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## Modeling Cable and Guide Channel Interaction in a High-Strength Cable-Driven Continuum Manipulator

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## Abstract

This paper presents several mechanical models of a high-strength cable-driven dexterous manipulator designed for surgical procedures. A stiffness model is presented that distinguishes between contributions from the cables and the backbone. A physics-based model incorporating cable friction is developed and its predictions are compared with experimental data. The data show that under high tension and high curvature, the shape of the manipulator deviates significantly from a circular arc. However, simple parametric models can fit the shape with good accuracy. The motivating application for this study is to develop a model so that shape can be predicted using easily measured quantities such as tension, so that real-time navigation may be performed, especially in minimally-invasive surgical procedures, while reducing the need for hazardous imaging methods such as fluoroscopy.

## **Index Terms**

Cable-driven robots; kinematics; medical robots and systems; snake-like robots; underactuated robots

## I. Introduction

Small, dexterous manipulators are becoming increasingly popular for minimally-invasive surgical procedures [1]–[7]. This paper is about a new type of high-strength cable-driven

This paper has supplementary downloadable material available at http://ieeexplore.ieee.org provided by the authors. This includes animations in gif format, showing comparison of model and experiment for eight separate test cases. Contact matt.moses@jhu.edu for further questions about this work.

flexible manipulator for medical applications [8]–[10]. This is a small snake-like device that can be hand held, or mounted to a larger robot for positioning (see Fig. 1). The manipulator is designed with a large open lumen expressly for the procedure of removing diseased bone during hip revision surgeries, without necessitating removal of a well-fixed acetabular implant [8], although it is well suited for other types of minimally-invasive surgical interventions. Visualizing the shape of the manipulator while it is inside the body is thus an important goal. In addition, a good understanding of the relation between design variables and manipulator stiffness) is important for designing new and improved systems.

Many existing cable-driven manipulators are described well by constant curvature models [11]. While the constant curvature assumption describes our manipulator with good accuracy at small deflections, the constant curvature prediction does not match the observed shape of our robot at higher cable tensions and deflections. One of the main advantages of our new design is that it is expressly intended for operation at high forces and deflections, so it is important to develop a predictive model that describes the manipulator behavior in these regimes of operation.

The purpose of this paper is to investigate the interaction between the actuating cables and their guide channels, in order to determine the extent to which this interaction is responsible for deviation from the "ideal" constant curvature model. Further, we take a step toward developing a model which predicts the shape of the manipulator given only material properties and the time-dependent tension on the drive cables. While specifically driven by our own particular manipulator design, this study can be applied to a broad range of cable-actuated flexible manipulators.

Friction between cable and manipulator plays an important role in manipulator behavior. While friction and stiffness effects have been relatively well-studied in the context of cabledriven robotic hands [13]–[16], friction is not considered in many works on cable-driven continuum manipulators, both pioneering and recent [17]-[20]. A number of studies investigate the frictional interaction between a snake-like medical device and the surrounding tissue, e.g., [21], but few studies explicitly address the friction of internal actuation cables. Some recent exceptions are [22]–[25]. In [22], cable friction is studied in depth, and shown to be at the root of hysteretic behavior shown in cable-driven surgical tools. A detailed lumped-parameter model using Dahl friction is presented in [23] and [24], which accurately predicts the curvature and tip location, after bending, of a manipulator constructed with a Teflon backbone and Teflon-coated cables. In this paper, we use a similar dynamic model with standard Coulomb friction to accurately predict the curvature and tip location, after bending, of a manipulator made with a Nitinol backbone and stainless steel cables. Extending beyond [23], [24], we present a geometric analysis of the free motion of the cable within the channels, a discussion of the separate roles of cable and backbone elasticity in manipulator stiffness, and we report experimental data on the effect of cable friction on time-dependent manipulator shape.

In Section II, basic kinematic and quasi-static dynamic model is presented. This section addresses the kinematic relation between cable length and manipulator shape, including

motion of the cable within the guide channel. Section III builds on the kinematic model from Section II to investigate how manipulator shape and stiffness is affected by cable tension and material properties. Section IV addresses the problem of parameterizing the nonconstant curvature manipulator shape, as well as effects of friction between cable and guide channel on the manipulator's shape and time-dependent behavior.

#### II. Kinematics of the Manipulator and Cables

This section presents a basic kinematic model of the manipulator and the drive cables (see Fig. 2).

#### A. Backbone Kinematics

The manipulator is designed to be used with a supporting robot (see Fig. 1) that actuates the base platform. For this reason, only the flexible backbone portion of the manipulator is considered in the kinematic model. The manipulator is made of two nested tubes with alternating slots cut in the sides. There are fourteen slots on each side of the tube, resulting in 27 flexure joints. Because of the symmetrical construction and small deflection of the flexure joints, the joints are well described by a simple kinematic revolute joint equipped with a torsional spring (see Section III-D for additional discussion). At each revolute joint, the stiffness of the flexure is represented by a nonlinear spring. Joint *i* connects rigid links *i* and i - 1. The base link (link 0) is fixed. The forward kinematics, written in product of exponentials form, are

$$g_{st}(\theta) = \exp\left(\hat{\xi}_1\theta_1\right) \exp\left(\hat{\xi}_2\theta_2\right) \dots \exp\left(\hat{\xi}_{27}\theta_{27}\right) g_{st_0}$$

where  $g_{st}$  is the current configuration of the end-effector,  $\theta$  is a 27 × 1 vector of joint angles, and  $g_{st_0}$  is the initial end-effector configuration. See the Appendix for a brief review of this notation, and [26] for detailed coverage. The initial configuration of each link is given by the 4 × 4 matrix representation

$$g_{si_0} = \left[ \begin{array}{cc} I & q_i \\ 0 & 1 \end{array} \right]$$

where  $q_i = (0, 0, (i-1)p)^T$ . The base link is link 0; links 1 through 26 are the repeated subunits of the compliant backbone; and the end-effector is link 27, that is,  $g_{st} = g_{s27}$ . The spacing between links is p = 1.21 mm. The configuration of the *ith* link is obtained by

$$g_{si}(\theta) = \prod_{k=1}^{i} \exp\left(\hat{\xi}_k \theta_k\right) g_{si_0}$$

The axis of rotation of each joint is in the direction of  $\omega_i = (0, 1, 0)^T$ . The joint twists are given by

$$\xi_{i} = \begin{pmatrix} -\omega_{i} \times q_{i} \\ \omega_{i} \end{pmatrix} = \begin{pmatrix} (1-i)p \\ 0 \\ 0 \\ 0 \\ 1 \\ 0 \end{pmatrix}$$

The manipulator Jacobian  $J_{st}^s$  relates joint velocities to the spatial velocity of the endeffector

$$V_{st}^{s} = ((\dot{g}_{st}) g_{st}^{-1})^{\vee} = (J_{st}^{s}(\theta)) \dot{\theta} .$$
(1)

 $J_{st}^{s}(\theta)$  is a 6 × 27 matrix. For a general redundant manipulator,  $J_{st}^{s}$  can have rank of at most 6. However, since all of the joints in this manipulator are parallel,  $J_{st}^{s}$  in fact has rank of at most 3, corresponding to the three degrees of freedom with which the end-effector can move in the plane. For most configurations  $g_{st}$  there are many corresponding  $\theta$ . The set of  $\theta$  corresponding to a given configuration is called the *self-motion manifold*. The manipulator can move freely among these  $\theta$ , given that the infinitesimal motions satisfy  $(J_{st}^{s}(\theta)) \dot{\theta} = 0$ . With the end-effector of the manipulator held fixed, the internal spring torque at the joints moves the manipulator along the self-motion manifold to a configuration of minimum energy.

#### **B. Cable Kinematics**

The manipulator is actuated by two cables, one threaded through each side. The length of each cable can be treated as a kinematic configuration variable, and the same technique of forward kinematics and Jacobians can be applied to gain useful insight into the functioning of the manipulator.

Referring to Fig. 2, consider  $L_{Ri}$ , the length of the right cable, in between link i - 1 and i + 1. The length of this cable segment is the length of the vector between cable entrance points in link i - 1 and link i + 1. For consistency with the rest of the discussion the cable is represented as a line segment in three dimensions, written as a vector in  $\mathbb{R}^3$ , even though the manipulator is assumed to lie in the *xz* plane. Writing this vector relative to a frame fixed on link *i*, the length of the cable segment is

$$L_{Ri} = \left\| \left[ \begin{array}{c} c_{i+1}r_x - s_{i+1}r_z \\ 0 \\ p - s_{i+1}r_x - c_{i+1}r_z \end{array} \right] - \left[ \begin{array}{c} c_ir_x - s_ir_z \\ 0 \\ s_ir_x + c_ir_z \end{array} \right] \right\|$$

where  $c_i = \cos(\theta_i)$  and  $s_i = \sin(\theta_i)$ .

Each joint angle has limited range of motion (approx. 0.14 rad) so it is reasonable to make the assumptions that  $c_i = 1$  and  $s_i = \theta_i$ . This simplifies the aforementioned to

$$L_{Ri} \approx \| \begin{bmatrix} (\theta_i - \theta_{i+1})r_z \\ 0 \\ p - (\theta_i + \theta_{i+1})r_x - 2r_z \end{bmatrix} \|$$
  
=  $((\theta_i - \theta_{i+1})^2 r_z^2 + (p - (\theta_i + \theta_{i+1})r_x - 2r_z)^2)^{1/2}$ .

The fact that  $r_z \ll r_x$  and  $r_z \ll p$  ( $r_z = 0.2$  mm, p = 1.21 mm,  $r_x = 2.5$  mm) allows us to neglect the  $r_z^2$  term, and  $L_{R_i}$  further simplifies to

$$L_{Bi} \approx L_{Bi} = p - (\theta_i + \theta_{i+1})r_x - 2r_z$$

where  $L_{RSi}$  is the approximated length of the cable segment under the aforementioned simplifying assumptions, and  $L_{RS}$  is likewise the approximated length of the entire cable. The total length of the right cable is then given by

$$\begin{split} L_{R} &\approx L_{RS} = L_{R_{\text{offset}}} + p - (\theta_{1})r_{x} \\ &- 2r_{z} + p + 2r_{z} + p - (\theta_{2} + \theta_{3})r_{x} - 2r_{z} + p + 2r_{z} + p - (\theta_{4} + \theta_{5})r_{x} \\ &- 2r_{z} + \ldots + p - (\theta_{26} + \theta_{27})r_{x} - 2r_{z}. \end{split}$$

The length of the right cable for a given set of joint angles is

$$L_{\scriptscriptstyle R}(\theta) \approx L_{\scriptscriptstyle RS}(\theta) {=} L_{\scriptscriptstyle R0} {-} r_x {\sum_{i=1}^{27}} \theta_i. \tag{2}$$

The same procedure gives the length of the left cable

$$L_{L}(\theta) \approx L_{LS}(\theta) = L_{L0} + r_{x} \sum_{i=1}^{27} \theta_{i}$$
 (3)

where  $L_{R0}$  and  $L_{L0}$  are constant terms containing the length of the cable offset (external to the flexible portion of the manipulator) and a sum of other dimensions.  $L_{R0} = L_{Roffset} + 27p - 2r_z$ , and  $L_{L0} = L_{Loffset} + 27p - 2r_z$ . For convenience we write the vector of cable length as  $L(\theta) = (L_R(\theta), L_L(\theta))^T$ .

#### **C. Cable Coupling Matrix**

The coupling matrix [26]  $P(\theta)$  relates joint velocities  $\dot{\theta}$  to change in cable length

$$\dot{L} = \begin{bmatrix} \dot{L}_{R} \\ \dot{L}_{L} \end{bmatrix} = P^{T}(\theta) \ \dot{\theta}$$
(4)

where

$$P^{T}(\theta) = \begin{bmatrix} \frac{\partial L_{R}}{\partial \theta_{1}} & \frac{\partial L_{R}}{\partial \theta_{2}} & \cdots & \frac{\partial L_{R}}{\partial \theta_{27}} \\ \frac{\partial L_{L}}{\partial \theta_{1}} & \frac{\partial L_{L}}{\partial \theta_{2}} & \cdots & \frac{\partial L_{R}}{\partial \theta_{27}} \end{bmatrix}_{2 \times 27}$$

In general the coupling matrix may depend on  $\theta$ . However, using the simplified expressions for cable length (2), (3) results in

$$P^{T} \approx r_{x} \begin{bmatrix} -1 & -1 & \dots & -1 \\ 1 & 1 & \dots & 1 \end{bmatrix}_{2 \times 27}$$
(5)

This simplified coupling matrix accurately describes many cable-driven manipulators with a flexible backbone. While the structure of  $P^T$  is trivial, it provides helpful insight into the manipulator. Starting at an initial  $\theta_0$ , any motion  $\dot{\theta}$  in the right nullspace of  $P^T$  will leave the cable lengths unchanged. By inspection, rank( $P^T$ ) = 1, so there is a 27 – 1 = 26 dimensional subspace of motions in  $\theta$  for which cable lengths will remain constant. By inspection of the manipulator kinematics, we see that the sum of the joint angles is in fact the angular orientation of the end-effector. The cable self-motion manifold is therefore any motion for which the orientation of the end-effector remains constant.

#### **D.** Comparison of Cable Length Models

In the physical manipulator, the actuating cables are free to move sideways in the guide channels, and the manipulator shape becomes bistable at high counter-tensions. A high-accuracy model of cable length for this manipulator was presented in [9]. This model uses Dijkstra's algorithm to find the shortest route by which a cable of nonzero thickness can pass through a channel. In this section, we compare the simple cable length model to the higher accuracy one of [9].

The ideal "zero channel width" length  $L_{Ri}$  represents the length of an ideal zero-thickness cable through a zero-thickness channel. This quantity is an upper bound to the length of a cable segment that is free to move in a wider channel, since widening the channel will only increase the number of possible shorter paths. The error between  $L_{Ri}$  and the simplified  $L_{RSi}$  is shown in Fig. 3.

The shortest possible length of an intersegment cable length  $L_{DMi}$  can be found as illustrated in Fig. 4. A lower bound on the total length of a cable free to move in a larger channel  $L_{DM}$ is found by summing the minimum possible lengths of all segments,  $L_{DM} = \sum_i L_{DMi}$ . Fig. 5 shows the maximum error between the simple cable model  $(L_{RSi})$  with the upper  $(L_{Ri})$  and lower  $(L_{DMi})$  bounds of the cable lengths predicted by the method of [9].

The error between the ideal "zero channel width" length  $L_{Ri}$  and the simple approximated length  $L_{RSi}$  is very small—on the order of micrometers across the whole operating range (see Fig. 3). The error between  $L_{RSi}$  and the bounds of method [9] is somewhat higher (see Fig. 5), but in both cases the error is small compared to the total length of the cable (see Fig. 6). In many cases, this small error can be neglected and the simple length model can be used. However, at inflection points, that is, places along the manipulator where the joint angles  $\theta_i$ change sign, it is important to consider these very small changes in cable length. Cases where joint angles change sign occur if there is an "S-bend" shape in the manipulator, and they also occur even when the manipulator is in the shape of a single circular arc, bending from one side to the other [9]. As the manipulator straightens, the most stable shape becomes a slightly buckled "S-bend" rather than a perfectly straight line. While these shape changes are rather small, they are accompanied by very rapid accelerations as the taught cable snaps from one side of a channel to another. Predicting and understanding these bifurcations of equilibrium state are important for a complete picture of manipulator behavior, but this topic is left for future work. The error analysis provides confidence that the conclusions we draw using the simplified length model and coupling matrix will closely approximate the true behavior of the physical manipulator in regions where bifurcations do not occur.

## III. Quasi-Static Dynamic Model, Cable Tension, and Manipulator Stiffness

Equipped with the manipulator and cable Jacobians, we are now ready to address how the manipulator, in the absence of cable friction, responds to cable tension and to forces applied to the end-effector. While there are sophisticated methods for solving the equilibrium configurations of elastic chains (e.g., [27]), we take a more basic approach. Recall the familiar robot dynamics [26]

 $\tau = M(\theta)\ddot{\theta} + C(\theta, \dot{\theta}) \dot{\theta} + N(\theta, \dot{\theta}).$ 

Several simplifying assumptions are made to reduce this general equation of motion to a set of decoupled first order differential systems. We are primarily interested in the equilibrium configuration, so we assume the movement of the manipulator is slow. Hence, the contributions of the Coriolis terms are negligible, and the joints are decoupled, i.e.,  $C(\theta, \dot{\theta})\dot{\theta} \approx 0$  and M is diagonal. The constitutive behavior of the flexible links is modeled as a nonlinear spring in parallel with a viscous damper. These terms appear in N, i.e.  $N(\theta, \dot{\theta}) = K_{\theta}(\theta)\theta + B\dot{\theta}$ . These same assumptions are used successfully to model continuum manipulators in [23], [24].

The matrix  $K_{\theta}(\theta)$  represents the elasticity of the flexures in addition to the obstruction to movement caused by collisions between joint segments. Referring to Fig. 2, note that the presence of a collision between link i-1 and i+1 is determined by both  $\theta_i$  and  $\theta_{i+1}$ . Precise conditions for collisions are cumbersome to calculate, but a simplifying assumption is to limit each  $\theta_i$  to a value such that collisions are guaranteed not to occur. The flexure elasticity and intersegment collisions are modeled using a dual rate torsional spring for each joint,  $K_{\theta}$  $(\theta) = \text{diag}(k_i(\theta_i))$  where

$$k_i(\theta_i) = \begin{cases} k_1, & \text{if}|\theta_i| \le \theta_{\text{thresh}} \\ k_2 + \theta_{\text{thresh}}(k_1 - k_2)/\theta_i, & \text{if} \, \theta_i > \theta_{\text{thresh}} \\ k_2 - \theta_{\text{thresh}}(k_1 - k_2)/\theta_i, & \text{if} \, \theta_i < -\theta_{\text{thresh}}. \end{cases}$$

The corresponding joint torque induced by this nonlinear spring constant,  $\tau_i(\theta_i) = k_i(\theta_i)\theta_i$ , is shown in Fig. 7.

The aforementioned assumptions result in a system of n decoupled second order equations

$$\tau = M\ddot{\theta} + K_{\theta}(\theta)\theta + B\dot{\theta}$$

which can be further reduced to a first order system by assuming that damping is dominant in the dynamic behavior

$$\tau = K_{\theta}(\theta)\theta + B \theta .$$
 (6)

While this is not an accurate assumption to make for modeling high-speed transient behavior of the manipulator, it is entirely reasonable in this case because the primary quantity of interest is the equilibrium system configuration.

The torque induced on the joints  $\tau_{ind}$  is due to the combined effects of the applied wrench at the end effector and tension on the cables. Recall that the transpose of  $J_{st}^s$  relates a general spatial wrench,  $F_{st}$  to the induced torques on the joints of the robot

$$\tau_{\text{wrench}} = (J_{st}^s(\theta))^T F_s. \quad (7)$$

The dimensions of  $(J_{st}^s(\theta))^T$  are 27 × 6. Wrenches in the right nullspace of  $(J_{st}^s(\theta))^T$ , i.e.,  $F_s$  such that  $(J_{st}^s(\theta))^T F_s = 0$  correspond to wrenches that are supported entirely by the structure of the manipulator, independent of joint torques. Because of the planar design of the manipulator,  $(J_{st}^s(\theta))^T$  is of rank at most 3. Basis vectors for the right nullspace (written in body coordinates) are  $F_1 = (0, 1, 0, 0, 0, 0)^T$ ,  $F_2 = (0, 0, 0, 1, 0, 0)^T$ , and  $F_3 = (0, 0, 0, 0, 0, 0, 0, 1)^T$ . That is,  $\tau_{wrench} = 0$  for any

$$F_{s} = \left(Ad_{g_{st}^{-1}}\right)^{T} (a_{1}F_{1} + a_{2}F_{2} + a_{3}F_{3}).$$
(8)

The adjoint transformation  $Ad_g$  is reviewed in the Appendix. In special configurations  $(J_{st}^s(\theta))^T$  can drop rank to less than 3, such as when  $\theta_i = 0$  for i = 0...26. In this case, the rank of  $(J_{st}^s(\theta))^T$  is 2 and the additional basis vector is  $F_4 = (0, 0, 1, 0, 0, 0)^T$ .

In a manner similar to (7), the coupling matrix relates cable tension to the induced joint torques [26]

$$\tau_{\text{cable}} = P(\theta)T$$
 (9)

where  $T = (T_R, T_L)^T$  is the vector of cable tensions, with  $T_L$  and  $T_R$  being the left and right cable tensions, respectively. It is assumed  $T_L$ ,  $T_R = 0$ .

One further point regarding (9) should be mentioned. Inspection of (5) shows that cable tension results in the same torque being uniformly induced on all joints. In the next section, it is shown that this is consistent with the familiar constant-curvature assumption that is often used to model continuum manipulators.

The full quasi-static equations of motion are given by incorporating the induced torques (7) and (9) into the simplified dynamics model (6) and setting the applied joint torque  $\tau$  to zero

$$0 = K_{\theta}(\theta)\theta + B \theta + \tau_{\text{wrench}} + \tau_{\text{cable}}$$

This equation can be solved iteratively for the joint angles  $\theta(t)$  by integrating

$$\dot{\theta} = -\frac{1}{B} (K_{\theta}(\theta)\theta + (J_{st}^{s}(\theta))^{T} F_{s} + P(\theta)T)$$
(10)

from an initial  $\theta(t_0)$ . This results in the equilibrium shape of the manipulator given a timedependent application of end-effector loading and cable tensions. One example is shown in Fig. 8. The manipulator starts from  $\theta = 0$  and then a force and moment load are applied. Different equilibrium solutions are reached depending on the order of application of force and moment, even though the final loading conditions are the same. Note that applying a pure moment to the end-effector, i.e.  $F_s = a(0, 0, 0, 0, 1, 0)^T$  is exactly equivalent to applying cable tension,  $T_L - T_R = 0$ . In each case torques are induced uniformly over all joints. For this reason, the example shown in Fig. 8 is also valid if a cable tension were applied instead of a moment to the end-effector.

#### A. Constant-Curvature Regime

In the absence of external load, i.e., when  $F_s = 0$ , the equilibrium configuration is

$$\theta = -K_{\theta}(\theta)^{-1} P(\theta)^{T}$$
. (11)

Using the simplified coupling matrix generated with small angle approximations (5), and assuming the joints are in their linear regime, i.e.  $k_i(\theta_i) = k_1$ , the equilibrium configuration is

$$\theta_i = -\frac{r_x}{k_1}(T_L - T_R).$$

Every joint angle is the same, so the joint axes are located along the arc of a circle with radius  $R = p/\theta_i$ . Thus, the model is consistent with a constant curvature approximation, and the curvature is given by

$$\kappa = \frac{1}{R} = \frac{r_x}{pk_1} |T_L - T_R|.$$
(12)

#### **B. Cable Length as Control Inputs**

The equations of motion in (10) use cable tension *T* as a control input. In some cases it is convenient to consider the length of the cables as the control input instead. Let  $u(t) = (u_R(t), u_L(t))^T$  be the controlled positions of the ends of the cables. Cable elasticity is modeled as a single spring element (external to the manipulator) in series with an ideal inelastic cable (see Fig. 9). Elongation of the cables, defined as the difference between the actual and the unstretched lengths, is given by

$$d = L(\theta) - L_0 + u \quad (13)$$

where  $d = (d_R, d_L)^T$  and  $L_0 = (L_{R0}, L_{L0})^T$ . In general, the cable tension is some function of the elongation

$$T(\theta, u) \equiv T(d) = \begin{bmatrix} T_R(d_R) \\ T_L(d_L) \end{bmatrix}.$$
 (14)

The expressions for T and d can be substituted directly into (10) for the equations of motion using cable lengths as the control input. It is assumed that u is chosen such that  $d_R$ ,  $d_L$ ,  $T_R$ ,  $T_L = 0$ .

In the upcoming discussion on stiffness, it is necessary to find the change in cable tension in response to a small perturbation in the manipulator shape. This can be done by linearizing the tension about some operating point  $d_0 = L(\theta_0) - L_0 + u_0$ . Noting that for a fixed control input  $u_0$ , an infinitesimal displacement in joint angles  $\delta\theta$  results in an elongation of  $\delta d = P^T$   $(\theta_0) \delta\theta$ , the linearized tension is

$$T(\theta_0 + \delta\theta, u_0) = T(\theta_0, u_0) + \begin{bmatrix} \frac{\partial T_R}{\partial d_R} & 0\\ 0 & \frac{\partial T_L}{\partial d_L} \end{bmatrix} P^T(\theta_0) \delta\theta$$
(15)

where the partial derivatives are evaluated at  $\theta_0$ ,  $u_0$ . When the cables are modeled as linear springs, this expression simplifies to

$$T(\theta_0 + \delta\theta, u_0) = T(\theta_0, u_0) + K_c P^T(\theta_0) \delta\theta$$
$$= K_c \left( L(\theta_0) - L_0 + u_0 + P^T(\theta_0) \delta\theta \right)$$
(16)

where  $K_c = \text{diag}(k_c)$  is the tensile spring constant of the cables. A simple approximation for cable spring constant is

$$k_c = \left(\frac{EA}{L_0}\right)$$

where *E* is Young's modulus and *A* is the cable cross-sectional area. Approximate values for the stainless steel cables used in this manipulator are  $E = 1.8 \times 10^{11} \text{ N/m}^2$ ,  $A = 5.1 \times 10^{-8} \text{ m}^2$ ,  $L_0 = 2.2 \times 10^{-1} \text{ m}$ , resulting in  $k_c = 4.1 \times 10^4 \text{ N/m}$ .

#### C. Manipulator Stiffness

Once the manipulator attains an equilibrium configuration, it will to some degree maintain that configuration in response to disturbances from the environment. This behavior can be modeled as a general stiffness matrix. Because a high stiffness is desirable, it is useful to consider the contributions of various design parameters to this stiffness. In this section, we address the effects that various design parameters have on manipulator stiffness.

1)  $K_m$ —Jointspace Stiffness—Assume that we start with an unloaded manipulator and then move the cables to some location specified by  $u_0$ . The manipulator will attain an equilibrium configuration  $\theta_0$ . We assume that  $\theta_0$  places the backbone flexures in their linear regime, i.e.  $K_{\theta}(\theta_0) = K_{\theta} = \text{diag}(k_1)$ . Holding  $u_0$  constant, suppose the backbone is displaced by an infinitesimal virtual motion  $\delta\theta$ . The resultant torque appearing on the flexure joints is due to deformation of both the backbone and the cables. This torque is given by

$$\tau = K_{\theta}\theta_0 + K_{\theta}\delta\theta + P(\theta_0 + \delta\theta)T(\theta_0 + \delta\theta, u_0).$$
(17)

It is a reasonable assumption that  $P(\theta_0 + \delta \theta) = P(\theta_0) = P$ , as it was shown in Section II-C that a constant *P* matrix is a very close approximation of the true  $P(\theta)$ . Expanding (17) and making use of (16) we have

$$\tau = K_{\theta}\theta_0 + K_{\theta}\delta\theta + P(\theta_0)T(\theta_0, u_0) + P(\theta_0)K_cP^T(\theta_0)\delta\theta.$$

Note from (11) that  $K_{\theta} \theta_0 = -P(\theta_0) T(\theta_0, u_0)$ , so we have now

$$\tau = (K_{\theta} + Pk_c P^T) \,\delta\theta = K_m \delta\theta$$

The quantity  $K_m$  is a 27 × 27 symmetric positive definite matrix representing the combined stiffness of the backbone and cables, in jointspace coordinates. The matrix  $PK_cP^T$  is equal to a constant times the matrix of ones, that is, every element is equal to  $2k_c r_x^2$ . Further,  $PK_cP^T$  has rank one, and the eigenvector corresponding to its only nonzero eigenvalue is  $(1, 1, 1, ..., 1)^T$ .

**2)**  $K_w$ —Workspace Stiffness—We define workspace stiffness as a natural extension to a linear spring, except for full 3-D motion, as described in [15], [16], [28]

$$F_s = K_s V_{st}^s \delta t$$

where  $F_s$  is a six-element wrench applied to the end-effector,  $\delta t$  is an infinitesimal timestep,  $V_{st}^s \delta t$  is an infinitesimal displacement of the end-effector, composed of a small rotation about an arbitrary axis and small translations along three orthogonal axes.  $K_s$  is a 6 × 6 stiffness matrix.

For the purposes of this discussion, only the "flexible" degrees of freedom are considered in the stiffness analysis. The manipulator is assumed to be effectively rigid in response to forces of the form in (8). Stiffness resisting forces in the directions of  $F_1$ ,  $F_2$ ,  $F_3$  will not be considered, as the stiffness in these is much higher than that in the intentionally flexible modes. With this assumption, parameterizing the end effector configuration as  $X = (x, z, \phi)^T$ , the Jacobian mapping between these coordinates and the joint velocities is

$$\dot{X} = \begin{bmatrix} \dot{x} \\ \dot{z} \\ \dot{\phi} \end{bmatrix} = \begin{bmatrix} J_x(\theta) \\ J_z(\theta) \\ J_{\phi}(\theta) \end{bmatrix} \dot{\theta} = J_X(\theta) \dot{\theta}$$
(18)

where  $J_X$  is a 3 × 27 matrix. Note that the Jacobian  $J_X$  is neither the spatial nor body manipulator Jacobian [see (27)]. The distinction between these three types of Jacobian mappings is described in detail in [26]. The *k*th entry of  $J_X$  is

$$J_{xk} = p \sum_{i=k}^{27} \cos \left( \sum_{j=1}^{i} \theta_j \right).$$
(19)

The kth entry of  $J_z$  is

$$J_{zk} = -p \sum_{i=k}^{27} \sin \left( \sum_{j=1}^{i} \theta_j \right)$$
(20)

where *p* is the spacing between links (see Section II-A), and the *k*th entry of  $J_{\phi}$  is simply equal to 1, i.e.  $J_{\phi k} = 1$ .

Assume that we start with an unloaded manipulator and then move the cables to some location specified by  $u_0$ . The manipulator will attain an equilibrium configuration  $\theta_0$ . We assume that rank  $J_X(\theta_0) = 3$  and that  $\theta_0$  places the backbone flexures in their linear regime, i.e.  $K_{\theta}(\theta_0) = K_{\theta} = \text{diag}(k_1)$ . We are interested in determining the  $3 \times 3$  stiffness matrix  $K_w$  that describes the reactive force of the end-effector in response to infinitesimal disturbances in position

$$F = \begin{bmatrix} f_x \\ f_z \\ \tau_\phi \end{bmatrix} = \begin{bmatrix} k_{xx} & k_{xz} & k_{x\phi} \\ k_{zx} & k_{zz} & k_{z\phi} \\ k_{\phi x} & k_{\phi z} & k_{\phi\phi} \end{bmatrix} \begin{bmatrix} \delta x \\ \delta z \\ \delta \phi \end{bmatrix} = K_w \delta X.$$

Holding  $u_0$  constant, suppose the end-effector is displaced by an infinitesimal motion  $\delta X$ . In jointspace coordinates, this deflection by  $\delta X$  corresponds to an infinitesimal motion  $\delta \theta$  away from the equilibrium joint angles  $\theta_0$ . We can solve for  $\delta \theta$  by minimizing the potential energy of the mechanical system, subject to the constraint  $\delta X = J_X(\theta_0)\delta\theta$ . Taking the zero point of the potential energy to be the value at equilibrium (i.e.,  $\theta = \theta_0$ ) the potential energy is given by

$$V = \frac{1}{2} \delta \theta^T K_m \delta \theta.$$
 (21)

The constrained minimization problem can then be formulated as a linear system of equations using the method of Lagrange multipliers

$$\begin{bmatrix} K_m & J_X^T \\ J_X & 0 \end{bmatrix} \begin{bmatrix} \delta\theta \\ \lambda \end{bmatrix} = \begin{bmatrix} 0 \\ \delta X \end{bmatrix}$$
(22)

where  $\lambda$  is a 3 × 1 vector of Lagrange multipliers. Provided that the three rows of  $J_X$  are linearly independent, the matrix of constraint equations can be inverted as

$$\begin{bmatrix} K_m & J_X^T \\ J_X & 0 \end{bmatrix}^{-1} = \begin{bmatrix} K_m^{-1} - K_m^{-1} J_X^T K_w J_X K_m^{-1} & K_m^{-1} J_X^T K_w \\ K_w J_X K_m^{-1} & -K_w \end{bmatrix}$$

where

$$K_w = (J_X K_m^{-1} J_X^T)^{-1}.$$
 (23)

With the inverted matrix, we may solve (22) for  $\delta\theta$ 

$$\delta\theta = (K_m^{-1} J_X^T K_w) \delta X. \quad (24)$$

Substituting (24) into (21) verifies that the potential energy as a function of  $\delta X$  is

$$V = \frac{1}{2} \delta X^T K_w \delta X \tag{25}$$

and the reactive force at the end effector is therefore

$$F = K_w \delta X.$$
 (26)

**3)** Stiffness Discussion—The structures of  $K_m$  and  $K_w$  provide us with useful insight into manipulator stiffness. When cable elasticity is linear, the stiffness does not depend on cable tension. This runs somewhat counter to our intuition, because we are familiar that tensioning both the agonist and antagonist cables usually results in a stiffening of the manipulator (a person can demonstrate this with their own limbs). However, it is known that this stiffening effect of cable counter-tension is in fact due to the cables themselves exhibiting a nonlinear spring constant [29], [30].

When cable stiffness is nonlinear, the manipulator stiffness can be controlled, to some extent, by changing the elongation of the cables. For example, if cable stiffness is

proportional to the square of the elongation,  $T(d) = K(d_R^2, d_L^2)^T$ , then the cable stiffness matrix depends linearly on d,  $K_c = \text{diag}(k_c d_R, k_c d_L)$ . A change in  $K_c$  will therefore influence  $K_m$  and  $K_w$ . However, the effect is limited by the span of  $PK_cP^T$ . Motions in jointspace that lie within the nullspace of  $PK_cP^T$  are not affected by the cables at all. The consequence of this is that there is a two-dimensional subspace of motions in the workspace (corresponding in this case to translations in x and z that leave  $\phi$  unchanged) for which stiffness is independent of  $K_c$  and determined only by  $K_{\theta}$ . The addition of more actuation cables to the manipulator increases the rank of P and allows stiffness control in additional degrees of freedom. For example, additional pairs of cables can be attached to adjacent pairs of segments, evenly spaced along the length of the manipulator. This is essentially the same as concatenating multiple two-cable-actuated manipulators together end-to-end.

An intuitive counting argument illustrates how many cables are needed to control joint angles as well as stiffness. As derived in [31] and clearly summarized in [32], a  $n \times n$ stiffness matrix K is determined by n(n + 1)/2 parameters. If n joint torques must also be controlled, then to fully control the joint angles and full stiffness matrix, a minimum of n(n + 3)/2 cables are required. Thus, at least nine cables are needed to fully control  $(x, z, \phi)$  and  $K_w$ . As pointed out in [32] it is common to neglect off-diagonal terms in  $K_w$ , which results in a requirement of 2n cables. Assuming we are concerned only with the diagonal terms of the stiffness matrix, at least six cables would be needed to control all three degrees of freedom  $(x, z, \phi)$  and their corresponding stiffnesses  $(k_{XX}, k_{ZZ}, k_{\phi\phi})$ . Since the manipulator only has two cables, it is clear that at most two of these six quantities are controllable. The two variables that are fully controllable are the angle of the tip  $\phi$  and the corresponding stiffness  $k_{\phi\phi}$ .

#### D. Validity of Modeling Flexures as Revolute Joints

It is common to model flexure joints in a segmented continuum manipulator as kinematic revolute joints (e.g., [10], [23], [24], [33], [34]), however in certain circumstances it is necessary to model additional degrees of freedom in the system. For example, in [19] the axial compression of a continuum manipulator is modeled in addition to the bending mode of deformation. In some applications, individual flexures may need to be modeled as more than one kinematic joint. For example, a planar parallel manipulator described in [35] required three of its nine flexure joints to be modeled as a combination of revolute and prismatic joints. In order to support our assertion that the flexure joints in our manipulator are accurately modeled as revolute joints, a segment of the manipulator was studied using finite element analysis (FEA) in Solidworks (see Fig. 10). Representative material parameters for Nitinol were chosen (E = 60GPa, v = 0.3) and typical loads were applied. The case of axial compression results in negligible deformation (maximum 0.015 mm). When compared with a segmented model made of rigid links connected by revolute joints, deformation due to a pure moment applied at the tip is consistent with constant joint angles,  $\theta_{1...4} = 0.129$  rad, while that due to a transverse load at the tip is consistent with uniformly decreasing joint angles,  $\theta_{1...4} = (0.096, 0.074, 0.053, 0.032)$  rad. All three cases produce results consistent with our model, indicating that no additional degrees of freedom need to be incorporated.

In addition to validating the kinematic assumptions, the same FEA results can be used to estimate an equivalent torsional stiffness for the flexure joints. Simple statics is used to determine the torque induced at each joint given the loading conditions. Joint stiffness is then estimated as the ratio of torque to angle. Values derived from the FEA results are compared with values derived from experimental measurements in Table I. The "cable-driven bending" experiments and the method for determining the associated stiffness are described in Section IV-B4. In the case of the "0.2-N transverse load" experiment, a transverse load was applied to the tip of an isolated manipulator backbone without cables, the tip deflection was measured using calibrated image processing, and then the model of (10) was used to find a value of  $k_1$  such that the model tip deflection matched the observed. The discrepancy between the two experimental cases may be due in part to differences in dimension and material properties between the two different manipulator backbones used in the test. The discrepancy between FEA results and between FEA and experiment may be due in part to inaccurate material properties and excessive deformations. Improving the accuracy of these estimates is a topic for future work.

## IV. Modeling Shape at High Strain and Nonnegligible Cable Friction

The model presented in Section III works well for small loads and deflections, but it is not accurate at higher deflections. This section includes cable friction in a model that more accurately predicts behavior at high deflections. The problem can be divided into two parts. First, we want to find a way to represent the shape of the manipulator with a small number of parameters. Second, we want to relate the shape parameters to quantities that are easy to measure, namely cable length and cable tension.

A representative dataset underlying the construction of this model is shown in Fig. 11. For this experiment, a light counter-tension was maintained on the right cable, while tension on the left cable was gradually increased and then decreased. An automated actuation system (for additional details see [8]) was used to control the manipulator during the test. The manipulator was moved slowly between sample points, and at rest when the sample data were taken. The typical time between sample points was about five seconds. At each sample point an image of the manipulator was also recorded. The cable tension was recorded with load cells (Honeywell 060-1426-04) that have an accuracy of  $\pm 0.11$ N. The cable displacement was controlled with a linear actuator (Haydon Kerk 35H4A-12) with a resolution of 0.0079mm per motion step.

The data from eight different experiments are uploaded as supplemental material to this paper. Fig. 11 shows results of one such test. Tension and displacement for the left cable is plotted for each sample point. Fig. 12 shows one example configuration during the test. This image shows the manipulator from underneath, so the "left" cable is actually on the right side of the image, and the corresponding angles  $\theta_i$  for the configuration shown are negative.

There are three regimes of bending during the experiment. In regime I the cable is tightening, and the manipulator bends in a predictable manner. This corresponds to a roughly linear increase in cable displacement (although the slope will increase as notches begin to close at high deflections, and there can be a small discontinuity if the manipulator passes

through  $\theta = 0$ ). In regime II, the tension in the left cable is decreasing, but the cable length is remaining close to constant and the overall shape of the manipulator is not changing dramatically. There is a sudden transition from regime II to III, when the cable to starts to move a substantial amount and the manipulator returns to its equilibrium shape. This behavior is repeatable over many tests and follows a similar pattern for different values of counter-tension, as can be seen by viewing the supplemental information.

#### A. Parameterizing Manipulator Shape

Each image taken during a test is processed by manually digitizing the location of corner points on the notches of the manipulator body. The location of these "notch points" is used as a baseline ground truth for the shape of the manipulator. The maximum error in point location using the manual digitization process is 0.15 mm [10]. Fifty-eight notch points ( $Q_i$ ) are selected, numbered sequentially from the lower left, as seen in Fig. 12, and proceeding anticlockwise around the manipulator (for additional details see [8], [10]). A variety of parameterized shape models can then be fit to the 58 selected points (additional shape models are presented in [10]).

A constant curvature assumption is not adequate to describe manipulator shape at high deflection, as shown in Fig. 12. Virtual notch points  $(\tilde{Q}_i)$  generated using the kinematic model described in Section II are overlaid on actual images of the manipulator. An error metric  $||Q_i - \tilde{Q}_i||$  between the model and the manually digitized points is defined by the sum of Euclidean distance from each model point to the actual location of its corresponding point. A "best fit" set of  $\theta_i$  can be defined by minimizing this error metric.

Fig. 12 shows the manipulator for the last point in regime I. While the constant curvature model is fairly close, a "linear model" given by  $\theta_j = a + bj$  provides a much closer agreement. Fig. 13 shows the pointwise error for all 58 notch points and all 18 samples in regime I, for both the constant curvature and linear model.

Fig. 14 shows pointwise error for the linear model in bending regimes II and III. The arrows in Figs. 13 and 14 indicate the progression of time through the sample set. Over the complete range of configurations, the maximum error using the linear fitting is 0.35 mm, compared to 3.36mmfor the constant curvature fitting.

In regime I, the manipulator starts with close to constant curvature. As tension is increased, the joints at the base begin to deflect more than those toward the tip. In regime II, as tension is released, the manipulator remains deflected overall, but the reduction in *b* shows that the shape is becoming more circular. The base joints, which are deflected most and therefore experiencing the highest force from the flexure joints, move before the other joints that are toward the tip. In regime III the manipulator becomes almost circular as the last of the tension in the cable is released.

#### **B. Physics-Based Model of Bending With Friction**

Fig. 15 shows the best-fit shape parameters a and b for the manipulator at each tension during the bending-unbending test shown in Fig. 11. This section presents a simple physical

model that captures this behavior. The manipulator is modeled as a series of rigid segments connected by pin joints equipped with torsion springs.

First order dynamics (i.e., overdamped motion) is assumed because we are primarily interested in the equilibrium state of the manipulator. The tension cable and backbone segment are lumped together, with the cable tension propagation dependent on a simple state machine model that represents sliding and static friction between cable and segment. Relative to a given segment, the cable has three states: sliding toward base, sliding toward tip, sticking. Fig. 16 shows the assumed geometry of the cables. Importantly, each cable is parallel to the link opposite it. This assumption is consistent with the geometry discussed in Section II-B. The dynamic modeling algorithm has four steps. Initially, the manipulator is placed at  $\theta_j = 0$ , all tension values are set to zero, and at each segment the cable is sticking. The simulation begins when tension is applied to the first two segments.

1) Step 1—Tension Propagation—The normal force acting on the cable at each segment is

$$N_j = T_{j,\text{out}} \sin(\theta_{j+1}) + T_{j,\text{in}} \sin(\theta_j).$$

For each segment, progressing from the base (j = 1) to the tip (j = 27) the tension is propagated as

$$T_{j,\text{out}} = \begin{cases} T_{j,\text{in}} - \mu_d |N_j|, & \text{cable sliding toward base} \\ T_{j,\text{in}} + \mu_d |N_j|, & \text{cable sliding toward tip} \\ & \text{unchanged}, & \text{cable sticking} \end{cases}$$

where  $\mu_d$  is the coefficient of sliding friction between cable and backbone segment.

**2)** Step 2—First Order Dynamics—Summing the dominant forces and moments on a segment provides the first order dynamic model

$$\dot{\theta}_{j} = \begin{cases} (r_{x}(T_{26,\text{out}} - T_{25,\text{out}}) - k(\theta_{27})\theta_{27})/b, & j = 27\\ (r_{x}(T_{j,\text{out}} - T_{j,\text{in}}) + \tau_{js})/b, & j \text{ even}\\ (r_{x}(T_{j,\text{in}} - T_{j,\text{out}}) + \tau_{js})/b, & j \text{ odd} \neq 27 \end{cases}$$

where  $\tau_{is} = k(\theta_{i+1})\theta_{i+1} - k(\theta_i)\theta_i$  and *b* is a first order damping coefficient.

3) Step 3—Determine Cable Speed—Given the joint angle speeds  $\dot{\theta}_i$  computed in Step 2, the velocity of the cable through each segment is given by

$$v_{cj} = \begin{cases} -r_x \sum_{k=j+1}^{k=27} \dot{\theta}_k, & j \text{ even} \\ r_x \sum_{k=j+1}^{k=27} \dot{\theta}_k, & j \text{ odd} \end{cases}$$

where j = 26 and a positive  $v_{cj}$  indicates the cable is sliding through segment j toward the base.

**4) Step 4—Determine Stick/Slip State**—A simple state machine determines the transitions between three conditions for the cable stick/slip state in each segment. For the *j*th segment, if the present state is "cable sliding toward base" or "cable sliding toward tip" and the following conditions are met:

$$\begin{array}{ll} |v_{cj}| &< v_{\rm thr} \\ |T_{j,{\rm out}} - T_{j,{\rm in}}| &< \mu_s |N_j| \end{array}$$

the new state is "cable sticking". If the present state is "cable sticking" and the following condition is met:

$$T_{j,\text{in}} - T_{j,\text{out}} > \mu_s |N_j|$$

the new state is "cable sliding toward base." If the present state is "cable sticking" and the following condition is met:

$$T_{j,\text{out}} - T_{j,\text{in}} > \mu_s |N_j|$$

the new state is "cable sliding toward tip." In all other cases the cable state is unchanged. The parameter  $v_{cj}$  is a threshold speed below which the cable is assumed to stick, and  $\mu_s$  is the static coefficient of friction,  $\mu_s > \mu_d$ .

Performance of the model is compared with measured data in Figs. 18 and 19. The only input to this model is cable tension. A value of b = 0.03 Nms/rad is used for the damping term (the results are not especially sensitive to this parameter). The friction terms are  $\mu_d = 0.29$  and  $\mu_s = 0.30$ , with  $v_{\text{thr}} = 5 \times 10^{-5}$  m/s. A piecewise linear torsional spring was used, shown in Fig. 17. The spring rate in the linear region where  $|\theta_i| < 0.07$  rad is equal to 0.4829 Nm/rad. The midrange spring rate for  $0.07 \quad |\theta_i| < 0.14$  rad models the "softening" behavior of the superelastic nitinol in high strain. This was necessary for obtaining a good fit to the data in the bending regime.

The model performs well in the bending regime across many different tests using different counter-tensions, as can be seen in the animated plots in the supplemental material. The model captures some of the hysteresis behavior, as seen clearly in Fig. 18, however the actual manipulator exhibits a much stronger hysteresis.

## V. Conclusion

In Section II, a kinematic model of the backbone and cables was presented. The effect of side to side motion of the cable within the guide channel was considered geometrically, and it was shown that the upper and lower bounds of cable length are very closely approximated by a simple coupling matrix  $P^T$  of rank 1. While the cable channel width plays an important role in shape bifurcations at high tension, the effects are not important when the joint angles  $\theta_i$  are all of the same sign.

In Section III, it was shown that, in general, stiffness is dependent on material properties and not cable tension, at least when neglecting cable friction. Cables with nonlinear stiffness characteristics can allow adjustable stiffness, but only in workspace directions that do not lie in the nullspace of the coupling matrix. Neither the stiffness properties of the backbone nor the cable contribute to the manipulator's deviation in shape from constant curvature.

In Section IV, it was shown that the manipulator shape was well described by the relation  $\theta_j = a + bj$ . Experimental data taken on the manipulator during bending and unbending showed three distinct regimes of behavior (I–III). A physics-based model was presented which works well for the bending in regime I. While the model captures some of the hysteresis shown in the actual robot, it underestimates the duration of sticking of the cables. Additional work is required to develop an improved friction model that captures this behavior.

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## Appendix

A detailed treatment of this material is given in [26]. This Appendix reviews the background material and notation that is of most importance for the study presented in this paper. Given an element  $g \in SE(3)$  represented by a 4 × 4 matrix



where *R* is a rotation matrix and *p* is a translation vector, the  $6 \times 6$  adjoint matrix is

$$Ad_g = \left[ \begin{array}{cc} R & \hat{p}R \\ 0 & R \end{array} \right].$$

The "hat"<sup> $\wedge$ </sup> and "vee"  $\vee$  operators are used to rearrange vector and matrix representations. Hat is used to map 3 × 1 vectors  $\omega = (\omega_1, \omega_2, \omega_3)^T$  to 3 × 3 skew matrices as

$$\hat{\omega} {=} \left[ \begin{array}{ccc} 0 & -\omega_3 & \omega_2 \\ \omega_3 & 0 & -\omega_1 \\ -\omega_2 & \omega_1 & 0 \end{array} \right]$$

and  $6 \times 1$  twists  $\xi = (v, \omega)^T$  to  $4 \times 4$  matrix elements of *se*(3) as

$$\hat{\xi} = \left[ \begin{array}{cc} \hat{\omega} & v \\ 0 & 0 \end{array} \right].$$

The vee performs the inverse operations of the hat, i.e.,  $(\hat{\omega})^{\vee} = \omega$  and  $(\hat{\xi})^{\vee} = \xi$ . The spatial manipulator Jacobian is constructed by

$$J_{st}^{s}(\theta) = [\xi_{1} \ \xi_{2}' \ \dots \ \xi_{27}']$$

where

$$\xi_i' = Ad_{h_i}\xi_i$$

and

$$h_i = \prod_{k=1}^{i-1} \exp\left(\hat{\xi}_k \theta_k\right)$$

The body manipulator Jacobian is related to the spatial Jacobian by the adjoint

$$J_{st}^{s}(\theta) = Ad_{g_{st}(\theta)} J_{st}^{b}(\theta).$$
(27)

Suppose the end-effector configuration is  $g_{st} = g(R_{st}, p_{st})$  and we represent a load applied to the manipulator end-effector as two three-element vectors, expressed relative to the world

frame: a force  $f = (f_x, f_y, f_z)^T$  and a torque  $\tau = (\tau_x, \tau_y, \tau_z)^T$  (see Fig. 8). These quantities must be expressed in the end effector frame in order to represent them as a body wrench  $F_t$ 

$$F_t = \left[ \begin{array}{c} R_{st}^T f \\ R_{st}^T \tau \end{array} \right].$$

The spatial wrench  $F_s$  used in (7) is related to the body wrench by the adjoint transpose

$$F_s = \left(Ad_{g_{st}^{-1}}\right)^T F_t \quad (28)$$

where

$$\left(Ad_{g_{st}^{-1}}\right)^{T} = \left[\begin{array}{cc} R_{st} & 0\\ \hat{p}_{st}R_{st} & R_{st} \end{array}\right].$$

(28) then evaluates to

$$F_s = \left[ \begin{array}{c} f \\ \hat{p}_{st} f + \tau \end{array} \right].$$

Importantly, note that neither the body wrench  $F_t$  nor the spatial wrench  $F_s$  are equal to  $(f^T, \tau^T)^T$ .



## Fig. 1.

High strength dexterous manipulator with large open lumen [8] mounted on a LARS surgical robot. The LARS is a 7-DOF robot designed to assist in surgical procedures [12].



#### Fig. 2.

Geometry for determining cable length. Link 0 is the base, and link numbering proceeds sequentially up the manipulator. The links shown in the inset are links 3, 4 and 5 (i = 4). Positive values of  $\theta$  correspond to the configuration shown.



## Fig. 3.

Contour map of the error between the ideal "zero channel width" length of one cable segment  $L_{R\dot{p}}$  and the approximate length  $L_{RS\dot{p}}$  where error  $= L_{Ri} - p + (\theta_i + \theta_{i+1})r_x + 2r_z$ . The error is less than five micrometers over nearly the entire range of operation (nominal segment length is 810 micrometers). Contour units are micrometers.



## Fig. 4.

Lower bound on the length of a cable that is free to move within a channel can be found by summing the minimum possible length of individual segments.  $L_{DMi}$  is the minimum possible length of a cable segment between two sections. The minimum length of a cable segment within a manipulator section is simply  $p + 2r_z$ . A cable segment can move sideways within a channel by a distance of  $\pm (r_{channel} - r_{cable})$ .



## Fig. 5.

Contour map of the worst case error between the simple length model and upper and lower bounds of the high-accuracy length model reported in [9]. The plot shows  $\max(//L_{RSi} - L_{Ri})//, //L_{RSi} - L_{DMi}//)$ . Contour units are micrometers.



## Fig. 6.

Contour map of intersegment cable length using the simplified expression  $L_{RSi} = p - (\theta_i + \theta_{i+1})r_x - 2r_z$ . Compare these values to the errors shown in Figs. 3 and 5. Contour units are micrometers.





Nonlinear dual rate spring used to model flexure elasticity and intersegment collisions.  $k_1 = 0.36$  Nm/rad,  $k_2 = 10$  Nm/rad,  $\theta_{\text{thresh}} = 0.135$  rad.



## Fig. 8.

Path-dependent solutions to the quasi-static dynamic model. In solution A, the moment is applied first, followed by the force. In solution B, the force is applied first and the moment second.





Each cable is modeled as an inelastic ideal cable in series with a spring. The position of the spring ends,  $u_L$  and  $u_R$ , are the control inputs.



## Fig. 10.

FEA results for deformation of a manipulator segment in response to different applied loads. Color scale indicates magnitude of displacement. Left: 10 N axial compression, deformation scale = 100. Center 0.1 Nm moment, deformation scale = 1. Right: 10 N transverse load, deformation scale = 1.



## Fig. 11.

Cable tension and displacement during bending and unbending of the manipulator. A constant counter-tension of 0.45 N is applied to the opposite cable. There are three regimes of operation (I–III).



## Fig. 12.

Image of the actual manipulator bending under high cable tension. Left: Virtual notch points  $(\tilde{Q}_i)$  for the best fit constant curvature fitting  $(\theta_j = \text{cons})$  are shown overlaid in red dots. Right:  $(\tilde{Q}_i)$  using the fitting  $\theta_j = a + bj$ .





Best fit joint angles using constant curvature (left) and linear (right) assumptions, for bending regime I.





Best fit joint angles using linear curvature assumption, for regime II (left) and regime III (right).



## Fig. 15.

Manipulator shape parameters corresponding to the data in Fig. 11. The fitting model is  $\theta_j = a + bj$ . Arrows indicate progression with time. See supplemental material for additional test results.



#### Fig. 16.

Right: Free body diagrams of link segments in a model that considers cable friction. Cable tension *T*, backbone compression force *F*, torsional spring constant *k*, joint angle  $\theta$ , and radius to cable  $r_x$ .









Comparison of measured and predicted backbone joint angles. Top: model prediction. Bottom: best fit to image data using two parameter model  $\theta_j = a + bj$ .





Comparison of measured and predicted shape parameters. Solid lines: measured data. Dashed lines: predicted. See supplemental material for more results.

## TABLE I

## Estimated Joint Stiffness Values

Configuration	k <sub>1</sub> (Nm/rad)
Experiment, 0.2-N transverse load	0.36
Experiment, cable-driven bending	0.48
FEA, 10-N transverse load	0.57
FEA, 0.1-Nm moment	0.77