

# Dynamics and mechanism of ultrafast water–protein interactions

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**Protein hydration is essential to its structure, dynamics, and function, but water–protein interactions have not been directly observed in real time at physiological temperature to our awareness. By using a tryptophan scan with femtosecond spectroscopy, we simultaneously measured the hydration water dynamics and protein side-chain motions with temperature dependence. We observed the heterogeneous hydration dynamics around the global protein surface with two types of coupled motions, collective water/side-chain reorientation in a few picoseconds and cooperative water/side-chain restructuring in tens of picoseconds. The ultrafast dynamics in hundreds of femtoseconds is from the outer-layer, bulk-type mobile water molecules in the hydration shell. We also found that the hydration water dynamics are always faster than protein side-chain relaxations but with the same energy barriers, indicating hydration shell fluctuations driving protein side-chain motions on the picosecond time scales and thus elucidating their ultimate relationship.**

hydration shell dynamics | protein side-chain motion | water-driven relaxation | coupled fluctuation | tryptophan scan

**W**ater–protein interactions are critical to protein structural stability and flexibility, functional dynamics, and biological activities (1, 2). Various methods such as neutron scattering (3), NMR (4), laser spectroscopy (5, 6), and molecular dynamics (MD) simulations (7) have been used to reveal protein surface hydration and coupled water–protein dynamics on different time and length scales. Hydration water molecules often participate in various protein functions and their motions even directly “control” protein fluctuations (2, 8). Frauenfelder et al. recently proposed a unified model for protein dynamics (8): large-scale protein motions are slaved to the fluctuations of bulk solvent and controlled by solvent viscosity while internal protein motions are slaved to the fluctuations of the hydration shell and controlled by hydration water. However, direct measurements of such coupled fluctuations at physiological temperature are challenging as a result of the ultrafast nature of water motions, and therefore most studies are indirect or at low temperature (3, 4). Here, we used a tryptophan (W) scan to probe global surface hydration (9) and used femtosecond spectroscopy to follow hydration water motions and local side-chain fluctuations in real time. With temperature dependence, we systematically measured their dynamics and thus finally elucidate their ultimate relationship.

## Results and Discussion

**Tryptophan Scan and Femtosecond Fluorescence Spectroscopy.** We used DNA polymerase IV (Dpo4) (10), a hyperthermal enzyme without a single tryptophan residue, as a model protein, and designed 10 tryptophan mutants, one at a time, to probe four different domains (Fig. 1A). We performed systematic measurements of tryptophan fluorescence intensity changes with wavelength and time and thereby constructed 3D fluorescence profiles (SI Appendix, SI Note 1). One example of an R176W mutant is shown in Fig. 1B. Based on the methodology we developed to follow the Stokes shifts of tryptophan with time (11, 12), we derived the solvation correlation functions for the 10 mutants. Fig. 1C shows our typical results

for three mutants with three exponential decays ( $\tau_{1S}$ ,  $\tau_{2S}$ , and  $\tau_{3S}$ , Y12W and R176W) or two exponential decays ( $\tau_{2S}$  and  $\tau_{3S}$ , Y108W) from hundreds of femtoseconds to tens of picoseconds. Simultaneously, we measured tryptophan side-chain relaxation through its fluorescence anisotropy dynamics (Fig. 1D, *Inset*), and the typical result is shown in Fig. 1D for R176W with four exponential decays. The initial component ( $\tau_{1C}$ ) is from ultrafast internal conversion of two concurrently excited electronic states ( $^1L_a$  and  $^1L_b$ ) in  $\sim 70$  fs (13), and the last decay ( $\tau_T$ ) represents the protein tumbling motion in  $\sim 20$  ns. The two middle decay components ( $\tau_{2W}$  and  $\tau_{3W}$ ) result from the local wobbling relaxations on picosecond time scales. For nine mutants, we measured protein solvation dynamics and local tryptophan relaxations at seven different temperatures from 1 °C to 60 °C (SI Appendix, Figs. S1–S4).

**Protein Surface Hydration Dynamics.** Fig. 2 shows all obtained solvation correlation functions in terms of relaxation energies (Fig. 2B) and times (Fig. 2C–E) for the 10 mutants at room temperature. Similar to the patterns that observed before in numerous proteins (9, 14–17), the relaxation energies are dominantly from hydration water molecules and the solvation dynamics are highly heterogeneous around the global protein surface (17). In Fig. 2B, the 10 mutants can be divided into two categories, exposed and buried, as usual with the dividing steady-state emission peak ( $\lambda_{\text{peak}}$ ) of tryptophan at  $\sim 338$  nm. The total solvation energy ( $E_{\text{tot}}$ ) shows a sigmoid-type increase along  $\lambda_{\text{peak}}$  when the probe moves to the protein surface with more exposure to hydration water molecules. For the exposed probes, we observed the solvation dynamics on three distinct time scales of  $\tau_{1S}$  in 300–400 fs (Fig. 2C),  $\tau_{2S}$  in 3–6 ps (Fig. 2D), and  $\tau_{3S}$  in 50–120 ps (Fig. 2E), corresponding to three relaxation energies of  $E_1$ ,  $E_2$ , and  $E_3$ , respectively. For the buried probes, we did not detect first ultrafast solvation dynamics ( $\tau_{1S}$ )

## Significance

Hydration water around a protein is fundamental to protein's property and function. How hydration water interacts with protein at their interface has not been fully understood. A universal slaving model that water drives protein fluctuations has been proposed, but direct experimental observation is difficult and challenging. Here, we directly measured hydration water dynamics and protein side-chain relaxations with temperature dependence. With extensive data, we conclude that the surface hydration-shell fluctuations drive protein side-chain motions on the picosecond time scales and thus play a critical role in protein dynamics.

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The authors declare no conflict of interest.

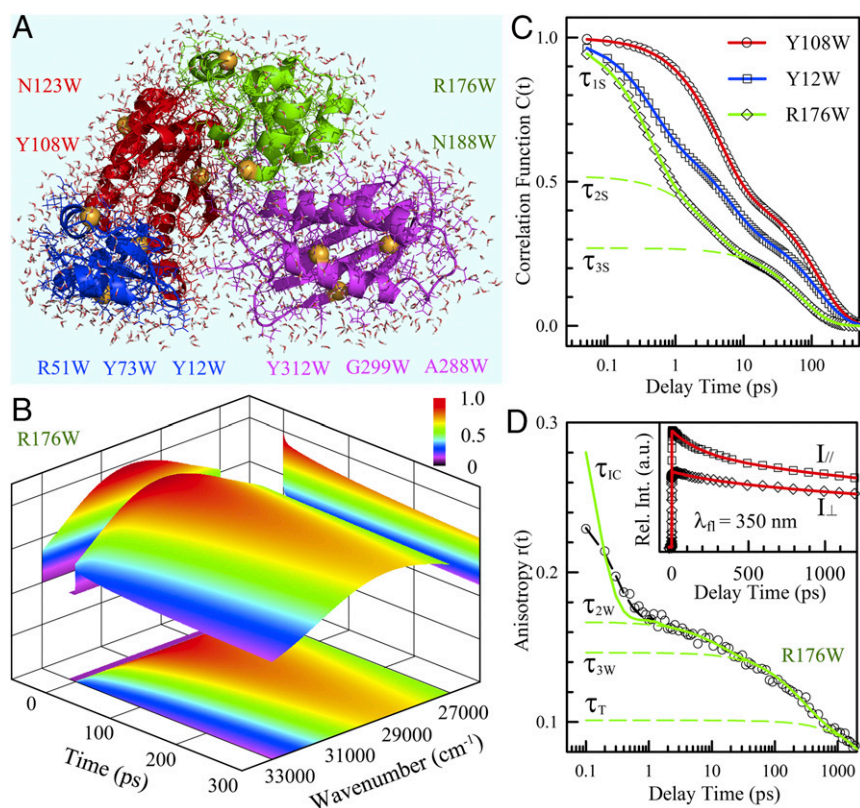
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**Fig. 1.** Hydration dynamics and protein side-chain motions detected by tryptophan. (A) A snapshot of MD simulations of solvated Dpo4 (Protein Data Bank ID 2RDI) shown with water molecules within 5 Å to the protein surface. Dpo4 consists of four domains, including thumb (green), palm (red), finger (blue), and little finger (magenta). The yellow spheres indicate the mutation sites replaced by tryptophan. (B) Constructed 3D FRES of tryptophan (R176W) at room temperature. (C) Typical solvation correlation functions of three mutants show the very different dynamics when the probe moves from the buried (Y108W) to exposed (R176W) locations. The symbols are the derived experimental data, and the solid lines are exponential fit. The dashed lines represent three exponential decay components ( $\tau_{1S}$ ,  $\tau_{2S}$ , and  $\tau_{3S}$ ) for the mutant R176W. (D) The anisotropy decay dynamics show four decay components corresponding to internal conversion ( $\tau_{IC}$ ), wobbling ( $\tau_{2W}$ ,  $\tau_{3W}$ ), and tumbling ( $\tau_T$ ) motions. The symbols are experimental data and the dashed black line is the best fit with the internal conversion model (13), and the solid line is after deconvolution from the instrument response. (Inset) Two parallel and perpendicular fluorescence transients used to calculate the anisotropy decay.

and observed only two solvation dynamics on picosecond scales ( $\tau_{2S}$  and  $\tau_{3S}$ ).

Considering the dipole-dipole interactions between tryptophan and neighboring water molecules mainly within  $\sim 10$  Å from our MD simulations (*SI Appendix*, Fig. S5), we found that for the buried probes the 10-Å water molecules are dominantly within 7 Å to the protein surface (Y108W and Y73W in Fig. 2A). Thus, the buried probes dominantly detected one to two inner layers of water molecules in the hydration shell, and the observed two relaxations are mainly from these inner-layer water molecules at the interface. For the exposed probes, besides the inner-layer hydration water response, we also observed  $\sim 30\%$  of 10-Å water molecules beyond 7 Å to the protein surface. Thus, the observed ultrafast bulk-type water relaxation (Fig. 2C) must be from those water molecules in the outer layers of the hydration shell (G299W in Fig. 2A). Consistent with the observations in apomyoglobin (16) and other proteins (18, 19), as well as theoretical work and recent MD simulations (20–24), the water motion in a few picoseconds (Fig. 2D) represents the collective water-network reorientational relaxation at the water-protein interface and significantly slows down, compared with the motions of bulk-type water in the outer layers, as a result of structured water-network collectivity near the protein. The water relaxation in tens of picoseconds results from the cooperative water-protein rearrangements, i.e., the coupled interfacial water-protein restructuring, and directly reflects such cooperativity of coupled water-protein fluctuations (Fig. 2E). Also, similar to our previous results observed in apomyoglobin (16), these water

motions are correlated with the local protein's chemical identities, topological roughness, higher-order structures, and tertiary intra-protein interactions, revealing various water behaviors in the inner hydration layers with wide heterogeneity (Fig. 2D, a–d, and Fig. 2E, e–g). For example, the hydration water near the densely charged protein surface (G299W) relaxes slowly, and the water molecules trapped at the active site and confined in a nanospace (Y12W) show the longest relaxation (Fig. 2D and E and *SI Appendix*, Figs. S5 and S6A).

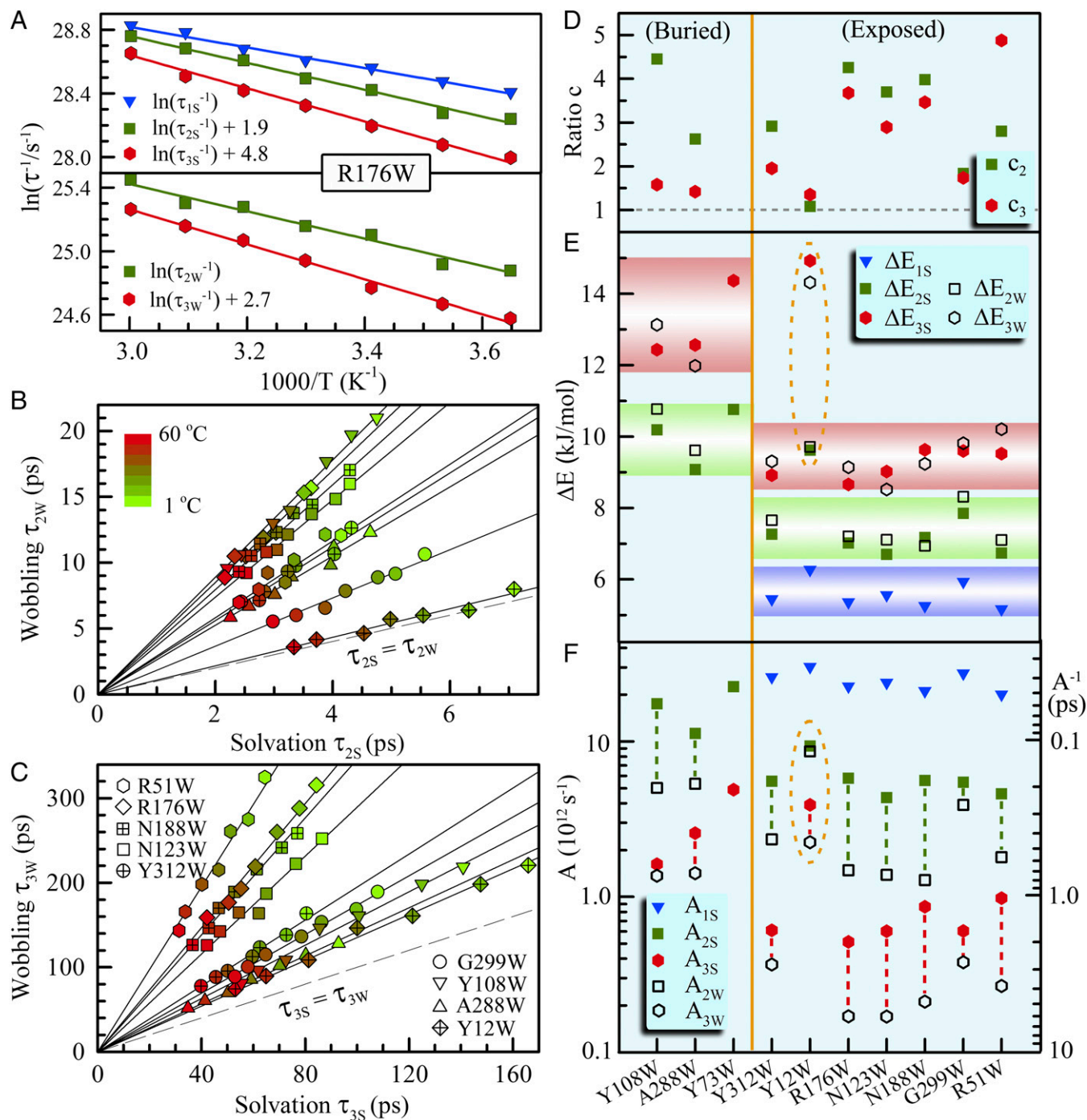
**Origin of Coupled Water-Protein Fluctuations.** The observed interfacial water motions must interact with the local protein. The observed two relaxations of tryptophan (Fig. 1D) on the picosecond scales should couple with two hydration water motions. Thus, we measured these dynamics with temperature changes from 1 °C to 60 °C for the nine mutants (*SI Appendix*, Figs. S7–S9), and Fig. 3A shows a typical result for the mutant R176W. We plotted the correlations of  $\tau_{2W}$  with  $\tau_{2S}$  in Fig. 3B and of  $\tau_{3W}$  with  $\tau_{3S}$  in Fig. 3C. Strikingly, the two sets of data are all linearly correlated through the original point (0,0). Such 18 linear correlations for all of the 9 mutants strongly show that both relaxations of hydration water and tryptophan side-chain are intrinsically connected and must be from one origin. Furthermore, all correlated straight lines lie above the equal line of  $\tau_{iW} = \tau_{iS}$  ( $i = 2,3$ ), i.e., all the tryptophan relaxations are always slower than the neighboring hydration water motions (Fig. 3B and 3C). This observation clearly suggests that such coupled motions are driven by hydration water relaxations, confirming the slaving model





activation barriers ( $\Delta E_{2S} \sim \Delta E_{2W}$  and  $\Delta E_{3S} \sim \Delta E_{3W}$ ) within the average SD of 0.4–0.6 kJ/mol and differ only in their prefactors, showing the two relaxations of hydration water and side chain from the same origin. (iii) For the exposed positions, the three barriers of hydration water relaxations are within  $5.6 \pm 0.3$ ,  $7.1 \pm 0.4$ , and  $9.2 \pm 0.5$  kJ/mol for  $\Delta E_{1S}$ ,  $\Delta E_{2S}$ , and  $\Delta E_{3S}$  with the corresponding prefactors in time ( $A^{-1}$ ) of 40 fs,

200 fs, and 1.5 ps, respectively. For the buried positions, the latter two barriers increase to  $9.9 \pm 0.6$  kJ/mol and  $13.6 \pm 0.6$  kJ/mol, also with greater corresponding prefactors. The barrier for the ultrafast motion of the outer-layer water is similar to that of free tryptophan in bulk water (*SI Appendix, Fig. S11*), indicating a similar relaxation (22, 26). However, the barriers for the inner-layer coupled motions



**Fig. 3.** Temperature dependence of hydration water and protein side-chain dynamics. (A) The dynamics of hydration water and tryptophan side chain change with temperature as shown for R176W, following an Arrhenius relation in the experimental temperature range of 1–60 °C (1, 10, 20, 30, 40, 50, and 60 °C). (B and C) Shown are the linear correlations between water solvation and tryptophan relaxation times for the observed two dynamics. The solid lines are linear fits for each mutant. The dashed lines are equal relations between two coupled water/side-chain relaxations. Note that the hydration dynamics are always faster than the tryptophan relaxations. (D) Shown are the ratios ( $c_i$ ;  $i = 2, 3$ ) of two coupled relaxations, which are also the slopes of the straight lines in B and C. (E and F) Show are the derived activation energies and prefactors by the Arrhenius relation. Hydration water and tryptophan side chain have the same barriers, indicating that the coupled relaxations are from one origin. Note that the exposed probes give the three same barriers.



## Materials and Methods

**Protein Preparation.** *Sulfolobus solfataricus* Dpo4 gene was cloned into the NdeI and XhoI sites of pET22b, and a His<sub>6</sub> tag was attached to the C terminal. The mutant plasmids were prepared by site-directed mutagenesis. Each plasmid was individually transformed into *Escherichia coli* BL21(DE3). All mutants were expressed and purified following procedures reported previously (34). The purified proteins were stored in a storage buffer solution containing 50 mM Tris-HCl (pH 7.5), 200 mM NaCl, 5 mM MgCl<sub>2</sub>, 1 mM DTT, 0.5 mM EDTA, and 50% (vol/vol) glycerol. For fluorescence spectroscopy experiments, the sample was diluted and concentrated again to remove the glycerol. The final concentration of the protein sample was ~1 mM, and the final buffer condition was 50 mM Tris-HCl (pH 7.5), 400 mM NaCl, 5 mM MgCl<sub>2</sub>, 1 mM DTT, 0.5 mM EDTA, and less than 5% (vol/vol) glycerol. L-tryptophan was purchased from Sigma-Aldrich and used as received. The tryptophan was dissolved in the same buffer as the protein, and the final concentration was 3 mM.

**Femtosecond Fluorescence Emission Spectroscopy.** All femtosecond-resolved fluorescence measurements were carried out by using the up-conversion method as described before (35). The pump wavelength was centered at 295 nm generated by a series of nonlinear mixing and doubling in 0.2-mm-thick barium borate (BBO) crystals (type I) with a repetition rate of 1 kHz. The pump pulse energy was attenuated to ~100 nJ before being focused into the motor-controlled rotating sample cell. Use of 295-nm wavelength to excite the sample minimized tyrosine absorption in Dpo4, and the observed fluorescence was es-

entially from the excited tryptophan emission. The fluorescence emission was collected by a pair of parabolic mirrors and then mixed with a gating pulse (800 nm) in a 0.2-mm-thick BBO crystal through a nonlinear configuration. The up-converted signal ranging from 221 to 253 nm was detected by a photomultiplier coupled with a double-grating monochromator. A single fluorescence transient at the fixed wavelength or a femtosecond-resolved emission spectrum (FRES) at the fixed delay time will be obtained. The instrument response time under the current noncollinear geometry is approximately 400 fs as determined from the water Raman signal around 327 nm. The pump-pulse polarization was set at the "magic angle" (54.7°) with respect to the acceptance axis (vertical) of the up-conversion crystal, and the gating-pulse polarization was set parallel to this axis. For all fluorescence anisotropy measurements, the pump-pulse polarization was rotated parallel or perpendicular to the acceptance axis to obtain the parallel ( $I_{\parallel}$ ) and perpendicular ( $I_{\perp}$ ) signals, respectively. The sample cell was mounted in the center of an aluminum block, and its temperature was controlled by a scientific solution bath (VWR). The lifetime of each mutant and tryptophan was determined by time-correlated single-photon counting.

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