

## High-Pressure NiAs-Type Modification of FeN

William P. Clark, Simon Steinberg, Richard Dronskowski, Catherine McCammon, Ilya Kuperenko, Maxim Bykov, Leonid Dubrovinsky, Lev G. Akselrud, Ulrich Schwarz, and Rainer Niewa\*

**Abstract:** The combination of laser-heated diamond anvil cells and synchrotron Mössbauer source spectroscopy were used to investigate high-temperature high-pressure chemical reactions of iron and iron nitride  $\text{Fe}_2\text{N}$  with nitrogen. At pressures between 10 and 45 GPa, significant magnetic hyperfine splitting indicated compound formation after annealing at 1300 K. Subsequent *in situ* X-ray diffraction reveals a new modification of FeN with NiAs-type crystal structure, as also rationalized by first-principles total-energy and chemical-bonding studies.

Since the pioneering research into binary nitrides,<sup>[1]</sup> this field of chemistry has grown into a rich diversity of sub-classes.<sup>[2–9]</sup> One of the families that has been investigated more intensely is that of iron nitrides. Several different phases are already known,<sup>[10]</sup> which range from the most iron-rich phases, such as  $\alpha''\text{-Fe}_{16}\text{N}_2$  or  $\alpha'\text{-Fe}_8\text{N}$  through  $\gamma'\text{-Fe}_4\text{N}$ ,  $\varepsilon\text{-Fe}_3\text{N}_{1\pm x}$  and  $\zeta\text{-Fe}_2\text{N}$  to  $\gamma''\text{-FeN}$ ,<sup>[11]</sup> which represents the most nitrogen-rich phase to date. However, in 2011, the existence of nitrogen-richer FeN<sub>2</sub> was predicted.<sup>[12]</sup> Iron nitrides are used as coatings for steel-based materials combining hardening with protection against corrosion. More recent studies have shown that iron nitrides can be used as catalysts for the production of hydrocarbons<sup>[13]</sup> and as a potential treatment of cancer cells.<sup>[14]</sup> Moreover, iron nitrides are of relevance for the geosciences as they may be an important constituent of the Earth's core. Since the interaction of iron and nitrogen at higher pressures has seen little

attention, there remains much to be investigated with regard to properties and phase relations.<sup>[15–17]</sup>

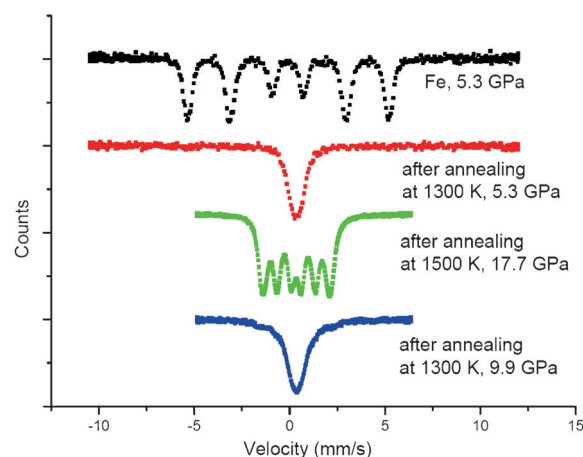
In 1993, the much debated cubic FeN was reported.<sup>[11,18–20]</sup> The published data supports either the rock salt structure,<sup>[21–24]</sup> or the zinc blende arrangement.<sup>[25,26]</sup> An independent prediction stated that cubic FeN would actually prefer the zinc blende type, while  $\text{FeN}_x$  ( $x = 0.5–0.7$ ) would adopt a rock salt arrangement,<sup>[25]</sup> but results remain controversial. Single-phase thin films of  $\gamma''\text{-FeN}$  adopt a zinc blende pattern.<sup>[20]</sup> Herein, we report on synthesis, magnetic properties, and electronic structure of NiAs-type bulk FeN. For synthesis at elevated temperatures and pressures, we chose the diamond anvil cell technique combined with laser heating. Phase transformations were probed with synchrotron source Mössbauer spectroscopy which serves as an ideal tool to investigate electronic configuration and chemical state of iron.<sup>[27–29]</sup>

Two different types of experiments were carried out, one using  $^{57}\text{Fe}$  and the other using  $\zeta\text{-}^{57}\text{Fe}_2\text{N}$ , both with  $\text{N}_2$  as reactant and pressure medium. Samples heated up to 1300 K at pressures below 10 GPa form an iron nitride with a composition close to  $\zeta\text{-Fe}_2\text{N}/\varepsilon\text{-Fe}_3\text{N}_{1.4}$ .<sup>[30]</sup> At pressures higher than 10 GPa, a new phase was observed after annealing, which gave moderate magnetic hyperfine splitting, as shown in Figure 1 (green curve). The magnetic splitting remains essentially up to 45 GPa, suggesting a quite robust

[\*] M. Sc. W. P. Clark, Prof. Dr. R. Niewa  
Institut für Anorganische Chemie, Universität Stuttgart  
Pfaffenwaldring 55, 70569 Stuttgart (Germany)  
E-mail: rainer.niewa@iac.uni-stuttgart.de  
Dr. S. Steinberg, Prof. Dr. R. Dronskowski  
Institut für Anorganische Chemie  
RWTH Aachen (Germany)  
Dr. C. McCammon, Dr. M. Bykov, Prof. Dr. L. Dubrovinsky  
Bayerisches Geoinstitut  
Universität Bayreuth (Germany)  
Dr. I. Kuperenko  
ESRF, Grenoble (France)  
Dr. L. G. Akselrud, Priv.-Doz. Dr. U. Schwarz  
Max-Planck-Institut für Chemische Physik fester Stoffe  
Dresden (Germany)

Supporting information for this article can be found under:  
<https://doi.org/10.1002/anie.201702440>.

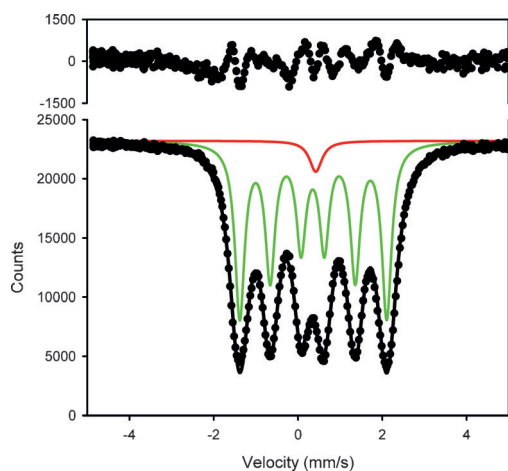
© 2016 The Authors. Published by Wiley-VCH Verlag GmbH & Co. KGaA. This is an open access article under the terms of the Creative Commons Attribution Non-Commercial License, which permits use, distribution and reproduction in any medium, provided the original work is properly cited, and is not used for commercial purposes.



**Figure 1.**  $^{57}\text{Fe}$ -Mössbauer spectra at various pressures before and after heat treatment of iron–nitrogen mixtures. Black: Initial  $^{57}\text{Fe}$  in  $\text{N}_2$  pressure medium at 5.3 GPa prior to laser heating. Red: After laser heating at 5.3 GPa the spectrum indicates formation of  $\zeta\text{-}^{57}\text{Fe}_2\text{N}/\varepsilon\text{-}^{57}\text{Fe}_3\text{N}_{1+x}$ . Green: Multiplet of the new phase after laser heating. Blue: Spectrum indicating the reformation of  $\zeta\text{-}^{57}\text{Fe}_2\text{N}/\varepsilon\text{-}^{57}\text{Fe}_3\text{N}_{1+x}$  after laser heating at pressures below 10 GPa.

magnetic order. Upon heating at a pressure below 10 GPa, the new phase transforms back into  $\zeta$ -Fe<sub>2</sub>N/ $\epsilon$ -Fe<sub>3</sub>N<sub>1.4</sub>.

There are a distinct major and a minor component observed by comparing the <sup>57</sup>Fe-Mössbauer spectra from all samples annealed above 1000 K and 10 GPa (Figure 2 and



**Figure 2.** <sup>57</sup>Fe-Mössbauer spectra of  $\zeta$ -<sup>57</sup>Fe<sub>2</sub>N in N<sub>2</sub> pressure medium, after annealing, at 17.7 GPa.

Figure S1 in the Supporting Information). The major component of all spectra consists of a magnetic sextet, while the minor component represents a non-magnetic singlet (Table 1). After annealing, magnetic components have similar

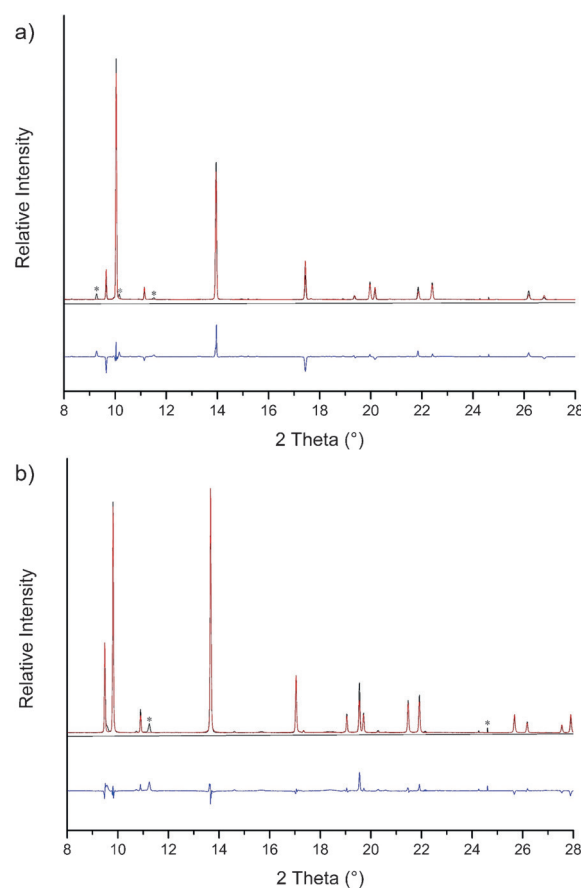
**Table 1:** Refined <sup>57</sup>Fe-Mössbauer spectra parameters before annealing and of the major component after annealing.[a]

Sample	$\delta$ [mm s <sup>-1</sup> ]	$\Delta E_Q$ [mm s <sup>-1</sup> ]	$\Delta E_M$ (T)
<sup>57</sup> Fe in N <sub>2</sub> (0.1 MPa, not annealed)	$-0.073 \pm 0.005$	$0.011 \pm 0.009$	$32.645 \pm 0.037$
<sup>57</sup> Fe in N <sub>2</sub> (11.9 GPa, annealed)	$0.305 \pm 0.012$	$0.015 \pm 0.019$	$11.553 \pm 0.063$
$\zeta$ - <sup>57</sup> Fe <sub>2</sub> N in N <sub>2</sub> (0.1 MPa, not annealed)	$0.441 \pm 0.023$	–	–
$\zeta$ - <sup>57</sup> Fe <sub>2</sub> N in N <sub>2</sub> (17.7 GPa, annealed)	$0.356 \pm 0.002$	$0.012 \pm 0.003$	$10.814 \pm 0.014$

[a]  $\delta$  represents the center shift,  $\Delta E_Q$  the quadrupole splitting, and  $\Delta E_M$  the hyperfine field. Spectra were fit using MossA software.<sup>[32]</sup>

hyperfine parameters: a small magnetic hyperfine field (ca. 11 T) and a moderate center shift (ca. 0.35 mm s<sup>-1</sup>), suggesting the presence of a new phase. The minor component singlet with a center shift of 0.4 mm s<sup>-1</sup> corresponds to the center shift of  $\zeta$ -Fe<sub>2</sub>N.<sup>[30]</sup>

X-ray diffraction patterns were obtained at pressures of 13.3 GPa, 4.4 GPa, and 0.1 MPa using synchrotron radiation. Patterns of the new phase were indexed on basis of a hexagonal unit cell, and Rietveld refinements indicate a NiAs-type arrangement (Figure 3 a, b and Table S1).<sup>[31]</sup> Most

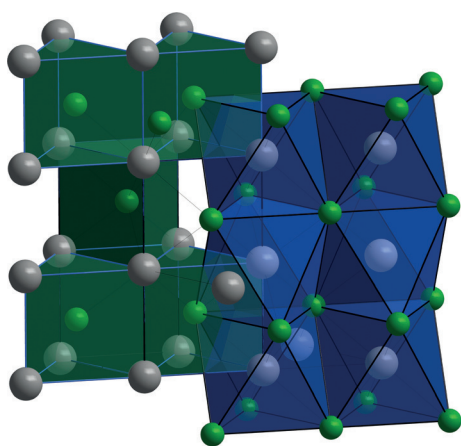


**Figure 3.** Crystal-structure refinements of NiAs-type FeN based on full diffraction profiles at a) 13.3 GPa and b) 0.1 MPa. The observed pattern is shown in black, the calculated one in red, and the difference in blue. The marked reflections (\*) are attributed to secondary phases (see text).

patterns contain extra reflections, which are mainly attributed to N<sub>2</sub> or residual  $\zeta$ -Fe<sub>2</sub>N. Despite the small number of experiments for pressure–volume correlation, a least squares fit of a Murnaghan-type equation of state to the experimental data results in the reasonable values  $V_0 = 33.78(3) \text{ \AA}^3$  and  $B_0 = 198(8) \text{ GPa}$  with  $B_0'$  fixed to 4.

Crystal structure refinements based on full diffraction profiles measured at various pressures and sample positions indicate a broad homogeneity range for the new phase Fe<sub>x</sub>N with  $x$  in the estimated range of between 0.60(5) and 1.0(1). In the following we will denote the phase as FeN.

NiAs-type FeN consists of a hexagonal closed packing of nitrogen, with the iron occupying all octahedral voids for  $x = 1$ . These FeN<sub>6</sub> octahedra form infinite chains by face sharing along the  $c$  direction (Figure 4). To our knowledge, this is not only the first face-sharing arrangement of FeN<sub>6</sub> octahedra in a nitride, but also the single NiAs-type 3d transition-metal nitride. Although some have been predicted, such as VN<sup>[33]</sup> and MoN,<sup>[34]</sup> they remain elusive and only heavier homologues, such as  $\delta$ -NbN<sup>[35]</sup> and TaN,<sup>[36]</sup> have been successfully synthesized. Comparison of the interatomic iron distances of NiAs-type FeN and  $\epsilon$ -iron with the same iron arrangement reveals<sup>[37]</sup> that  $d(\text{Fe}–\text{Fe})$  along the  $c$  axis varies only slightly (2.467 Å versus 2.442 Å) whereas along the  $a$  axis there is



**Figure 4.** Crystal structure of NiAs-type FeN, showing the face-sharing condensation of the iron-centered octahedra (blue), as well as the trigonal-prismatic coordination environment of the nitride ions (green).

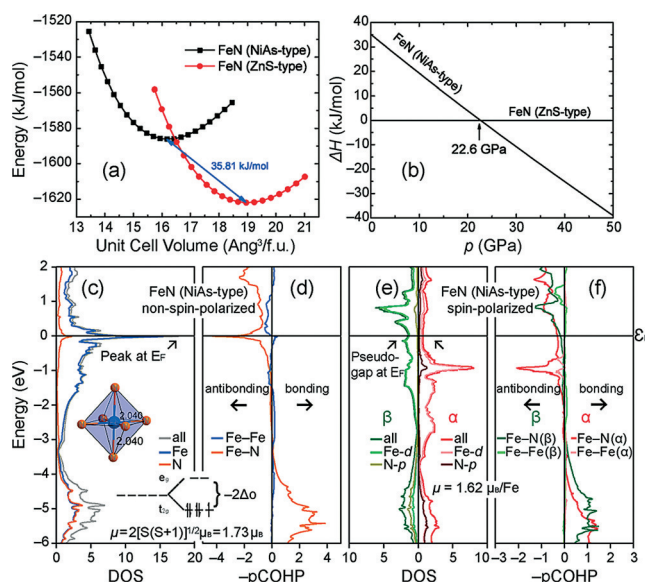
a larger difference: In the nitride, the Fe–Fe distance is significantly larger than in  $\epsilon$ -iron (2.737 Å versus 2.473 Å). Typical Fe–N distances range from 1.8 Å to 2.0 Å, but in NiAs-type FeN this distance is at the upper limit of 2.0 Å and hardly changes with pressure. FeN shows a large  $c/a$  ratio of about 1.80 (Table 2) which is typical for NiAs-type compounds of transition metals with Group V or VI elements.

**Table 2:** The  $c/a$  ratios of NiAs-type FeN, under pressure and ambient pressure, with related NiAs-type compounds.

NiAs-Type Compound	$c/a$ ratio
NiAs	1.3920 <sup>[41]</sup>
FeN (0.1 MPa)	1.7911
FeN (13.3 GPa)	1.8015
$\delta$ -NbN	1.8665 <sup>[35]</sup>
TiS	1.9515 <sup>[42]</sup>
VP	1.9560 <sup>[43]</sup>

An insight into the electronic structure of NiAs-type Fe<sub>x</sub>N for  $x = 1$  was gained by examining band structure and density of states (DOS), while a bonding analysis used the projected crystal orbital Hamilton populations (–pCOHP), a variant of the COHP technique<sup>[38]</sup> and their integrated values (–IpCOHP; Computational Details are provided in the Supporting Information). Features at the Fermi level,  $E_F$ , in the non-spin-polarized DOS and –pCOHP curves (Figure 5c,d) indicate an electronically unfavorable situation as  $E_F$  falls into a maximum of the DOS curves<sup>[39]</sup> with strongly antibonding Fe–N interactions; such attributes at  $E_F$  however, may also suggest that the material alleviates this unfavorable situation by approaching a magnetic state,<sup>[40]</sup> as indicated by Mössbauer spectroscopy.

A comparison of the total energies between various magnetic models reveals that the antiferromagnetic NiAs-type FeN model is just 0.6 kJ mol<sup>–1</sup> higher in energy than the ferromagnetic model, which, accordingly, tends to be preferred. Because the Fermi level falls into a pseudogap of the spin-polarized DOS curves (Figure 5e and Figure S2), an



**Figure 5.** a) Energy–volume curves and b) relative enthalpies as functions of pressure for NiAs-type as well as ZnS-type FeN at 0 K; c), d) non-spin-polarized DOS and –pCOHP curves; e), f): spin-polarized DOS and –pCOHP curves of NiAs-type FeN.

electronically favorable situation is inferred for magnetic NiAs-type FeN. This outcome is in good agreement with the Mössbauer spectra in which the magnetic sextet clearly denotes the presence of a magnetic ground state (Figure 2). The theoretical magnetic moment of 1.62  $\mu_B$ /Fe follows a spin-only rule-derived value for low-spin Fe<sup>3+</sup> ( $t_2^5 e_g$  configuration), for which the ligand-field stabilization energy is maximized (Figure 5c). The presence of such a low-spin-configuration for the iron atoms corresponds well to the observed Fe–N distances of approximately 2.01 Å, which is expected based on the covalent radii of nitrogen (0.71 Å) and low-spin iron (1.32 Å).<sup>[44]</sup> An alternative calculus using Shannon's ionic radii of N<sup>3–</sup> (1.54 Å) and low-spin Fe<sup>3+</sup> (0.55 Å) yields a similar Fe–N distance of 2.09 Å.<sup>[45]</sup> The states near and below the Fermi level stem primarily from Fe d as well as N p atomic orbitals (Figure 5e). Although antibonding Fe–N interactions are evident below  $E_F$  in the spin-polarized –pCOHP curves of both magnetic models (Figure 5f and Figure S3), the integrated values (Table S2) of the Fe–N –pCOHP (–IpCOHP) are indicative of a net bonding character for these contacts.

Previous research on the electronic and vibrational properties revealed that the zinc blende-type structure of FeN maintains dynamically stable.<sup>[25,46]</sup> To evaluate the stability of NiAs-type iron nitride, we determined the energy–volume curves and pressure-dependent relative enthalpies at absolute zero (Figure 5a,b). Because crystal structure data, Mössbauer spectra and electronic structure calculations suggest a ferromagnetic ground state for NiAs-type FeN while an antiferromagnetic ground state is proposed for ZnS-type iron nitride,<sup>[11]</sup> spin-polarized band structure computations were used. A comparison reveals that the total energy of the ZnS modification is lower by almost 36 kJ mol<sup>–1</sup>. At the atomic scale, the preference to adopt the ZnS-type

rather than NiAs-type modification originates from shorter Fe–N contacts corresponding to larger  $-I_p\text{COHP}$  values in ZnS-type FeN, that is, the modification is favored owing to a larger number of strong heteroatomic contacts. However, the values of the pressure-dependent relative enthalpies imply a phase transition from the ZnS-type to the NiAs-type for pressures  $\geq 22.6$  GPa. Additionally, investigations of the vibrational properties for NiAs-type FeN (Figure S4) point to dynamic stability.

In summary, a new high-pressure phase has been discovered in the Fe–N system by combining diamond anvil cell laser heating techniques with synchrotron-source Mössbauer spectroscopy and in situ X-ray diffraction. Theoretical investigations confirm that NiAs-type FeN exhibits a magnetic ground state. Although the modification is less stable than the ZnS-type form at ambient conditions, the application of pressure provides access to this modification at elevated temperatures.

### Acknowledgements

We acknowledge the European Synchrotron Radiation Facility for provision of synchrotron radiation facilities. We thank Rudolf Rüffer and the staff members for using beamline ID18 under experiment CH-4383. The allocation of CPU time through the IT Center of RWTH Aachen University is gratefully acknowledged. S.S. is grateful for a Liebig Stipend from the Fonds der Chemischen Industrie e.V. (FCI).

### Conflict of interest

The authors declare no conflict of interest.

**Keywords:** high-pressure chemistry · Mössbauer spectroscopy · nitrides · structure elucidation

**How to cite:** *Angew. Chem. Int. Ed.* **2017**, *56*, 7302–7306  
*Angew. Chem.* **2017**, *129*, 7408–7412

- [1] a) R. Juza, H. Hahn, *Z. Anorg. Allg. Chem.* **1938**, 239, 282; b) R. Juza, K. Langer, K. Von Benda, *Angew. Chem. Int. Ed. Engl.* **1968**, *7*, 360; *Angew. Chem.* **1968**, *80*, 373.
- [2] W. Schnick, *Angew. Chem. Int. Ed. Engl.* **1993**, *32*, 806; *Angew. Chem.* **1993**, *105*, 846.
- [3] G. V. Vajenine, G. Auffermann, Y. Prots, W. Schnelle, R. K. Kremer, A. Simon, *Inorg. Chem.* **2001**, *40*, 4866.
- [4] G. Auffermann, Y. Prots, R. Kniep, *Angew. Chem. Int. Ed.* **2001**, *40*, 547; *Angew. Chem.* **2001**, *113*, 565.
- [5] J. von Appen, M. W. Lumeij, R. Dronskowski, *Angew. Chem. Int. Ed.* **2006**, *45*, 4365; *Angew. Chem.* **2006**, *118*, 4472.
- [6] A. F. Young, C. Sanloup, E. Gregoryanz, S. Scandolo, R. J. Hemley, H. K. Mao, *Phys. Rev. Lett.* **2006**, *96*, 1.
- [7] J. C. Crowhurst, A. F. Goncharov, B. Sadigh, J. M. Zaug, D. Aberg, Y. Meng, V. B. Prakapenka, *J. Mater. Res.* **2008**, *23*, 1.
- [8] K. Niwa, D. Dzivenko, K. Suzuki, R. Riedel, I. Troyan, M. Eremets, M. Hasegawa, *Inorg. Chem.* **2014**, *53*, 697.
- [9] V. S. Bhadram, D. Y. Kim, T. A. Strobel, *Chem. Mater.* **2016**, *28*, 1616.
- [10] H. A. Wriedt, N. A. Gokcen, R. H. Nafziger, *Bull. Alloy Phase Diagr.* **1987**, *8*, 355.
- [11] K. Suzuki, H. Morita, T. Kaneko, H. Yoshida, H. Fujimori, *J. Alloys Compd.* **1993**, *201*, 11.
- [12] M. Wessel, R. Dronskowski, *Chem. Eur. J.* **2011**, *17*, 2598.
- [13] Z. Yang, S. Guo, X. Pan, J. Wang, X. Bao, *Energy Environ. Sci.* **2011**, *4*, 4500.
- [14] Y. Namiki, S. Matsunuma, T. Inoue, S. Koido, A. Tsubota, Y. Kuse, N. T., in *Nanocrystal*, InTech, **2011**, p. 349.
- [15] U. Schwarz, A. Wosylus, M. Wessel, R. Dronskowski, M. Hanfland, D. Rau, R. Niewa, *Eur. J. Inorg. Chem.* **2009**, 1634.
- [16] R. Niewa, D. Rau, A. Wosylus, K. Meier, M. Wessel, M. Hanfland, R. Dronskowski, U. Schwarz, *J. Alloys Compd.* **2009**, *480*, 76.
- [17] R. Niewa, D. Rau, A. Wosylus, K. Meier, M. Hanfland, M. Wessel, R. Dronskowski, D. A. Dzivenko, R. Riedel, U. Schwarz, *Chem. Mater.* **2009**, *21*, 392.
- [18] D. M. Borsa, D. O. Boerma, *Hyperfine Interact.* **2003**, *151*, 31.
- [19] E. Andrzejewska, R. Gonzalez-Arrabal, D. Borsa, D. O. Boerma, *Nucl. Instrum. Methods Phys. Res. Sect. B* **2006**, *249*, 838.
- [20] I. Jouanny, P. Weisbecker, V. Demange, M. Grafoute, O. Pena, E. Bauer-Grosse, *Thin Solid Films* **2010**, *518*, 1883.
- [21] H. Shimizu, M. Shirai, N. Suzuki, *J. Phys. Soc. Jpn.* **1998**, *67*, 922.
- [22] A. Filippetti, W. E. Pickett, *Phys. Rev. B* **1999**, *59*, 8397.
- [23] A. Houari, S. F. Matar, M. A. Belkhir, M. Nakhil, *Phys. Rev. B* **2007**, *75*, 1.
- [24] Y. Kong, *J. Phys. Condens. Matter* **2000**, *12*, 4161.
- [25] B. Eck, R. Dronskowski, M. Takahashi, S. Kikkawa, *J. Mater. Chem.* **1999**, *9*, 1527.
- [26] P. Lukashev, W. R. L. Lambrecht, *Phys. Rev. B* **2004**, *70*, 1.
- [27] L. C. Ming, W. A. Bassett, *Rev. Sci. Instrum.* **1974**, *45*, 1115.
- [28] I. Kupenko, L. Dubrovinsky, N. Dubrovinskaia, C. McCammon, K. Glazyrin, E. Bykova, T. B. Ballaran, R. Sinmyo, A. I. Chumakov, V. Potapkin, A. Kantor, R. Rüffer, M. Hanfland, W. Crichton, M. Merlini, *Rev. Sci. Instrum.* **2012**, *83*, 124501.
- [29] V. Potapkin, A. I. Chumakov, G. V. Smirnov, J. P. Celse, R. Rüffer, C. McCammon, L. Dubrovinsky, *J. Synchrotron Radiat.* **2012**, *19*, 559.
- [30] V. Ksenofontov, S. Reiman, M. Waldeck, R. Niewa, R. Kniep, P. Gütllich, *Z. Anorg. Allg. Chem.* **2003**, *629*, 1787.
- [31] Crystal structure for FeN from Rietveld refinement: hexagonal, space group  $P6_3/mmc$  (194); Fe in 2a; N in 2c: at 13.3 GPa  $a = 2.737(3)$  Å;  $c = 4.933(5)$  Å;  $V = 32.01$  Å<sup>3</sup>; at 4.4 GPa  $a = 2.775(7)$  Å;  $c = 4.982(1)$  Å;  $V = 33.24$  Å<sup>3</sup>; at 0.0001 GPa  $a = 2.800(2)$  Å;  $c = 5.015(5)$  Å;  $V = 34.05$  Å<sup>3</sup>. Further details on the crystal structure investigation(s) may be obtained from the Fachinformationszentrum Karlsruhe, 76344 Eggenstein-Leopoldshafen, Germany (fax: (+49)7247-808-666; e-mail: crysdta@fiz-karlsruhe.de), on quoting the depository numbers CSD-432891 (0.1 MPa), CSD-432893 (4.4 GPa), CSD-432892 (13.3 GPa).
- [32] C. Prescher, C. McCammon, L. Dubrovinsky, *J. Appl. Crystallogr.* **2012**, *45*, 329.
- [33] C. Ravi, *Calphad* **2009**, *33*, 469.
- [34] E. Zhao, J. Wang, Z. Wu, *Phys. Status Solidi B* **2010**, *247*, 1207.
- [35] N. Terao, *J. Less-Common Met.* **1971**, *23*, 159.
- [36] J. C. Fontbonne, A. Gilles, *Rev. Int. Hautes Temp.* **1969**, 181.
- [37] T. Takahashi, W. A. Bassett, H.-K. Mao, *J. Geophys. Res.* **1968**, *73*, 4717.
- [38] a) R. Dronskowski, P. E. Blöchl, *J. Phys. Chem.* **1993**, *97*, 8617; b) V. L. Deringer, A. L. Tchougréeff, R. Dronskowski, *J. Phys. Chem. A* **2011**, *115*, 5461; c) S. Maintz, V. L. Deringer, A. L. Tchougréeff, R. Dronskowski, *J. Comput. Chem.* **2013**, *34*, 2557; d) S. Maintz, V. L. Deringer, A. L. Tchougréeff, R. Dronskowski, *J. Comput. Chem.* **2016**, *37*, 1030.
- [39] a) V. L. Deringer, M. Lumeij, R. P. Stoffel, R. Dronskowski, *Chem. Mater.* **2013**, *25*, 2220; b) S. Steinberg, J. Brgoch, G. J. Miller, G. Meyer, *Inorg. Chem.* **2012**, *51*, 11356.

- [40] G. A. Landrum, R. Dronskowski, *Angew. Chem. Int. Ed.* **2000**, *39*, 1560; *Angew. Chem.* **2000**, *112*, 1598.
- [41] N. Alsén, *Geol. Foeren. Stockholm Foerh.* **1925**, *47*, 19.
- [42] M. Onoda, H. Wada, *J. Less-Common Met.* **1987**, *132*, 195.
- [43] N. Schönberg, W. G. Overend, A. Munthe-Kaas, N. A. Sørensen, *Acta Chem. Scand.* **1954**, *8*, 226.
- [44] B. Cordero, V. Gómez, A. E. Platero-Prats, M. Revés, J. Echeverría, E. Cremades, F. Barragán, S. Alvarez, *Dalton Trans.* **2008**, 2832.
- [45] R. D. Shannon, *Acta Crystallogr. Sect. A* **1976**, *32*, 751. Because no Shannon radius  $r$  is provided for a sixfold coordinated  $\text{N}^{3-}$  anion, its ionic radius has been approximated based on a fourfold coordinated one ( $^{[4]}r$ ) using the expression  $^{[4]}r = ^{[6]}r(4/6)^{1/8}$ . Further details are provided in: R. Dronskowski, *Computational Chemistry of Solid State Materials*, Wiley-VCH Publishers, Weinheim, **2005**.
- [46] H. R. Soni, V. Mankad, S. K. Gupta, P. K. Jha, *J. Alloys Compd.* **2012**, *522*, 106.

Manuscript received: March 8, 2017

Revised manuscript received: April 10, 2017

Version of record online: May 18, 2017