SCIENTIFIC REPORTS

Received: 07 March 2016 Accepted: 16 May 2016 Published: 16 June 2016

OPEN Acoustic one-way metasurfaces: **Asymmetric Phase Modulation of** Sound by Subwavelength Layer

Xue Jiang^{1,2}, Bin Liang^{1,2}, Xin-ye Zou^{1,2,3}, Jing Yang^{1,2}, Lei-lei Yin⁴, Jun Yang⁵ & Jian-chun Cheng^{1,2}

We theoretically design and numerically demonstrate an acoustic one-way metasurface, which is a planar and acoustically subwavelength layer behaving like a nearly-reflectionless surface with arbitrary wave-steering capability for incident wave impinging on one side, while virtually blocking the reversed wave. The underlying mechanism is based on an asymmetric phase modulation by coupling a phase array and a near-zero-index medium. We exemplify a metastructure-based implementation by combining the hybrid metastuctures and labyrinthine structures. Moreover, the performance of the proposed implementation is demonstrated via three distinct phenomena of anomalous refraction, wave splitting and conversion of propagation wave to surface wave. Our findings may offer more possibilities for sound manipulation and improve the application potential of acoustic artificial devices in situations such as ultrasonic imaging and therapy.

The recent emergence of acoustic metamaterials provides acoustic properties unavailable in nature¹⁻¹⁰, and has enabled design of various conceptual devices with special wave-manipulation functionalities such as invisibility and cloaking²⁻³, sub-diffraction imaging lens^{5,6}, acoustic field rotator⁷, etc. In particular, acoustic metasurfaces¹⁰⁻¹⁸ have aroused increasing interests due to the ability to control the transmitted and reflected phase of waves to generate diverse functionalities with a planar layer of subwavelength thickness. The existing designs of metasurfaces, however, only yield completely symmetric manipulations for incident wave impinging on the two opposite sides of the metasurfaces, which poses a fundamental challenge for the development of wave-steering devices. The possibility of breaking through the barrier of how to realize asymmetric acoustic transmission has attracted extensive interests¹⁹⁻²⁸ and is believed to have the potential to improve the quality of medical ultrasound imaging^{20,27}. Despite the recent progresses in the acoustic one-way devices developed in both linear²⁴⁻²⁸ and nonlinear regimes¹⁹⁻²³, the resulting "acoustic diodes", which generally have a bulky size, are only designed to possess asymmetric transmission of energy flux but unable to provide free manipulation on the transmitted wavefront via rational design of their structural parameters. The great potential of combining the capability of one-way manipulation and the flexibility of wavefront steering have not yet been explored, which should have deep implications for acoustic devices, acoustic-based applications and the field of acoustics in general.

In this article, we theoretically design and numerically realize an acoustic one-way metasurface (AOM) capable of freely tailoring the transmitted wave for incident wave impinging on one side but prohibiting the reversed wave from passing. Unlike the previous metasurface schemes giving rise to identical phase profile for incident wave from two sides, on which a versatile but unavoidably symmetric sound manipulation is based, the current AOM design has an inherently different mechanism that generates an asymmetric phase modulation along two opposite directions. Such a unique ability is realized by using a combination of acoustic phase array (PA) with the flexibility of free phase-steering and near-zero-index material (ZIM) with the important feature of high phase selectivity and phase preservation. We have also demonstrated a practical implementation of our design by artificial structures, in which a hybrid metastucture¹⁶ with full phase control is employed to freely steer the transmitted

¹Key Laboratory of Modern Acoustics, MOE, Institute of Acoustics, Department of Physics, Nanjing University, Nanjing 210093, P. R. China. ²Collaborative Innovation Center of Advanced Microstructures, Nanjing University, Nanjing 210093, P. R. China. ³State Key Laboratory of Acoustics, Chinese Academy of Sciences, Beijing 100190, P. R. China. ⁴Imaging Technology Group, Beckman Institute, University of Illinois at Urbana-Champaign, Urbana, Illinois 61801, USA. ⁵Key Laboratory of Noise and Vibration Research, Institute of Acoustics, Chinese Academy of Sciences, Beijing 100190, P. R. China. Correspondence and requests for materials should be addressed to B.L. (email: liangbin@nju.edu.cn) or J.C.C. (email: jccheng@nju.edu.cn)



Figure 1. (a) Schematic of the proposed AOM design which is a combination of PA and ZIM as shown in the inset. (b) Schematic of the asymmetric phase manipulation of the AOM in PD and ND cases. Inset: a zoom-in of a part of the AOM.

wave¹⁶ and a labyrinthine structure^{27,29} is used to effectively mimic a ZIM³⁰. Performances of the proposed AOM are demonstrated by realizing three distinctive phenomena of anomalous refraction, splitting acoustic wave and converting the propagating wave into surface wave. Combining the capability of asymmetric transmission and flexibility of manipulating the wave profiles together, the realization of AOM allows for more versatile manipulation of sound and opens up possibilities for the design of asymmetric manipulation devices that may find applications in various situations such as ultrasonic imaging and therapy.

Results

The schematic diagram of the AOM is shown in Fig. 1(a), which is characterized by a subwavelength, planar configuration and the special wave-steering functionality that only works for incident wave from one particular side. The proposed AOM is a combination of acoustic PA and ZIM as shown in the inset of Fig. 1(a). In the current study, we define the propagating directions of the wave normally incident on the ZIM and PA as the positive direction (PD) and negative direction (ND) respectively. Only the acoustic plane wave incident from PD is permitted to transmit through the AOM and makes up the desired wavefront profile, while the wave propagating from ND is prohibited, as represented by the purple and blue arrows in Fig. 1(a), respectively. The PA offers a great deal of freedom and high efficiency in modulating the phase of transmitted wave flexibly that is, however, entirely symmetric in PD and ND. The extremely large phase velocity provided by the ZIM, on the other hand, leads to the near-zero critical angle and intrinsically high selectivity of the ZIM²⁷. These result in two important features that are crucial for the effectiveness of the AOM: the high phase selectivity that only allows incident wave with uniform phase distribution in the transverse direction to penetrate while blocking the wave with non-uniform phase distribution, and the phase preservation capability that maintains the propagating phase of incident wave unchanged when passing through the planar ZIM and causes the transmitted wave and incident wave to share exactly the same phase profile. For the proposed AOM composed of the two components possessing their respective important features mentioned above, it can therefore be theoretically predicted that the AOM would behave differently in steering the phase of the incident wave in opposite directions, as will be proven later. This should be understood as the underlying mechanism responsible for the realization of one-way wavefront manipulation, which is inherently different from the mechanisms of the existing metasurface designs that provide symmetric phase responses only. This asymmetric phase manipulation ability of the AOM is well performed in Fig. 1(b). When the acoustic plane wave normally incident along the PD, the wave will be permitted to transmit through the ZIM due to the uniform phase profile which can be preserved until the wave leaves this part. Then the interaction with the PA will freely steer the phase to construct the desired phase profile $\phi(y)$, as represented by the purple dots in Fig. 1(b). In the ND, contrarily, the uniform phase profile of the incident wave has been disturbed and appended with phase $\phi(y)$ by PA before reaching the ZIM, leading to the blockage and reflection of the incident wave. The reflective wave in the output plane of the metasurface will be appended with phase $2(\phi(y) + kd)$ as indicated by the blue dots in Fig. 1(b), where d is the thickness of the PA in the propagation direction. As a consequence, a unidirectional transmission of acoustic waves is expected in the proposed system, and the wavefront of the transmitted wave in the PD can be arbitrarily engineered by designing different $\phi(y)$ shown in Fig. 1(b).

In practice, the PA and ZIM, the two key components characterized by flexible phase control and extremely large phase velocity respectively, need to be implemented by real acoustical materials. Next we will demonstrate a practical implementation of the proposed scheme of AOM for airborne sound, as shown in Fig. 2. In this study we use three layers of labyrinthine structure designed by coiling up space and the hybrid metastructure comprising four Helmholtz cavities and a straight pipe as the ZIM and PA, respectively. In these two employed structures, the walls are regarded as rigid with a thickness of λ /500, where λ is the sound wavelength. Other geometric parameters are shown in Fig. 2. Through adjusting the height of Helmholtz cavities one can manipulate the phase of the





transmitted wave in the 2π range freely with almost-unity transmission¹⁶. And it has also been proven that the labyrinthine metastructure is able to behave as a ZIM by delicately adjusting its geometric paramters^{27,29}. As a result, this implementation of AOM has a total thickness as thin as 0.57λ and a completely planar configuration. These important features would be particularly significant for practical applications due to the potential to minimize the size of devices and improve the compatibility to couple with the surroundings. Furthermore, it is worth pointing out that the general scheme proposed here for designing AOM are not exclusively restricted to these two particular configurations, but can also implemented by employing other candidates with tunable phase response and near-zero index to further improve the performance of the resulting device. This particular implementation composing of a hybrid metastructure and a labyrinthine structure for simplicity, has the optimal performance near the resonance frequency, where these two components are delicately designed to yield the near-zero index and the phase-modulating ability respectively. Beyond the designed working band, the performance of the implemented one-way metasurfaces will be affected, leading to reduction in the contrast factor. It is worth stressing however our proposed scheme of designing a one-way metasurface by coupling a phase array and a ZIM is general and should be irrelevant to the operating frequency. A broadband functionality can be achieved in practical implementations as long as specific kinds of metamaterials with broadband or non-dispersive near-zero indices are employed.

We have carried out a series of numerical simulations to verify our theoretical predictions as well as demonstrate the performance of the resulting AOM. As particular examples, three distinctive phenomena are chosen and will be exemplified in what follows: acoustic anomalous refraction, wave splitting and conversion of propagation wave to surface wave. The AOMs for these three examples have the same thickness 0.57λ , which are comprised of a layer of hybrid metastructure and three layers of labyrinthine structure.

Acoustic anomalous refraction is the phenomenon that the transmitted wave deviates from the incident direction when the wave impinges on a surface³¹. And the anomalous emergent angle, θ_t , can be determined by designing the specific phase distribution $\phi(y)$ on the surface, which is governed by the generalized Snell's law³¹:

$$\theta_{\rm t} = \arcsin\left(\sin\theta_{\rm i} + \frac{\lambda}{2\pi}\frac{d\phi}{dy}\right) \tag{1}$$

where θ_i is the incident angle, and the phase gradient term $(\lambda/2\pi)(d\phi/dy)$ determines the anomalous emergent angle θ_i . Here we choose a normally incident plane wave with θ_i being zero, and the emergent angle θ_i is selected to be $\pi/4$. Figure 3(a) is the desired asymmetric phase profiles of the AOM in this situation, which can be realized by appropriately designing the geometric parameter of each unit cell of the hybrid metastructure mimicking the PA. The pressure patterns in PD and ND cases are shown respectively in Fig. 3(b,c), in which the black arrows denote the corresponding incident directions, the green arrow indicates the theoretical emergent angle and the red arrow is the reflection direction. As shown in Fig. 3(b), when the AOM is illuminated normally by a plane wave propagating in PD, full transmission and preserved phase distribution are permitted by the ZIM, and the designed wavefront profile of the transmitted wave is formed by the modulation effect of the PA. Excellent agreement can be observed between the theoretical prediction of the anomalous transmitted direction (green arrow) and the simulation result. On the contrary, the phase of the incident plane wave in the ND will be disturbed and attached with an additional phase by the PA first. Therefore, the ND incident wave is reflected back as shown in Fig. 3(c). For quantitatively evaluating the discrepancy between the transmissions in PD and ND case, a parameter of contrastive factor, defined as $10 \log(I_{PD}/I_{ND})$ with $I_{PD}(I_{ND})$ being the acoustic intensity in PD (ND) case, is introduced. The acoustic intensities are calculated by integrating the intensity in the corresponding surfaces perpendicular to the propagation direction. The numerical results reveals that the contrastive factor in this situation reaches 12.5 dB, which verifies the high efficiency of the distinct unidirectional property of the proposed AOM.

Then we present the realization of splitting acoustic beam based on the proposed scheme, with the plane wave being confined and guided to two symmetric directions with respect to the incident direction^{32,33} in the PD situation and reflected back in the ND situation. The asymmetric phase profiles of the AOM in PD and ND cases are shown in Fig. 3(d), and the corresponding acoustic pressure fields in PD and ND cases are displayed in Fig. 3(e,f), in which the black, green and red arrows indicate the incident, theoretical transmitted and reflected directions respectively. Acoustic plane wave incident from PD is able to pass through and shifted into two symmetric beams





as expected. However, in ND case the incident wave is prohibited to penetrate and is subject to a total reflection. In such a case, a contrastive factor of 13.2 dB is obtained. The asymmetric splitting ability of the AOM can be applied in constructing one-way beam-splitting prisms, which may have wide applications in practice.

Finally, we consider the conversion of acoustic propagating wave into surface wave with the AOM, which is the wave bounded at the surface and evanescent on the transmitted side³⁴. The transmitted wave with the angle $\theta_t = \pi/2$ can be obtained according to Eq. (1), and the desired phase profile is illustrated in Fig. 3(g). Similarly, the PD incident wave is allowed to transmit through the ZIM and converted to the surface wave, while the ND incident wave is prohibited. The corresponding sound pressure levels (SPLs) in PD and ND cases are shown in Fig. 3(h,i), respectively. The results clearly demonstrate that the acoustic field is obviously enhanced on the surface and attenuates quickly away from the interface, revealing the nearly perfect conversion of the propagating wave into the surface mode. Physically, in the PD case, the discrete phase shifts provided by the PA along the *y* direction provide extra momenta to compensate the momentum mismatch between the propagating wave and the surface wave on the metasurface, resulting in the high efficiency conversion. The asymmetric transmission property of the AOM is verified by the 12.1 dB contrast of the sound intensity between the opposite cases. The asymmetric conversion with the AOM could act as a bridge to link the propagation wave and surface wave, which would lead to both the theoretical interests and practical applications.

Discussion

The design of an AOM has been proposed, which is capable of manipulating the wavefront of transmitted wave flexibly in only one direction but blocking the transmission in the opposite direction, and demonstrated a practical implementation by metamaterials via numerical simulations. This intriguing characteristics results from a distinct ability of asymmetric phase manipulation achieved by coupling PA with ZIM which are implemented by employing hybrid metastructures and labyrinthine units respectively. As a few examples, we have investigated the acoustic anomalous refraction, splitting wave and converting the propagating wave into surface wave to demonstrate the performances of the proposed AOM. The subwavelength thickness and planar configuration of the particular implementation would contribute to the practical applications. We need to emphasize that the proposed AOM has not broken the reciprocity principle and the time-reversal symmetry is still valid in this system, which limits the one-way transmission functionality to the normal incidence case. Due to the high selectivity of the ZIM, the oblique incident wave will be totally reflected from the surface as well. With the capability of realizing asymmetric transmission and wavefront manipulation simultaneously with high efficiency, the proposed AOM offers a framework in integrating the unidirectional propagation with versatile wave manipulation, which may find applications in a great variety of scenarios such as ultrasound imaging or therapy where special control of acoustic transmission is always highly desired.

Method

Throughout the paper, the numerical simulations are performed by using finite element method based on COMSOL Multiphysics software. The background medium is air, the mass density and sound speed of it is 1.21 kg/m³ and 343 m/s. The material of the AOM in simulation is chosen to be acrylonitrile butadiene styrene (ABS) plastic, whose mass density and sound speed are 1200 kg/m³ and 2200 m/s, respectively. Perfectly matched layers (PMLs) are utilized to eliminate the reflected waves by the outer boundaries.

References

- 1. Li, J. & Chan, C. T. Double-negative acoustic metamaterial. Phys. Rev. E 70, 055602 (2004).
- 2. Zhu, X. F. et al. Acoustic cloaking by a superlens with single-negative materials. Phys. Rev. Lett. 106, 014301 (2011).
- 3. Zhang, S., Xia, C. G. & Fang, N. Broadband acoustic cloak for ultrasound waves. Phys. Rev. Lett. 106, 024301 (2011).
- 4. Mei, J. et al. Dark acoustic metamaterials as super absorbers for low-frequency sound. Nat. Commun. 3, 756 (2012).
- 5. Li, J. et al. Experimental demonstration of an acoustic magnifying hyperlens. Nature Mater. 8, 931–934 (2009).
- 6. Zhu, J. et al. A holey-structured metamaterial for acoustic deep-subwavelength imaging. Nature Phys. 7, 52–55 (2011).
- 7. Jiang, X. et al. Ultra-broadband absorption by acoustic metamaterials. Appl. Phys. Lett. 105, 243505 (2014).
- 8. Jiang, X. et al. Broadband field rotator based on acoustic metamaterials. Appl. Phys. Lett. 104, 083510 (2014).
- 9. Jing, Y., Xu, J. & Fang, N. Numerical study of a near-zero-index acoustic metamaterial. Phys. Lett. A 376, 2834–2837 (2012).
- 10. Li, Y. et al. Acoustic focusing by coiling up space. Appl. Phys. Lett. 101, 233508 (2012).
- 11. Li, Y. et al. Reflected wavefront manipulation based on ultrathin planar acoustic metasurfaces. Sci. Rep. 3, 2546 (2013).
- 12. Shen, C., Xu, J., Fang, N. & Jing, Y. Anisotropic Complementary Acoustic Metamaterial for Canceling out Aberrating Layers. *Phys. Rev. X* 4, 041033 (2014).
- 13. Ma, G. C. et al. Acoustic metasurface with hybrid resonances. Nat. Mater. 13, 873-878 (2014).
- 14. Zhao, J. J., Li, B. W., Chen, Z. N. & Qiu, C. W. Redirection of sound waves using acoustic metasurface. Appl. Phys. Lett. 103, 151604 (2013).
- Li, Y. et al. Experimental Realization of Full Control of Reflected Waves with Subwavelength Acoustic Metasurfaces. Phys. Rev. Appl. 2, 064002 (2014).
- Li, Y., Jiang, X., Liang, B., Cheng, J. C. & Zhang, L. K. Metascreen-Based Acoustic Passive Phased Array. Phys. Rev. Appl. 4, 024003 (2015).
- 17. Zhao, J., Li, B. W., Chen, Z. N. & Qiu, C. W. Manipulating acoustic wavefront by inhomogeneous impedance and steerable extraordinary reflection. *Sci. Rep.* **3**, 2537 (2013).
- 18. Zhao, J. *et al.* Manipulation of acoustic focusing with an active and configurable planar metasurface transducer. *Sci. Rep.* **4**, 6257 (2014).
- 19. Liang, B., Yuan, B. & Cheng, J. C. Acoustic diode: rectification of acoustic energy flux in one-dimensional systems. *Phys. Rev. Lett.* **103**, 104301 (2009).
- 20. Liang, B. et al. An acoustic rectifier. Nat. Mater. 9, 989 (2010).
- 21. Popa, B. I. & Cummer, S. A. Non-reciprocal and highly nonlinear active acoustic metamaterials. Nat. Commun. 5, 3398 (2014).
- 22. He, Z. J. et al. Asymmetric acoustic gratings. Appl. Phys. Lett. 98, 083505 (2011).
- 23. Fleury, R. et al. Sound isolation and giant linear nonreciprocity in a compact acoustic circulator. Science 343, 516-9 (2014).
- 24. Li, X. F. *et al.* Tunable Unidirectional Sound Propagation through a Sonic-Crystal-Based Acoustic Diode. *Phys. Rev. Lett.* **106**, 084301 (2011).
- 25. Cicek, A., Kaya, O. A. & Ulug, B. Refraction-type sonic crystal junction diode. Appl. Phys. Lett. 100, 111905 (2012).
- 26. Hwan Oh, J. *et al.* Inverted bi-prism phononic crystals for one-sided elastic wave transmission applications. *Appl. Phys. Lett.* **100**, 213503 (2012).
- 27. Li, Y. et al. Unidirectional acoustic transmission through a prism with near-zero refractive index. Appl. Phys. Lett. 103, 053505 (2013).
- 28. Zhu, Y. F. et al. Acoustic one-way open tunnel by using metasurface. Appl. Phys. Lett. 107, 113501 (2015).
- 29. Liang, Z. X. & Li, J. Extreme Acoustic Metamaterial by Coiling Up Space. Phys. Rev. Lett. 108, 114301 (2012).
- 30. Zhao, J., Chen, Z. N., Li, B. & Qiu, C. W. Acoustic cloaking by extraordinary sound transmission. J. Appl. Phys. 117, 214507 (2015).
- Yu, N. F. et al. Light Propagation with Phase Discontinuities: Generalized Laws of Reflection and Refraction. Science 334, 333–337 (2011).
- 32. Rahm, M. *et al.* Optical design of reflectionless complex media by finite embedded coordinate transformations. *Phys. Rev. Lett.* **100**, 063903 (2008).
- Ma, Y. G., Wang, N. & Ong, C. K. Application of inverse, strict conformal transformation to design waveguide devices. J. Opt. Soc. Am. A. 27, 968–72 (2010).
- 34. Sun, S. et al. Gradient-index meta-surfaces as a bridge linking propagating waves and surface waves. Nat. Mater. 11, 426–431 (2012).

Acknowledgements

This work was supported by the National Basic Research Program of China (973 Program) (Grant Nos. 2011CB707900 and 2012CB921504), National Natural Science Foundation of China (Grant Nos 11174138,

11174139, 11222442, 81127901, and 11274168), NCET-12-0254, and a project funded by the Priority Academic Program Development of Jiangsu Higher Education Institutions.

Author Contributions

X.J., B.L. and X.Y.Z. performed the analytical and numerical simulations. X.J. and B.L. conceived the idea and wrote the manuscript. B.L. and J.C.C. conceived and supervised the manuscript. J.Y., L.L.Y. and J.Y. contributed in the analysis. All authors contributed to the discussions.

Additional Information

Competing financial interests: The authors declare no competing financial interests.

How to cite this article: Jiang, X. *et al.* Acoustic one-way metasurfaces: Asymmetric Phase Modulation of Sound by Subwavelength Layer. *Sci. Rep.* **6**, 28023; doi: 10.1038/srep28023 (2016).

This work is licensed under a Creative Commons Attribution 4.0 International License. The images or other third party material in this article are included in the article's Creative Commons license, unless indicated otherwise in the credit line; if the material is not included under the Creative Commons license, users will need to obtain permission from the license holder to reproduce the material. To view a copy of this license, visit http://creativecommons.org/licenses/by/4.0/