



Published in final edited form as:

Magn Reson Med. 2016 June ; 75(6): 2265–2277. doi:10.1002/mrm.25729.

Biexponential Longitudinal Relaxation in White Matter: Characterization and Impact on T_1 Mapping with IR-FSE and MP2RAGE

James A. Rioux¹, Ives R. Levesque^{1,2,3}, and Brian K. Rutt^{1,*}

¹Department of Radiology, Stanford University, Stanford, California, USA

²Medical Physics Unit, McGill University, Montreal, Quebec, Canada

³Research Institute of the McGill University Health Centre, Montreal, Quebec, Canada

Abstract

Purpose—Magnetization transfer in white matter (WM) causes biexponential relaxation, but most quantitative T_1 measurements fit data assuming monoexponential relaxation. The resulting monoexponential T_1 estimate varies based on scan parameters and represents a source of variation between studies, especially at high fields. In this study, we characterized WM T_1 relaxation and performed simulations to determine how to minimize this deviation.

Methods—To characterize biexponential relaxation, four volunteers were scanned at 3T and 7T using inversion recovery fast spin echo (IR-FSE) with 13 inversion times (TIs). Three volunteers were scanned with IR-FSE using TIs chosen by simulations to reduce T_1 deviation, and with MP2RAGE.

Results—At 3T, the biexponential relaxation has a short component of $T_1 = 48$ ms (9%) and a long component of $T_1 = 939$ ms. At 7T the short component is $T_1 = 57$ ms (11%) and the long component is 1349 ms (89%). For IR-FSE, acquiring four TIs with a minimum of 150 ms (3T) or 200 ms (7T) yielded monoexponential T_1 estimates that match the long component to within 10 ms. For MP2RAGE, significant differences (90 ms at 3T, 125 ms at 7T) remain at all parameter values.

Conclusion—Many T_1 mapping sequences yield robust estimates of the long T_1 component with suitable choice of TIs, allowing reproducible, sequence-independent T_1 values to be measured. However, this is not true of MP2RAGE in its current implementation.

Keywords

quantitative MRI; T_1 mapping; inversion recovery; MP2RAGE; white matter

*Correspondence to: Brian K. Rutt, Ph.D., Richard M. Lucas Center for Imaging, Stanford University, Radiology Department, 1201 Welch Road PS064, Stanford, CA 94305. brutt@stanford.edu.

INTRODUCTION

The primary goal of quantitative MRI is to obtain accurate, spatially resolved maps of tissue parameters and to relate these quantities to changes associated with disease. For such measurements to be reproducible and comparable across sites, platforms, and sequences, they must be independent of all contributions to the MRI signal aside from the parameter of interest. For example, a map of the longitudinal relaxation time T_1 must not be influenced by other relaxation effects (eg, T_2 , T_2^*), diffusion, or inhomogeneity in the B_0 and B_1 fields. It should also be independent of the specific scan parameters used during acquisition. Removing such influences is critical to enabling quantitative MRI for basic research and clinical applications.

In this study, we consider the application of T_1 mapping in the white matter (WM) of human brain, which may aid in the identification of pathologies caused by multiple sclerosis (1,2), Alzheimer disease (3), and other neurological diseases. A key prerequisite for the detection of diseased tissue is the accurate characterization of healthy, normal tissue. In the case of WM, however, there is no strong consensus in the literature about what constitutes a “baseline” healthy control T_1 value. For example, Table 1 lists a number of reported WM T_1 values at 7T. Significant variation is evident; a majority of values cluster around 1100–1200 ms, but some are much longer [eg, (4,12)], and one is much shorter (7).

Some of this variation is sequence-based. Stikov et al. (13) have shown that, at 3T, sequences based on variable flip angle acquisition such as DESPOT1 (4,14) tend to overestimate T_1 values compared with inversion recovery fast spin echo (IR-FSE) methods, whereas sequences based on the Look–Locker acquisition method (15) tend to underestimate T_1 . It can be assumed that such deviations will only be larger at 7T, because one of their main causes is error in B_1 estimation (13), and that B_1 variations will begin to impact IR-FSE as well. Individual variation and changes due to aging or anatomical location (eg, frontal versus callosal WM) are also expected (16). Nevertheless, large differences in reported T_1 values obtained with similar sequences [eg, (6,9) or (8,12)] are concerning from the viewpoint of quantitative MRI and indicate that some variation in the T_1 measurement is not yet fully understood or accounted for.

We hypothesized that biexponential relaxation is a major source of variation in T_1 measurements in WM. Such biexponential relaxation is established in the literature, primarily in applications involving magnetization transfer (MT), but its impact on quantitative T_1 imaging has remained largely unexplored. We first quantified the relaxation time and signal fraction of the biexponential components in WM at both 3T and 7T using a 13-point IR-FSE series. However, because collecting sufficient data to allow such characterization is impractical for clinical applications, we sought methods to quickly and accurately estimate the long T_1 component, with minimal influence from the short component. We demonstrate that, for sequences that acquire multiple inversion times (TIs) to fit with a monoexponential model, the apparent T_1 value is shorter than the long T_1 component by an amount that depends mainly on the minimum TI, and this difference can be minimized by selecting an appropriate set of TIs. The resulting measurement of the long T_1 value should be consistent across sites and methods, and independent of scan parameters.

However, the effect of the short T_1 component on other T_1 mapping methods such as MP2RAGE (10) cannot be mitigated as easily and will remain an important source of variation in such measurements.

THEORY

Biexponential T_1 Relaxation

The presence of biexponential T_1 relaxation in tissues such as muscle (17), cartilage (18), gray matter, and WM (19) has been acknowledged for some time. Although early studies of human WM at 1.5T (20) found that this biexponential T_1 relaxation did not contribute significantly to signal evolution, it is more easily detected and quantified at higher field strengths. Prantner et al. (21) tested several hypotheses that could explain the biexponential T_1 behavior observed in rat brains at 4.7T and 11.7T. These mechanisms included sequence or scanner artifact, blood flow, and multiple exchanging or nonexchanging compartments. It was concluded that MT between exchanging spin populations was the most likely mechanism for this biexponential behavior.

This is consistent with MT models such as that outlined by Dortch et al. (19), based on earlier work (22–24). The two spin populations in these models are often assumed to be hydrogen protons associated with free water and macromolecules, respectively, with rates of exchange k_{fm} and k_{mf} between pools. The evolution of the longitudinal magnetization following a radiofrequency (RF) pulse can be modeled with coupled Bloch equations:

$$\begin{aligned} dM_f/dt &= -(M_f - 1) \times R_{1f} - k_{fm}(M_f - M_m) \\ dM_m/dt &= -(M_m - 1) \times R_{1m} - k_{mf}(M_m - M_f). \end{aligned} \quad [1]$$

Here, M_f and M_m are the time-varying longitudinal magnetizations of the pools, normalized to their values at thermal equilibrium, and $R_1 = 1/T_1$. The solution of these equations (19) for the longitudinal magnetization $M(t)$ of the visible protons (also normalized to its equilibrium value) now contains two relaxation rates, one with short T_1 ($T_{1S} = 1/R_{1S}$), the other with long T_1 ($T_{1L} = 1/R_{1L}$):

$$M(t) = 1 + b_s \exp(-t \times R_{1S}) + b_L \exp(-t \times R_{1L}), \quad [2]$$

where

$$\begin{aligned} 2R_{1S} &= R_{1f} + R_{1m} + k_{fm} + k_{mf} + \sqrt{(R_{1f} - R_{1m} + k_{fm} - k_{mf})^2 + 4k_{fm}k_{mf}} \\ 2R_{1L} &= R_{1f} + R_{1m} + k_{fm} + k_{mf} + \sqrt{((R_{1f} - R_{1m} + k_{fm} - k_{mf})^2 + 4k_{fm}k_{mf})} \end{aligned}$$

and

$$\begin{aligned} b_s &= ((M_f(0) - 1) \times (R_{1f} - R_{1L}) + (M_f(0) - M_m(0)) \times k_{fm}) / (R_{1S} - R_{1L}) \\ b_L &= -((M_f(0) - 1) \times (R_{1f} - R_{1S}) + (M_f(0) - M_m(0)) \times k_{fm}) / (R_{1S} - R_{1L}). \end{aligned}$$

For most common experimental parameters (ie, an RF pulse that fully or almost fully inverts the free pool), both coefficients $b_{S,L}$ are negative and the normalized magnetization $M(t)$ starts from a negative value and relaxes toward 1. For clarity when discussing the relative contributions of the two biexponential components, we may also refer to normalized weights $w_{S,L} = b_{S,L}/(b_S + b_L)$.

The observed biexponential relaxation times T_{1L} and T_{1S} represent a mixture of the relaxation times of the underlying pools (T_{1f} and T_{1m}), as well as the exchange constants k_{mf} and k_{fm} . Reported values for these relaxation times in WM (11,19) indicate that the pools are in an intermediate exchange regime. The long component T_{1L} is weighted more heavily by T_{1f} and T_{1m} but also contains some contribution from exchange effects. In contrast, the short component T_{1S} is dominated by exchange effects, and in the intermediate exchange regime it is significantly shorter than T_{1L} while remaining observable with achievable TIs. As an example of typical values, Dortch et al. (11) found $T_{1L} = 1370$ ms ($R_{1L} = 0.73$ s⁻¹) at 7T. Based on their assumption $R_{1f} = R_{1m}$ and their reported values of the pool size ratio (PSR = 17.6%) and exchange constants ($k_{mf} = 14.5$ s⁻¹, $k_{fm} = \text{PSR} * k_{mf} = 2.55$ s⁻¹), the short relaxation time they observed must have been $T_{1S} \approx 60$ ms. The relative weights of the components $w_{S,L}$ are also not identical to the population fractions of the two pools; for example, they are influenced by $M_f(0)$ and $M_m(0)$, the magnetizations of each pool immediately after the inversion pulse. Using the same assumptions as reported in Dortch et al. (11), we can infer an observed $w_S = 0.09$ and $w_L = 0.91$ such that $w_S/w_L \approx \text{PSR}$.

Biexponential Relaxation and IR-FSE

Although the effects of exchange have been considered in the context of quantitative MT (qMT) sequences and others such as spoiled gradient echo imaging (25), this exchange model applies to even the common gold standard sequences for T_1 mapping, inversion recovery spin echo (IR-SE), or IR-FSE, and according to Equation [1], the T_1 relaxation observed by these sequences should be biexponential. However, this biexponential relaxation is typically neglected, and IR-FSE data are processed and analyzed assuming monoexponential recovery. This yields an apparent T_1 , which we refer to as T_1^* . For the simple case of a three-parameter fit to three acquired TIs (TI_1, TI_2, TI_3), and given certain assumptions about the values of T_{1S} and T_{1L} compared with these TIs (see Appendix), T_1^* can be approximated as

$$T_1^* = T_{1L} - [T_{1L}^2 / (TI_2 - TI_1)] \times b_S / b_L \times [\exp(-TI_1 / T_{1S}) / \exp(-TI_1 / T_{1L})] \quad [3]$$

T_1^* will therefore be shorter than the long T_1 component by an amount that depends on T_{1L} , T_{1S} , the coefficients b_S and b_L , and the TIs at which the signal is sampled. Because the deviation from T_{1L} increases as T_{1L}^2 , and the T_{1L} of most tissues increases with field strength, the absolute magnitude of the error term will also increase with field strength. The dependence of T_1^* on the TIs can be reduced by choosing $TI_1 \gg T_{1S}$. In that case, to good approximation, $T_1^* = T_{1L}$ and it becomes valid to use a monoexponential model such as

$$M_z(TI)=c_1+c_2\exp(-TI/T_1^*). \quad [4]$$

As an alternative interpretation, once $TI_1 \gg T_{1S}$, the short component no longer contributes to signal change during the acquisition and can be modeled as a constant that is subsumed into c_1 in Equation 4. Note that the magnetization of the exchanging pool that gives rise to T_{1S} is not necessarily fully relaxed after time TI_1 , but because this pool's signal is not observed directly, the single exponential model is still valid.

However, this choice of TI_1 is in opposition to the heuristic (26,27) of selecting the first TI to be as short as possible, which maximizes the dynamic range of the sampled signal intensity and improves the precision of the measurement. Reducing the difference between T_1^* and T_{1L} will necessarily involve increasing the variance in the resulting measurement. For this reason, it is important to have accurate knowledge of T_{1S} values in WM, such that parameters can be chosen to balance gains in accuracy with loss of precision.

Biexponential Relaxation and MP2RAGE

An emerging method for fast three-dimensional (3D) T_1 mapping is MP2RAGE (10,28,29), which offers inherent robustness against M_0 , T_2^* , and most B_1 effects, making it a particularly promising candidate for high-field quantitative imaging. MP2RAGE acquires two images following a single inversion pulse, and instead of fitting a model to the acquired data, uses a lookup table computed from the signal equations to convert the intensity of a composite image into a T_1 value. This composite image is obtained from the acquired (complex) images using

$$S_{MP2}=S(TI_1) \times S(TI_2)/(S(TI_1)^2+S(TI_2)^2). \quad [5]$$

The image intensities $S(TI_1)$ and $S(TI_2)$ have a complicated dependence on T_1 , α , repetition time (TR), and other parameters that can be computed based on the Bloch equations (10). This makes MP2RAGE sensitive to the choice of parameters but also allows T_1 map quality to be optimized for a desired range of T_1 values while minimizing B_1 sensitivity (29).

The MP2RAGE signal equations assume a single spin pool and monoexponential T_1 recovery. Including a second exchanging spin pool further complicates the signal equations and makes straightforward estimation of T_1 far more challenging. To conceptually explore the general mechanism by which biexponential relaxation impacts the estimated T_1 in MP2RAGE, consider an MP2RAGE signal consisting of independent (nonexchanging) contributions from long and short T_1 components:

$$S(TI_n)=b_L \times S(TI_n, T_{1L})+b_S \times S(TI_n, T_{1S}). \quad [6]$$

This simplification assumes that the RF pulse train does not significantly modify the apparent relaxation due to exchange (ie, that $T_{1S,L}$ and $b_{S,L}$ have the same values as they would under free relaxation). In practice, exchange effects are likely to produce further deviations from this simplified behavior. However, even in this simple case of unperturbed

relaxation during the RF pulse train, it can be shown that MP2RAGE will still yield an underestimation of T_1 . If both TIs are long enough to ensure complete relaxation of the short component, but not long enough to allow full relaxation of the long component, then

$$S(TI_n) = b_L \times S(TI_n, T_{1L}) + b_S. \quad [7]$$

As with the IR-FSE case, the signal is that observed in the absence of any biexponential relaxation, plus a constant value to which the short T_1 component has relaxed. The expression for the MP2RAGE signal becomes

$$S_{MP2} = (S(TI_1) + b_S) \times (S(TI_2) + b_S) / ((S(TI_1) + b_S)^2 + (S(TI_2) + b_S)^2). \quad [8]$$

The constant but unknown contribution of the short T_1 component's coefficient will not cancel from this expression, and though it may represent only a fraction of the desired signal, it can introduce significant errors into the computation of T_1 . This can also be simulated once the biexponential relaxation has been characterized.

METHODS

Experiments

To better characterize relaxation behavior in WM, data were acquired from four healthy volunteers (male, $n = 3$; female, $n = 1$; mean age, 31 ± 1 y) at 7T and from four healthy volunteers (male, $n = 2$; female, $n = 2$; mean age, 30 ± 2 y) at 3T. Three of the volunteers were scanned at both field strengths. All scans were conducted after obtaining informed consent compliant with our institutional review board. 7T images were acquired on a GE Discovery MR950 scanner (GE Healthcare, Waukesha, Wisconsin, USA) using a 32-channel transmit/receive head coil (Nova Medical, Wilmington, Massachusetts, USA). 3T data were acquired on a GE MR750 (GE Healthcare) scanner with an eight-channel receive-only head coil (InVivo, Gainesville, Florida, USA).

In each volunteer, a series of 13 IR-FSE images was acquired with TI = 10, 20, 35, 55, 85, 125, 200, 350, 600, 1000, 1600, 2500, and 4000 ms. Each image was of a single 1.5-mm-thick axial slice centered on the thalamus and basal ganglia acquired using the following parameters: matrix = 128×128 ; field of view = 19.2 cm; echo train length = 8; TR = 6000 ms; bandwidth = 25 kHz. The inversion was performed with an adiabatic hyperbolic secant (HSn) pulse (30) using the following parameters: pulse width = 16 ms; bandwidth = 990 Hz; HSn = 1.0; $\beta = 2.66$; adiabatic threshold = $8.45 \mu\text{T}$; peak $B_1 = 14.1 \mu\text{T}$ (the same parameters were used at both field strengths). Echoes were collected in centric order to maximize T_1 contrast and reduce T_2 contrast. Acquisition time for each 2D image was 1:42 for a total of 22 min.

In three additional volunteers at 7T (male, $n = 1$; female, $n = 2$; mean age, 32 ± 2 y), a four-point IR-FSE time series (single 1-mm-thick slice; matrix = 192×192 ; TI = 200, 600, 1500, 4000 ms; TR = 6000 ms; scan time = 2:30 per image) was collected to compare with an MP2RAGE volume centered on the same slice. The MP2RAGE parameters [based on

Marques et al. (10) but adjusted to improve robustness to flip angle errors] were as follows: MP2RAGETR = 7500 ms; TR = 7.4 ms; TI₁/TI₂ = 1000/3300 ms; α₁/α₂ = 5°/4°. A 180 × 180 × 180 matrix was collected at 1-mm isotropic resolution using a centrically ordered phase encoding scheme (31) that decouples the number of slices from the number of acquired phase encode lines per inversion, which was 200. The hyperbolic secant adiabatic inversion pulse was 16 ms, bandwidth of 913 Hz, HS_n = 1.0, β = 4.5, adiabatic threshold = 8.35 μT, and peak B₁ = 20.9 μT. The scan time was 10 min with 2× acceleration via parallel imaging.

Analysis

For each volunteer, a region of interest (ROI) that covers most of the WM in the acquired slice was defined using a region-growing algorithm in MATLAB (Mathworks, Natick, Massachusetts, USA), starting from seed pixels in WM of the IR-FSE image with TI = 600 (Fig. 1a). This image was used because the WM signal is approximately nulled and the white/gray contrast is highest. Seeds were manually chosen until most of the WM was included. The resulting ROI was eroded with a 3 × 3 filter to remove pixels on the white/gray boundary, giving a mask such as that shown in Figure 1b.

For the 13-point IR-FSE time series, T₁ was calculated at each pixel in the ROI by fitting a biexponential recovery model to the entire time series, as well as monoexponential fits to the entire series, and subsets of the data with 4–12 TIs at successively larger minimum TI values (eg, only the images with TI ≥ 200 ms, as shown in Fig. 1c). Examples of these fits are shown in Figure 2. Monoexponential fits were performed with a reduced-dimension nonlinear least squares fit (32) in MATLAB. This algorithm uses a four-parameter model that accounts for effects such as imperfect RF pulses and finite TR, so fitting only three TIs is generally not feasible. The biexponential fit uses a modified version of the model from (32) with separate parameters for the long and short components, and a standard Levenberg–Marquardt nonlinear least-squares fit.

Unless otherwise specified, parameter values obtained from the various fits are reported as the mean ± standard deviation of a Gaussian function fitted to the histogram of that parameter over the entire ROI. To quantify whether the biexponential fit represents a significant improvement over a monoexponential fit to the entire series, the F statistic was computed for each pixel in the ROI according to

$$F = \frac{(RSS_1 - RSS_2) \times f_2}{RSS_2 \times (f_1 - f_2)}, \quad [9]$$

where RSS_{*i*} is the residual sum-of-squares of either the single-exponential (*I* = 1) or biexponential (*I* = 2) fit to all 13 TI values at that pixel, and *f_i* is the number of degrees of freedom in the fit. Using the associated cumulative distribution function, *F* can then be converted to a *P* value for evaluation of significance.

T₁ maps were created from MP2RAGE composite images (Eq. 6) using the signal equations given in Marques et al. (10), modified to account for the centric phase encode order. The

same procedure as described above was used to delineate and compare ROIs with corresponding IR-FSE T_1 maps; an example is shown in Figure 1d.

Simulations

To assess the impact of a short T_1 component on a wide range of IR-FSE acquisition schemes, simulations were performed in which a biexponential T_1 recovery defined according to Equation [1] was sampled at various TIs and fitted with the monoexponential function in Equation [4] to obtain T_1^* and the deviation $T_{1L} - T_1^*$. Experimentally determined values for T_{1L} , T_{1S} , and w_S at 3T and 7T were used to define the biexponential recovery curve. TI values were geometrically spaced between a minimum and maximum TI. The maximum TI was set to 4000 ms to match experimental data, and the minimum TI varied between 10 and 600 ms. The number of TI values in each set varied between four and 12. This allowed direct comparison between simulated results and T_1 values from subsets of experimental data. Gaussian noise with a standard deviation of 2% or 1% of the equilibrium magnetization was added to the recovery curves to simulate images with signal-to-noise ratio (SNR) of 50 and 100, respectively. Five thousand repetitions were performed at each noise level and set of TI values, allowing the variance σ^2 to be computed. To evaluate the optimum combination, we used the figure of merit $(T_{1L} - T_1^*)^2 + \sigma^2$, which is the mean squared error (MSE) of T_1^* relative to T_{1L} .

T_1 mapping with MP2RAGE involves a lookup table based on the signal equations, which relates the intensity of the ratio image in Equation [5] to a T_1 value. However, if the signal behavior is modified by the presence of a short T_1 component, as in Equation [7], the lookup will generally be incorrect (Fig. 3). To assess the extent of this deviation, MP2RAGE signals were simulated for a range of TIs TI_1 (200–1000 ms at 3T and 200–1900 ms at 7T) and TI_2 (2000–6000 ms at 3T and 7T). Combinations where $TI_2 - TI_1$ was not large enough to accommodate the required number of TR intervals were disallowed, as were combinations in which the lookup table was non-monotonic over a range of interest centered on T_{1L} (400–1400 ms at 3T and 800–1800 ms at 7T). This ensures that the lookup yields a unique T_1 value for a given intensity. The following parameters were chosen to match experiments: TR = 7.4 ms; TS = 7500 ms between inversions; 200 phase encode lines per inversion; $\alpha_1/\alpha_2 = 5^\circ/4^\circ$.

For each combination of TI_1 and TI_2 , the actual MP2RAGE signal in the presence of the short T_1 component was calculated using Equations [5] and [6]. The lookup table for all of the T_1 values in the range of interest is generated assuming monoexponential relaxation, and the resulting signal is converted into a T_1^* using that lookup table (see Fig. 3 for an example). As with the IR-FSE simulations, 5000 repetitions were performed with 2% noise added to the biexponential $S(TI_n)$, such that the variance and MSE of the lookup table result T_1^* can be computed.

RESULTS

Experimental

The various T_1 relaxation times and short T_1 component weight w_S obtained from the 13-point IR-FSE series in a typical volunteer at 7T are illustrated in Figure 4. There was clear evidence of biexponential relaxation, with a short component comprising approximately 11% of total signal. If all of the data from this biexponential recovery curve are fitted with a monoexponential model, as shown by the red dotted lines in Figure 2, the result is a fit that gives a T_1^* significantly less than the true value of the long T_1 component. For the volunteer whose data are shown in Figure 4, the T_1^* obtained from the monoexponential fit to all data (1161 ms, Fig. 4d) was shorter than T_{1L} (1342 ms, Fig. 4a) by 181 ms. Figure 4e shows that this can be remedied by fitting to only those data with sufficiently long TIs ($TI > 200$ ms in this case).

All of the results from each of the volunteers at both field strengths are given in Table 2. In addition to those parameters illustrated in Figure 4, Table 2 also provides the percent difference in T_1 between monoexponential (T_1^*) and biexponential (T_{1L}) fits to all data, and the percentage of voxels in the ROI for which the biexponential fit performed significantly better ($P < 0.05$ or $P < 0.01$) than the monoexponential fit. All volunteers at both field strengths consistently showed biexponential relaxation, with a larger effect at 7T (mean difference of 13% and significantly better biexponential than monoexponential fit in 96% of voxels at $P = 0.05$) than at 3T (mean difference of 6%, significantly better biexponential fit in 45% of voxels).

The WM T_1 values determined by MP2RAGE in three volunteers at 7T are given in Table 3, along with the corresponding T_1 values calculated by a four-point IR-FSE sequence with $TI_{\min} = 200$ ms, and the percent difference between them. Whereas the four-point IR-FSE measurement was consistent with the biexponential T_{1L} obtained from the 13-point series, MP2RAGE measurements showed an underestimation of T_{1L} of a degree similar to monoexponential IR-FSE fits with short minimum TIs (17% underestimation of T_{1L}).

Simulations

As shown graphically in Figure 2 and numerically in Figure 4, omitting the short TI values from the analysis of IR-FSE data resulted in a monoexponential T_1^* fit that was in far better agreement with the long component T_{1L} . By comparing the results of fits with experimental data with increasingly longer minimum TI values, it can be seen that T_1^* steadily approached T_{1L} , though the variance in the measurement also tended to increase (Fig. 5a,b). The corresponding simulated data show the same behavior at 3T and 7T.

The simulations also demonstrated that acquiring more than four TIs has no significant impact on the deviation $T_{1L} - T_1^*$ except when the minimum TI was very short (data not shown). Additional TIs did result in decreased variance, but once this variance was normalized by the total scan time, this effect was also minimal, and the largest contribution to the variance was the minimum TI. Results for the entire range of simulated minimum TI values, with four TIs and 2% added noise, are shown in Fig. 5c and 5d (7T and 3T, respectively). As expected, the deviation diminished as TI_{\min} increased, while the variance

increased. The mean square error was lowest, with a minimum TI of 150 ms at 7T and 120 ms at 3T. If only 1% noise was added to simulated measurements instead (data not shown), the optimum TI_{\min} values were 200 ms and 150 ms at 7T and 3T, respectively.

Corresponding results for the errors introduced in MP2RAGE T_1 measurements are shown in Figure 6. Unlike IR-FSE, the difference between T_1^* and T_{1L} was never removed entirely. The difference decreased as both TIs increased, and though this increased the variance (Fig. 6b), the optimal MSE occurred when TI_1 and TI_2 were maximized. However, even at very long TIs, the difference at 3T could not be reduced below 90 ms with the parameters tested, and at 7T the minimum difference was 125 ms. Other image parameters, such as flip angle, TR, and TS, did not have a significant effect on these results (data not shown). With the parameters used experimentally at 7T (marked with an asterisk in Fig. 6a), the anticipated difference was approximately 200 ms, consistent with experimental observations (Table 3).

DISCUSSION

Biexponential T_1 relaxation is a known characteristic of many tissues, including WM, and the resulting short T_1 component is commonly measured in quantitative MT applications and used to compute parameters of interest. However, because the effect of T_{1S} at lower field strengths is on the order of a few percent, its impact on quantitative T_1 mapping has been largely ignored. Only now that T_1 mapping is becoming more common at very high field strengths is the true impact of biexponential relaxation becoming clear. Because accurate measurement of both biexponential components may be too time-consuming for most clinical applications, strategies to reduce the impact of T_{1S} are a more practical way to make T_1 measurements reproducible and independent of scan parameters.

It is worth emphasizing that T_{1S} does not correspond directly to the relaxation time of a particular population of protons, but is an apparent relaxation time constant whose value is governed by exchange effects. This may be one reason why the effects of biexponential T_1 behavior are not always recognized. For example, Kingsley et al. (33) noted that the quality of monoexponential fits to T_1 recovery data in WM would be greatly improved by the presence of a second T_1 component on the order of 50 ms that comprises approximately 10% of signal. The authors of that study believed a component having such properties was unlikely from a biological viewpoint, but these values are consistent with an interpretation of biexponential relaxation based on magnetization transfer, and in good agreement with both the present results and with previously reported qMT data (11,19).

At 7T, our IR-FSE data show the mean difference between T_1^* and T_{1L} to be 13% when T_1^* was calculated using a monoexponential fit to all available data. Because this data set had a minimum TI of 10 ms, this represents a near worst-case scenario in terms of deviation. If this T_1^* of 1153 ms is compared with T_1 values in WM reported in the literature (Table 1), it is clear that many of those measurements may have been affected by this deviation to varying degrees. In addition, some of the variation in reported WM T_1 values may be due to differences in minimum TIs, the selection of which varies significantly from protocol to protocol. In contrast, our $T_{1L} = 1349$ ms agrees with the value reported by Dortch et al. (11), who also used a biexponential analysis that should yield results that are independent of the

TIs acquired. It is our assertion that this is a more reliable estimate of the true T_1 of free water in WM.

The deviation from T_{1L} is also observable at 3T, albeit to a lesser degree. This is consistent with Equation [2], which predicts smaller deviations ($T_{1L} - T_1^*$) as T_{1L} decreases. Rather than the 13% deviation in T_1 observed at 7T, the effect at 3T is on the order of 6%. Both T_1^* and T_{1L} are within the range of accepted WM T_1 values at 3T [eg, (13)]. In addition, the proportion of WM voxels in which biexponential fits provide a significant benefit was much larger at 7T (90%–96%) compared with 3T (20%–45%). This finding is consistent with data at 1.5T (20) showing that biexponential fits provide little extra information at low fields and helps illustrate why short T_1 effects have not previously been considered significant for quantitative T_1 mapping.

The role of magnetic field inhomogeneity in the increased T_1^* deviation at high fields cannot be ignored. As B_0 increases, flip angles become less uniform and inversion pulses become less efficient, though these effects can be accounted for with appropriate models (32), and we attempted to minimize these effects on our acquisition by using overdriven adiabatic inversion pulses and centric phase encoding strategies that acquired the center of k-space after as few RF excitations as possible. Nevertheless, some of the observed deviation will depend on RF pulse parameters. For instance, as illustrated in Equation [1], changes in the magnetization of the pools due to, for example, incomplete saturation of the macromolecular pool will affect the initial amplitude of the observed relaxation components b_S and b_L , and therefore the observed deviation in T_1^* . For the pulses used in this study, we calculated that the saturation of the macromolecular pool [S_r in the notation of Pike (34), which is analogous to the value S_m reported by Dortch et al. (11)] varies from 0.16 to 0.31 (ie, macromolecular magnetization is reduced by 69%–84% after a single HSn pulse). This value depends on the pulse shape and power characteristics as well as the line shape (Gaussian, Lorentzian) and T_2 value ($T_{2m} = 10$ – $20 \mu\text{s}$) used to model the macro-molecular pool. The saturation effect of the pulses is significant, and slight differences between the HSn pulses used will affect the observed b_S and b_L to some degree. Though we have calculated the impact of S_r variations to be secondary to variations in macromolecular PSR, the influence of the inversion pulse remains a source of variability to be studied and controlled. However, we contend that not all of the observed T_1^* deviation can be attributed to RF inhomogeneity and pulse characteristics alone, because for a given value of w_S , our theoretical analysis shows a strong field strength dependence of T_1^* independent of these RF pulse effects.

To reduce the influence of T_{1S} on IR-FSE measurements, the minimum TI can be increased, but this will necessarily increase the variance in the measurement. Although this trade-off can be optimized by minimizing the mean squared error, some deviation from T_{1L} may remain, depending on the SNR of the underlying data. For example, in the simulations with 2% noise, the remaining deviation at the optimal minimum TI is 20 ms at 3T, and >40 ms at 7T. If higher accuracy is desired, the minimum TI must be further increased (eg, to 150 ms at 3T and 200 ms at 7T) to reduce the deviation to acceptable levels (eg, <10 ms). Additional scan time may then be needed to compensate for the increased variance.

Although this strategy has been presented here in the context of IR-FSE, it applies to some other methods as well. For example, the 3-TI MPRAGE approach described by Liu et al. (35) acquires 3 MPRAGE images at different TIs, with a ratio of differences used to create a T_1 map free of B_0 , B_1 , and T_2 effects. This method is a promising candidate for fast 3D T_1 mapping and can be used to obtain T_1 maps with larger brain coverage and faster acquisition times than IR-FSE methods (36). If all TIs are long enough that the short component has fully relaxed, the resulting T_1 map will be unaffected by T_{1S} . Indeed, Liu et al. suggest a minimum TI of 150 ms for this reason, a value that is in agreement with our recommendations.

However, the influence of T_{1S} cannot be removed from all sequences in this way. In particular, the T_1^* computed by MP2RAGE is significantly different from T_{1L} regardless of the TIs chosen. Although MP2RAGE T_1 values in WM were very consistent across the three volunteers imaged in this study, the T_1 given by MP2RAGE is significantly less than the T_1 obtained from a four-point IR-FSE series with optimized minimum TI, and the latter measurement is in good agreement with the biexponential T_{1L} value obtained using a 13-TI series. Although choosing very long TIs in MP2RAGE reduces this deviation somewhat, this comes at the cost of increased variance and scan time and will never eliminate the effect completely. This is a direct consequence of the lookup expression used for MP2RAGE, and only by exploring different approaches to ratio-based combination of images can this bias be potentially addressed. It is possible that this property of MP2RAGE has not been recognized previously because the T_1 values reported by MP2RAGE have often been in agreement with the T_1^* obtained from IR-FSE and similar measurements made with typically minimum short TIs. Expansion of the MP2RAGE signal equation to include the effects of two-site exchange may provide methods for increasing the accuracy of MP2RAGE T_1 measurements, though the collection of only two TIs may still limit the extent to which the biexponential relaxation can be modeled.

The impact of biexponential relaxation on other quantitative MRI methods must also be considered. For example, MR fingerprinting with inversion recovery balanced steady state free precession (37) uses a database spanning a range of potential parameters to attempt to match the acquired signal. If this database does not properly model biexponential relaxation, the apparent T_1 may also not reflect the underlying signal behavior. T_1 quantification using variable flip angle methods like DESPOT1 (12,14) may also be susceptible to deviations, though in a different way than described here due to the lack of inversion preparation.

This study was performed with the assumption of biexponential relaxation caused by two-site exchange between a free and a macromolecular pool. It is possible that a more accurate model would include components corresponding to more pools of varying MR visibility in different exchange regimes (eg, myelin water). While studies of T_1 relaxation using multicomponent DESPOT analysis (38) indicate that the myelin pool is likely in fast exchange with the free water pool, implying that the myelin water would primarily contribute to T_{1L} , this hypothesis should be explored further. T_1 analysis of three or more components may be challenging with the methods presented herein. It has been shown (39) that five TIs are sufficient to determine qMT parameters assuming a two-pool model, but

more points and higher SNR are likely necessary to reliably discriminate additional components.

Because T_{1S} arises from an exchange process between spin populations that may be modified by pathology, it is quite possible that T_{1S} contains information relevant to this pathophysiology and could itself be a desirable target for quantitative MRI. In this study, the goal was to reduce the influence of T_{1S} on the measurement of T_{1L} , but direct mapping of T_{1S} is also feasible. In addition to the method used in this study, MT methods have been used to estimate T_{1S} and w_S in 3D volumes with scan times on the order of 20 min (11), though this requires multiple TIs and provides limited spatial resolution (2–3 mm). It is possible that two consecutive T_1 mapping scans—one with parameters chosen to increase sensitivity to T_{1S} and one made insensitive to T_{1S} —could be combined to create simultaneous T_{1S} and T_{1L} maps with fewer total acquisitions. Correlations between T_{1S} and other measures of intercompartment exchange or macro-molecular content could then be explored more fully.

CONCLUSIONS

Variability between methods is a key obstacle to clinical implementations of quantitative T_1 mapping. The influence of biexponential T_1 relaxation on quantitative T_1 images has been overlooked because of its relative unimportance at low field strengths. However, as T_1 mapping at high field becomes common, it becomes critical to recognize these effects that, left unchecked, can introduce significant parameter-dependent variation in measurements of T_1 . While differences between the long T_1 component and the apparent T_1 in WM are on the order of 6% at 3T, we have shown that this difference can be as large as 13% at 7T.

For sequences such as IR-FSE in which data are fitted to a monoexponential model, this deviation between T_{1L} and T_1^* can be reduced through judicious selection of TIs. Optimal performance in terms of mean-squared error is obtained by acquiring four TIs with a minimum TI in the 100–150 ms range at 3T and 125–200 ms at 7T. The specific value of TI_{\min} should be guided by the SNR of the underlying data to achieve a suitable balance of increased variance and reduced bias. This strategy can also be employed for some other T_1 mapping methods such as 3-TI MPRAGE (36). In contrast, the $T_{1L} - T_1^*$ deviation present in MP2RAGE cannot be mitigated in the same way, and remains a significant source of error even at very long TIs. This must be kept in mind when considering the application of this particular sequence to T_1 mapping of the brain, particularly at high fields.

Acknowledgments

Grant sponsor: National Institutes of Health; Grant numbers: P41-EB015891; 1U54-A151459; 1P50-A114747; S10RR026351-01A1; Grant sponsor: GE Healthcare.

References

1. Vrenken H, Geurts JGG, Knol DL, et al. Whole-brain T_1 mapping in multiple sclerosis: global changes of normal-appearing gray and white matter. *Radiology*. 2006; 240:811–820. [PubMed: 16868279]

2. Manfredonia F, Ciccarelli O, Khaleeli Z, Tozer DJ, Sastre-Garriga J, Miller DH, Thompson AJ. Normal-appearing brain T₁ relaxation time predicts disability in early primary progressive multiple sclerosis. *Arch Neurol*. 2007; 64:411–415. [PubMed: 17353385]
3. Gouw AA, Seewann A, Vrenken H, van der Flier WM, Rozemuller JM, Barkhof F, Scheltens P, Geurts JJG. Heterogeneity of white matter hyperintensities in Alzheimer's disease: post-mortem quantitative MRI and neuropathology. *Brain*. 2008; 131:3286–3298. [PubMed: 18927145]
4. Li, TQ.; Deoni, SC. Fast T₁ Mapping of the Brain at 7T with RF Calibration Using Three Point DESPOT1 Method. Proceedings of the 14th Annual Meeting of ISMRM; Seattle, Washington, USA. 2006. p. 2643
5. Rooney WD, Johnson G, Li X, Cohen ER, Kim SG, Ugurbil K, Springer CS. Magnetic field and tissue dependencies of human brain longitudinal ¹H₂O relaxation in vivo. *Magn Reson Med*. 2007; 57:308–318. [PubMed: 17260370]
6. Ikonomidou, VN.; van Gelderen, P.; de Zwart, JA.; Duyn, JH. T₁ Measurements at 7T with Applications to Tissue Specific Imaging. Proceedings of the 14th Annual Meeting of ISMRM; Seattle, Washington, USA. 2006. p. 920
7. Wright, PJ.; Peters, A.; Brookes, M.; Coxon, R.; Morris, P.; Francis, S.; Bowtell, R.; Gowland, P. T₁ Measurements for Cortical Grey Matter, White Matter and Sub-cortical Grey Matter at 7T. Proceedings of the 14th Annual Meeting of ISMRM; Seattle, Washington, USA. 2006. p. 921
8. Wright PJ, Mouglin OE, Totman JJ, et al. Water proton T₁ measurements in brain tissue at 7, 3, and 1.5T using IR-EPI, IR-TSE, and MPRAGE: results and optimization. *Magn Reson Mater Phy*. 2008; 21:121–130.
9. Grinstead, J.; Rooney, W. Fast T₁ Mapping in Human Brain Using Inversion Recovery EPI with GRAPPA at 3T and 7T. Proceedings of the 16th Annual Meeting of ISMRM; Toronto, Ontario, Canada. 2008. p. 3084
10. Marques JP, Kober T, Krueger G, van der Zwaag W, van de Moortele PF, Gruetter R. MP2RAGE, a self bias-field corrected sequence for improved segmentation and T₁-mapping at high field. *NeuroImage*. 2010; 49:1271–1281. [PubMed: 19819338]
11. Dortch RD, Moore J, Li K, Jankiewicz M, Gochberg DF, Hirtle JA, Gore JC, Smith SA. Quantitative magnetization transfer imaging of human brain at 7T. *NeuroImage*. 2013; 64:640–649. [PubMed: 22940589]
12. Dieringer MA, Deimling M, Santoro D, Wuerfel J, Madai VI, Sobesky J, von Knobelsdorff-Brenkenhoff F, Schulz-Menger J, Niendorf T. Rapid parametric mapping of the longitudinal relaxation time T₁ using two-dimensional variable flip angle magnetic resonance imaging at 1.5 Tesla, 3 Tesla, and 7 Tesla. *PLoS One*. 2014; 9:e91318. [PubMed: 24621588]
13. Stikov N, Boudreau M, Levesque IR, Tardif CL, Barral JK, Pike GB. On the accuracy of T₁ mapping: searching for common ground. *Magn Reson Med*. 2014;1002/mrm.25135
14. Deoni SC, Peters TM, Rutt BK. High-resolution T₁ and T₂ mapping of the brain in a clinically acceptable time with DESPOT1 and DES-POT2. *Magn Reson Med*. 2005; 53:237–241. [PubMed: 15690526]
15. Look DC, Locker DR. Time saving in measurement of NMR and EPR relaxation times. *Rev Sci Instrum*. 1970; 41:250–251.
16. Cho S, Jones D, Reddick WE, Ogg RJ, Steen RG. Establishing norms for age-related changes in proton T₁ of human brain tissue in vivo. *Magn Reson Imaging*. 1997; 15:1133–1143. [PubMed: 9408134]
17. Balaban RS, Chesnick S, Hedges K, Samaha F, Heineman FW. Magnetization transfer contrast in MR imaging of the heart. *Radiology*. 1991; 180:671–675. [PubMed: 1871277]
18. Wolff SD, Chesnick S, Frank JA, Lim KO, Balaban RS. Magnetization transfer contrast: MR imaging of the knee. *Radiology*. 1991; 179:623–628. [PubMed: 2027963]
19. Dortch RD, Li K, Gochberg DF, Welch EB, Dula AN, Tamhane AA, Gore JC, Smith SA. Quantitative magnetization transfer imaging in human brain at 3T via selective inversion recovery. *Magn Reson Med*. 2011; 66:1346–1352. [PubMed: 21608030]
20. Kjaer L, Thomsen C, Henriksen O. Evaluation of biexponential relaxation behaviour in the human brain by magnetic resonance imaging. *Acta Radiologica*. 1989; 30:433–437. [PubMed: 2775605]

21. Prantner AM, Bretthorst GL, Neil JJ, Garbow JR, Ackerman JJH. Magnetization transfer induced biexponential longitudinal relaxation. *Magn Reson Med*. 2008; 60:555–563. [PubMed: 18759367]
22. Edzes HT, Samulski ET. The measurement of cross-relaxation effects in the proton NMR spin-lattice relaxation of water in biological systems: hydrated collagen and muscle. *J Magn Reson*. 1978; 31:207–229.
23. Henkelman RM, Huang X, Xiang QS, Stanisz GJ, Swanson SD, Bronskill MJ. Quantitative interpretation of magnetization transfer. *Magn Reson Med*. 1993; 29:759–766. [PubMed: 8350718]
24. Hazelwood CF, Chang DC, Nichols BL, Woessner DE. Nuclear magnetic resonance transverse relaxation times of water protons in skeletal muscle. *Biophys J*. 1974; 14:583–606. [PubMed: 4853385]
25. Ou X, Gochberg DF. MT effects and T1 quantification in single-slice spoiled gradient echo imaging. *Magn Reson Med*. 2008; 59:835–845. [PubMed: 18302249]
26. Weiss GH, Ferretti JA. Optimal design of relaxation time experiments. *Prog Nucl Magn Reson Spectrosc*. 1988; 20:317–355.
27. Ogg RJ, Kingsley PB. Optimized precision of inversion-recovery T₁ measurements for constrained scan time. *Magn Reson Med*. 2004; 51:625–630. [PubMed: 15004808]
28. van de Moortele PF, Auerbach EJ, Olman C, Yacoub E, Ugurbil K, Moeller S. T₁ weighted brain images at 7 Tesla unbiased for Proton Density, T₂* contrast and RF coil receive B₁ sensitivity with simultaneous vessel visualization. *NeuroImage*. 2009; 46:432–446. [PubMed: 19233292]
29. Marques JP, Gruetter R. New developments and applications of the MP2RAGE sequence—focusing the contrast and high spatial resolution R1 mapping. *PLoS One*. 2013; 8:e69294. [PubMed: 23874936]
30. Tannus A, Garwood M. Adiabatic pulses. *NMR Biomed*. 1997; 10:424–434.
31. Saranathan M, Tourdias T, Bayram E, Ghanouni P, Rutt BK. Optimization of white-matter-nulled magnetization prepared rapid gradient echo (MP-RAGE) imaging. *Magn Reson Med*. 2014; 76:1002–1009. [PubMed: 25298]
32. Barral JK, Gudmundson E, Stikov N, Etezadi-Amoli M, Stoica P, Nishimura DG. A robust methodology for in vivo T₁ mapping. *Magn Reson Med*. 2010; 64:1057–1067. [PubMed: 20564597]
33. Kingsley PB, Ogg RJ, Reddick WE, Steen RG. Correction of errors caused by imperfect inversion pulses in MR imaging measurements of T₁ relaxation times. *Magn Reson Imaging*. 1998; 16:1049–1055. [PubMed: 9839989]
34. Pike GB. Pulsed magnetization transfer contrast in gradient echo imaging: a two-pool analytic description of signal response. *Magn Reson Med*. 1996; 36:95–103. [PubMed: 8795027]
35. Liu JV, Bock NA, Silva AC. Rapid high-resolution three-dimensional mapping of T₁ and age-dependent variations in the non-human primate brain using magnetization-prepared rapid gradient-echo (MPRAGE) sequence. *NeuroImage*. 2011; 56:1154–1163. [PubMed: 21376814]
36. Rioux, JA.; Levesque, IR.; Saranathan, M.; Rutt, BK. 10-Minute High-Resolution Whole-Brain T1 Mapping: A Comparison of Three Candidate Methods. Proceedings of the 22nd Annual Meeting of ISMRM; Milan, Italy. 2014. p. 4320
37. Ma D, Gulani V, Seiberlich N, Liu K, Sunshine JL, Duerk JL, Griswold MA. Magnetic resonance fingerprinting. *Nature*. 2013; 495:187–193. [PubMed: 23486058]
38. Deoni SCL, Rutt BK, Arun T, Pierpaoli C, Jones DK. Gleaning multi-component T₁ and T₂ information from steady-state imaging data. *Magn Reson Med*. 2008; 60:1372–1387. [PubMed: 19025904]
39. Li K, Zu Z, Xu J, Janve VA, Gore CJ, Does MD, Gochberg DF. Optimized inversion recovery sequences for quantitative T1 and magnetization transfer imaging. *Magn Reson Med*. 2010; 64:491–500. [PubMed: 20665793]

APPENDIX A

Consider a monoexponential fit to three data points using, for example, an IR-FSE sequence with inversion times TI_1 , TI_2 and TI_3 . The most general model that can be fitted reliably to such data using a nonlinear least-squares method has three parameters:

$$S(TI)=a+b \times \exp(-TI/T_1). \quad [A1]$$

While analytically characterizing such a nonlinear fit is difficult in general, for the specific case of a three-parameter fit to three data points, the solution should be unique. Therefore, a solution found by alternative means should be identical, given the same input data. One alternate method for computing T_1 (35) is by forming the expression

$$L=\frac{S(TI_3)-S(TI_1)}{S(TI_3)-S(TI_2)}=\frac{\exp(-TI_3/T_1)-\exp(-TI_1/T_1)}{\exp(-TI_3/T_1)-\exp(-TI_2/T_1)}. \quad [A2]$$

This expression depends only on T_1 , and while it is generally not possible to solve directly for T_1 in terms of the TIs and signal intensities, a simple lookup can be performed based on the known TIs to compute T_1 for any given value of L .

Now consider the case in which the underlying relaxation is actually biexponential, of the form

$$S(TI)=a+b_L \exp(-TI/T_{1L})+b_S \exp(-TI/T_{1S}). \quad [A3]$$

The relative contributions of the two T_1 components, T_{1L} and T_{1S} , are given by b_L and b_S , respectively. An expression analogous to Equation [A2] that includes this biexponential behavior is

$$L=\frac{b_L \exp(-TI_3/T_{1L}) \pm b_S \exp(-TI_3/T_{1S}) - b_L \exp(-TI_1/T_{1L}) - b_S \exp(-TI_1/T_{1S})}{b_L \exp(-TI_3/T_{1L}) + b_S \exp(-TI_3/T_{1S}) - b_L \exp(-TI_2/T_{1L}) - b_S \exp(-TI_2/T_{1S})}. \quad [A4]$$

Two assumptions can be made to greatly simplify this expression. First, for most experimental parameters, TI_2 and TI_3 are much longer than T_{1S} , such that $\exp(-TI_3/T_{1S}) = \exp(-TI_2/T_{1S}) \approx 0$. We will also assume that TI_3 is long compared with T_{1L} and $\exp(-TI_3/T_{1L}) \approx 0$. This is decreasingly justified at higher fields, but as a first-order approximation this assumption is worth considering. The lookup expression for biexponential relaxation is then approximately

$$L_{bi}=(b_L \exp(-TI_1/T_{1L})+b_S \exp(-TI_1/T_{1S}))/b_L \exp(-TI_2/T_{1L}). \quad [A5]$$

The corresponding expression for purely monoexponential relaxation has also been simplified to

$$L_{\text{mono}} = \exp(-TI_1/T_1) / \exp(-TI_2/T_1) \\ = L_{\text{bi}} - b_s \exp(-TI_1/T_{1s}) / b_L \exp(-TI_2/T_{1L}). \quad [\text{A6}]$$

If the monoexponential lookup function L_{mono} is used to determine T_1 for data that actually exhibit biexponential relaxation, the lookup will be incorrect by an amount given by

$$\Delta L = b_s \exp(-TI_1/T_{1s}) / b_L \exp(-TI_2/T_{1L}). \quad [\text{A7}]$$

This can be converted to units of T_1 using $T_1 = L \, dT_1/dL = L/(dL/dT_1)$. Using $L = L_{\text{mono}}$ and

$$dL/dT_1 = \frac{(TI_1 - TI_2) \times \exp(-TI_1/T_1)}{T_1^2 \exp(-TI_2/T_1)}. \quad [\text{A8}]$$

we find the following expression for the change in T_1^* introduced by the short component:

$$\Delta T_1^* = -[T_{1L}^2 / (TI_2 - TI_1)] \times b_s / b_L \times [\exp(-TI_1/T_{1s}) / \exp(-TI_1/T_{1L})]. \quad [\text{A9}]$$

Though it is acknowledged that this expression relies on the validity of the assumptions concerning the magnitudes of TI_2 and TI_3 compared with T_{1s} and T_{1L} , this expression does give some intuition into the increased influence of biexponential relaxation at high field. In addition, the general behavior described by Equation [A9] should be extensible to larger numbers of TIs.

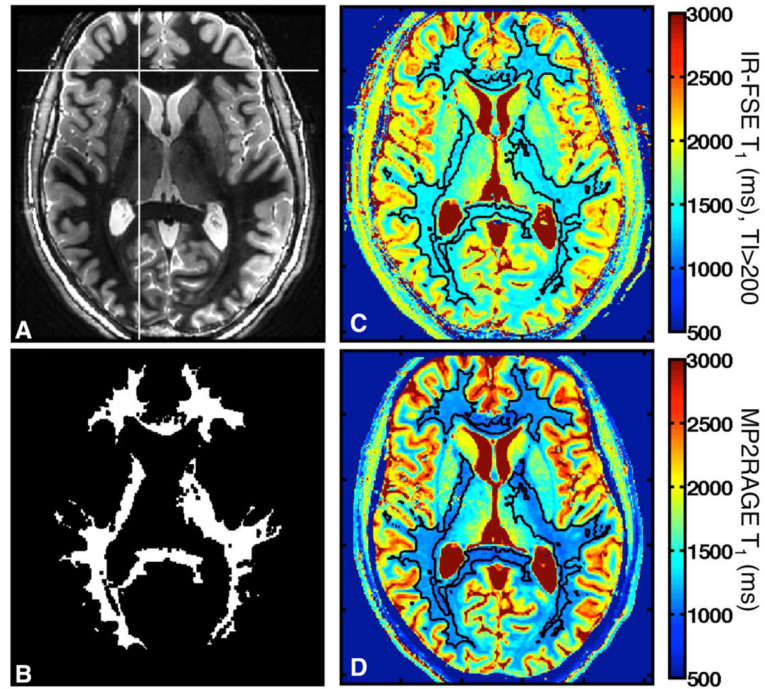


FIG. 1.

Generation of WM mask for T_1 analysis. (a) Example IR-FSE image with $TI = 600$ ms and one seed voxel used for ROI generation (crosshairs). (b) Resulting WM ROI after erosion. (c) Corresponding IR-FSE T_1 map ($TI > 200$ ms), with the ROI outlined in black. (d) Corresponding MP2RAGE T_1 map.

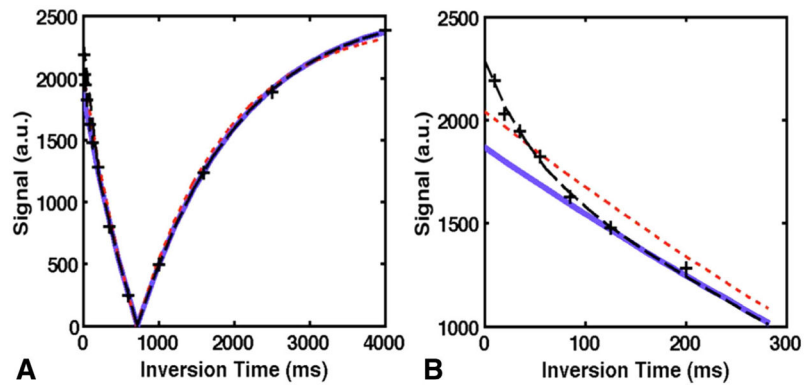


FIG. 2. Comparison of monoexponential and biexponential fits. **(a)** Fits to data at the voxel shown in Figure 1a; black dashed line = biexponential fit, red dotted line = monoexponential fit to all data, solid blue line = monoexponential fit to TI > 200 ms. **(b)** Inset showing early TI values only.

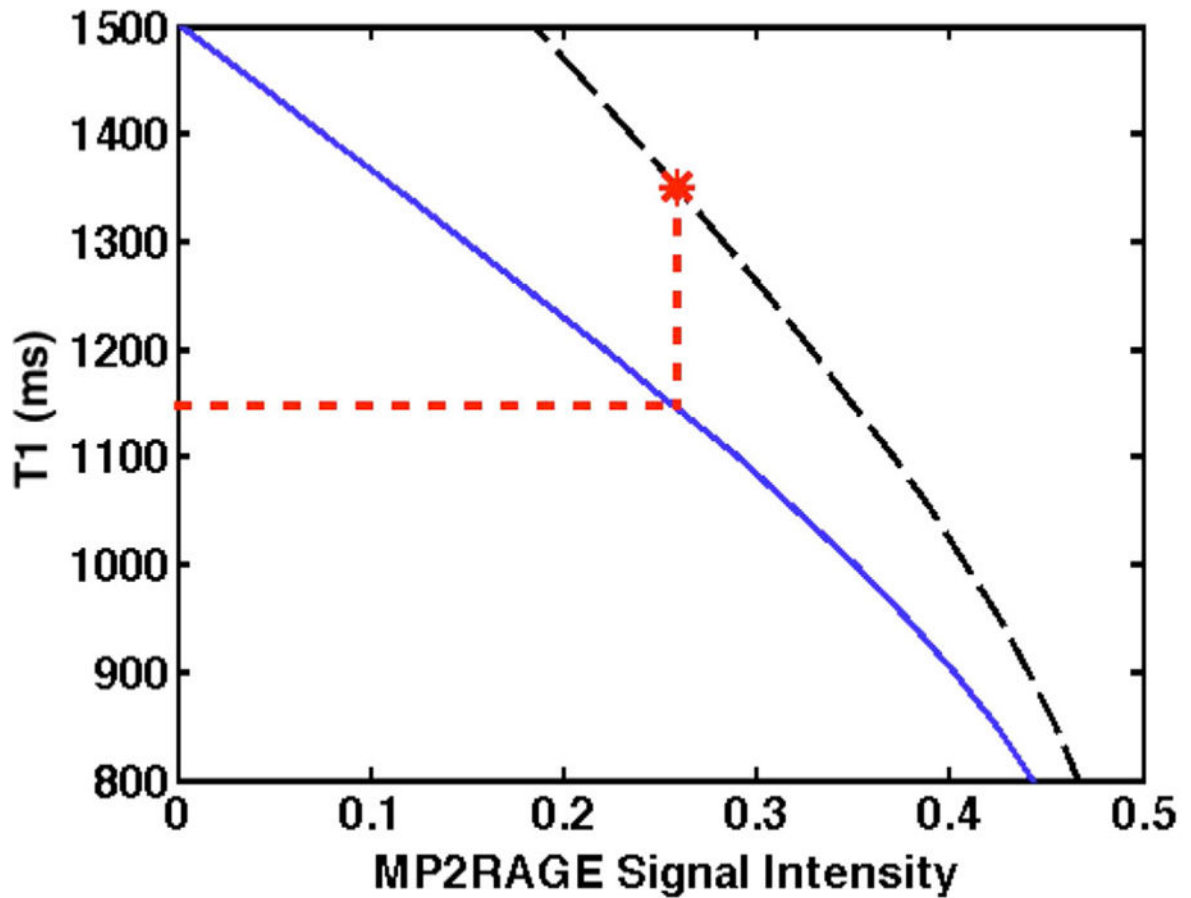


FIG. 3.

Example of an incorrect MP2RAGE T_1 lookup in the presence of a short T_1 component. The correct lookup table, based on biexponential relaxation for a range of T_{1L} values with T_{1S} and w_S constant, is indicated by the black dashed line, with a particular T_1 value of interest ($T_{1L} = 1350$ ms) marked by the asterisk. The lookup table assuming monoexponential relaxation is indicated by the blue solid line. If a lookup (red dotted line) is performed using the monoexponential table and the highlighted signal intensity, the resulting T_1^* (1146 ms) is much less than T_{1L} .

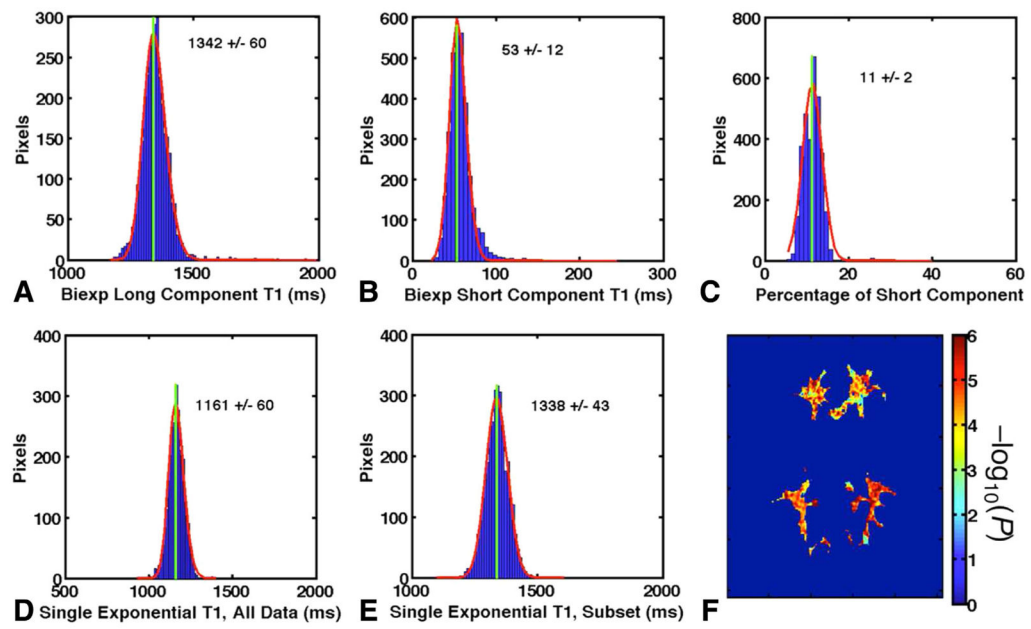


FIG. 4. Results from a typical WM ROI (shown in Fig. 1b). (a) Gaussian fit to histogram of biexponential long component T_{1L} . (b) Biexponential short component T_{1S} . (c) Short component weight w_s . (d) Monoexponential T_1^* using all data. (e) Monoexponential T_1^* using only $TI > 200$ ms. (f) P value map of entire WM ROI. For ease of visualization, $-\log_{10}(P)$ is displayed such that $P = 0.05$ corresponds to a value of 1.3 and $P = 0.01$ corresponds to a value of 2.

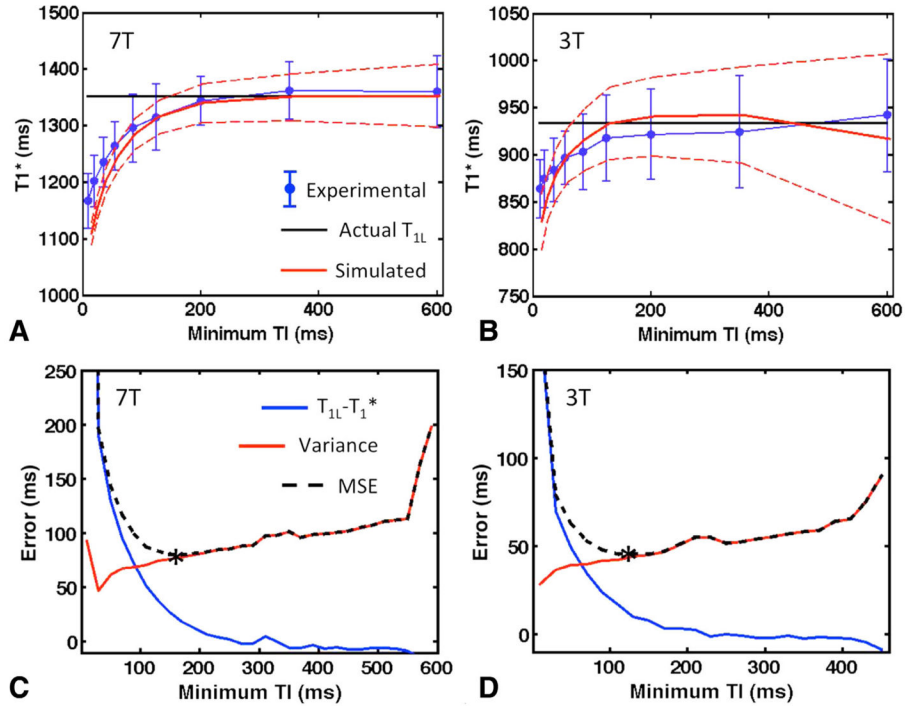


FIG. 5. Behavior of T_1^* in IR-FSE acquisitions with increasing minimum TI. **(a, b)** Comparison of simulated and experimental data at 7T **(a)** and 3T **(b)**. Blue dots are experimental results of monoexponential T_1^* fits (mean \pm standard deviation throughout ROI from subject #1 at 7T and subject #2 at 3T). The solid black line represents the T_{1L} used as input to the simulations, and the red dashed lines are T_1^* fits to simulated data with the same minimum TI and number of TIs (mean \pm standard deviation over 5000 repetitions). **(c, d)** Simulated mean deviation $T_{1L} - T_1^*$, variance σ^2 in the fitted T_1^* values, and $MSE = (T_{1L} - T_1^*)^2 + \sigma^2$ as a function of minimum TI at 7T **(c)** and 3T **(d)**. Mean and variance were measured over 5000 Monte Carlo repetitions with 2% noise added to the data. The optimal TI for each field strength is marked with an asterisk.

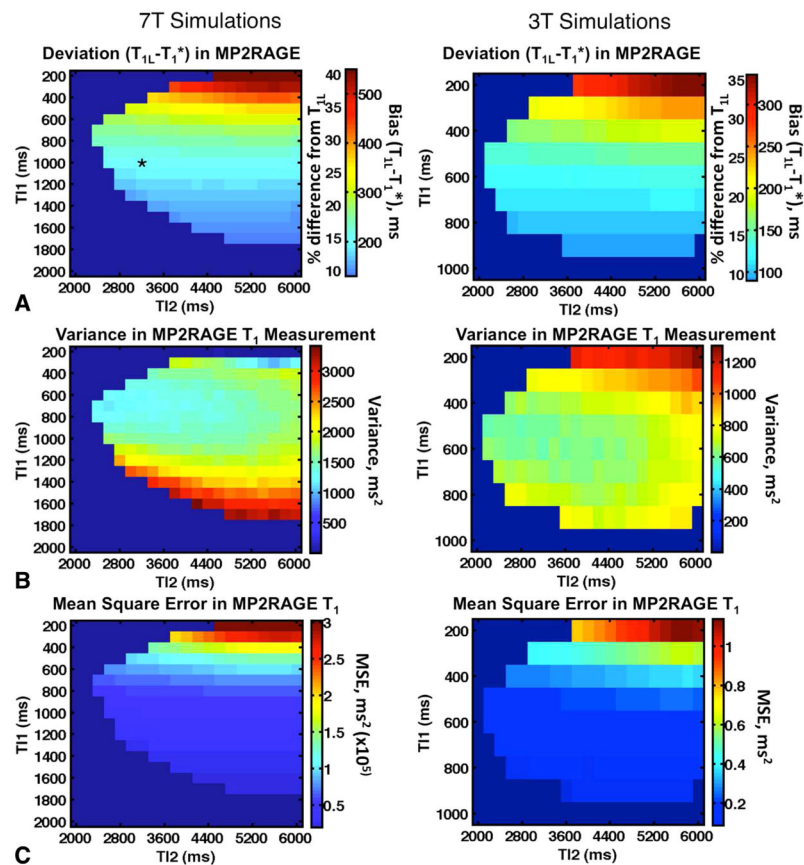


FIG. 6.

(a) Simulated deviation in MP2RAGE T_1 measurement ($T_{1L} - T_1^*$) as a function of TIs T_{11} and T_{12} , at 7T (top) and 3T (bottom). Values shown are the mean of 5000 Monte Carlo repetitions. The combination used for experimental comparison of MP2RAGE with IR-FSE is marked with an asterisk. Combinations that were disallowed due to ambiguous lookups or unfeasible parameters appear in dark blue. (b) Variance σ^2 in the T_{1app} values at 2% added noise. (c) $MSE = (T_{1L} - T_1^*)^2 + \sigma^2$.

Table 1Reported Measurements of T_1 in WM at 7T

Study	Reported WM T_1 (ms)	Region	Sequence and Parameters
Li and Deoni (4)	1500 ± 100	Not stated	DESPOT1, flip angle = $5^\circ/10^\circ/20^\circ$
Rooney et al. (5)	1220 ± 36	Frontal	Look-Locker, 32 TIs, 20 ms to 10 s
Ikonomidou et al. (6)	1357 ± 22	Not stated	IR-EPI, TIs 30–8000 ms
Wright et al. (7)	1017 ± 86	Cortical	IR-FSE, 9 TIs, 120–2000 ms
Wright et al. (8) ^a	~1250?	Cortical	IR-EPI, 9 TIs, 120–4000 ms
Wright et al. (8)	1130 ± 100	Cortical	MPRAGE, 8 TIs, 160–2100 ms
Grinstead and Rooney (9)	1180 ± 40	Not stated	IR-EPI, TI values not reported
Marques et al. (10)	1150 ± 60	Frontal	MP2RAGE, TI = 1000/3300 ms
Marques et al. (10)	1110 ± 60	Corpus callosum	MP2RAGE, TI = 1000/3300 ms
Dortch et al. (11)	1372	Global mean	Selective IR qMT, 15 TIs, 6 ms to 2s
Dieringer et al. (12)	1284 ± 22	Global mean	IR-FSE, 8 TIs, 60–5000 ms
Dieringer et al. (12)	1855 ± 141	Global mean	2D VFA, flip angle = $3^\circ/8^\circ/25^\circ$

Abbreviations: DESPOT1, driven equilibrium single point observation of T_1 ; IR-EPI, inversion recovery echo planar imaging; IR-FSE, inversion recovery fast spin echo; MPRAGE, magnetization prepared rapid acquisition of gradient echo; MP2RAGE, magnetization prepared 2 rapid acquisition gradient echoes; qMT, quantitative magnetization transfer; VFA, variable flip angle.

^aThe measurement using IR-EPI is not reported in the text but is inferred from Figure 5 in Wright et al. 2008.

Table 2

T₁ Relaxation in Volunteers at 7T and 3T Based on Biexponential and Monoexponential Fits to 13 IR-FSE Images

Volunteer	T _{1L} (ms)	T _{1S} (ms)	Short T ₁ Weight, w _S	T ₁ * (ms)	% Difference (T ₁ */T _{1L} - 1) ^a	% Voxels	
						P < 0.05	P < 0.01
7T							
#1	1342 ± 60	53 ± 12	11% ± 2%	1161 ± 60	-13%	99.5%	98.8%
#2	1347 ± 64	51 ± 19	10% ± 2%	1173 ± 39	-13%	94.8%	80.0%
#3	1331 ± 77	61 ± 12	11% ± 2%	1106 ± 95	-17%	91.1%	86.9%
#4	1374 ± 46	64 ± 14	12% ± 2%	1162 ± 38	-18%	97.2%	93.3%
Mean ^b	1349 ± 18	57 ± 6	11% ± 1%	1151 ± 30	-15% ± 3%	96% ± 4%	90% ± 8%
3T							
#1	949 ± 54	54 ± 40	7% ± 2%	906 ± 43	-4%	50.8%	27.5%
#2	928 ± 37	47 ± 48	10% ± 2%	864 ± 26	-7%	43.0%	14.8%
#3	926 ± 79	49 ± 64	9% ± 3%	869 ± 60	-6%	43.0%	19.3%
#4	952 ± 67	42 ± 43	8% ± 2%	887 ± 40	-7%	44.3%	20.2%
Mean ^b	939 ± 14	48 ± 5	9% ± 1%	882 ± 19	-6% ± 1%	45% ± 4%	20% ± 5%

Values are for an ROI encompassing all of the WM in an axial slice through the basal ganglia. Subjects #2-4 are the same volunteer at both field strengths.

^a Gives the error in the fitted T₁* seen in monoexponential fits to all data compared with the biexponential long component T_{1L}.

^b Mean ± standard deviation of the mean across all volunteers at that field strength.

Table 3Comparison of WM T₁ Values in MP2RAGE and IR-FSE at 7T

Subject	IR-FSE T ₁ (ms)	MP2RAGE T ₁ (ms)	% Difference
#5	1380 ± 66	1128 ± 69	-18%
#6	1354 ± 50	1125 ± 75	-17%
#7	1351 ± 66	1130 ± 56	-16%
Mean ^a	1362 ± 16	1127 ± 3	-17% ± 1%

^aMean ± standard deviation of the data for all volunteers.

Author Manuscript

Author Manuscript

Author Manuscript

Author Manuscript