

Surface Plasmon Resonance Effect in Inverted Perovskite Solar Cells

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This work reports on incorporation of spectrally tuned gold/silica (Au/SiO₂) core/shell nanospheres and nanorods into the inverted perovskite solar cells (PVSC). The band gap of hybrid lead halide iodide (CH₃NH₃PbI₃) can be gradually increased by replacing iodide with increasing amounts of bromide, which can not only offer an appreciate solar radiation window for the surface plasmon resonance effect utilization, but also potentially result in a large open circuit voltage. The introduction of localized surface plasmons in CH₃NH₃PbI_{2.85}Br_{0.15}-based photovoltaic system, which occur in response to electromagnetic radiation, has shown dramatic enhancement of exciton dissociation. The synchronized improvement in photovoltage and photocurrent leads to an inverted CH₃NH₃PbI_{2.85}Br_{0.15} planar PVSC device with power conversion efficiency of 13.7%. The spectral response characterization, time resolved photoluminescence, and transient photovoltage decay measurements highlight the efficient and simple method for perovskite devices.

1. Introduction

Solar cells using hybrid lead halide perovskite as light harvester have evidenced significant progress in the past few years.^[1–4] Currently, the certified efficiency for perovskite solar cells (PVSC) is 20.1% and expected to improve further.^[5] This success is closely associated with the photoelectrical properties of the specific light absorber, such as CH₃NH₃PbX₃ (X = I, Br, Cl), exhibiting an appropriate band gap, small exciton binding energy, and long and balanced ambipolar charge transport.^[6,7] Most efficient perovskite devices rely on the p-i-n heterojunction structure, in which TiO₂ and 2,2',7,7'-tetrakis (*N,N*-di-*p*-methoxyphenylamine)-9,9'-spirobifluorene (spiro-MeOTAD) are

the widely used electron and hole transporting materials, owing to their good optical transparency and well band alignment with respect to CH₃NH₃PbI₃.^[8,9] Recently, an emerged inverted PVSC, has attracted intensive interests, in which the photo-induced holes are collected at the front transparent conductive glass substrate after photoexcitation.^[10–12] Perovskite devices will be certainly benefited from new materials and device configurations, and thus, resolving some obstacles existing in the conventional cells, such as the current–voltage scan hysteresis.^[13,14] Recently, the elimination of photocurrent hysteresis in the inverted CH₃NH₃PbI₃-based solar cells has been demonstrated by different groups.^[15,16]

Poly(3,4-ethylenedioxythiophene) poly(styrene-sulfonate) (PEDOT:PSS) has been used as hole selective contact in inverted

perovskite cells, achieving outstanding power conversion efficiency (PCE) of about 15%.^[17,18] Meanwhile, NiO-based PVSCs have also shown growing attention, due to its high stability and electrical conductivity.^[19–21] For example, a faster charge transfer between NiO and charge transfer materials than that of TiO₂ has been revealed with scanning electrochemical microscopy.^[21] Chen et al. first reported the replacement of PEDOT:PSS with a thin NiO hole selective layer, resulting in a PCE of 7.8%.^[11] We applied a reactive magnetron sputtered NiO ultrathin layer to promote the efficiency up to 9.8%.^[22] Subsequently, Chen et al. proposed a hybrid interfacial layer with an important “dual blocking effect,” resulting in PCE of ≈13.5%.^[23] Recently, CH₃NH₃PbI₃-based solar cell with 15.4% PCE was demonstrated by using solution-processed Cu-doped NiO hole selective interlayer.^[24] The outstanding performance was attributed to an improved electrical conductivity of NiO. An applicability of Cu:NiO_x with large band gap (*E*_g) perovskite [CH₃NH₃PbI_{3–x}Br_x] solar cell has also been pointed out in this work. The PVSC devices using CH₃NH₃PbI_{3–x}Br_x with large *E*_g encounter less potential loss and exhibit high photovoltage of 1.13–1.16 V. However, a relative low photocurrent of 8–12 mA cm^{–2} was obtained due to a less utilization of photons in the wavelength range from 670 to 780 nm. More recently, a remarkable efficiency of 17.3% was achieved on devices using pulsed laser deposited NiO.^[25] The (111)-oriented nanostructured NiO film with good optical transparency plays a key role in the efficient extraction of holes and the prevention of electron leakage.

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DOI: 10.1002/adv.201500312

Optical manipulation has been successfully explored to improve light utilization efficiency in various photovoltaic devices, such as using folded device architecture, aperiodic dielectric stack, diffraction grating, and plasmon resonant metallic nanostructure.^[26–29] Among them, the introduction of metallic nanostructures as localized surface plasmon resonance (LSPR) into photovoltaic devices has been regarded as the most efficient and simple approach.^[30] Actually incorporation of Au@SiO₂ core-shell nanoparticles into PVSCs has first reported in 2013.^[31] The authors showed the evidence for the reduced exciton binding energy in the perovskite absorber through photoluminescence study. It is known that the LSPR properties originate from collective oscillation of their electrons in response to optical excitation. However, in the previous report the photon absorption of metallic nanostructures can be negligible when considering the broad visible absorption spectrum and high adsorption coefficient of CH₃NH₃PbI₃ active layer in devices.^[31] Therefore, the application of LSPR effect in perovskite devices has been long underestimated in this community. In line with this, the solar cells need to be structured so that light remains trapped inside to increase the photocurrent generation.

In this work, we report on an inverted CH₃NH₃PbI_{3-x}Br_x-based PVSC employing a hybrid interfacial layer of “compact NiO_x/meso-Al₂O₃” in combination with nanostructured Au nanoparticles. The CH₃NH₃PbI_{3-x}Br_x absorbers with valence band (VB) ≈5.4 eV are specially selected to provide an optimized optical utilization, and well energy alignment with that of NiO to minimize energy loss. Most importantly, the optical band gap E_g of CH₃NH₃PbI_{3-x}Br_x can be intentionally tuned wide enough, leaving a suitable spectrum window for the Au nanostructures. Therefore, an expected LSPR excitation in the range of 740–860 nm can be promoted by the resonant interaction between electromagnetic field of incident photons and surface charge oscillation of the Au nanorods (NRs). It should be noted that in the wavelength range, the CH₃NH₃PbI_{3-x}Br_x active layer does not absorb any photons. Therefore, we could intentionally investigate SPR effect for efficiency improvement of inverted PVSCs using Au nanoparticles. By imbedding a thin layer of Au NRs in the active layer of the PVSCs, the PCE was improved by a maximum value of 13.7%, which is attributed to enhanced local electric field originate from SPR effect.

2. Results and Discussion

Figure 1a depicts the typical inverted PVSC device configuration in this study. To incorporate metal nanoparticles into the perovskite active layer, ethanol solution containing Au@SiO₂ was added to the Al₂O₃ colloid solution at a range of concentrations prior to porous alumina film deposition. The Au@SiO₂ core-shell metal-dielectric nanoparticles were prepared with a three-step synthesis process described elsewhere.^[32] X-ray diffraction (XRD) analysis in Figure S1 (Supporting Information) exhibits a common pattern of bulk gold as well as a broad peak at around 25° corresponding to amorphous silica. Au@SiO₂ nanoparticles with different AR can be obtained by adjusting either Au seeds or the content of the growth solution.^[33] The corresponding absorption characterization of these Au nanostructures is shown in Figure S2 (Supporting Information). Herein, Au@SiO₂ nanorods with AR of 3.8 were chosen on account of its most red-shift absorption peak (≈785 nm). As depicted in Figure 1b, the designed optical manipulation will take full advantage of the incident photons when CH₃NH₃PbI_{3-x}Br_x absorber and Au@SiO₂ NRs are assembled in such a device. As most of photons in the overlap region could be cut off by light absorbers, a CH₃NH₃PbI_{3-x}Br_x perovskite with onset of absorption band from 786 nm (1.58 eV) to 704 nm (1.78 eV) has been synthesized and tested in this study. The light harvesting in the Au@SiO₂ device is indistinguishable from the one without Au@SiO₂ nanostructures due to a low loading of Au@SiO₂ NRs (2.0 wt%) in the Al₂O₃ film.

The CH₃NH₃PbI_{3-x}Br_x absorber was deposited by spun stacking of double layers of PbI₂ and CH₃NH₃I+CH₃NH₃Br as illustrated in Figure 2a. Details of the perovskite film preparation procedure can be found in the Experimental Section. Figure 2b shows the scanning electron microscopy (SEM) surface image of as-prepared CH₃NH₃PbI_{3-x}Br_x film on the compact-NiO/meso-Al₂O₃ layer. The grain size increases fast to ≈1 μm in a 25 min annealing process. It has been reported that the big grain size CH₃NH₃PbI_{3-x}Br_x has less inner grain boundaries, which is benefited for reduction of charge recombination.^[34,35] The EDX mapping result shows that “Au,” “Br,” and “I” elements are homogeneously distributed in the meso-Al₂O₃ layer (Figure 2b), reflecting a good evidence of the existence of Au NPs and CH₃NH₃PbI_{3-x}Br_x materials.

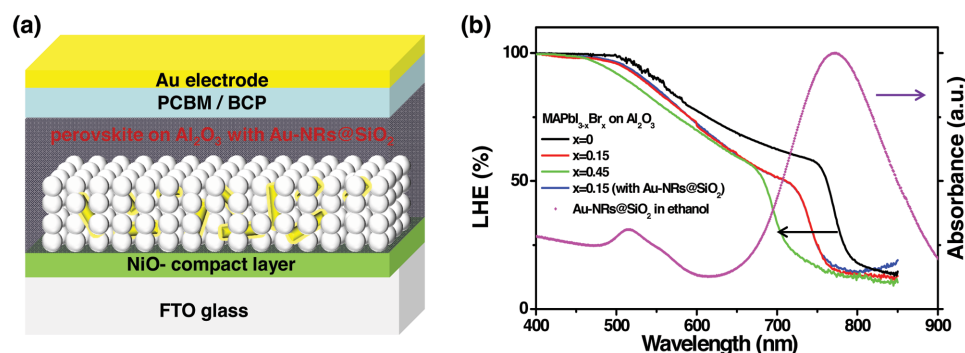


Figure 1. a) Illustration of device structure labeled with different components. b) Light harvesting efficiency (LHE) of CH₃NH₃PbI_{3-x}Br_x absorber (with different Br content) deposited on the same thickness of compact-NiO/meso-Al₂O₃ layer and UV–visible spectroscopy in ethanol for Au@SiO₂ NRs with AR of 3.8. The LHE of CH₃NH₃PbI_{3-x}Br_x coated compact-NiO/meso-Al₂O₃ layer with Au@SiO₂ NRs is also presented.

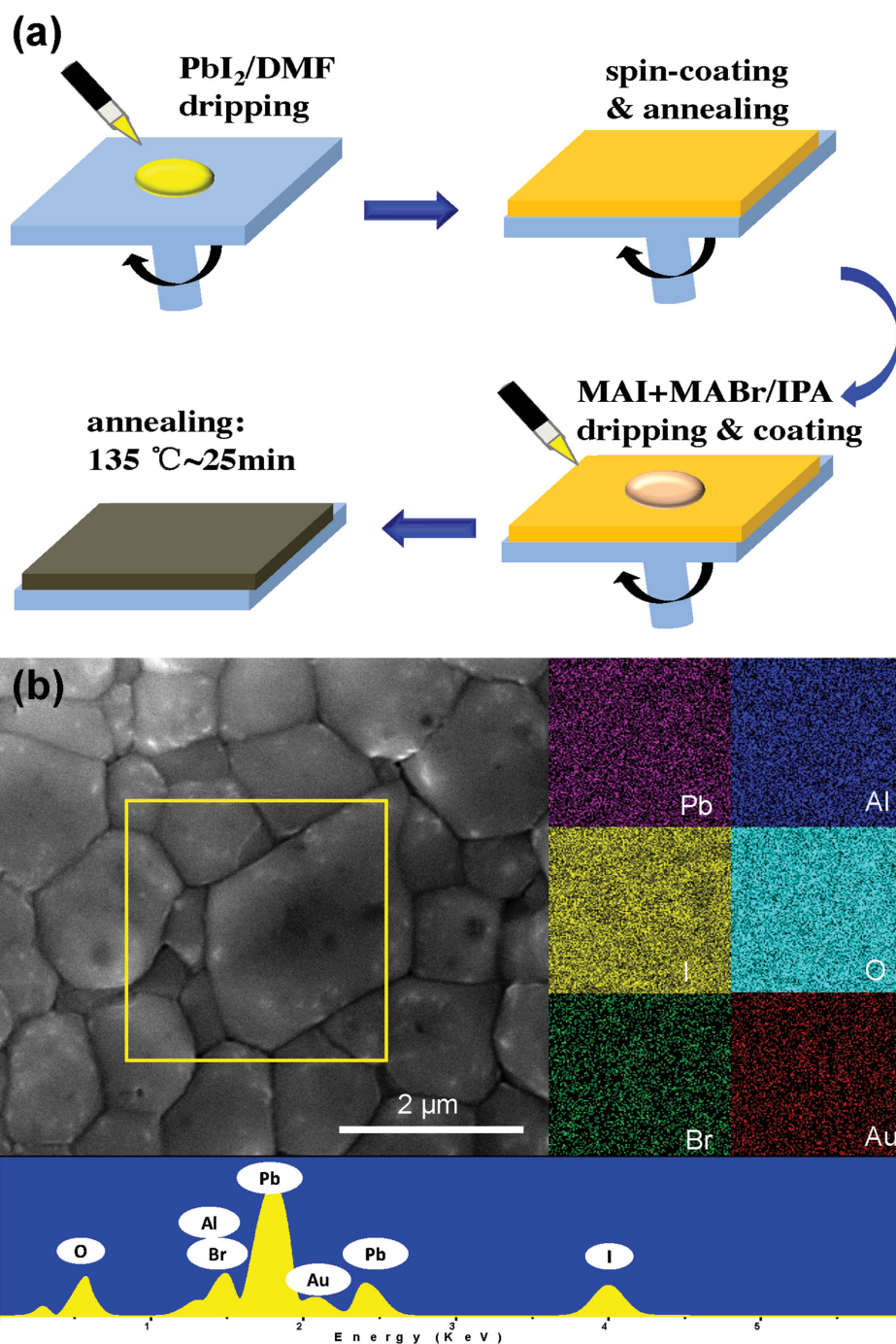


Figure 2. a) Schematics of the preparation approach to $\text{CH}_3\text{NH}_3\text{PbI}_{3-x}\text{Br}_x$ perovskite film. b) Top view SEM image of $\text{CH}_3\text{NH}_3\text{PbI}_{3-x}\text{Br}_x$ coated meso- Al_2O_3 film with Au-NRs and EDX mapping results of the yellow square region.

The evaluation of validity and rationality of our strategy on optical manipulation was first carried out with consideration of (1) Au NRs with various AR values, and (2) $\text{CH}_3\text{NH}_3\text{PbI}_{3-x}\text{Br}_x$ with different Br contents. The absorption spectra of $\text{CH}_3\text{NH}_3\text{PbI}_{3-x}\text{Br}_x$ absorber can be controlled through tuning $\text{CH}_3\text{NH}_3\text{I}/\text{CH}_3\text{NH}_3\text{Br}$ molar ratio during the two-step spin-coating process. The used Au nanospheres with diameter of 41 nm (AR = 1, Figure 3a) and Au-NRs of about

42 nm length and 11 nm width (AR = 3.8, Figure 3b) were coated with an approximate 1–2 nm SiO_2 shell. The absorption peaks are observed at ≈ 541 and ≈ 785 nm for the two Au nanostructure samples in ethanol (Figure S2, Supporting Information), respectively. After adding Au-nanosphere (AR = 1) into the meso- Al_2O_3 film, the photocurrent was increased by 12% from 13.6 to 15.3 mA cm^{-2} for the $\text{CH}_3\text{NH}_3\text{PbI}_{3-x}\text{Br}_x$ ($x = 0.15$)-based perovskite device. Keeping the same concentration of the

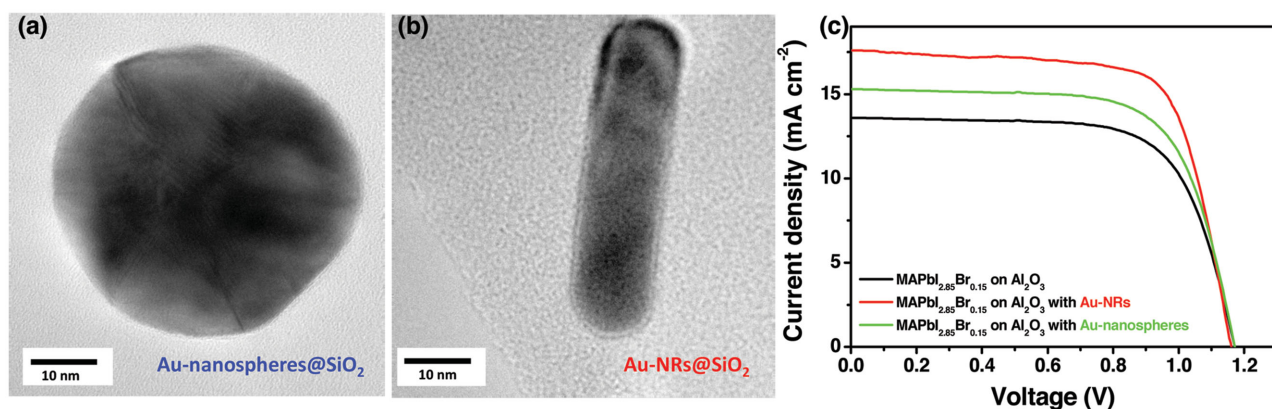


Figure 3. TEM images of the synthesized core-shell SiO_2 coated a) Au-nanospheres@ SiO_2 (AR = 1) and b) Au@ SiO_2 NRs (AR = 3.8). c) Photocurrent–voltage characterization of $\text{CH}_3\text{NH}_3\text{PbI}_{2.85}\text{Br}_{0.15}$ -based PVSCs with and without Au nanostructures.

Au additives, the device with Au-NRs (AR = 3.8) showed even higher photocurrent of 17.5 mA cm^{-2} (Figure 3c). Compared to that of Au-NRs, the J_{sc} enhancement of perovskite devices with Au-nanospheres has been suppressed, mainly due to an optical trapping competition raised from high extinction coefficient $\text{CH}_3\text{NH}_3\text{PbI}_{3-x}\text{Br}_x$ ($x = 0.15$) absorber around 541 nm (Figure 1b).

The photovoltaic performance for devices using compact-NiO/meso- Al_2O_3 with/without Au@ SiO_2 NRs (AR = 3.8)/ $\text{CH}_3\text{NH}_3\text{PbI}_{3-x}\text{Br}_x$ / [6,6]-phenyl C_{61} -butyric acid methyl ester (PCBM)/bathocuproine (BCP)/Au configuration by varying Br content in perovskite was further investigated (Figure S3, Supporting Information). Indeed, as expected, the perovskite devices in presence of Br component exhibited higher V_{OC} than the pristine $\text{CH}_3\text{NH}_3\text{PbI}_3$ -based counterparts (Figure S3b, Supporting Information), confirming superior energy level alignment between $\text{CH}_3\text{NH}_3\text{PbI}_{3-x}\text{Br}_x$ and NiO. The photocurrent of PVSC devices with adding Au@ SiO_2 NRs was higher than the control devices (Figure S3a, Supporting Information), which originates from the reutilization of photon in longer wavelength region around 785 nm. Increasing the Br portion in the $\text{CH}_3\text{NH}_3\text{PbI}_{3-x}\text{Br}_x$ decreases the device's photocurrent. The decrease of photocurrent in the Br portion

of $x = 0-0.45$ was accompanied by an augmentation of the V_{OC} compensating the decrease in the fill factor. As a result, the overall power conversion efficiency reaches a maximum value of 13.5% at $x = 0.15$ over 15 piece of devices. The statistical data for device performance are shown in Figure S3a (Supporting Information). Therefore, the $\text{CH}_3\text{NH}_3\text{PbI}_{2.85}\text{Br}_{0.15}$ absorber in combination with Au@ SiO_2 NRs (AR = 3.8) takes the most advantage of the optical manipulation design. The work function analysis by photoelectron spectroscopy for the spin-coated $\text{CH}_3\text{NH}_3\text{PbI}_{3-x}\text{Br}_x$ perovskite film in line with the cell energy levels are depicted in Figure S4 (Supporting Information).

Some photovoltaic experiments were conducted to evaluate the performance of devices by varying the concentration of Au nanoparticles. It was found that the Au NRs concentration played a vital role on the devices performance. Figure S5 (Supporting Information) presents the photovoltaic parameters for the $\text{CH}_3\text{NH}_3\text{PbI}_{2.85}\text{Br}_{0.15}$ -based perovskite device as the concentration of Au@ SiO_2 NRs (AR = 3.8) varies from 0 to 4.2 wt%. The introduction of Au@ SiO_2 NRs results in an increased photocurrent, reaching a maximum value of 17.4 mA cm^{-2} at 2.0 wt% of Au@ SiO_2 NRs. The V_{OC} is remarkably kept similar for various devices (Figure S5b, Supporting Information). Figure 4a compares the photocurrent density–voltage ($J-V$)

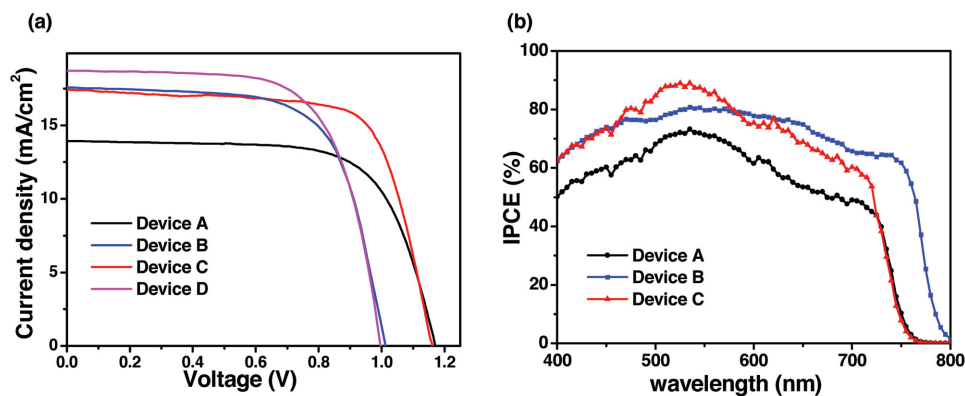


Figure 4. a) Representative $J-V$ curves for devices using Al_2O_3 -only (devices A and B) and Al_2O_3 film incorporated with Au@ SiO_2 NRs (devices C and D) with $\text{MAPbI}_{2.85}\text{Br}_{0.15}$ (devices A and C) or MAPbI_3 (devices B and D) absorbers measured under AM1.5 simulated sunlight (100 mW cm^{-2} irradiance). b) IPCE spectra of corresponding devices. MA: CH_3NH_3^+ cation.

Table 1. Photovoltaic parameters of the compact-NiO/meso-Al₂O₃/CH₃NH₃PbI_{3-x}Br_x ($x = 0$ and $x = 0.15$)/PCBM/BCP/Au configuration inverted PVSCs, with/without 2.0 wt% incorporation of Au@SiO₂ NPs.

Device architecture	V_{oc} [V]	J_{sc} [mA cm ⁻²]	FF	PCE [%]	
Device A	CH ₃ NH ₃ PbI _{2.85} Br _{0.15} on Al ₂ O ₃	1.17	13.9	0.66	10.7
Device B	CH ₃ NH ₃ PbI ₃ on Al ₂ O ₃	1.01	17.2	0.65	11.3
Device C	CH ₃ NH ₃ PbI _{2.85} Br _{0.15} on Al ₂ O ₃ with Au@SiO ₂ NRs	1.16	17.4	0.68	13.7
Device D	CH ₃ NH ₃ PbI ₃ on Al ₂ O ₃ with Au@SiO ₂ NRs	0.99	18.7	0.66	12.2
Device E	CH ₃ NH ₃ PbI _{2.85} Br _{0.15} on Al ₂ O ₃ with Au@SiO ₂ nanospheres	1.16	15.3	0.65	11.5

curves of the CH₃NH₃PbI₃ or CH₃NH₃PbI_{2.85}Br_{0.15} based PVSC devices W/O 2.0 wt% incorporation of Au@SiO₂ NRs within the same thickness meso-Al₂O₃ film. The photovoltaic parameters are tabulated in Table 1. Device A (with Al₂O₃-only using CH₃NH₃PbI_{2.85}Br_{0.15}) exhibits the PCE of 10.7% with a J_{SC} of 13.9 mA cm⁻², a V_{OC} of 1.17 V, and a fill factor of 0.66. The replacement of CH₃NH₃PbI_{2.85}Br_{0.15} with CH₃NH₃PbI₃ results in a significant increase in J_{SC} , achieving a PCE of 11.3% for device B. The addition of Au@SiO₂ NRs in devices C and D result in a significant increase in J_{SC} . Device C with compact-NiO/meso-Al₂O₃ with Au@SiO₂ NRs/CH₃NH₃PbI_{2.85}Br_{0.15}/PCBM/BCP/Au configuration showed the highest PCE of 13.7% with a V_{OC} of 1.16 V, a J_{SC} of 17.4 mA cm⁻², and a fill factor of 0.68. Devices A and C show higher V_{OC} s compared to devices B and D, indicating CH₃NH₃PbI_{3-x}Br_x has better energy alignment with the deep VB of NiO. Considering the similar configuration in the four devices, the augmented output photovoltage and current could be contributed to the synergy effect from Au@SiO₂ NRs and CH₃NH₃PbI_{2.85}Br_{0.15} absorber. Figure 4b presents the incident photon-to-current conversion efficiency (IPCE) for the corresponding devices. Even though the difference of light harvesting capability can be ignorable for the CH₃NH₃PbI_{2.85}Br_{0.15}-based devices with and without Au NRs (Figure 1b), an enhanced photocurrent was unambiguously observed for the former. We conclude that the photocur-

rent improvement is caused by the LSPR effect of Au@SiO₂ NRs, rather than the scattering effect from metal nanoparticles.

Photoluminescence (PL) measurements on CH₃NH₃PbI_{2.85}Br_{0.15} coated NiO/meso-Al₂O₃ layer with and without the addition of Au@SiO₂ nanoparticles (AR = 1 and 3.8) are performed to further probe the influence of the LSPR effect. The excitation wavelength was set at 435 nm in order to avoid any external electromagnetic radiation on Au nanoparticles (absorption peak at 530 nm for that with AR of 1, and absorption peak of 785 nm with AR of 3.8). Figure 5a shows time-integrated PL spectrum at room temperature, revealing a significant reduction in the signal for the samples incorporating Au NRs over six samples. Figure 5b shows the result of time-resolved PL measurement at the emission peak (760 nm). An accelerated PL quenching was observed for the sample with the Au NRs. The time constant was evaluated to be 10.9 and 12.1 ns for the sample with Au nanoparticles of AR = 3.8 and AR = 1, respectively, indicating both of them play similar role during the charge separation process. A longer time constant of 24.7 ns was observed for the film without Au nanostructures. The enhanced PL quenching could be contributed from ionization of the excitons and enhanced charge separation, which is related to the LSPR effect of Au NRs.^[36,37]

The interfacial charge-recombination in the CH₃NH₃PbI_{2.85}Br_{0.15}-based PVSCs was further investigated with

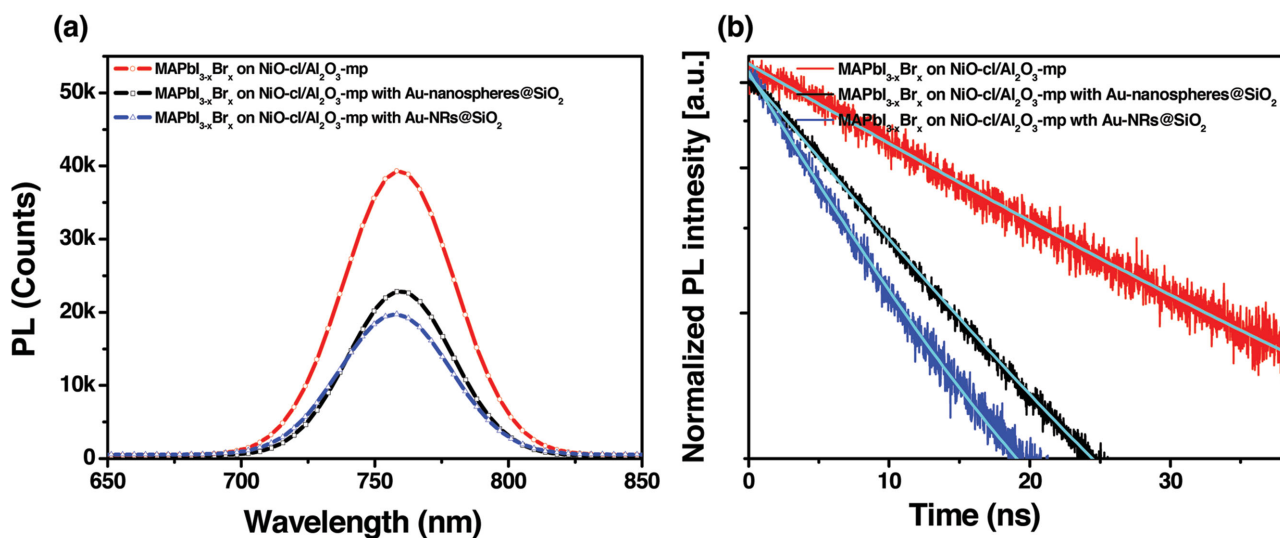


Figure 5. Photoluminescence study: a) Time-integrated spectra and b) time-resolved PL decays (detected at 760 nm) for MAPbI_{3-x}Br_x perovskite coated on meso-Al₂O₃ film with and without incorporating Au-nanospheres@SiO₂ (AR = 1) or Au@SiO₂ NRs (AR = 3.8).

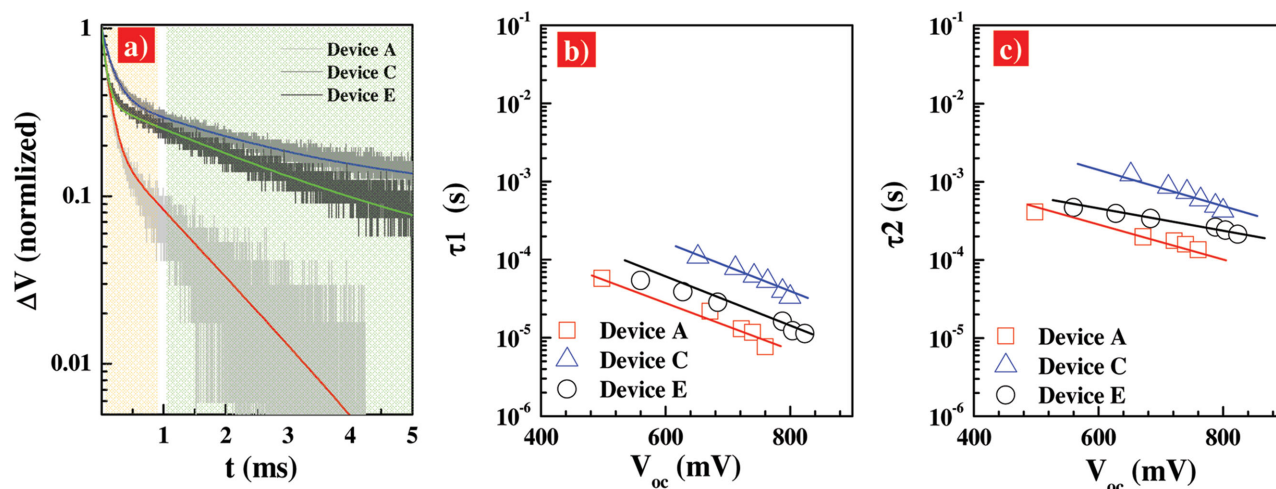


Figure 6. a) Transient photovoltage decay (TPD) data for device A, C in Table 1 and device E with compact-NiO/meso-Al₂O₃/CH₃NH₃PbI_{2.85}Br_{0.15}/PCBM/BCP/Au configuration with 2.0 wt% incorporation of Au-nanospheres@SiO₂, measured under 0.5% sun illumination. And charge recombination lifetime, b) τ_1 and c) τ_2 , as a function of the open-circuit voltage, derived from the double exponential fitting of TPD curves.

transient photovoltage decay (TPD) measurements, which is presented in **Figure 6a**. This technique is widely employed for the determination of charge lifetimes (τ_n) in sensitized solar cells.^[38,39] Figure S6 (Supporting Information) depicts the typical voltage transient dynamics for all devices responding to the light perturbation. The recombination kinetics in these inverted PVSCs are yet to be fully understood, however, we observed biphasic decays with double exponential functions (Figure 6a). The two time constants (τ_1 and τ_2) are suggestive of the presence of two distinguished populations of generated carriers and their recombination independently. Figure 6b,c presents the extracted charge lifetime by fitting photovoltage decay results as a function of the open-circuit voltage for various devices. This result is different with that observed for dye-sensitized solar cells.^[40,41] The lifetime constant in the range of 0.1–10 ms depending on the sensitizers and electrolytes can be ascribed to the interfacial charge recombination between electrons in the dye-covered TiO₂ and holes in the redox mediators. Therefore, for the perovskite devices, the charge population bearing a short lifetime (τ_1 , in the range of microseconds in Figure 6b) might be attributable to the charge confined within the perovskite bulk layer recombining with defect traps or PCBM. And we associate the longer voltage decay component (τ_2 , in the range from 0.1 to 10 ms in Figure 6c) to the electrons recombining with holes in NiO near the perovskite/charge selective layer interface. An increased lifetime was observed when Au nanoparticles were used. This is contrast to the findings of LSPR effect in dye-sensitized solar cell. A shorter interfacial charge recombination lifetime of about milliseconds was usually observed for the dye-sensitized devices when the metallic nanostructures were incorporated.^[41] This observation was explained by a triggered interfacial charge recombination between the metallic nanoparticles and their neighbor mediators.^[42] Previous investigation on dye-sensitized solar cells revealed that the augmented photocurrent could be contributed both from the scatter effect of metallic nanoparticles and/or the boosted charge separation efficiency when gold NRs was used.^[42] As shown in Table 1, there is about 3.5 mA cm⁻² enhancement in the J_{SC} for device C compared with device A.

The charge lifetime for device A is about ten times longer than that of device C. Therefore, this result indicates that the augment in J_{SC} for device A could be contributed to the longer lifetime. It is noted that, herein, the perovskite devices with Au @SiO₂ NRs (AR = 3.8) show the longer charge lifetime than that of devices with Au nanosphere (AR = 1). A long charge lifetime guarantees effective charge collection efficiency, thus the device output photocurrent. This result indicates the aspect ratio of Au nanoparticles could play critical role in the photo-induced enhanced dissociation of excitons. The similar PL quenching of the films with Au nanostructures with various ARs (Figure 5) indicates the same charge generation rate, however, the device C (with Au NRs) shows an order of magnitude higher in charge lifetime than the device E (with Au nanospheres).

3. Conclusion

In summary, we demonstrated the incorporation of Au@SiO₂ nanorod into the inverted CH₃NH₃PbI_{2.85}Br_{0.15} perovskite solar cells could significantly increase the device photocurrent, which could be contributed to the LSPR excitation in the range of 740–860 nm by the resonant interaction. An increased interfacial charge recombination lifetime in line with effective charge collection was further observed in the inverted CH₃NH₃PbI_{2.85}Br_{0.15} perovskite solar cells containing Au nanorods. As a result, a considerably higher PCE of up to 13.7% was achieved. This optical manipulation strategy can be extendable to a broad community of thin film solar cells.

4. Experimental Section

Materials and Sample Preparation: All solvents and reagents, unless otherwise stated, were of analytically pure quality and used as received. PbI₂ (99%), Al₂O₃ nanoparticles dispersion in isopropanol (<50 nm, 20 wt%), nickel acetylacetonate (95%), super dehydrated solvents of dimethyl formamide (DMF), isopropanol, and chlorobenzene were all purchased from Sigma Aldrich. Synthesis process of Au@SiO₂ NPs can be referred to the Supporting Information.

Fabrication of $\text{CH}_3\text{NH}_3\text{PbI}_{3-x}\text{Br}_x$ -Based Perovskite Solar Cells: The ethylammonium lead iodide ($\text{CH}_3\text{NH}_3\text{PbI}_3$) was prepared according to a previous work.^[22] $\text{CH}_3\text{NH}_3\text{PbI}_{3-x}\text{Br}_x$ was deposited with a modified two-step spin-coating method.^[13] First, the p-type selective interlayer, NiO_x was deposited onto the precleaned and prepatterned transparent glass/fluorine-doped tin oxide (FTO) coated glasses (Pilkington TEC 15 $\Omega \square^{-1}$) with a spray pyrolysis method. For transparent meso- Al_2O_3 film preparation, 2.0 g Al_2O_3 nanoparticles dispersion in isopropanol (20 wt%) was added to 12.0 g α -terpinol and 2.0 g ethanol solution of ethyl cellulose (10 wt%), the mixture was then treated in ultrasonic bath for 30 min and string for another 30 min to form a homogenous paste. 5 wt% Au@SiO_2 NRs ethanol solution were mixed with different concentrations (from 0 to 4.2 wt% of Au@SiO_2 NRs/ Al_2O_3). The paste was subsequently spin-coated and sintered at 550 °C for 30 min. A concentration of 460 mg mL⁻¹ of PbI_2 in DMF solution was then spin-coated onto $\text{NiO}/\text{Al}_2\text{O}_3$ layer (3000 rpm for 30 s). MAI and MABr (MA: CH_3NH_3^+) were mixed with different MAI/MABr molar ratio in 2-propanol at 0.3 M and then spin-coated on dry PbI_2 layer at room temperature. The formation of continuous, compact $\text{CH}_3\text{NH}_3\text{PbI}_{3-x}\text{Br}_x$ perovskite films will be completed after annealing at 100 °C for 25 min. The surface morphology of $\text{CH}_3\text{NH}_3\text{PbI}_{3-x}\text{Br}_x$ film was characterized by the cross section images of SEM. The PCBM (≈ 60 nm) and BCP (< 10 nm) were spin-coated inside an argon-filled glove box. Au layer (65 nm) was thermally deposited on the substrate inside a vacuum chamber (10^{-6} Torr). The active area of the device is 0.16 cm².

Characterization: A xenon light source solar simulator (450 W, Oriel, model 9119) with AM 1.5G filter (Oriel, model 91192) was used to give an irradiance of 100 mW cm⁻² at the surface of the solar cells. Various irradiance intensities from 0.01 to 1.0 sun can be provided with neutral wire mesh 50 attenuators, and the light intensity was calibrated with a standard silicon solar cell. The current–voltage characteristics of the devices under these conditions were obtained by applying external potential bias to the devices and measuring the generated photocurrent with a Keithley model 2400 digital source meter. The J – V characteristics were recorded by reverse scan or forward scan with a scan rate of 50 mV s⁻¹. A similar data acquisition system was used to control the IPCE measurement. SEM images were obtained using FEI Nova NanoSEM 450. XRD results were acquired using Phillips X'Pert PRO. The time-resolved and steady-state PL measurements were recorded with Edinburgh instruments (FLSP920 spectrometers). The excitation light source was a femtosecond laser centered at 435 nm, operated at a power of 10 mW.

Transient photovoltage decay measurements were performed on all the cells using a ring of red LED (Lumiled) controlled by a fast solid-state switch. The pulse widths were 2 ms. An array of InGaN diodes (Lumiled) supplied the white bias light, transients were measured at different white light intensities via tuning the voltage applied to the bias diodes. The voltage output was recorded on an oscilloscope directly connected with the cells.

Supporting Information

Supporting Information is available from the Wiley Online Library or from the author.

Acknowledgements

Financial support from the Director Fund of the Wuhan National Laboratory for Optoelectronics (WNLO), the 973 Program of China (Grant Nos. 2014CB643506 and 2013CB922104) is gratefully acknowledged. The authors thank the Analytical and Testing Centre at the Huazhong University of Science and Technology (HUST) for performing characterization of various samples.

Received: September 17, 2015

Published online: January 21, 2016

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