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## The first CO<sup>+</sup> image: I. Probing the H<sup>+</sup>/H<sub>2</sub> layer around the ultracompact H<sub>II</sub> region Mon R2

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### Abstract

The CO<sup>+</sup> reactive ion is thought to be a tracer of the boundary between a HII region and the hot molecular gas. In this study, we present the spatial distribution of the CO<sup>+</sup> rotational emission toward the Mon R2 star-forming region. The CO<sup>+</sup> emission presents a clumpy ring-like morphology, arising from a narrow dense layer around the HII region. We compare the CO<sup>+</sup> distribution with other species present in photon-dominated regions (PDR), such as [CII] 158 μm, H<sub>2</sub> S(3) rotational line at 9.3 μm, polycyclic aromatic hydrocarbons (PAHs) and HCO<sup>+</sup>. We find that the CO<sup>+</sup> emission is spatially coincident with the PAHs and [CII] emission. This confirms that the CO<sup>+</sup> emission arises from a narrow dense layer of the H<sup>+</sup>/H<sub>2</sub> interface. We have determined the CO<sup>+</sup> fractional abundance, relative to C<sup>+</sup> toward three positions. The abundances range from 0.1 to 1.9 × 10<sup>-10</sup> and are in good agreement with previous chemical model, which predicts that the production of CO<sup>+</sup> in PDRs only occurs in dense regions with high UV fields. The CO<sup>+</sup> linewidth is larger than those found in molecular gas tracers, and their central velocity are blue-shifted with respect to the molecular gas velocity. We interpret this as a hint that the CO<sup>+</sup> is probing photo-evaporating clump surfaces.

## Keywords

Astrochemistry – (ISM); Photo-dominated regions (PDRs) – HII region – Radio lines; ISM – Individual; Monoceros R2

## 1 Introduction

Reactive ions are destroyed in almost every collision with H and H<sub>2</sub> and recombine rapidly with e<sup>-</sup>. These compounds present enhanced abundances toward the hot layers of the photon-dominated regions (PDRs), where the far ultraviolet (FUV) field is only partially attenuated and maintains high abundances of the parent species C<sup>+</sup> and S<sup>+</sup> (Sternberg & Dalgarno 1995). In particular, Sternberg & Dalgarno (1995) predict a high CO<sup>+</sup> abundance at the HI/H<sub>2</sub> interface ( $A_V \approx 1$  mag) of dense PDRs, where it is mainly produced by the C<sup>+</sup> + OH → CO<sup>+</sup> + H reaction. So far, the CO<sup>+</sup> ion has been detected in several PDRs; such as the M17SW, Orion Bar, NGC7027, NGC7023 (Latter et al. 1993; Stoerzer et al. 1995; Fuente & Martín-Pintado 1997; Fuente et al. 2003), G29.96, MonR2 (Rizzo et al. 2003) and S140 (Savage & Ziurys 2004). However, all the detections have been obtained after long integrations toward a single position and they lack the information on the spatial distribution.

The Mon R2 star-forming region, located at 830 pc (Herbst & Racine 1976), contains an ultracompact (UC) HII region surrounded by a series of PDRs with different physical conditions (Pilleri et al. 2013; Treviño-Morales et al. 2014). The main PDR, corresponding to IRS 1 (hereafter IF), is irradiated by a high UV field of  $G_0 > 10^5$  (in units of the Habing flux; Habing 1968), and presents high densities ( $> 10^5$  cm<sup>-3</sup>) and kinetic temperatures ( $T_k \approx 600$  K; Berné et al. 2009). A second PDR, associated with the molecular peak MP2, is detected 40'' north from IF, and shows chemical properties similar to those found in low- to mid-UV irradiated PDRs (Ginard et al. 2012). Due to its proximity and physical conditions, Mon R2 turns to be an excellent candidate to study the HI/H<sub>2</sub> interface. CO<sup>+</sup> is thought to be a good PDR tracer, and its distribution is potentially an excellent diagnostic tool to learn about the physical structure of these regions.

In this paper, we present a study of the CO<sup>+</sup> ( $J=2-1$ ) transition line toward Mon R2 and compare its spatial distribution with *Spitzer* data reported by Berné et al. (2009), *Herschel* data from Pilleri et al. (2012) and Ossenkopf et al. (2013) and the HCO<sup>+</sup> and H<sup>13</sup>CO<sup>+</sup> molecules from Treviño-Morales et al. (2014).

## 2 Observations and data reduction

We observed 2' × 2' maps of CO<sup>+</sup> transition lines ( $J=2-1$  at 235.380 GHz, 235.789 GHz and 236.062 GHz) using the IRAM-30m telescope (Pico Veleta, Spain). The observations were performed using the EMIR receiver with the Fast Fourier Transform Spectrometer (FTS) at 200 kHz of resolution. Throughout this paper, we use the main-beam brightness temperature ( $T_{MB}$ ) as intensity scale. The data were reduced using standard procedures with the CLASS/GILDAS package (Pety et al. 2005). The three lines were detected, but only the 236.0625 GHz has sufficient signal-to-noise ratio (S/N) for good imaging. In order to improve the S/N, we smoothed the native observation to an angular resolution of 16'' (see

Fig. 1) and to a spectral resolution of  $1 \text{ km s}^{-1}$ . In the final datacube,  $\text{CO}^+$  has linewidths of  $6\text{--}8 \text{ km s}^{-1}$  and intensity peaks of  $60\text{--}200 \text{ mK}$  (with  $\text{rms} \sim 20 \text{ mK}$ ). The main  $\text{CO}^+$  line is located close to a bright  $^{13}\text{CH}_3\text{OH}$  line (at  $236.0628 \text{ GHz}$ ). However, we dismiss the idea of a possible blending in the  $\text{CO}^+$  line, as in the spectral line survey conducted toward Mon R2 (Treviño-Morales 2016) we do not find  $^{13}\text{CH}_3\text{OH}$  emission at any frequency. Moreover, the main compound  $\text{CH}_3\text{OH}$  is not detected at the positions where  $\text{CO}^+$  is bright.

### 3 Results

In Fig. 1 we compare the  $\text{CO}^+$  spatial distribution with other species. *Panel A* shows the  $\text{H}^{13}\text{CO}^+$  (3–2) line emission (black contours) tracing the molecular gas (Treviño-Morales et al. 2014), that is distributed around the  $\text{CO}^+$  emission.  $\text{CO}^+$  presents a clumpy structure, where the main  $\text{CO}^+$  clumps seem to have a counterpart in the  $\text{H}^{13}\text{CO}^+$  (3–2) emission. However, the peaks of  $\text{CO}^+$  are located  $\approx 5\text{--}10''$  closer to the  $\text{HII}$  region, likely tracing an inner layer of the region. We find that  $\text{CO}^+$  emission appears surrounding the  $\text{HII}$  region with its intensity peak at the offset  $[0'', -7'']$ , very close the IF position (see *Panel B* of Fig. 1). Moreover, the two most intense  $\text{CO}^+$  clumps are correlated with the  $\text{H}_2$  emission, in an area where the density is presumably larger. *Panels C* and *D* show the PAHs and the  $[\text{CII}]$  emission, respectively (black contours). The  $\text{CO}^+$  emission present a clumpy ring-like distribution, spatially coincident with the PAHs emission. The  $\text{CO}^+$  secondary clumps are associated with the PAHs emission peaks, but not the most intense one. *Panel D* shows a comparison of the  $\text{CO}^+$  spatial distribution with the emission of its chemical precursor  $[\text{CII}]$ . These species are spatially associated, with the main difference being the location of the peaks:  $\text{CO}^+$  has its intensity peak to the south of the IF position, while the  $[\text{CII}]$  peak is located to the west. Therefore, we interpret that the  $\text{CO}^+$  is found toward the densest area of the region (where all the molecular gas piles up, e. g., the  $\text{H}^{13}\text{CO}^+$  spatial distribution) as expected since the critical densities of  $\text{H}^{13}\text{CO}^+$  ( $n_{\text{cr}} \sim 10^5 \text{ cm}^{-3}$ ) and  $\text{CO}^+$  ( $n_{\text{cr}} \sim \text{a few } 10^5 \text{ cm}^{-3}$ ; Stäuber & Bruderer 2009) are larger than that of  $[\text{CII}]$  ( $n_{\text{cr}} \sim \text{a few } 10^3 \text{ cm}^{-3}$ ; Goldsmith et al. 2012). It is worth noting that the  $\text{CH}^+$  molecule is also related to the  $\text{CO}^+$  and  $[\text{CII}]$  chemistry. When comparing their spatial distribution, we find that the  $\text{CO}^+$  emission also coexist with  $\text{CH}^+$  but, as  $[\text{CII}]$ , its intensity peak is located to the west of the IF position (Pilleri et al. 2014).

In order to better understand the spatial distribution, we did intensity cuts with a position angle of  $45^\circ$  throughout the IF position and cutting the ring-like structure seen in the  $\text{CO}^+$  emission at the south-east and north-west (pink dashed line in left panel of Fig. 2). The intensity cuts for the species  $\text{CO}^+$ ,  $\text{H}^{13}\text{CO}^+$ ,  $[\text{CII}]$ ,  $\text{H}_2$  and PAHs are shown in *Panels A* to *F* of Fig. 2. The molecular gas as traced by  $\text{H}^{13}\text{CO}^+$  shows emission between the offset  $-20''$  and  $0''$  (corresponding to the area between  $[20'', -20'']$  and  $[0'', 0'']$  in the map) with no emission associated with the  $\text{HII}$  region. The most intense emission of  $\text{CO}^+$  comes from this region, but there is also  $3\sigma$ -level emission associated with the  $\text{HII}$  region (gray-shaded area in the map). Note, however, that the  $\text{CO}^+$  peak is closer to the  $\text{HII}$  region than the  $\text{H}^{13}\text{CO}^+$  peak. The  $[\text{CII}]$  emission is very intense and extended in the whole area, and its intensity cut shows two emission peaks: one coincident with the  $\text{CO}^+$  peak, and a second one (the brightest one) on the opposite edge of the ring-like structure (to the north-west). The PAHs emission is weaker than  $[\text{CII}]$  but their intensity cuts are quite similar. Finally, the  $\text{H}_2 \text{ S}(3)$

line only shows emission over  $3\sigma$  between the offset  $-10''$  and  $0''$ . Summarizing we can see a trend of spatial segregation, with the  $\text{H}^{13}\text{CO}^+$  tracing the outer layers (far from the  $\text{H}\text{II}$  region), then  $\text{CO}^+$  and  $[\text{CII}]$  and PAHs peaking closer to the ionized gas, and finally the  $\text{H}_2$  emission tracing a hotter layer close to the  $\text{H}\text{II}$  region (*Panel F* of Fig. 2).

The circles in the left panel of Fig. 2 mark the positions where we compare the spectral profiles of  $\text{CO}^+$ ,  $[\text{CII}]$ ,  $\text{H}^{13}\text{CO}^+$  and  $\text{HCO}^+$  (*Panel G*). The positions correspond to (i) the ionization front (IF at offset  $[0'', 0'']$ , red), (ii) the offset  $[-10'', 24'']$  (close to the PDR3 in Berné et al. 2009, black), and (iii) the MP2 position (offset  $[0'', 40'']$ , blue) that is known to be a PDR with a low UV field and high density (Ginard et al. 2012). In Fig. A.1, we present a comparison of the  $\text{CO}^+$  line profile (in red) with the  $[\text{CII}]$  (Ossenkopf et al. 2013),  $^{13}\text{CO}$  (10–9) and  $\text{CO}$  (9–8) lines at IF (Pilleri et al. 2012). We find that  $\text{CO}^+$  presents larger linewidth than  $^{13}\text{CO}$ ,  $\text{H}^{13}\text{CO}^+$  and  $\text{HCO}^+$  ( $\approx 7 \pm 0.7 \text{ km s}^{-1}$  vs  $3\text{--}5 \text{ km s}^{-1}$ ). The  $\text{CO}^+$  linewidth, however, is comparable to that of carbon recombination lines ( $\approx 6 \text{ km s}^{-1}$ ; Treviño-Morales 2016). Similarly, the  $\text{CO}^+$  line has a similar profile to that of the  $[\text{CII}]$  and the high excitation  $\text{CO}$  (9–8) line. The  $\text{CO}^+$  and  $[\text{CII}]$  spectra are closely similar in the line shape and line widths; this similarity may be well explained if  $\text{CO}^+$  lies in the dissociation front, between the ionized gas and the molecular gas. We find weak  $\text{CO}^+$  emission toward the  $\text{H}\text{II}$  region that might come from the back and front walls of the PDR.

On the basis of reactive ions ( $\text{CH}^+$ ,  $\text{OH}^+$ , and  $\text{H}_2\text{O}^+$ ) observations, Pilleri et al. (2014) constructed a schematic view of the Mon R2 geometry. They find that the emission of the high density molecular gas seems to come from the back side (relative to the observer) of the  $\text{H}\text{II}$  region. This is confirmed by the detection of the cluster at infrared wavelengths (i. e., not obscured by in-front molecular gas). In this scenario, one expects that an expansion of the  $\text{H}\text{II}$  region is indicated by an excess of redshifted emission in the observed lines. However, while we find hints of the molecular tracers ( $\text{H}^{13}\text{CO}^+$ ,  $\text{CO}$ ) to be skewed toward red-shifted velocities, the PDR tracers ( $\text{CO}^+$ ,  $[\text{CII}]$ ) seem to have a spectral profile skewed toward blue-shifted velocities (Figs. 2 and A.1). In order to quantify this, we have measured the amount of emission in each spectral line that is blue-shifted and red-shifted with respect to the systemic velocity ( $10 \text{ km s}^{-1}$  for IF and PDR3, and  $8 \text{ km s}^{-1}$  for MP2; Treviño-Morales et al. 2014). For the IF position, we find  $55 \pm 11\%$  of the emission of PDR tracers to be blue-shifted, and only  $33 \pm 6\%$  for the molecular tracers. Similarly, for MP2 and PDR3 we find  $44 \pm 9\%$  and  $55 \pm 11\%$  of blue-shifted emission for the PDR tracers and  $18 \pm 4\%$  and  $39 \pm 8\%$  for the molecular tracers. In general, we have  $51 \pm 10\%$  of the PDR-tracer emission to be blue-shifted, and only  $30 \pm 6\%$  for the molecular tracers. Considering that the molecular gas is located behind the UC  $\text{H}\text{II}$  region, the difference between the PDR and molecular tracers can be explained if the PDR is formed by dense condensations that are being photo-evaporated. In this case, the photo-evaporated gas (PDR tracers) would be ejected toward us, and therefore, blue-shifted with respect to the molecular gas. This effect is also visible in Fig. A.1, where  $[\text{CII}]$  presents its velocity peak at  $\sim 9 \text{ km s}^{-1}$  (blue-shifted) and the emission associated with the molecular gas is red-shifted. Finally, we do not find significant velocity gradients ( $< 1 \text{ km s}^{-1}$ ) between the  $\text{CO}^+$  clumps, suggesting low levels of turbulence between them, and consistent with the low expansion velocity found by Fuente et al. (2010) and Pilleri et al. (2012).

## 4 CO<sup>+</sup> and HCO<sup>+</sup> fractional abundances

On the basis of the spectra presented in Figs. 2 and A.1, we calculated the CO<sup>+</sup>, HCO<sup>+</sup> and C<sup>+</sup> column density ( $N$ ) in ranges of 1 km s<sup>-1</sup> at the three selected positions (see Appendix B). We assume optically thin emission and a beam filling factor of 1 for all the species. The  $M[\text{CO}^+]$  values have been calculated assuming a Boltzmann distribution of rotational levels with  $T_{\text{ex}} = 18$  K. This value of the excitation temperature is based on calculations by Stäuber & Bruderer (2009) and the detection of several CO<sup>+</sup> lines towards the Orion Bar (Cuadrado et al. 2015, S. Cuadrado, priv. communication). We used the MADEX large velocity gradient (LVG; Cernicharo 2012) and RADEX code (van der Tak et al. 2007), to derive  $M[\text{H}^{13}\text{CO}^+]$  and  $M[\text{C}^+]$ . According to the physical conditions derived by Berné et al. (2009) from the H<sub>2</sub> ground state rotational lines, we assume  $T_{\text{k}} = 600$  K,  $n_{\text{H}} = 4 \times 10^5$  cm<sup>-3</sup> for the IF and  $T_{\text{k}} = 300$  K,  $n_{\text{H}} = 4 \times 10^4$  cm<sup>-3</sup> for PDR3. Unfortunately, the MP2 position is out of the *Spitzer* map; for this position, we make a reasonable guess of  $T_{\text{k}} = 300$  K and  $n_{\text{H}} = 2 \times 10^5$  cm<sup>-3</sup>. The assumption of optically thin emission is not valid for the HCO<sup>+</sup> (3–2) line at velocities close to cloud systemic velocity. Thus, between 8–11 km s<sup>-1</sup>, we derived the  $M[\text{HCO}^+]$  using the rarer isotopologue line H<sup>13</sup>CO<sup>+</sup> (3–2) and assuming <sup>12</sup>C/<sup>13</sup>C = 50 (Treviño-Morales et al. 2014). The derived H<sup>13</sup>CO<sup>+</sup> excitation temperatures (~10–16 K) toward the IF and MP2 are in agreement with those measured by Treviño-Morales et al. (2014). The collisional rate coefficients for HCO<sup>+</sup> and H<sup>13</sup>CO<sup>+</sup> are taken from Flower (1999). To derive the  $M[\text{C}^+]$  at the IF position, between 6–12 km s<sup>-1</sup>, we used the rarer isotopologue line [<sup>13</sup>CII]. A significant fraction of the [CII] emission is expected to come from the atomic layer. The calculated C<sup>+</sup>–H collisional rates are similar to those with H<sub>2</sub> within a factor 1.3 (Wiesenfeld & Goldsmith 2014, Barinovs et al. 2005). Hence, the relevant parameter regarding collisional excitation is the number of particles, either H or H<sub>2</sub>. In our calculations we assume that the hydrogen is in molecular form which implies an uncertainty of a factor of 2 in the assumed density. Taking into account that we are well over the critical density of the [CII] 158 μm line, this translates into an uncertainty of <30% in the  $M[\text{C}^+]$ . The obtained  $M[\text{C}^+]$  are in good agreement with those obtained by Ossenkopf et al. (2013).

We have computed the  $M[\text{CO}^+]/M[\text{HCO}^+]$  ratio. We find values between 0.01–0.1 toward IF, with the highest values in the velocity wings. Note that we have not detected CO<sup>+</sup> toward the PDR with lower UV field (MP2) with a significant upper limit of  $M[\text{CO}^+]/M[\text{HCO}^+] < 0.008$ . Toward PDR3, we obtain values of  $M[\text{CO}^+]/M[\text{HCO}^+] \sim 0.006$ , i. e., a factor of  $\approx 2$  lower than IF. C<sup>+</sup> is known to be a good probe of the skin ( $A_{\text{v}} < 4$  mag) of PDRs. Because of the similar spatial distribution and velocity profiles between the CO<sup>+</sup> and C<sup>+</sup>, we used C<sup>+</sup> to estimate the absolute fractional abundance of CO<sup>+</sup>. We can safely assume that almost all the carbon is in C<sup>+</sup> in the region from which CO<sup>+</sup> is coming. Assuming a carbon elemental abundance of 10<sup>-4</sup> (Ossenkopf et al. 2013, Wakelam & Herbst 2008), we derive  $X[\text{CO}^+]$  between a few 10<sup>-11</sup> to  $\sim 1.9 \times 10^{-10}$  toward both, IF and PDR3. It is worth noting that the beam of CO<sup>+</sup>, HCO<sup>+</sup>, H<sup>13</sup>CO<sup>+</sup>, [<sup>13</sup>CII] and [CII] are quite similar, thus the calculated  $M[\text{CO}^+]/M[\text{HCO}^+]$  and  $X[\text{CO}^+]$  values are not affected by beam filling factors. Our results are in good agreement with the model predictions presented by Sternberg & Dalgarno (1995), for  $G_0 \sim 5 \times 10^5$  and  $n_{\text{H}} \sim 10^6$  cm<sup>-3</sup>. More recently, the models of Spaans & Meijerink (2007) predict



the  $X[\text{CO}^+]$  in PDRs, for  $n_{\text{H}} = 10^5 \text{ cm}^{-3}$  and  $G_0 = 10^{3.5}$ . However, the high cosmic ray ionization rate (about 100 times larger than the Galactic value) prevents us from a direct comparison with Mon R2. In general, the production of  $\text{CO}^+$  seems to depend on the temperature of the gas. Stäuber & Bruderer (2009) suggest that  $X[\text{CO}^+]$  of about  $10^{-11}$  are only reached in gas with  $T_{\text{k}} \approx 300 \text{ K}$ . As  $T_{\text{k}}$  depends on  $G_0$  and  $n_{\text{H}}$ , it is expected that the production of  $\text{CO}^+$  only occurs in regions with  $n_{\text{H}} \approx 2 \times 10^4 \text{ cm}^{-3}$  and  $G_0 \approx 10^3$ . The densities and  $G_0$  measured in IF and PDR 3 are in good agreement with these values, as IF presents  $G_0 \sim 5 \times 10^5$  and  $n_{\text{H}} \sim 5 \times 10^4 \text{ cm}^{-3}$  (Rizzo et al. 2003; Fuente et al. 2010) and PDR3 presents  $G_0 \sim 3.7 \times 10^4$  and  $n_{\text{H}} \sim 3.7 \times 10^4 \text{ cm}^{-3}$  (Berné et al. 2009).

## 5 Discussion and summary

We present a  $\text{CO}^+$  map toward Mon R2 star-forming region. This is the first map ever reported of this reactive ion. The spatial distribution of  $\text{CO}^+$  consists of a ring-like structure (similar to PAHs), tracing the layer between the  $\text{H II}$  region and the molecular gas. The maps reveal a clumpy structure in the hot layer of the mainly atomic gas. Previous works (Young et al. 2000; Goicoechea et al. 2016) suggest that there exist fragmentation in the photodissociation front. This is, there are not uniform layers between the  $\text{H II}$  region and the molecular cloud, but they present clumps that allow the radiation to penetrate deeper into the cloud. In this scenario where the PDR is conformed by a series of clumps, the emission of PDR tracers would be related to the external layers of dense clumps being photo-evaporated by the UV radiation. Despite the moderate angular resolution of our observations, we find hints that favor this scenario: (a) the spatial distribution of the  $\text{CO}^+$  as observed in the higher-angular resolution ( $11''$ ) map (*Panel A* of Fig. 1) suggests that the  $\text{CO}^+$  emission is coming from the illuminated surface of the  $\text{H}^{13}\text{CO}^+$  clumps, and b) the excess of blue-shifted emission seen for the PDR tracers in comparison with the molecular tracers. Considering that the molecular gas is located behind the UC  $\text{H II}$  region (Pilleri et al. 2014) and the chemical segregation, the difference in velocity between tracers can be explained if the PDR is formed by dense condensations that are being photo-evaporated. In this case, the photo-evaporated gas (PDR tracers) would be ejected toward us, and therefore, blue-shifted with respect to the molecular gas. Future higher angular resolution observations will help to confirm or discard this scenario.

Finally, we have determined  $X[\text{CO}^+]$  in three positions. Toward IF we derive an abundance of a few  $10^{-11}$ , in agreement with chemical model predictions (Sternberg & Dalgarno 1995) for  $n_{\text{H}} \sim 10^6 \text{ cm}^{-3}$  and  $G_0 \sim 5 \times 10^5$ . Toward MP2 we do not detect  $\text{CO}^+$  emission with an upper limit to the  $\text{CO}^+$  abundance of  $< 4 \times 10^{-11}$ . Abundances of  $10^{-11}$ – $10^{-10}$  had been previously observed in PDRs with  $G_0 > 10^3$  Habing field (M17SW: Latter et al. 1993, Stoerzer et al. 1995; Orion Bar: Fuente & Martín-Pintado 1997; NGC7023: Fuente et al. 2003; G29.96–0.02: Rizzo et al. 2003). The non-detection of  $\text{CO}^+$  in MP2, together with the abundances found in the other PDRs, suggest that the production of  $\text{CO}^+$  only occurs in dense regions with high radiation fields. High UV fields ( $G_0 > 10^3$ ) and  $n_{\text{H}} (> 2 \times 10^4 \text{ cm}^{-3})$  are required to achieve gas temperatures  $\approx 300 \text{ K}$  that are necessary to produce high abundances of OH in the external layer of the PDR ( $A_{\text{V}} \sim 1 \text{ mag}$ ; Stäuber & Bruderer 2009). This is also consistent with the non-detection of  $\text{CO}^+$  in the Horsehead PDR (Goicoechea et al. 2009), where the UV field is  $G_0 \sim 100$  and presents chemical properties similar to MP2

(Ginard et al. 2012). A counter-example that challenge this interpretation could be the detection of CO<sup>+</sup> toward S140 where the incident UV field is estimated to be  $G_0 \sim 100\text{--}300$  (Savage & Ziurys 2004). However, the number of CO<sup>+</sup> detections is scarce and the statistics is not enough to draw firm conclusions between the relation of the CO<sup>+</sup> and the physical properties ( $n_{\text{H}}$  and  $G_0$ ) of PDRs. A larger sample of objects need to be studied, including maps to characterize and understand the spatial distribution of the CO<sup>+</sup> in different environments.

## Appendix

Refer to Web version on PubMed Central for supplementary material.

## Acknowledgements

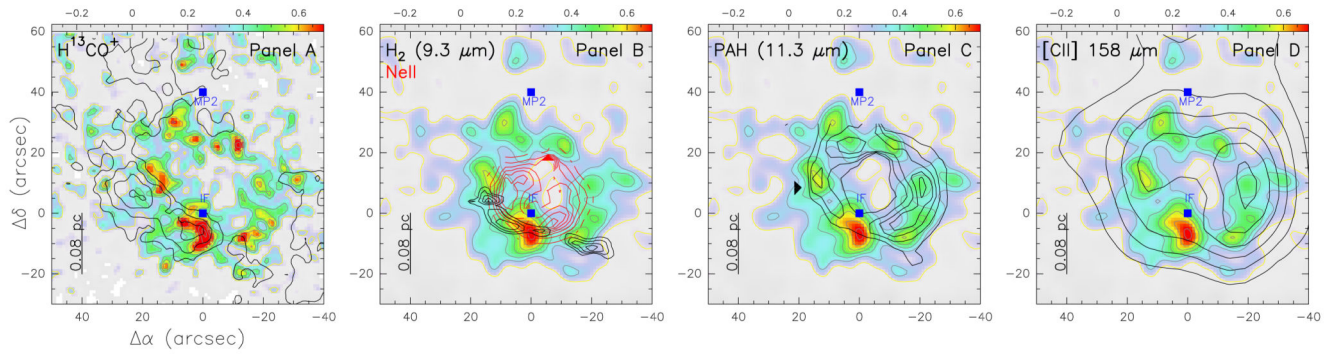
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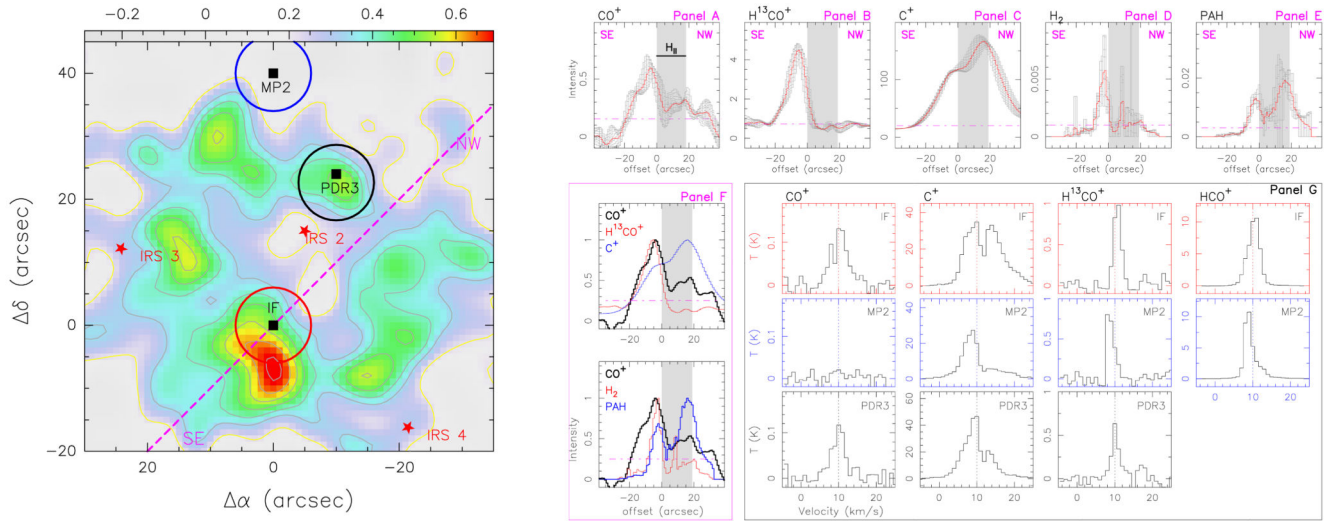
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**Fig. 1.**

Color image shows the integrated emission (in  $\text{K km s}^{-1}$ ) of the  $\text{CO}^+$  line, with the original ( $11''$ , *Panel A*) and smoothed ( $16''$ , *Panel B* to *D*) angular resolution. In *Panels A* to *D*, the gray contour levels range from 40% to 100%, in steps of 10% of the intensity peak, where the lower contour level corresponds to a  $S/N = 5\sigma$ . The yellow contour indicates the  $3\sigma$  emission. The blue squares indicate the IF and MP2 positions, where IF corresponds to  $\alpha(\text{J2000}) = 06\text{h}07\text{m}46.2\text{s}$ ,  $\delta(\text{J2000}) = -06^\circ 23' 08.3''$ . *Panel A* shows the  $\text{H}^{13}\text{CO}^+$  (3–2) emission (black contours) tracing the molecular gas (Treviño-Morales et al. 2014). *Panel B* shows the [NeII] emission (red contours) tracing the HII region and the emission of the  $\text{H}_2$  S(3) rotational line at  $9.7 \mu\text{m}$  (black contours). Black contours in *Panel C* show the PAHs ( $11.3 \mu\text{m}$ ) emission. *Panel D* shows the [CII] emission at  $158 \mu\text{m}$  (black contours, Pilleri et al. 2014). The *Spitzer* data are explained in Berné et al. (2009).

**Fig. 2.**

**Left:** CO<sup>+</sup> integrated emission (see Fig. 1). The pink-dashed line indicates the direction of the intensity cuts. The color circles indicate the positions where spectra were extracted.

**Right:** *Panel A to E* show the intensity cuts (red continuum lines) of CO<sup>+</sup>, H<sup>13</sup>CO<sup>+</sup>, [CII], H<sub>2</sub> and PAHs. In these panels, the gray continuum lines indicates the errors in the cuts. *Panel F* shows the comparison of the intensity cuts scaled to unity. In *Panels A to F*, the pink dot-dashed indicate the 3σ level for each species and the gray area indicates the position of the UC HII region. *Panel G* shows the CO<sup>+</sup>, [CII], H<sup>13</sup>CO<sup>+</sup> and HCO<sup>+</sup> spectra at IF, MP2 and PDR3. The color of the boxes are related to this positions (circles in left panel).