriers to removal of Ca²⁺ from the binding pocket. This event is rare, and so potentials of mean force calculations were implemented using steered molecular dynamics (SMD) (26, 27). We wished to eliminate the effects of water friction from this calculation, and so a previously described, bidirectional algorithm was used (28). The water molecules and ions in the solution add resistance and drag on the Ca²⁺, in the form of dissipative work in both the forward and reverse directions (28). The dissipative work could then be removed from the calculated work to isolate the barrier to removal of Ca²⁺.

The concept behind SMD is that a rare event at equilibrium may be simulated by applying an external force to the system such that a desired event occurs on computationally accessible time scales. In our case, the force is created by pulling on a pseudoatom that is attached to the Ca^{2+} with a pseudospring (27). The force that the spring feels is converted into work. The ensemble average work can be equated to the free energy through the application of Jarzynski's equality (29).

Table 1 shows that the calculated total free energy required to remove the Ca^{2+} completely from the binding pocket of the R92L mutation was larger than the calculated free energy of Ca^{2+} removal from the WT, whereas the R92W mutation had a smaller



Fig. 5. Coordinating oxygen distances. The mean coordinating oxygen distances for the N-lobe binding pocket. The square points represent the ensemble average distance, and the diamond points represent the mean distance for the 5-ns simulations. Delta oxygen 2 (OD2), gamma oxygen (OG), carbonyl oxygen (O), delta oxygen 1 (OD1), epsilon oxygen 1 (OE1), and epsilon oxygen 2 (OE2) refer to the position of the oxygen of the amino acid. The 69 OG coordination for the WT and 5-ns WT is the distance of the transferable intermolecular potential 3 point model (TIP3P) water oxygen.

Table 1. SMD calculated work: Total calculated forward and reverse work, free energy, and percentage of cTnl E32 interactions for the WT, R92L, and R92W from SMD

Interaction parameters	WT, kcal/mol	cTnT R92L, kcal/mol	cTnT R92W, kcal/mol
$\langle W \rangle_f$	1,190.13	1,185.92	1,098.57
$\langle W \rangle_r$	1,047.07	1,025.77	983.27
ΔG	71.53	80.07	57.40
cTnI E32 interaction, %	60	90	10

calculated free energy than the WT's Ca^{2+} removal. However, the three systems had similar results for the work required to remove the Ca^{2+} only from the coordinating oxygens of the binding pocket, as shown in Fig. 6. The difference in the calculated work was due to the interaction of the Ca^{2+} and the E32 of cTnI. The negatively charged carboxylate side chain of the cTnI E32 was able to interact with the positively charged Ca^{2+} as it was pulled from the binding pocket. The R92 mutations altered the positioning of the N-terminus region of the cTnI. The E32 carboxylate of the R92L mutation was positioned closer to the direct center above the binding pocket, whereas the E32 carboxylate of the R92W mutation was positioned farther away from the direct center above the binding pocket (Fig. S2).

Our hypothesis that E32 of cTnI partially controls Ca^{2+} dissociation suggests an immediate test. Replacement of this residue with neutral Ala should increase the rate of dissociation for all species (mutant and WT), the overall rates should be more similar, and the rate of R92L with the E32A additional mutation should be roughly comparable to the dissociation in the single mutant R92W. Computations and experiments to test this hypothesis directly are shown in Fig. 7. Note that the substitution essentially abrogates the second Ca^{2+} interaction in the pulling computations. Any residual difference in rate is due to steric effects not sampled in the computation.

It should be noted that the free energy values obtained are not the exact free energy barriers to actual Ca^{2+} exit. Equality via the Jarzynski theorem only obtains in the limit of "slow" pulling of infinitely many realizations (27, 29). Using higher velocity "pulling" allows the generation of computational results in a



Average Work

Fig. 6. SMD calculated work. The calculated forward work at each step of the Ca^{2+} pulling is shown. The first peak is representative of the Ca^{2+} work to break free from the coordinating oxygens of the binding pocket, whereas the second peak is where the cTnI E32 oxygens interact with the Ca^{2+} .



Fig. 7. SMD (A–C) and dissociation kinetic (*D*) measurements for WT, R92L, and R92W cTnT with the further replacement of E32A. This mutation tests the hypothesis that E32 interaction with Ca²⁺ dominates the contribution to work in the 25-Å region of the SMD plots.

finite time, but also leads to increased displacement of the spring. The method has been used before with similar results (30, 31). On the other hand, at a slower velocity, the force applied to the pseudoatom pulled the N-lobe of the cTnC along with the Ca²⁺. The pulling force would result in deformation of the N-lobe before the Ca²⁺ was released. We chose a velocity of the pulling force that was the slowest with which the deformation of the N-lobe was minimal.

Flexibility and Position of cTnl Relative to Ca²⁺ Binding. Given the empirical results obtained from the SMD calculations (i.e., that changes in apparent Ca²⁺ dissociation rates are caused by differential interactions with cTnI that, in turn, alter the free barrier to Ca^{2+} escape), we probed the atomic-level differences in this part of the complex to determine the root cause of this effect. First, we address the lower rate and higher free energy barrier found in the R92L mutation. We detected important changes in flexibility of the N-terminal region of cTnI induced by the cTnT mutation. The root mean square fluctuation (RMSF) of the cTnI (Fig. 8) and cTnC (Fig. S1) was calculated. The N-terminus end of the cTnI is largely unstructured until the start of an α -helix at residue 43, which leads to the region where α -helices from both the cTnI and cTnT interact (IT arm). Examination of the RMSF in the unstructured N-terminus region of cTnI shows significant differences for the R92L mutant compared with the WT and R92W complexes. The RMSF is both larger in magnitude and rising in the region around residues 10-25 for R92L compared with both WT and R92W. This change causes the N-terminal section of the protein to be more prone to follow the motion of the double-charged cation because of electrostatic forces, thus causing a greater probability of interaction.

The R92W mutation has a slight increase in k_{off} , and this increase is caused by a static structural mechanism. The N-terminus of the R92W mutation is slightly more mobile than the WT (Fig. 8), but the dominant effect shown in Fig. 9 is that the average position of the R92W compared with the WT is shown to shift, on average, away from the Ca²⁺ binding pocket. It does not have the flexibility of the R92L mutation, however, and so although the R92L mutant is shown in the static representation as being also located away from the binding pocket, due to the increased mobility, the region can sample more conformational space. One might ask what causes the static and dynamic structural changes. This answer is found in changes in the structure of the IT arm region of the complex. As shown in Fig. S3 and Table S2, both the angle between the helices and pseudodihedral angles of the same



Fig. 8. RMSF of the cTnl for the 1-ns ensemble. The WT is shown in black, the R92L is shown in red, and the R92W is shown in green. The light blue region of the graph represents the N-terminus region. The N-terminus end contains the largest differing RMSF among the models.

helices of cTnT change upon mutation of cTnT, and, in turn, this change is transmitted to the corresponding helices of cTnI. Fig. S4 shows representative results from SMD, demonstrating the differences in cTnI E32 Ca²⁺ association. This transmission of structural and dynamic information may be viewed as the converse of such information flow upon Ca²⁺ binding.

Discussion

This paper reports the application of a newly developed atomistic model of the CTF to two cTnT mutations that exhibit differential effects on calcium homeostasis in vivo. The mutations are far from the site of calcium binding, and, in fact, our primary result is that complex allosteric interactions are responsible for the differences observed. The measured Ca²⁺ dissociation rates can be explained by the variation of the cTnI E32 interaction. If the WT results are used as the benchmark, the increased Ca^{2+} dissociation rate of the R92W can be explained by the weaker interaction of the Ca²⁺ and the cTnI, which is, in turn, explained by cTnI displacement to the outer edge of the binding pocket. The decreased Ca²⁺ dissociation rate of the R92L can be explained by the increased interaction of the Ca^{2+} and cTnI. The RMSF of the N-terminus region of the cTnI demonstrates the location and manner of how and where the mutations are altering interactions of the WT. This alteration is in a location that is relevant to the Ca²⁺ binding pocket and the release of Ca²⁺, even though the mutations occur over 100 Å away. Whereas the in vitro Ca²⁺ dissociation rate experiment gave numerical values to the effects of the mutations, the atomistic model illuminated the mechanism by which the mutations cause the observed effects on Ca²⁺ dynamics.

The computational results give insight into the function of cTnC. In simulation production runs, the Ca^{2+} distances from the coordinating oxygens remained constant throughout the dynamics. Similarly, the SMD experiment showed that the same work was needed to remove the Ca^{2+} from the binding pocket in the WT and the mutants. We and others have previously shown that mutations can have long-range effects on the CTF (17, 18, 20, 21, 32, 33). These results demonstrate both the challenge and promise of molecular intervention in such genetic alterations. The effect that causes dysfunction is found remote from the site of mutation. It is important to note that the development of targeted therapeutics is strongly advanced by the development and application of this model. In complex machines, such as the sarcomere, the strong binding of agents is not the only goal; one needs to have a target

for the binding. Without the type of work presented in this paper, one would never know how a mutation in cTnT affects the function of cTnC, and the therapeutic options available would remain the same calcium-channel blocking agents used for 50 years. We now have the potential to begin to target the actual molecularlevel insult to sarcomeric function.

It is important to note that one of the most confounding issues in the management of human genetic cardiomyopathies is the inability to use genotype to either inform or prognosticate clinical phenotype [recently reviewed by van der Velden et al. (16)]. This significant limitation persists despite 20 years of concerted efforts by many groups. It has been argued that HCM is the most "biophysical" of disorders, with mutations occurring in proteins with known physiological roles and extensive functional data. The thin filament, in particular, is an experimentally "tractable" system; thus, significant advances have been made regarding molecular mechanism (reviewed in 5). More recently, however, it has become clear that the widely used steady-state and biochemical approaches do not fully differentiate between the known clinical phenotypes, suggesting that further, high-resolution approaches are needed to provide mechanistic insight. The two human cTnT mutations studied here are representative of this issue. The R92W and R92L cTnT missense mutations cause dissimilar patterns of ventricular remodeling in mice and men (21, 34, 35). In animal models, they also exhibit marked mutation-specific effects on Ca²⁺ homeostasis and myocyte energetics despite near-identical effects on the Ca²⁺ sensitivity of force generation (36, 37). The results we describe here, for the first time to our knowledge, provide the level of molecular biophysical understanding required to develop a robust mechanistic paradigm for the differential phenotypes observed in the human disorder. Moreover, given that any thin filament mutation can be tested in a rigorous, high-throughput manner using these computational models, it is highly likely that they can eventually be incorporated into existing predictive HCM algorithms, and thus greatly extend the biophysical quality of these approaches.

Materials and Methods

Detailed materials and methods are provided in *SI Materials and Methods*, including the CTF model building methodology, mutation models, MD and SMD settings, coordinating oxygen distance calculation, RMSF calculations, energy calculations, in vitro CTF reconstruction, and stopped-flow kinetics techniques.

Neither our experimental model nor our computational model includes myosin. Thus, there is no chemical reaction of ATP hydrolysis, and this level of control is simply not present. Although the physiological role of myosin heads in cardiac relaxation is an important factor, the measurement of Ca^{2+} dissociation rates in the presence of S1 heads is challenging. Previous groups have measured the calcium dissociation rates in reconstituted thin filaments plus myosin



Fig. 9. MD ensemble average structures of the 1-ns simulations. The view of the N-lobe binding pocket perpendicular to the vector from which the Ca^{2+} was pulled out of the pocket during the SMD simulations is shown. The red represents the cTnC EF-hand motif of residues 53 through 88. The blue represents the cTnI residues 12 through 38. The silver sphere is the Ca^{2+} . The small red, cyan, and white spheres represent the oxygen, carbon, and hydrogen atoms of the cTnI E32 side chain, respectively. The black arrows point at the distance of the Ca^{2+} and the cTnI E32 side chain. WT (A), the R92L (B), and the R92W (C) are shown.

S1 heads (23, 38) and reported rates that are similar to the rate of cardiac relaxation (~15 s⁻¹). However, each of these systems has notable limitations. For example, the binding of S1 heads to the reconstituted thin filament is not regulated due to the lack of an organized thick filament, and thus cannot account for specific binding sites. Additionally, within the I-band of the cardiac sarcomere, the thin and thick filaments do not overlap, and thus represent a region that is governed by the calcium release rate of the thin filament alone. Furthermore, there are only two possible states in these experiments: myosin-bound and myosin-unbound. Calcium cycling plays a role in the equilibrium between unbound and weakly bound myosin; thus, the effects of altered

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dissociation rates in this transition may be overlooked. Because we cannot directly account for these different possibilities in either the in vitro measurements or our computational model, we have chosen to focus on the thin filament alone.

ACKNOWLEDGMENTS. We thank Mark McConnell (University of Arizona) for development of the analysis program for the stopped-flow kinetics. This research was financially supported by NIH Grant R01 HL107046 (to J.C.T. and S.D.S.). M.R.W. was supported via Training Grant GM084905 and S.J.L. received support from Grant T32 HL007249. J.C.T. acknowledges the support of the Steven M. Gootter Foundation.

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