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## Engineering of inorganic nanoparticles as magnetic resonance imaging contrast agents

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### Abstract

Magnetic resonance imaging (MRI) is a highly valuable non-invasive imaging tool owing to its exquisite soft tissue contrast, high spatial resolution, lack of ionizing radiation, and wide clinical applicability. Contrast agents (CAs) can be used to further enhance the sensitivity of MRI to obtain information-rich images. Recently, extensive research efforts have been focused on the design and synthesis of high-performance inorganic nanoparticle-based CAs to improve the quality and specificity of MRI. Herein, the basic rules, including the choice of metal ions, effect of electron motion on water relaxation, and involved mechanisms, of CAs for MRI have been elucidated in detail. In particular, various design principles, including size control, surface modification (*e.g.* organic ligand, silica shell, and inorganic nanolayers), and shape regulation, to impact relaxation of water molecules have been discussed in detail. Comprehensive understanding of how these factors work can guide the engineering of future inorganic nanoparticles with high relaxivity. Finally, we have summarized the currently available strategies and their mechanism for obtaining high-performance CAs and discussed the challenges and future developments of nanoparticulate CAs for clinical translation in MRI.

### 1. Introduction

Diagnostic imaging techniques to detect lesion information (*e.g.*, type, location, and stage) are currently playing diverse roles in the combat against serious diseases such as cancer.<sup>1-3</sup> Among all existing imaging techniques, magnetic resonance imaging (MRI) is a popular choice because of its exquisite soft tissue contrast, high spatial resolution, lack of ionizing radiation, unlimited signal penetration depth, and wide clinical applicability.<sup>4</sup> Contrast agents (CAs) can be used to further improve the sensitivity and detection capability of MRI. The use of CAs in MRI has greatly expanded the possible conditions that may be explored,

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#### Conflicts of interest

There are no conflicts to declare.

offering physicians unprecedented noninvasive visualization of many ailments.<sup>5,6</sup> Most clinically available CAs are chemical agents, traditionally gadolinium ( $\text{Gd}^{3+}$ ) chelates. Generally, these  $\text{Gd}^{3+}$  chelates have very short circulation lifetimes in the body, work in a nonspecific manner, and have certain limitations as molecular CAs for long-time tracking. Another disadvantage of clinically used  $\text{Gd}^{3+}$  chelates is their relatively low proton relaxation efficiency, which means that high doses are necessary for imaging (*e.g.*, mM concentrations of  $\text{Gd}^{3+}$  chelates at target areas).<sup>7</sup> However, high-dose injection of  $\text{Gd}^{3+}$  chelates may be harmful to the body as it has been reported that the released free  $\text{Gd}^{3+}$  ions may cause severe nephrogenic system fibrosis syndrome *in vivo*.<sup>8</sup> Therefore, the imaging performances of clinically available CAs to date have been far from optimal.

Nanoparticles naturally lend themselves to be used as CAs in MRI, where magnetically responsive atoms are essential.<sup>9</sup>  $T_1$  or spin–lattice contrast may be enhanced through the use of paramagnetic centers such as gadolinium<sup>10,11</sup> or manganese ions,<sup>12,13</sup> whereas iron oxide-based nanosystems are often used for  $T_2$ -weighted<sup>14–16</sup> or spin–spin imaging. Unlike other molecular imaging agents, these nanoparticles are not directly imaged in the resulting MRI image; in fact, the metal ions perturb the magnetic properties of protons around them, which are the actual target of a traditional MRI scan.<sup>6</sup> Among various types of nanoscale imaging agents, inorganic nanoparticle (NP)-based CAs (hereafter, “NPs” refers to inorganic nanoparticles) have received great attention and have been regarded as promising imaging probes. This is due to their unique electronic, magnetic and optical properties at the nanoscale, facile intracellular uptake and accumulation, and excellent *in vivo* stability.<sup>17</sup> The high surface areas of these NPs also greatly enhance their chemical reactivity, allowing for easy surface modification with biological and bioactive materials. These characteristics have caused NPs to emerge as promising candidates for next-generation imaging CAs.<sup>18</sup>

The contrast behavior of NPs strongly depends on their perturbation of magnetic properties of nearby protons, which can be tuned through a series of factors (*e.g.*, size, surface, and shape) and various innovative approaches or strategies. Great efforts have been made for the rational design and manufacture of NPs with high relaxivity to elevate imaging sensitivity of CAs and even impart MRI with molecular imaging capabilities for better monitoring of biological processes (*e.g.*, tumor microenvironment, cell tracking, and cellular interactions). This has resulted in the wide expansion of this research area, leading to the formation of sub-branches such as ultra-high field MRI,<sup>19–21</sup> multi-modal imaging,<sup>22–24</sup> and incorporation of theranostic techniques into MRI.<sup>25</sup>

Both the rational design and engineering of nanoparticulate CAs play important roles in this field, which may, in the meantime, broaden biomedical applications of MRI. However, the complexity of this process is much greater than initially expected, calling for broad expertise in a multidisciplinary manner through chemistry, physics, biology, and engineering. The goal of this review is not to comprehensively summarize nanoparticulate CAs for MRI that have been reported so far,<sup>26</sup> as there are many reviews available that list the major types of NPs as CAs ( $\text{Gd}^{3+}$ -based  $T_1$  CAs,<sup>11,27</sup> magnetic  $T_2$  CAs,<sup>14,15,28–30</sup> and dual-  $T_1/T_2$  CAs<sup>9,31</sup>) or summarize a special application of CAs (*e.g.*, bioresponsive CAs,<sup>32–34</sup> and multimodal CAs<sup>24</sup>). Instead, this review intends to comprehensively discuss the main rules and mechanisms of nanoparticulate CAs, key factors influencing the relaxation of CAs, and

innovative strategies to engineer CAs with high relaxivity. As the Chinese aphorism goes, “It is better to teach a man to fish than to give him a fish,” we hope that this timely and comprehensive review can familiarize readers with the general principles of designing high-performance CAs, stimulate extensive further studies, and attract extensive attention from researchers worldwide.

## 2. Basic rules for MRI contrast agents

Since the mechanisms of MRI have been thoroughly described elsewhere,<sup>11,29</sup> herein, we will focus on the basic rules of CAs for MRI, an in-depth understanding of which is extremely important for formulating design principles of CAs. The contrast enhancement of MRI originates from the interaction between CAs and surrounding water protons, which shorten the longitudinal ( $T_1$ ) or transverse ( $T_2$ ) relaxation time of nearby water molecules. Usually,  $T_1$  CAs are applied to increase signal intensity, resulting in a positive contrast enhancement (brighter image) in  $T_1$ -weighted MRI, whereas  $T_2$  CAs can decrease signal intensity, providing a negative contrast enhancement (darker image) in  $T_2$ -weighted MRI. Relaxivity ( $r_i$ ,  $i = 1, 2$ ) is therefore used to evaluate the efficacy of CAs, which is defined as the slope of a plot of relaxation rate ( $1/T_i$ ,  $i = 1, 2$ ) versus concentration of CAs:

$$1/T_i = 1/T_{i0} + r_i C \quad (1)$$

where  $T_{i0}$  is the relaxation time without CAs and  $C$  is the concentration of CAs (mM). The values of  $r_1$ ,  $r_2$ , and  $r_2/r_1$  are typically used to evaluate MRI efficiency of nanoparticulate CAs. The transition metal ions ( $\text{Mn}^{2+}$ ,  $\text{Fe}^{2+}/\text{Fe}^{3+}$ ,  $\text{Co}^{2+}$ , and  $\text{Cu}^{2+}$ ) and lanthanide metal ions ( $\text{Eu}^{3+}$ ,  $\text{Gd}^{3+}$ ,  $\text{Ho}^{3+}$ , and  $\text{Dy}^{3+}$ ) are the major candidate relaxation centers that can be utilized to construct NPs with proper chemical compositions for enhanced contrast in MRI. The choice of a metal ion requires consideration of its electron motion, nanostructure-dependent effects, and involved relaxation mechanism to obtain high-relaxivity CAs.

### 2.1 Choice of ions

The selection of appropriate ions to construct nanoparticulate agents should obey certain rules. Inherently, one metal ion includes both the electron orbital motion and electron spin motion. The fast electron orbital motion is significantly higher than the speed of water proton relaxation, while the electron spin motion is slower and can be closely in tune with the slow water proton relaxation. Only the electron spin magnetic moment of ions can efficiently influence the longitudinal water proton relaxation ( $r_1$ ). Therefore,  $r_1$  will be higher if the total electron angular momentum ( $J$ ) contains a larger contribution from electron spin angular momentum ( $S$ ) of adjacent ions. On the other hand, it will be small if  $J$  consists only of electron orbital angular momentum ( $L$ ).<sup>35</sup> Table 1 summarizes the outer orbital, quantum parameters ( $S$ ,  $L$ , and  $J$ ) and effective magnetic moment ( $\mu_{\text{eff}}$ ) of candidate metal ions for MRI, from which we can determine suitable ions for chemical composition of CAs. For example, since  $\text{Gd}^{3+}$  has seven unpaired electrons and its  $J$  ( $J = 7/2$ ) value is completely derived from  $S$  ( $S = 7/2$ ,  $L = 0$ ), inorganic  $\text{Gd}^{3+}$ -based NPs (e.g.,  $\text{Gd}_2\text{O}_3$ ,  $\text{NaGdF}_4$ ,  $\text{GdF}_3$ ) have been widely explored as  $T_1$  CAs.<sup>36-49</sup> Notably,  $\text{Mn}^{2+}$  plays a highly important role in biological processes as a cofactor of enzymes and through neurotransmitter

release control; moreover,  $\text{Mn}^{2+}$  has five unpaired electrons ( $J = 5/2$ ,  $S = 5/2$ ,  $L = 0$ ), showing prominent contrast effects that have been shown to enable visualization of brain anatomic structure and its neuronal activity.<sup>50</sup> Therefore,  $\text{Mn}^{2+}$ -based NPs constitute a new imaging category and are promising candidates as  $T_1$  CAs.<sup>51-69</sup>

The  $r_2$  of a CA is largely proportional to the square of  $\mu_{\text{eff}}$  of the ion, as  $T_2$  CAs show reduced signal dependence on induced magnetic moment of these CAs.<sup>70</sup> Among all lanthanide ions,  $\text{Dy}^{3+}$  (together with  $\text{Ho}^{3+}$ ) has the highest effective magnetic moment ( $\mu_{\text{eff}} = 10.6 \mu_{\text{B}}$ ) and much attention has been paid to synthesizing  $\text{Dy}^{3+}$ - or  $\text{Ho}^{3+}$ -based NPs as  $T_2$  CAs,<sup>70-72</sup> especially for ultra-high field imaging.<sup>19-21,73</sup> Studies of D-glucuronic acid-coated ultra-small  $\text{Ln}_2\text{O}_3$  NPs ( $\text{Ln} = \text{Eu}, \text{Gd}, \text{Dy}, \text{Ho}, \text{or Er}$ ) show that  $\text{Gd}_2\text{O}_3$  NPs exhibit the highest  $r_1$  value, whereas  $\text{Dy}_2\text{O}_3$  NPs possess the highest  $r_2$  value among all studied NPs.<sup>74</sup> Unlike the lanthanides, iron is one of the most abundant metal elements in living organisms and deeply involved in various biological processes including cellular respiration by redox enzymes and oxygen transport by hemoglobin.<sup>75</sup> Magnetic iron oxide NPs (*e.g.*,  $\text{Fe}_3\text{O}_4$ ,  $\gamma\text{-Fe}_2\text{O}_3$ ,  $\text{Fe@Fe}_3\text{O}_4$ ,  $\text{MnFe}_2\text{O}_4$ , *etc.*) are therefore very popular  $T_2$  CAs that can produce negative imaging contrast.<sup>29,30,76-84</sup> The magnetization of paramagnetic NPs, such as Gd-based NPs and Mn-based NPs, directly depends on the total number of metal ions, and no magnetization is present in the absence of an external magnetic field. However, magnetic iron oxide NPs exhibit strong intrinsic magnetization, which may cause microscopic field inhomogeneity and accelerate the dephasing of protons for MRI contrast enhancement when an external magnetic field is applied.

In addition to relaxation impact from ion species, the ion valency also affects relaxation interaction between ions and water protons. Variable valencies (*e.g.*,  $\text{Mn}^{4+}/\text{Mn}^{2+}$  ions,<sup>54</sup>  $\text{Co}^{3+}/\text{Co}^{2+}$  ions,<sup>85</sup>  $\text{Eu}^{3+}/\text{Eu}^{2+}$  ions<sup>86</sup>) can endow these ions with tunable MRI performance under certain reducing or oxidative microenvironments, thus creating responsive MRI CAs for molecular imaging.<sup>32</sup>  $\text{Mn}^{4+}$ -based NPs are widely used as pH or redoxresponsive MRI CAs because  $\text{Mn}^{4+}$  ( $J = 9/2$ ,  $S = 3/2$ ,  $L = 3$ ) with negligible relaxation interactions can be easily reduced to  $\text{Mn}^{2+}$  ( $J = 5/2$ ,  $S = 5/2$ ,  $L = 0$ ), which shows high MRI contrast capability according to the rules mentioned above. Overall, understanding the basic principles of NPs such as ion species, ion valency, and magnetic properties will facilitate the engineering of nanoparticulate CAs with excellent MRI performance for various applications.

## 2.2 Mechanisms of contrast enhancement

The relaxation enhancement of water protons near a paramagnetic center generally follows the inner sphere, secondary sphere, and outer sphere mechanism (Fig. 1).<sup>87</sup> The water molecules directly coordinated to the paramagnetic center are responsible for inner sphere relaxation contribution, while bulk water molecules in the nearby environment affect outer sphere relaxation. In some cases, water molecules that are hydrogen bonded to chelators are regarded as the secondary sphere. The contrast enhancement of individual CAs closely depends on their chemical composition, structure, and surfacemodification.

For  $T_1$  CAs, on the basis of the classical Solomon–Bloembergen–Morgan (SBM) theory,<sup>35,87,88</sup> common strategies for enhancing relaxivity mainly include increasing the number of bound water molecules ( $q$ ), enhancing rotational correlation time ( $\tau_{\text{R}}$ ), and

optimizing water residence time ( $\tau_M$ ). Higher  $\tau_R$  values imply slower molecular tumbling and lead to increased proton relaxivity.  $\tau_M$  is the inverse of water exchange rate in the inner sphere ( $k_{ex} = 1/\tau_M$ ), which is favorable for proton relaxivity in the case of fast water exchange (high  $k_{ex}$ ). The relationship between  $\tau_R$  and  $\tau_M$  can be described by the following equation:<sup>89</sup>

$$1/\tau_{ci}=1/T_{ie}+1/\tau_R+1/\tau_M \quad i=1, 2 \quad (2)$$

where  $\tau_c$  is the correlation time parameter and  $T_{ie}$  ( $i = 1, 2$ ) is the electronic relaxation time. Since it is challenging to predict and experimentally alter  $T_{ie}$ , approaches to obtain highly efficient CAs commonly focus on the regulation of  $\tau_R$  and  $\tau_M$ .<sup>37</sup> By immobilizing  $Gd^{3+}$  on the surface or lattice of NPs, prolonging  $\tau_R$  will result in significantly enhanced  $r_1$ , which is a commonly used strategy because these nanoparticulate CAs can couple the tumbling rate of  $Gd^{3+}$  chelates to that of much larger species (*e.g.*, titanium, gold, and silica NPs).<sup>90-98</sup> Another widely used approach to enhance  $r_1$  is the optimization of  $\tau_M$ ,<sup>99,100</sup> which is readily accessible because nanoparticulate CAs inherently have relatively slow tumbling rates.

Different from  $T_1$  CAs that influence the relaxation of water proton through direct energy exchange between water and paramagnetic ions,<sup>101</sup> magnetized NPs can induce increased  $T_2$  relaxivity, which reduces phase coherence of water molecules in an inhomogeneous magnetic field generated by magnetized NPs.<sup>102-104</sup> Based on the quantum mechanical outer sphere theory,  $T_2$  relaxivity is given by eqn (3)<sup>105</sup>

$$1/T_2=(256\pi^2\gamma^2/405)V^*M_s^2r^2/D(1+L/r) \quad (3)$$

where  $r$  and  $M_s$  are the effective radius of magnetic NPs and saturation magnetization, respectively,  $\gamma$  is the proton gyromagnetic ratio,  $V^*$  is the volume fraction,  $D$  is the diffusivity of water molecules, and  $L$  is the thickness of an impermeable surface shell on the magnetic core. It should be noted that  $M_s$  value of a magnetic NP can be tuned by its high surface area, where the interior spins of magnetic NPs are usually aligned, while the spins of surface atoms are magnetically disordered. This is known as the spin canting effect, which is non-negligible because NPs have a large surface-to-volume ratio.<sup>29</sup> According to eqn (3),  $T_2$  relaxivity of magnetic NPs is highly dependent on effective radius of the magnetic core,  $M_s$ , and thickness of an impermeable surface.

### 3. Factors influencing relaxivity of MRI contrast agents

Recently, great advances in nanotechnology have stimulated the quick development of nanoparticulate CAs, which have two main traits of interest: the dominant quantum confinement effect and the large surface area. Size-dependent magnetic, electronic, and optical properties originating from dominant quantum confinement effects can be utilized in the rational design of high-performance MRI CAs at the nanoscale. Moreover, the large surface area of NPs offers increased chemical reactivity for surface modification with additional functional species, enabling enhanced magnetic relaxivity and higher-quality MRI

with reasonable detection limits. Typically, three techniques are applied for obtaining high-relaxivity CAs at the nanoscale: size control, surface modification, and shape regulation (Fig. 2).

### 3.1 Size effects

Inorganic NPs present a unique size-dependent effect on physical properties including surface plasmon resonance of metal nanoparticles,<sup>106,107</sup> fluorescence emission of semiconductor NPs,<sup>108</sup> and superparamagnetism or heating of magnetic NPs.<sup>109,110</sup> By controlling nucleation and growth, the size of NPs can be well controlled. The dependence of cellular uptake, biodistribution, and clearance rate on NP size has been widely investigated by many groups.<sup>111-115</sup> To date, several types of CAs have been developed for various applications in MRI, which can be classified according to their chemical composition (Gd-, Mn-, Fe-, Dy- or Ho-based CAs; Table 2). A deep understanding of the involved mechanisms between size effect and proton relaxation will undoubtedly benefit the rational design of  $T_1$  or  $T_2$  CAs for high-performance MRI. For  $T_1$  CAs, smaller-sized NPs usually generate greater contrast enhancement because their surface-to-volume ratio is the dominant contributing factor, which increases as the size of NPs is decreases. However,  $\tau_R$  will become much lower for large NPs and plays a dominant role in affecting relaxivity, causing larger NPs to possess higher  $r_1$  values. For  $T_2$  CAs, larger-sized magnetic NPs usually have higher magnetization than smaller-sized NPs, and therefore,  $r_2$  values increase as size increases. However, this trend breaks down when sizes of magnetic NPs become too large.

**3.1.1 Gd-based contrast agents**—The recently developed Gd<sup>3+</sup>-based NPs such as Gd<sub>2</sub>O<sub>3</sub>,<sup>43-45,116</sup> NaGdF<sub>4</sub>,<sup>36,38-41</sup> GdF<sub>3</sub>,<sup>46,117,118</sup> and Gd<sup>3+</sup>-doped NPs (*e.g.*, quantum dots,<sup>119-123</sup> upconversion NPs,<sup>124-126</sup> and carbon dots<sup>127,128</sup>) have been regarded as promising  $T_1$  CAs because of their controllable size and morphology. By optimizing the reaction parameters during both nucleation and growth, Johnson *et al.* synthesized size-tunable ultra-small NaGdF<sub>4</sub> NPs with uniform sizes of around 2.5 nm, 4 nm, 6.5 nm, and 8 nm (Fig. 3a) and found that the  $r_1$  value increased dramatically with decreasing size from 3.0 mM<sup>-1</sup> s<sup>-1</sup> for 8.0 nm NPs to 7.2 mM<sup>-1</sup> s<sup>-1</sup> for the smallest 2.5 nm ones (Fig. 3b).<sup>49</sup> Correlations of NPs' sizes and their  $r_1$  values showed that contributions from the surface-to-volume ratio increased with decreasing NP size, becoming the dominant contributing factor. Several other research groups also reported similar size effects of NaGdF<sub>4</sub> NPs on relaxivity.<sup>36,39</sup> However,  $\tau_R$  became much lower for larger NaGdF<sub>4</sub> NPs (>15 nm) and played a dominant role in affecting  $r_1$  value, as demonstrated by Gao and co-workers, who found  $r_1$  values of 5.7 mM<sup>-1</sup> s<sup>-1</sup> for 15 nm NaGdF<sub>4</sub> NPs, which increased to 8.78 mM<sup>-1</sup> s<sup>-1</sup> for 20 nm ones (Fig. 3c).<sup>41</sup> Inspired by this, we recently synthesized ultra-small sized and renal-clearable NaGdF<sub>4</sub> nanodots (~2 nm) with high  $r_1$  values of 8.93 mM<sup>-1</sup> s<sup>-1</sup>.<sup>40</sup> These nanodots exhibited excellent magnetic resonance angiography (MRA) performance, as the vascular signal was much stronger than that achieved using clinical MRI CAs (Fig. 3d). With the new CAs, even capillary vessels and atherosclerotic plaques could be clearly delineated (Fig. 3e).

The same trends were also observed for ultra-small Gd<sub>2</sub>O<sub>3</sub> NPs,<sup>45,129</sup> such that  $r_1$  values increased from 4.4 mM<sup>-1</sup> s<sup>-1</sup> (4.6 nm Gd<sub>2</sub>O<sub>3</sub>) to 8.8 mM<sup>-1</sup> s<sup>-1</sup> (2.2 nm Gd<sub>2</sub>O<sub>3</sub>).<sup>45</sup> Based on the fitting of Gaussian, Lorentzian, and log-normal functions, Lee *et al.* concluded that the optimal range of diameters for high-relaxivity Gd<sub>2</sub>O<sub>3</sub> NPs should be 1–2.5 nm with a theoretical maximum  $r_1$  value of 10.1–10.6 mM<sup>-1</sup> s<sup>-1</sup>.<sup>48</sup> In that study, they reported a large  $r_1$  value of 9.9 mM<sup>-1</sup> s<sup>-1</sup> of ultra-small Gd<sub>2</sub>O<sub>3</sub> NPs coated with D-glucuronic acid. Consistent with ultra-small NaGdF<sub>4</sub> NPs, such a large  $r_1$  value of ultra-small Gd<sub>2</sub>O<sub>3</sub> NPs was attributed to the large surface-to-volume ratio, as well as the cooperative induction of surface Gd<sup>3+</sup> ions for longitudinal relaxation. It is generally accepted that the  $r_1$  value should be as large as possible, and for high-performance T<sub>1</sub> MRI CAs, the  $r_2/r_1$  ratio should be close to unity, while a greater value of  $r_2/r_1$  will generally result in higher contrast enhancement efficiency for T<sub>2</sub> MRI CAs.<sup>9</sup> The  $r_2/r_1$  ratio of ultra-small Gd<sub>2</sub>O<sub>3</sub> NPs was estimated to be 1.06 ( $r_1 = 9.9$  mM<sup>-1</sup> s<sup>-1</sup>,  $r_2 = 10.5$  mM<sup>-1</sup> s<sup>-1</sup>), which seemed to satisfy this condition.

Biologically, these ultra-small sized agents can escape from recognition and can be captured by the liver and then be rapidly cleared through the kidneys, in which the basal lamina has pores of approximately 10 nm to filter metabolites (*e.g.*, ultrasmall sized NPs) from the bloodstream.<sup>38,48,130</sup> Such rapid renal clearance diminishes *in vivo* toxicity of NPs owing to the suppressed tissue accumulation and certainly facilitates their future clinical translation and application. In addition, ultrasmall Gd<sup>3+</sup>-based NPs can be integrated with additional components such as multifunctional nanoplateforms for non-targeted blood pool imaging,<sup>40</sup> targeted molecular imaging,<sup>41</sup> pH/H<sub>2</sub>O<sub>2</sub>- responsive MRI,<sup>131</sup> and brain tumor diagnosis and image-guided therapy.<sup>130</sup>

**3.1.2 Mn-based contrast agents**—As another major class of T<sub>1</sub> CAs, Mn-based NPs (*e.g.*, MnO NPs, Mn<sub>3</sub>O<sub>4</sub> NPs, and MnO<sub>2</sub> nanosheets) have gained extensive interest as GSH- or pH-responsive MRI CAs due to their readily changeable valency and easy release of manganese under acidic conditions.<sup>54,55,63,133</sup> The size-dependent magnetic properties of MnO and Mn<sub>3</sub>O<sub>4</sub> NPs with monodispersed sizes in the range of 6–22 nm were initially reported by Seo *et al.*<sup>134</sup> Subsequently, the size-dependent T<sub>1</sub> relaxation properties of MnO NPs were investigated by Hyeon and co-workers,<sup>69</sup> demonstrating higher T<sub>1</sub> signal enhancement for smaller MnO NPs than larger ones (Fig. 4a and b). The interaction between proton spins and surface Mn<sup>2+</sup> ions on NPs is regarded as the main dipolar interaction responsible for their relaxation properties. Clear T<sub>1</sub>-weighted MR images of the brain, liver, kidney, and spinal cord with a fine anatomical structure were obtained after administration of MnO NPs. Since MnO NPs spontaneously get oxidized to Mn<sub>3</sub>O<sub>4</sub> NPs in air,<sup>68</sup> An *et al.* reported measurements of T<sub>1</sub> and T<sub>2</sub> relaxation of spherical and tetrahedral Mn<sub>3</sub>O<sub>4</sub> NPs with varied uniform sizes.<sup>132</sup> In agreement with the previous results, the  $r_1$  value increased as the size of Mn<sub>3</sub>O<sub>4</sub> NPs decreased for both spherical (Fig. 4c) and tetrahedral NPs (Fig. 4d), and 3 nm-sized spherical NPs exhibited the highest  $r_1$  value of 2.38 mM<sup>-1</sup> s<sup>-1</sup>, again indicating that higher surface-to-volume ratios lead to higher relaxivity of ultrasmall NPs.

To further increase surface-to-volume ratio, Mn-based NPs were etched into hollow structures by selectively removing the core MnO phase from the Mn<sub>3</sub>O<sub>4</sub> shell phase, thus exposing more active Mn<sup>2+</sup> ions to water protons.<sup>68</sup> As a result, the hollow nanostructure

with a larger water-accessible surface area possessed a much higher  $r_1$  value of  $1.42 \text{ mM}^{-1} \text{ s}^{-1}$  compared to the solid ones ( $r_1 = 0.21 \text{ mM}^{-1} \text{ s}^{-1}$ ) at 3.0 T. Such hollow structures could be used as a drug delivery system with marked MRI contrast enhancement for the monitoring of drug release.<sup>139,140</sup>

**3.1.3 Fe-based contrast agents**—Iron oxide NPs are generally regarded to be safer than Gd- or Mn-based CAs because iron is essential in the human body, mostly stored as ferritin in the blood.  $\text{Fe}^{3+}$  has 5 unpaired electrons ( $J = 5/2$ ,  $S = 5/2$ ,  $L = 0$ ), the same as  $\text{Mn}^{2+}$ , making  $\text{Fe}^{3+}$  a promising candidate for  $T_1$  CAs.<sup>136,141-143</sup> Small-sized iron oxide NPs show enhanced  $T_1$  contrast effects due to the existence of a large number of  $\text{Fe}^{3+}$  ions on the surface, suppressing  $T_2$  relaxation by their small magnetic moment. The magnetic properties and  $r_1$  values of various small-sized iron oxide NPs from 1.5 to 12 nm have been investigated (Fig. 5a).<sup>135</sup> By canting the most surface spins (around 93.6%), 2.2 nm-sized iron oxide NPs with low magnetic moments exhibited a relatively high  $r_1$  value of  $4.78 \text{ mM}^{-1} \text{ s}^{-1}$ , while the  $r_1$  values of 12 nm NPs (nearly 38.6% of spins were canted, Fig. 5b) decreased to  $2.37 \text{ mM}^{-1} \text{ s}^{-1}$  (Fig. 5c). This phenomenon was further verified in cellular studies, in which significant  $T_1$  signal enhancement for 3 nm-sized NPs and attenuated  $T_1$  signal for cells incubated with 12 nm-sized iron oxide NPs were found (Fig. 5d). Another study from Zhang *et al.* also indicated similar size-dependences of  $r_1$  values based on ultra-small  $\text{MnFe}_2\text{O}_4$  NPs.<sup>136</sup> Recently, exceedingly small iron oxide NPs with a hydrodynamic size of 4.7 nm, consisting of approximately 3 nm inorganic cores and about 1 nm organic hydrophilic shell, were reported to have high  $T_1$  relaxivity ( $r_1 = 5.2 \text{ mM}^{-1} \text{ s}^{-1}$  at 1.5 T).<sup>142</sup> Excreted mostly within the first 2.5 h through the urine, the exceedingly small iron oxide NPs have been successfully applied for high-performance magnetic resonance angiography (MRA) imaging (Fig. 5e and f). Strong positive contrast from both the heart and blood vessels has been found, which decreased over time, while the signal from the bladder increased, indicating renal clearance.

Importantly, iron oxide NPs are also widely used as  $T_2$  CAs, where the magnetic moment of these NPs decreases rapidly as the size decreases due to reduced magnetic anisotropy and spin disordering on the surface of NPs.<sup>30,144</sup> The size-dependence of  $T_2$  contrast has been extensively investigated both theoretically and experimentally. Theoretically, two different regimes are predicted. The first is called the “motional averaging regime” (MAR) that describes  $T_2$  relaxation for relatively small iron oxide NPs,<sup>104</sup> where water diffusion near NPs occurs on a much faster timescale than the resonance frequency shift, resulting in increased  $r_2$  values with increasing NP size. However, this theory breaks down for larger NPs because the induced surrounding perturbing field around larger NPs is much stronger, and proton diffusion becomes a nondominant factor for signal decay. Therefore, the relaxation rate is independent of diffusion and does not significantly increase with further increase in size of NPs. This process is termed the “static dephasing regime” (SDR), first introduced by Yablonskiy and Haacke.<sup>145</sup> Importantly, this tendency has been validated on NPs of varied sizes synthesized by many different groups.

Generally, small iron oxide NPs synthesized by various research groups are in the MAR, where larger NPs possess higher magnetization and therefore larger  $r_2$  values. For example, such size effects on  $T_2$  contrast have been reported by Cheon and co-workers,<sup>79</sup> using small



and size-controlled Fe<sub>3</sub>O<sub>4</sub> NPs (Fig. 6a), where the variation of NPs size from 4 to 6, 9, and 12 nm resulted in  $r_2$  values of 78, 106, 130, and 218 mM<sup>-1</sup> s<sup>-1</sup>, respectively (Fig. 6b–d). Similar size effect on relaxation rates has also been reported by many other research groups.<sup>80,83,105,146–149</sup> Importantly, this size effect is a universal phenomenon, which is not limited to spherical NPs. For example, octapod iron oxide NPs with average edge lengths of 30 nm exhibited a large  $r_2$  value of 679.25 mM<sup>-1</sup> s<sup>-1</sup>, almost three times higher than that of the NPs of 20 nm in edge lengths.<sup>137</sup> However, according to the SDR theory, the  $r_2$  value reaches a plateau as the size continues to increase. It has been reported by Chen and co-workers that the size effect became less significant for single iron oxide NPs of 50 nm (Fig. 6e).<sup>138</sup> They found that  $r_2$  values increased rapidly from 173.37 to 203.86 and 239.98 mM<sup>-1</sup> s<sup>-1</sup> at 7 T, when the particle size increased from 8 to 23 and 37 nm, respectively (Fig. 6f), but  $r_2$  value was only slightly higher at 248.89 mM<sup>-1</sup> s<sup>-1</sup> for 65 nm-sized iron oxide NPs (Fig. 6g). This SDR theory has also been experimentally verified in cells loaded with iron oxide NPs,<sup>150</sup> as well as in an induced clustering process.<sup>151</sup>

**3.1.4 Dy- or Ho-based contrast agents**—Ultra-high field MRI, the future of next-generation MRI scanners, brings not only new opportunities for biomedical research, but also great challenges in developing exogenous CAs, since most currently used CAs are only suitable at relatively low magnetic field strengths (*e.g.*, 1.5 T or 3.0 T) but are inefficient at higher magnetic field strengths (*e.g.*, 7.0 T). Dysprosium (Dy<sup>3+</sup>) and holmium (Ho<sup>3+</sup>) are promising candidates for imaging in an ultra-high magnetic field due to their relatively short electronic relaxation time ( $\sim 10^{-13}$  s) and large magnetic moment ( $\sim 10.6 \mu_B$ ).<sup>71</sup> The size-dependent properties of NaDyF<sub>4</sub> NPs under ultra-high magnetic fields were first reported by van Veggle and co-workers,<sup>19</sup> demonstrating that  $T_2$  contrast was enhanced with increasing NP size and magnetic field strength (Fig. 7b and c), while  $r_1$  value remained almost unchanged (Fig. 7a).

In a follow-up study, van Veggle and co-workers investigated in detail the involved relaxation mechanisms of NaDyF<sub>4</sub> and NaHoF<sub>4</sub> NPs in the range of 9–40 nm based on the outer sphere theory (Fig. 7d).<sup>20</sup> The larger particle had a larger volume fraction, and a three-fold increase in core size resulted in  $r_2$  values that were 30 times larger. Further theoretical simulation results showed that Curie contribution became predominant under an increasing magnetic field. Recently, we have reported the size effect of NaHoF<sub>4</sub> NPs on magnetic relaxation (Fig. 7e).<sup>21</sup> Both experimental and theoretical findings revealed that Curie contribution to contrast enhancement was the principal mechanism when particle size was smaller than 7 nm (Fig. 7f and g), as the  $\varphi(\tau_D)$  value from Curie relaxation remarkably decreased with increasing diffusion correlation time ( $\tau_D$ ) from 1 ns (3 nm NaHoF<sub>4</sub>) to 5 ns (7 nm NaHoF<sub>4</sub>). The  $\varphi(\tau_D)$  value then slightly decreased as  $\tau_D$  became larger, explaining the opposite change of  $r_2$  values that was observed when NP size was less than 7.4 nm. As the size of NPs increased to 13 nm and 29 nm, the contribution of dipolar mechanism could not be neglected, and all relevant mechanisms mentioned above contributed to MRI contrast enhancement.

### 3.2 Surface effects

Apart from the size effect on MRI relaxation, surface properties of NPs also play an important role in MRI contrast enhancement based on interaction between NPs and water molecules, which happens primarily on the surface of NPs. Individual contributions from the inner sphere, secondary sphere, and outer sphere depends on the CA's structure and composition, which can be regulated by surface modifications of tuned composition and thicknesses. Most NPs developed for *in vivo* MRI applications could be stabilized and functionalized with coating materials such as organic surface ligands, dense or mesoporous silica shells, or an inorganic nanolayer (*e.g.*,  $\text{Fe}_x\text{O}_y$  or  $\text{NaGdF}_4$  layers). These modifications, here termed surface effects, can in turn impact the relaxation of water protons in various configurations, such as hydration, hydrogen binding, and diffusion. For  $T_1$  CAs, surface effects can enhance proximity of water to center paramagnetic ions by tuning the distance of water protons to the NP surface, generating greater contrast than small-molecule CAs alone. Such a process depends on the species, chain length, density, and functional group of organic ligands, as well as the thickness of silica shells or inorganic nanolayers. For  $T_2$  CAs, proper surface modification of magnetic NPs will impact intrinsic surface spin disordering, increase inhomogeneity in the local magnetic field, or induce much stronger magnetization, leading to improvement in  $T_2$  contrast effects.

**3.2.1 Organic surface ligands**—Organic surface ligands are critical for the synthesis of NPs, not only for controllable growth during synthesis but also for ensuring their biocompatibility for *in vivo* application. An early study by Lee *et al.* suggested that different types of organic surface moieties greatly impacted  $T_1$  contrast enhancement.<sup>60</sup> Organic ligands including mercaptosuccinic acid (MSA), poly(maleic anhydride-*alt*-1-octadecene) (PMAO), PMAO-PEG, and Pluronic PF127 were capped on MnO NPs. The MnO@MSA exhibited the largest  $r_1$  values of  $2.55 \text{ mM}^{-1} \text{ s}^{-1}$  at pH 7.4 and  $6.94 \text{ mM}^{-1} \text{ s}^{-1}$  at pH 5.0 among all the coatings, indicating that  $r_1$  relaxivity of the NPs depended on the surface ligands.

Furthermore, the addition of polyethylene glycol (PEG) to the surface (*i.e.*, PEGylation) of NPs has been widely applied to improve biocompatibility of NPs, prolong *in vivo* circulation, and reduce the capture by the reticuloendothelial system (RES).<sup>152-155</sup> Recently, Johnson *et al.* reported a compact micellization strategy that could achieve ultra-high  $r_1$  values (around  $80 \text{ mM}^{-1} \text{ s}^{-1}$  at 1.4 T), about 25 times higher than that of clinically used  $\text{Gd}^{3+}$  chelates, by regulating the surface ligands (Fig. 8a–c).<sup>39</sup> The compactly packed PEGylated lipids would bring the water protons closer to the surface of  $\text{NaGdF}_4$  NPs and thus lead to enhanced relaxivity. They also found that the relaxivity was independent of outer PEG chain length (Fig. 8d and e), whereas increasing the inner core  $\text{NaGdF}_4$  size resulted in lowered relaxivity values (Fig. 8f and g). Two important criteria to achieve ultra-high relaxivity with  $\text{Gd}^{3+}$ -based NPs were proposed: (i) extremely small core NPs (<5 nm) and (ii) surface coating strategies that can render compact and individually stabilized NPs in water.

After surface PEGylation of ultra-small  $\text{NaGdF}_4$  nanodots, our group further conjugated DTPA on the surface of NPs to capture potentially released  $\text{Gd}^{3+}$  ions, effectively preventing

their liberation.<sup>40</sup> The  $\text{Gd}^{3+}$  ions captured by DTPA could directly interact with water protons to further enhance relaxivity. Recently, Yan and co-workers presented another simple strategy to improve  $r_1$  values by taking advantage of surface ligands' strong interaction with water protons.<sup>36</sup>  $\text{NaGdF}_4$  NPs coated with poly(acrylic acid) (PAA) exhibited largely enhanced  $r_1$  values compared to those coated by polyethylenimine (PEI) or PEG (Fig. 8h), which resulted from the strong hydrogen-bonding capacity of PAA to water molecules according to their theoretical calculation (Fig. 8i).

By rationally designing the functional group of surface ligands, the transverse  $r_2$  value can also be regulated for iron oxide NPs. Zeng *et al.* reported the surface modification of PEG2000 ligands with three different functional groups: diphosphate (DP-PEG), hydroxamate (HX-PEG), and catechol (CC-PEG) on 3.6 nm and 10.9 nm  $\text{Fe}_3\text{O}_4$  NPs (Fig. 9a).<sup>149</sup> The  $r_2$  values increased with increasing sizes for all three ligand-modified  $\text{Fe}_3\text{O}_4$  NPs, and the HX- or CC-PEG modified NPs showed much higher  $r_2$  values than the unconjugated functional group of DP-PEG for both sized  $\text{Fe}_3\text{O}_4$  NPs (Fig. 9b), while the  $r_1$  value only changed slightly with different anchored groups and NP sizes (Fig. 9c). The  $\pi$ - $\pi$  conjugation and p- $\pi$  conjugation from the CC and HX groups were believed to create increased inhomogeneity in the local magnetic field around the underlying  $\text{Fe}_3\text{O}_4$  NPs when compared to the unconjugated DP group (which has no  $\pi$ - $\pi$  or p- $\pi$  conjugation), which effectively accelerated transverse relaxation of the water protons nearby.

Besides the surface groups, the ligand species and the methods of modification also affect  $T_2$  contrast enhancement. The species effects of surface ligands on magnetic properties were verified by Duan *et al.* by coating ultra-small  $\text{Fe}_3\text{O}_4$  NPs with different levels of hydrophilic ligands including hyperbranched PEI, PEG grafted PEI (PEG-*g*-PEI), and poly(maleic acid)-octadecene (PMO).<sup>148</sup> The hydrophilicity of surface ligands was found to greatly impact the intrinsic surface spin disordering of NPs for water proton relaxation, where the most hydrophilic PEI coated  $\text{Fe}_3\text{O}_4$  NPs showed the highest  $r_2$  value, nearly three times higher than that of the least hydrophilic PMO-coated NPs. Among different methods for surface modification of  $\text{MnFe}_2\text{O}_4$  NPs such as ligand exchange, encapsulation with polymer shells, and embedding them into micelles (Fig. 9d), embedding NPs into lipid micelles was found to be the most efficient method for greatly enhancing  $r_2$  relaxation.<sup>147</sup> The larger-sized core inside lipid micelles and the contribution from dephasing of proton spins of the lipid molecules due to high viscosity of these molecules were considered to be responsible for this phenomenon.

The accessibility of water molecules to the magnetic core of NPs depends on the thickness of surface ligands, which can influence  $T_2$  relaxation of NPs. Tromsdorf *et al.* found that increasing PEG chain length from PEG350 to PEG550 on the surface of  $\text{Fe}_3\text{O}_4$  NPs led to decreased  $r_2$  and almost unchanged  $r_1$  values due to aggregation of NPs.<sup>141</sup> However, slight changes were found for both  $r_1$  and  $r_2$  values when PEG chain length was increased from 550 to 2000. Another report by Tong *et al.* indicated that  $\text{Fe}_3\text{O}_4$  NPs had a critical PEG ligand size, at which the  $r_2$  value changed dramatically.<sup>105</sup> The coating thickness could be well controlled by varying the size of DSPE-PEG from 500 Da to 5000 Da (Fig. 9e). Interestingly, the  $r_2$  value of 14 nm  $\text{Fe}_3\text{O}_4$  NPs increased by 2.5-fold when PEG molecular weight decreased from 5000 to 1000 Da (Fig. 9f). It was suggested that 14 nm iron oxide

NPs and PEG1000 surface modification was the optimal combination of core size and coating thickness according to  $T_2$  relaxivity on a per-particle basis (Fig. 9g). However,  $r_2$  values did not increase further when PEG size decreased to 750 and 550 Da. Such effects of PEG chain length on  $T_2$  relaxivity can be partly attributed to changes in the shell/core ratio, *i.e.*,  $L/r$  in eqn (3), under the assumption that the coating was water impermeable.

**3.2.2 Surface silica coatings**—The growth of a dense or mesoporous silica shell on the surface of MRI CAs has been a commonly used strategy since silica exhibits excellent biocompatibility, and the active surface or pore volume of the shells can be used for conjugating or loading biological molecules, functional particles, and therapeutic agents.<sup>156-158</sup> Our group extensively investigated the effect of both dense (dSiO<sub>2</sub>) and mesoporous (mSiO<sub>2</sub>) silica shells and their thickness on the MRI relaxation of Gd<sup>3+</sup>-based upconversion nanoparticles (UCNPs) (Fig. 10a).<sup>159</sup> The thicknesses of dSiO<sub>2</sub> and mSiO<sub>2</sub> shell could be well controlled and uniformly coated on each particle (Fig. 10b and c). It was found that both  $r_1$  and  $r_2$  values were independent of thickness of mSiO<sub>2</sub> shell, the thickness of which could be readily controlled from 11.1 nm to 23.4 nm. The dominant outer sphere mechanism determined by the translational diffusion of water molecules was thought to be responsible for this phenomenon, and the diffusion of water molecules in the mesoporous channels was not affected by the increased thickness of the mSiO<sub>2</sub> shell (Fig. 10d). However, the diffusion of water molecules in the dSiO<sub>2</sub> shell, in which large pore channels did not exist, became much more restricted by thicker dSiO<sub>2</sub> shells, leading to longer translational diffusion time and thus larger  $r_1$  and  $r_2$  values. Therefore, the  $r_1$  value of UCNPs@dSiO<sub>2</sub> increased from 4.03 mM<sup>-1</sup> s<sup>-1</sup> to 9.7 mM<sup>-1</sup> s<sup>-1</sup> as the thickness of dSiO<sub>2</sub> shell increased from 2.6 nm to 16.9 nm, while  $r_2$  values almost doubled (from 57.58 mM<sup>-1</sup> s<sup>-1</sup> to 102.08 mM<sup>-1</sup> s<sup>-1</sup>, Fig. 10e).

In another study, both mSiO<sub>2</sub> and dSiO<sub>2</sub> shells were used to coat the surface of MnO NPs.<sup>160</sup> With these modifications, the  $r_1$  value decreased substantially from 0.17 mM<sup>-1</sup> s<sup>-1</sup> for MnO NPs to 0.07 mM<sup>-1</sup> s<sup>-1</sup> for MnO@dSiO<sub>2</sub>. For MnO@mSiO<sub>2</sub> nanostructures, the  $r_1$  value was similar to MnO NPs (0.16 mM<sup>-1</sup> s<sup>-1</sup>), which was attributed to the existence of worm-like nanochannels in these NPs, where increased water exchange rates could be achieved. Similar results were also reported by Kim *et al.*, who found that mSiO<sub>2</sub>-coated MnO NPs showed a much higher  $r_1$  value of 0.65 mM<sup>-1</sup> s<sup>-1</sup> than those of dSiO<sub>2</sub>-coated NPs ( $r_1 = 0.08$  mM<sup>-1</sup> s<sup>-1</sup>) and PEGylated MnO NPs ( $r_1 = 0.11$  mM<sup>-1</sup> s<sup>-1</sup>). The  $r_1$  value could be further increased to 0.99 mM<sup>-1</sup> s<sup>-1</sup> after acid etching of MnO NPs to generate a hollow nanostructure (*i.e.*, HMnO@mSiO<sub>2</sub>) (Fig. 10f and g).<sup>62</sup> Such a hollow structure allowed more interior Mn<sup>2+</sup> ions to interact with water molecules on the inner surface, resulting in a significantly higher  $r_1$  value than that of MnO NPs encapsulated with PEG phospholipids.<sup>69</sup> Therefore, intracranial transplantation of HMnO@mSiO<sub>2</sub>-labeled mesenchymal stem cells could be monitored by serial MRI over 14 days (Fig. 10h). Among a number of surface modifications strategies of NPs with organic ligands and silica, Lee *et al.* found that silica-coated MnO exhibited the highest stability at neutral pH and the slowest dissolution rate when dispersed in acetate buffer solutions and showed continuous contrast enhancement in the rat brain.<sup>60</sup>

**3.2.3 Inorganic surface nanolayers**—Coating an inorganic nanolayer (e.g., Fe<sub>2</sub>O<sub>3</sub>, Fe<sub>3</sub>O<sub>4</sub>, NaGdF<sub>4</sub>, and so forth) on the surface of NPs to induce stronger magnetization and ensure adequate exposure to water molecules nearby is an attractive strategy to obtain high-relaxivity  $T_2$  or  $T_1$  CAs. MR relaxation enhancement effects of  $\gamma$ -Fe<sub>2</sub>O<sub>3</sub> nanolayers were first reported by Xu and co-workers based on core/shell, yolk/shell, or hollow structures (Fig. 11a–d).<sup>77</sup> Field-dependent magnetization measurements showed that all these NPs were superparamagnetic (Fig. 11e), and the saturated magnetization ( $M_s$ ) of the  $\gamma$ -Fe<sub>2</sub>O<sub>3</sub> hollow nanostructure was the highest among these nanostructures (about 50.8 emu g<sup>-1</sup>). Additionally, both the growth of a 5 nm Fe<sub>x</sub>O<sub>y</sub> layer and self-assembly of ultra-small Fe<sub>3</sub>O<sub>4</sub> on the surface of UCNPs generated much higher  $r_2$  values than conventional iron oxide NPs.<sup>161,162</sup> Cheong *et al.* synthesized iron/iron oxide core/shell (Fe@Fe<sub>3</sub>O<sub>4</sub>) NPs by oxidizing the surface of iron nanocrystals in the air.<sup>163</sup> These core/shell Fe@Fe<sub>3</sub>O<sub>4</sub> NPs had a  $r_2$  value of 324 mM<sup>-1</sup> s<sup>-1</sup>, more than twice that of Fe<sub>3</sub>O<sub>4</sub> NPs (145 mM<sup>-1</sup> s<sup>-1</sup>, Fig. 11h and i). Such an increase in  $r_2$  value was attributed to the presence of a single-crystal iron core, which can induce much stronger magnetization than Fe<sub>3</sub>O<sub>4</sub> NPs.

The coating of a NaGdF<sub>4</sub> layer is an intriguing approach not only to suppress surface quenching of UCNPs for enhanced luminescence emission but also to endow NPs with MRI capabilities, creating a dual-modality imaging system.<sup>125,165,166</sup> Results from different groups have suggested that the presence of Gd<sup>3+</sup> ions on the surface of NPs might play a major role in accelerating relaxation, whereas the contribution from inner Gd<sup>3+</sup> ions could be neglected; however, no direct evidence has been found to support this hypothesis.<sup>46,48,49</sup> To investigate this issue, our group designed Gd<sup>3+</sup>-doped core NPs with multilayer coating of NaYF<sub>4</sub> and/or NaGdF<sub>4</sub>.<sup>164</sup> The  $r_1$  value was calculated to be 1.24 mM<sup>-1</sup> s<sup>-1</sup> for the core NPs, which decreased to nearly zero ( $r_1 = 0.037$  mM<sup>-1</sup> s<sup>-1</sup>) after coating a ~4 nm NaYF<sub>4</sub> layer on the core NPs, and recovered to 1.18 mM<sup>-1</sup> s<sup>-1</sup> after further coating with a NaGdF<sub>4</sub> layer, clearly indicating the different effects between the bulk and surface Gd<sup>3+</sup> ions in shortening  $T_1$  relaxation time.

To further validate the hypothesis that Gd<sup>3+</sup> ions in the inner crystal lattice contributed little to longitudinal relaxivity, Gd<sup>3+</sup>-free core NPs were coated with a NaGdF<sub>4</sub> layer of varied thicknesses (Fig. 12). With an increase in NaGdF<sub>4</sub> layer thickness from approximately 0.2 nm to 3.7 nm,  $r_1$  values decreased gradually from 6.18 mM<sup>-1</sup> s<sup>-1</sup> to 2.18 mM<sup>-1</sup> s<sup>-1</sup>, confirming that the main contribution to shortening  $T_1$  values was by surfacebound Gd<sup>3+</sup> ions on the NPs. Subsequently, these NaGdF<sub>4</sub> coated UCNPs were applied as dual-targeted nanoprobes by conjugating angiopep-2 peptide to cross the blood–brain barrier and target brain tumors for MRI-aided preoperative diagnosis and intraoperative identification of brain tumors by upconversion luminescence in an animal model.<sup>166</sup>

### 3.3 Shape effects

The shape effects of NPs can be as important as the size and surface effects in determining the properties of nanomaterials. The unusual shapes of these NPs could potentially dictate their interactions with single cells or whole organisms.<sup>167</sup> Shape-dependent relaxivity of NP-based MRI CAs has been revealed by many different groups, where shape anisotropy of magnetic NPs or shape confinement of Gd<sup>3+</sup>-based CAs can improve the relaxation of

nearby water molecules. Employing the highly controllable shape of spherical or cubic magnetic  $\text{CoFe}_2\text{O}_4$  NPs, Song *et al.* found that shape was an important factor in determining coercivity of magnetic NPs due to the varying effects of surface anisotropy.<sup>82</sup> Very strong localized magnetic field inhomogeneity was produced by shape anisotropy, and the generated magnetic field strength was found to be dependent on the shape of magnetic NPs due to demagnetization effect.<sup>168</sup> Shape-dependent enhancement of  $r_2$  values was also observed with superparamagnetic NPs.<sup>80,81,137</sup> For example, faceted irregular cobalt ferrite nanostructures showed increased  $r_2$  values when compared to spherical NPs of similar volume.<sup>81</sup> Zhao *et al.* synthesized branched octapod-shaped iron oxide NPs (edge length of 30 nm, Fig. 13a) with an ultra-high  $r_2$  relaxivity of  $\sim 679.3 \text{ mM}^{-1} \text{ s}^{-1}$ ,<sup>137</sup> which was a much more effective  $T_2$  CA for small tumor detection than conventional spherical iron oxide NPs (Fig. 13d and e) because of their unique morphology and larger effective boundary radius (Fig. 13c). The significant enhancement was attributed to an increase in the effective radius of the magnetic core, a parameter that is impacted by shape anisotropy and directly affects  $r_2$  values (Fig. 13b).<sup>80,137</sup>

Recently, Mohapatra *et al.* suggested that although spherical iron oxide NPs exhibited higher magnetization than iron oxide nanorods (NRs) of the same volume,<sup>80</sup> high relaxivity ( $r_2 = 608 \text{ mM}^{-1} \text{ s}^{-1}$ ) was obtained for 50 nm NRs (Fig. 13g and h), while the relaxivity of the same volume of spherical NPs (the volume of a 50 nm NR was equivalent to that of a 16 nm NP) was only  $297 \text{ mM}^{-1} \text{ s}^{-1}$ . This phenomenon was believed to be a result of the larger surface area of the anisotropic morphology for NRs. Using the software Comsol Multiphysics, the induced local field inhomogeneity surrounding the 50 nm NRs was calculated to be significantly stronger than that induced by 16 nm NPs (Fig. 13g and h). The higher outer sphere diameter of rod-shaped NPs resulted in a stronger local magnetic inhomogeneity with a larger number of water protons experiencing rapid dephasing, which gave a higher  $r_2$  value. On the other hand, spherical NPs exhibited a smaller volume of magnetic field variation. Therefore, water protons diffused slowly and thus dephased slowly, resulting in a much lower  $r_2$  value than NRs of equivalent volume (Fig. 13i).<sup>80</sup>

Importantly, the shape effects of NPs on  $r_2$  values cannot be translated to  $r_1$  values because mechanisms for enhancements in  $r_2$  and  $r_1$  are different. An early study from Hyeon and co-workers demonstrated the significant effects of shape on  $r_1$  values using spherical and tetrahedral  $\text{Mn}_3\text{O}_4$  NPs.<sup>132</sup> Recently, Rotz *et al.* described the surface conjugation of DNA-Gd onto gold nanostars and nanospheres<sup>95</sup> and found that DNA-Gd@stars exhibited remarkable  $r_1$  values up to  $98 \text{ mM}^{-1} \text{ s}^{-1}$ , much higher than those of DNA-Gd@spheres, where  $r_1$  values were 25 times larger than FDA-approved Gd chelates at  $25^\circ\text{C}$  (Fig. 14a and b). Nuclear magnetic relaxation dispersion (NMRD) analysis indicated that secondary-sphere relaxivity contributed significantly to the production of such an exceptionally high  $r_1$  value of DNA-Gd@stars, while secondary-sphere enhancements were not observed for DNA-Gd@spheres (Fig. 14c). The shape of gold nanostars with hydrophilic microenvironment played a meaningful role in the sequestration of water molecules around the  $\text{Gd}^{3+}$  complexes, which facilitated longer access to secondary-sphere water molecules and thereby generated greater contrast enhancement. Further studies from the same group reported that  $r_1$  depended on the number of branches on the nanostars, as  $r_1$  increased with increasing average number of branches (Fig. 14d).<sup>94</sup> Researches from many other groups

have also demonstrated that the shape of NPs can greatly impact their cellular uptake, biodistribution, clearance, and biocompatibility *in vivo*.<sup>169-173</sup> Thus, an in-depth understanding of the effects of shape on relaxation can enable the rational design of high-relaxivity CAs with outstanding performance for *in vivo* applications.

## 4. Strategies to achieve high relaxivity for MRI contrast enhancement

Size control, surface modifications (*e.g.*, organic surface ligands, silica shells, and inorganic nanolayers), and shape regulation at the nanoscale have been widely applied to regulate the behavior of CAs, resulting in largely accelerated relaxation for efficient  $T_1$  or  $T_2$  MRI contrast enhancement. A deep understanding of the involved nucleation, growth, and surface chemistry of NPs enables intelligent engineering of the key factors that affect relaxation, either independently or in combination. The development of such high-relaxivity CAs is crucial for molecular MRI applications such as early-stage disease detection,<sup>174-176</sup> gene expression monitoring,<sup>177,178</sup> drug release,<sup>140,179-181</sup> and cell tracking.<sup>182-185</sup> In addition to the traditionally investigated factors influencing the performance of CAs at the nanoscale, several novel strategies have been exploited to obtain high-relaxivity CAs, such as confining CAs in mesopores, doping with metal ions, self-assembly or disassembly, and so forth (Fig. 15).

### 4.1 Confinement in mesopores

The accessibility of water molecules to the paramagnetic center is highly important for the design of high-relaxivity CAs. Mesoporous silica nanoparticles (MSNs) provide an ideal platform for ready access of water molecules to the center due to their tunable pore structures and high surface areas. Dispersing CAs such as  $Gd^{3+}$  chelates,<sup>90-92</sup> MnO NPs,<sup>61,64,139</sup> and iron oxide NPs<sup>143</sup> into the mesopores of MSNs can efficiently prolong  $\tau_R$  and thus greatly increase relaxivity. For example, Lin *et al.* loaded Gd–Si–DTPA complexes into the hexagon-ordered nanochannels (~2.4 nm diameter) of MSNs, obtaining extremely high  $r_1$  values of  $10.2 \text{ mM}^{-1} \text{ s}^{-1}$  at 9.4 T and  $28.8 \text{ mM}^{-1} \text{ s}^{-1}$  at 3.0 T.<sup>91</sup> The geometrical confinements of  $Gd^{3+}$  chelates, gadofullerenes, and gadonanotubes within the porous structure of silicon microparticles (SiMPs) (Fig. 16a) were all reported by Ananta *et al.*,<sup>186</sup> demonstrating that such a strategy can successfully influence the paramagnetic behavior of  $Gd^{3+}$  ions, resulting in increased  $r_1$  values. By prolonging  $\tau_R$  and reducing the mobility of outer sphere water molecules,  $r_1$  values about 4–50 times larger than their original values without SiMPs were obtained (Fig. 16b).

Recently, the use of such geometrical confinement for enhanced relaxivity was further confirmed by Gao and co-workers.<sup>187</sup> Ultrasmall  $Gd_2O_3$  NPs with numerous coordinating water molecules were dispersed into worm-like mesoporous channels of MSNs, which, by increasing the tumbling time and creating a geometrically confined space, remarkably increased  $r_1$  value from  $14.93 \text{ mM}^{-1} \text{ s}^{-1}$  to  $45.08 \text{ mM}^{-1} \text{ s}^{-1}$  at 0.5 T. (Fig. 16c). Thus, confinement of CAs in a porous structure offers a new and widely applicable strategy for strong  $T_1$  contrast enhancement, where mesopores also hold the potential for drug loading and delivery for therapeutic applications.

As described above, loading small MnO or iron oxide NPs within the mesopores of MSNs was demonstrated to be more beneficial for relaxation enhancement than traditional silicacoated structures because such coating layers shielded the surface paramagnetic centers from interaction with the water molecules.<sup>139,143</sup> To further enhance the  $r_1$  value of MnO NPs to exceed that of clinical Gd<sup>3+</sup> chelates (*e.g.* Magnevist,  $r_1 = \sim 3.4 \text{ mM}^{-1} \text{ s}^{-1}$ ), we loaded small MnO NPs into hollow mesoporous silica nanoparticles (HMSNs) *via* an *in situ* redox reaction to further facilitate the interaction between manganese paramagnetic centers and water protons (Fig. 16d).<sup>61</sup> Using this strategy, we were able to obtain a maximum  $r_1$  value of  $8.81 \text{ mM}^{-1} \text{ s}^{-1}$ , which is more than twice that of commercial Gd<sup>3+</sup> chelates (Fig. 16e). In addition, the well-defined mesoporosity of HMSNs can enable efficient encapsulation of photosensitizers (*i.e.*, protoporphyrin) for therapeutic purposes, as well as further confinement of paramagnetic Mn ions based on metalloporphyrin chelating chemistry.<sup>188</sup> This mesoporous structure endowed these intriguing sensitizers with remarkably enhanced MRI performance ( $r_1 = 9.43 \text{ mM}^{-1} \text{ s}^{-1}$ ).

By a facile but efficient physical-vapor-infiltration method, ultra-small iron oxide NPs have also been confined to the mesoporous networks of MSNs.<sup>143</sup> In particular, the highly dispersed iron species in the mesoporous channels can be used as anchoring sites for anti-cancer drugs, such that gradual exposure of these iron paramagnetic centers along with controllable drug release from the mesopores enables monitoring of drug release *in vivo* through synchronously intensified MRI contrast. In another recent report, Zhang *et al.* incorporated both GdPO<sub>4</sub> and g-Fe<sub>2</sub>O<sub>3</sub> NPs in ordered mesoporous carbon composites, which exhibited satisfactory  $r_1$  and  $r_2$  values of 2.7 and  $183.7 \text{ mM}^{-1} \text{ s}^{-1}$ , respectively.<sup>189</sup>

## 4.2 Metal alloying and metal ion doping

The MRI performance of magnetic NPs is strongly dependent on their magnetic characteristics. Metal alloy-based NPs such as FeCo and FePt NPs are another class of CAs,<sup>77,146,190-194</sup> which possess magnetic spins that align in parallel to the external magnetic field, resulting in typically higher magnetic moments than traditional ferromagnetic NPs. For example, the magnetic moment of FeCo alloys is approximately  $2.4 \mu_B$  per magnetic atom, while that of Fe<sub>3</sub>O<sub>4</sub> NPs is approximately  $1.3 \mu_B$ .<sup>146</sup> These FeCo alloys have a comparably high  $M_s$  of 215 emu per gram of metal when coated with a graphite shell, leading to an exceptionally high  $r_2$  value of  $644 \text{ mM}^{-1} \text{ s}^{-1}$ , much larger than that of conventional Fe<sub>3</sub>O<sub>4</sub> NPs (usually about  $100 \text{ mM}^{-1} \text{ s}^{-1}$ ). High X-ray absorption of Pt ( $6.95 \text{ cm}^2 \text{ g}^{-1}$  at 50 keV) makes FePt alloy NPs promising candidates for MRI/CT dual-modality CAs.<sup>194</sup> By conjugating FePt NPs with an antibody targeting Her2/neu, *in vivo* dual-modality imaging was achieved in cancer lesions in an animal model. Recently, it was reported that FePt NPs could also be used as an infrared-active nanoplatform for four imaging modalities (optical/photoacoustic/CT/MRI), which may allow for multiscale imaging from the cellular level to wholebody imaging.<sup>192</sup>

According to eqn (3), the  $r_2$  value of magnetic NPs is proportional to their  $M_s$ . Thus, metal ion doping has been verified to be an effective method to engineer the magnetism of iron oxide NPs and obtain extremely high  $r_2$  values, as investigated in detail by Cheon and co-workers.<sup>84</sup> They synthesized various types of MFe<sub>2</sub>O<sub>4</sub> NPs in which Fe<sup>2+</sup> ions can be



readily replaced by other transition-metal dopants  $M^{2+}$  where  $M = \text{Mn, Ni, or Co}$ . The magnetic moments per unit of  $\text{MnFe}_2\text{O}_4$ ,  $\text{CoFe}_2\text{O}_4$ , and  $\text{NiFe}_2\text{O}_4$  were estimated to be  $5 \mu_B$ ,  $3 \mu_B$ , and  $2 \mu_B$ , respectively, leading to gradually decreasing  $M_s$  values from 110 emu per gram of metal for  $\text{MnFe}_2\text{O}_4$  NPs to 85 emu per gram of metal for  $\text{NiFe}_2\text{O}_4$  NPs (Fig. 17a). As a result,  $\text{MnFe}_2\text{O}_4$  NPs exhibited the highest  $r_2$  value of  $358 \text{ mM}^{-1} \text{ s}^{-1}$ , two-fold higher than that of conventional iron oxide NPs. Such an enhancement is significant for clinical applications as the dosages of nanoparticulate CAs can potentially be decreased. By further co-doping  $\text{Zn}^{2+}$  and  $\text{Mn}^{2+}$  ions in  $(\text{Zn}_x\text{Mn}_{1-x})\text{Fe}_2\text{O}_4$  NPs, the same group obtained NPs that exhibited an extremely high magnetization value ( $175 \text{ emu g}^{-1}$ ) and provided the largest  $r_2$  value ( $860 \text{ mM}^{-1} \text{ s}^{-1}$ ) among the CAs reported to date (Fig. 17c).<sup>83</sup> In this study, it was also found that the magnetism was dependent on amount of metal ion doping and  $x = 0.4$  doping of  $\text{Zn}^{2+}$  demonstrated the best contrast (Fig. 17b).

### 4.3 Self-assembly and disassembly

“Activatable” or “smart” nanoprobcs with tunable MRI relaxivity through self-assembly or disassembly in the presence of certain pathophysiological parameters can serve as high-relaxivity CAs, resulting in enhanced detection sensitivity of MRI for specific molecular or biological events. For example, such a self-assembly strategy allows for small NPs to circulate in the blood pool over a prolonged time duration, which can then be triggered in the tumor to self-assemble into a large structure for enhanced MRI. The trigger may include many physiological characteristics, such as molecular interactions,<sup>195</sup> enzymes,<sup>196</sup> receptors,<sup>197</sup> and redox status.<sup>198</sup> On the other hand, the disassembly strategy utilizes the different tumor microenvironments to release CAs preconfined in NPs, also achieving amplified MRI contrast. Both strategies are highly versatile and useful as they could be potentially used to monitor *in vivo* biological processes of the tumor microenvironment and improve the accuracy of cancer diagnosis or monitoring the therapeutic efficacy.

**4.3.1 Self-assembly**—Liang and co-workers reported the controlled self-assembly of  $\text{Gd}^{3+}$ -based NPs through the condensation reaction between 1,2-aminothiol and 2-cyanobenzothiazole.<sup>199</sup> The probe was activated through disulfide reduction by tris(2-carboxyethyl)phosphine or glutathione (GSH) in cells and generated intermediate 1,2-aminothiol groups to react with the cyano groups on 2-cyanobenzothiazole, which results in self-assembly of large  $\text{Gd}^{3+}$  NPs with enhanced relaxivity. Presumably due to an increased rotational correlation time, a more than 2-fold increase in  $r_1$  value was observed for the larger assembly when compared to its small molecule form of  $\text{Gd}^{3+}$  chelates. Through further optimization of self-assembly strategy by intermolecular macrocyclization chemistry, the same group obtained biothiol-specific  $\text{Gd}^{3+}$ -based CAs with amplified relaxivity after self-assembly, which were capable of detecting GSH concentrations with high sensitivity using MRI.<sup>200</sup> In addition to redoxactivatable self-assembly, Ye *et al.* reported novel caspase-3/7 activated  $\text{Gd}$ -based CAs, which could self-assemble into  $\text{Gd}^{3+}$  NPs upon reduction and activation by caspase-3/7, an important early biomarker for apoptosis.<sup>201</sup> Higher  $r_1$  relaxivity and longer retention time in apoptotic tumors were found with such self-assembling CAs, suggesting the potential for MRI of *in vivo* enzyme activity.

Self-assembly of small magnetic NPs into nanoclusters to increase effective magnetic size is also a widely used strategy for enhancing their magnetic properties and thus obtaining a higher  $r_2$  value. Using charge compensation and electrostatic adsorption between oppositely charged moieties, Berret *et al.* assembled 6.3 nm iron oxide NPs into larger nanoclusters, resulting in remarkably increased  $r_2$  values from  $39 \text{ mM}^{-1} \text{ s}^{-1}$  for single NPs to  $71 \text{ mM}^{-1} \text{ s}^{-1}$  for 70 nm nanoclusters and  $162 \text{ mM}^{-1} \text{ s}^{-1}$  for 170 nm nanoclusters.<sup>202</sup> The detection sensitivity of MRI using these NPs was thus substantially improved. A similar effect was also reported by Ai *et al.* in their first proof-of-concept study regarding clustering of monodispersed iron oxide NPs inside the core of polymeric micelles.<sup>203</sup> When compared to single core nanoparticles, multi-core nanoclusters with an average core size of 50 nm (self-assembled from 13 nm iron oxide NPs) provided much higher  $M_s$  values and subsequently exhibited higher performance in MRI contrast enhancement.<sup>204</sup>

Although larger sized iron oxide NPs present enhanced magnetic properties for MRI contrast enhancement, an ideal scenario is the self-assembly of small iron oxide NPs into larger nanostructures at the target sites (*e.g.*, tumor) for substantially amplified relaxation, since smaller NPs (below renal clearance threshold of approximately 10 nm) exhibit longer blood circulation half-life and hence are cleared more slowly from the blood pool than their larger counterparts. Mao and co-workers reported that ultrafine iron oxide NPs ( $\sim 3.5$  nm) with high  $T_1$  relaxivity exhibited prolonged *in vivo* circulation life time and could be self-assembled or clustered into larger ones in the tumor, thereby switching from bright  $T_1$  contrast to dark  $T_2$  contrast in MRI.<sup>205</sup> In another report, two complementary iron oxide NPs that were linked to a peptide substrate for enzymatic cleavage and a functional group (azide or alkyne) for click reaction<sup>196</sup> were able to self-assemble into a larger nanocluster network in the presence of matrix metalloproteinase enzymes, leading to approximately 160% enhancement of  $T_2$  signal intensity.

Recently, Liang and co-workers reported a caspase-3/7-instructed biorthogonal cyclization to self-assemble  $\text{Fe}_3\text{O}_4@1$  NPs in apoptotic cells and tumors for greatly enhanced  $T_2$  MRI (Fig. 18a).<sup>197</sup> Caspase-3/7-mediated self-assembly of  $\text{Fe}_3\text{O}_4@1$  NPs ( $r_2 = 185 \text{ mM}^{-1} \text{ s}^{-1}$ ) induced an approximately 65% increase of  $r_2$  value when compared to the control group ( $r_2 = 112 \text{ mM}^{-1} \text{ s}^{-1}$ ) (Fig. 18b). More importantly, *in vivo* MRI revealed that  $\text{Fe}_3\text{O}_4@1$  NPs provided specific  $T_2$  enhanced contrast in apoptotic tumors, whereas the control  $\text{Fe}_3\text{O}_4@1$ -Scr NPs showed no contrast enhancement (Fig. 18c). In another recent study, Gao and co-workers synthesized  $\text{Fe}_3\text{O}_4$  NPs modified with a responsive peptide sequence,<sup>198</sup> which enabled GSH-triggered self-assembly of  $\text{Fe}_3\text{O}_4$  NPs within the tumor microenvironment and the formation of aggregates through inter-particle cross-linking reactions (Fig. 18d and e). Both *in vitro* and *in vivo* experiments demonstrated that such self-assembly strategies substantially enhanced MRI performance (Fig. 18f).

**4.3.2 Disassembly**—In general, disassembly strategies switch the relaxivity of CAs from “off” to “on” after responding to pathophysiological parameters such as low pH value and high GSH concentration within the tumor, thereby maximizing the signal from region-of-interest and minimizing the signal from background tissues to improve specificity and sensitivity of CAs. For example, Viger *et al.* encapsulated a large number of ultra-small  $\text{Gd}_2\text{O}_3$  NPs in biodegradable polymer capsules, which showed extremely weak MRI contrast

due to effective silencing of both inner and outer sphere relaxation.<sup>131</sup> The  $r_1$  of the ultra-small  $\text{Gd}_2\text{O}_3$  NPs could be recovered after disassembling the large capsules under acidic conditions and/or in the presence of  $\text{H}_2\text{O}_2$ , which are both characteristic of many diseases such as cancer (Fig. 19a). In this study, encapsulation of ultra-small  $\text{Gd}_2\text{O}_3$  NPs within pH-responsive polymer particles resulted in deactivation of their contrast enhancement at pH 7.4. However, burst degradation was observed within 1 min at pH 6.5 (approximately the extracellular pH of diseased tissue), leading to an instantaneous 4.2-fold increase of  $r_1$  value, which remained constant thereafter (Fig. 19b). This concept was also verified using a peroxide-degradable polymer for encapsulation of ultra-small  $\text{Gd}_2\text{O}_3$  NPs, indicating a general strategy of disassembly (Fig. 19c).<sup>131</sup> A redox-modulated disassembly of  $\text{Gd}^{3+}$ -based nanosystem was developed by Martinelli *et al.*<sup>206</sup> The disassembly was achieved *via* the cleavage of S–S bridges in response to reducing environments, leading to the release of  $\text{Gd}^{3+}$  complexes and a subsequent significant change in relaxivity.

Similar to the above-mentioned strategies used for Gd-based NPs, the release of  $\text{Mn}^{2+}$  from the disassembled matrix at low pH sites was also employed to obtain high-relaxivity CAs for pH-responsive MRI.<sup>61,63</sup> For example, Kim *et al.* reported core–shell  $\text{Fe}_3\text{O}_4@ \text{Mn}_3\text{O}_4$  NPs, where  $\text{Mn}^{2+}$  ions released from the  $\text{Mn}_3\text{O}_4$  shell under reducing conditions allowed  $\text{Fe}_3\text{O}_4$  NPs to interact with water protons, thereby enhancing contrast of both  $T_1$  and  $T_2$ -weighted MRI.<sup>207</sup> By confining  $\text{Mn}^{2+}$  within pH-sensitive calcium phosphate NPs (PEGMnCaP NPs, Fig. 20a), Mi *et al.* found a remarkable increase of  $r_1$  from  $4.73 \text{ mM}^{-1} \text{ s}^{-1}$  to  $19.96 \text{ mM}^{-1} \text{ s}^{-1}$  after exposure to acidic environments simulating intratumoral pH, wherein the pH-sensitive CaP was disassembled (Fig. 20b).<sup>52</sup> The enhanced MRI signal by the disassembled CaP matrix and released  $\text{Mn}^{2+}$  was attributed to the increase in relaxivity after binding of  $\text{Mn}^{2+}$  with the protein. Such  $\text{Mn}^{2+}$ -triggered signal amplification could be further used for the precise imaging of tumors. After intravenous injection of PEGMnCaP NPs, enhanced MRI signal in tumor areas could be clearly visualized within 6 h after injection. Such a release of  $\text{Mn}^{2+}$  from the nanocomposites allowed these NPs to sensitively detect millimeter-sized metastatic tumors in the liver of mice, which were otherwise invisible (Fig. 20c). The signal of the gallbladder gradually increased due to hepatic clearance of PEGMnCaP. As hypoxic regions in tumors usually exhibit a lower local pH than surrounding tissues, the contrast in hypoxic regions was 20–30% higher than that in normally oxygenated regions within the tumors after intravenous injection of PEGMnCaP NPs (Fig. 20d and e).

By doping  $\text{Mn}^{2+}$  ions into HMSN (Mn–HMSNs), we have recently reported a strategy called “manganese extraction” to promote the disassembly of Mn–HMSNs and enhance MRI contrast (Fig. 20f).<sup>208</sup> Such a  $\text{Mn}^{2+}$  doping approach enabled accelerated disassembly of Mn–HMSNs in either mildly acidic or reducing microenvironments by breaking the Mn–O bond and subsequently triggering “manganese extraction,” which further accelerated breakage of Si–O–Si bonds in the framework and biodegradation of HMSNs (Fig. 20g and h). Such a fast disassembly of Mn–HMSNs enhanced the interaction probability of Mn paramagnetic centers with water molecules, allowing pH- and GSH-responsive MRI to be realized (Fig. 20i and j). The disassembly of Mn–HMSNs also accelerated the release of anticancer drugs to obtain tumor-sensitive theranostic functions.

#### 4.4 Other strategies

**Valency change**—In principle, any physical or chemical method to change the interaction between CAs and water molecules such as the number of bound water molecules ( $q$ ), rotational correlation time ( $\tau_R$ ), and water residence time ( $\tau_M$ ) of  $T_1$  CAs or to enhance magnetic properties of  $T_2$  CAs will undoubtedly result in enhanced MRI relaxivity of CAs. One major strategy for MRI contrast enhancement focuses on valency changes of ions, which has been widely applied through using redox-activated chemical exchange saturation transfer (CEST) agents that can be detected by MRI.<sup>85,86,209</sup> As previously discussed, the valency of ions affects electron spin that in turn influences relaxation of water protons. Owing to its pH- and redox-responsive properties,  $MnO_2$  can be reduced to  $Mn^{2+}$  by GSH or  $H_2O_2/H^+$ .<sup>51,53,210</sup> For example, Chu *et al.* reported that  $MnO_2$  nanosheets could serve as highly effective generators to trigger self-assembly of photosensitizers (sinoporphyrin sodium, DVDMS) after reduction by GSH and  $H_2O_2$  (Fig. 21a).<sup>211</sup> Such processes could be monitored by *in vitro* and *in vivo* MRI (Fig. 21b and c), which was validated by TEM imaging of *ex vivo* tumor tissues (Fig. 21d).

Research from our group also suggested that  $MnO_2$  nanosheets could generate oxygen to overcome the hypoxic environment of tumors by reacting with  $H_2O_2$ .<sup>53</sup> and the released  $Mn^{2+}$  could be used for pH-responsive MRI.<sup>55</sup> Zhao *et al.* designed a redox-responsive  $MnO_2$  nanosheet, in which manganese atoms were coordinated in an octahedral geometry to six oxygen atoms that prevented direct contact with the aqueous environment, therefore making no contributions to the alteration of relaxation of water protons.<sup>54</sup> Intracellular GSH could then reduce endocytosed  $MnO_2$  nanosheets, generating numerous  $Mn^{2+}$  ions for MRI with greatly enhanced  $r_1$  and  $r_2$  values.

**Oxygen vacancy**—To date, the strategy of increasing the number of bound water molecules has not been employed in inorganic  $Gd^{3+}$ -based NPs. Recently, we introduced oxygen vacancy into  $Gd^{3+}$ -based NPs by designing a novel  $Gd^{3+}$ -doped tungsten bronze (termed as  $PEG-Na_xGdWO_3$  nanorods).<sup>37</sup> For the first time, controllable oxygen vacancies in inorganic nanosized CAs were generated for accelerated proton relaxation, favorably resulting in a substantially enhanced  $r_1$  value due to the oxygen vacancies' natural affinity to oxygen atoms in water molecules.<sup>212</sup> Such a strategy was validated with oxygen-deficient  $PEG-Na_xGdWO_3$  nanorods, which exhibited extraordinarily high  $r_1$  values up to  $80 \text{ mM}^{-1} \text{ s}^{-1}$  at 0.7 T and  $32.1 \text{ mM}^{-1} \text{ s}^{-1}$  on a clinical 3.0 T scanner (Fig. 21e and f), offering excellent blood pool imaging and tumor diagnosis capabilities at a low CA dosage. NMRD profiles indicated that  $r_1$  value was apparently increased by changes in the distribution of the number of bound water molecules ( $q$ ) either through secondary- or outer-sphere effects (Fig. 21g).

**Distance regulation**—Recently, Choi *et al.* have introduced a distance-dependent magnetic resonance tuning system based on the interaction between a superparamagnetic 'quencher' and a paramagnetic 'enhancer', where  $T_1$  signal was tuned "on" or "off" depending on separation distance between the enhancer and quencher.<sup>213</sup> As the enhancer moved far away from the quencher ( $d_1 > d_c$ , where  $d_c$  is the critical separation distance), the electron spin fluctuation of the enhancer accelerated water proton relaxation, resulting in a

stronger  $T_1$  signal (“on” state); however, when the enhancer and quencher were in close proximity to one another ( $d_2 < d_c$ ), the resulting spin fluctuation became too slow to induce effective water proton relaxation, resulting in a low  $T_1$  signal “off” state (Fig. 22a). The feasibility of this strategy was demonstrated through the control of thickness of  $\text{SiO}_2$  shell, thereby separating the enhancer and quencher from 18 to 2 nm (Fig. 22b). The  $d_c$  was set to 7 nm, where the  $T_1$  relaxivity quenching efficiency was 50% (Fig. 22c). Measurements from X-band electron paramagnetic resonance and NMRD profiles revealed distance-dependent values of correlation time ( $\tau_c$ ) and  $T_{1e}$  (for details, please see eqn (2)), indicating that the enhancer electron spin fluctuation became slower as separation distance decreased. Furthermore, such a distance-dependent tuning effect of MRI signal was verified using three representative models of molecular interactions (*i.e.*, cleavage, binding, and folding/unfolding) *via* modular combinations of enhancers and quenchers through interactive linkers (Fig. 22d).

**Tuning inhomogeneity**—A novel strategy to greatly enhance  $r_2$  relaxivity by tuning magnetic field inhomogeneity of iron oxide nanoclusters has been recently reported by Chen and co-workers.<sup>214</sup> Using both experimental validations and theoretical simulations, the group found that the clustering of iron oxide NPs could enhance local field inhomogeneity due to reduced field symmetry. Clusters of mixed 5 nm plus 15 nm iron oxide NPs exhibited a much higher  $r_2$  value of  $533.4 \text{ mM}^{-1} \text{ s}^{-1}$ , when compared to clusters containing 5 nm NPs only ( $r_2 = 231.6 \text{ mM}^{-1} \text{ s}^{-1}$ ) or 15 nm NPs only ( $r_2 = 358.3 \text{ mM}^{-1} \text{ s}^{-1}$ ) (Fig. 23a), indicating the important roles of particle distribution and fraction in generating local field inhomogeneity. By artificially involving iron oxide NPs with heterogeneous geometries (*e.g.*, iron oxide cubes and plates to assemble clusters) (Fig. 23b and c), the levels of field inhomogeneity were further elevated when compared to iron oxide sphere clusters. The iron oxide cube clusters showed the highest  $r_2$  value ( $589.3 \text{ mM}^{-1} \text{ s}^{-1}$ ) due to the presence of spinpolarized cubes in nanoclusters (Fig. 23d). Landau–Lifshitz–Gilbert simulations confirmed their hypothesis (Fig. 23e and f), suggesting that regulating magnetic field inhomogeneity is a powerful way to establish correlations between magnetization and  $T_2$  relaxivity of iron oxide NPs.

## 5. Summary and outlook

The design and fabrication of high-relaxivity and ultra-sensitive CAs for MRI provides not only a great challenge to overcome, but also provides great opportunities to make significant contributions in the biomedical field. The imaging performances of commercially available MRI CAs have been far from optimal, which necessitates the development of novel nanosized CAs with large enough contrast effects for highly sensitive detection of lesions. Undoubtedly, controllable synthesis and facile engineering of inorganic nanoparticulate CAs and their excellent *in vivo* behavior as imaging agents will greatly deepen our understanding and subsequent effective diagnosis and treatment of diseases.

This review article provides a comprehensive summary of methods that have been investigated to construct highperformance CAs, including size control of Gd-, Mn-, Fe-, Dy and Ho-based NPs, surface modification with different ligands, coating of silica shells (dense or mesoporous shells) or inorganic nanolayers on NPs, and regulation of shape of

NPs for accelerated relaxation of water protons. Some widely applied strategies to achieve high relaxivity for MRI contrast enhancement such as confinement of CAs in mesopores, doping with metal ions, and self-assembly or disassembly have been discussed. In addition, other recently developed novel strategies including valency change, oxygen vacancy, and distance-dependent tuning have been briefly described. While there is no simple answer to the question of “how do these factors impact relaxivity of NPs,” several principles that can guide the engineering of future NPs as high-performance MRI CAs are as follows:

1. Size effect of NPs will determine whether surface-to-volume ratio or  $\tau_R$  is the dominant contributing factor for generating  $T_1$  contrast with any given NP. Meanwhile, size-dependent magnetization of magnetic NPs can be utilized in the design of  $T_2$  CAs with high  $r_2$  values.
2. Surface modification of NPs is indispensable not only for controlling NP growth but also for increasing compatibility and functionalization of NPs. The surface of NPs can be further adjusted to enhance or reduce MRI signal after size, which depends on coating species, hydrophilicity, density, and thickness of organic ligands, silica shells, and inorganic layers.
3. Shape confinement of  $Gd^{3+}$ -based CAs has recently been utilized to impact relaxation of water protons by engineering  $Gd^{3+}$  chelates into dense and hydrophilic microenvironments that are closely related to the shape of NPs. In the meantime, the magnetic field inhomogeneity enhanced by shape anisotropy of magnetic NPs can be artificially regulated to enhance  $T_2$  relaxivity by reducing symmetry of the induced field.
4. Biodistribution, pharmacokinetics, and metabolism of NPs are also mediated by factors such as particle size, surface, and shape. The successful engineering of NPs as high-performance imaging agents must consider their behavior *in vivo*. Ultrasmall sized NPs with higher relaxivity are promising imaging agents because they can be excreted by the kidneys on a short time scale.
5. Strategies to achieve high relaxivity for MRI contrast enhancement are mainly focused on tuning the interaction between NPs and nearby water protons, either enhancing or reducing the contrast of MRI. By further exploiting tumor microenvironment characteristics such as acidic pH, unevenness of blood flow, and hypoxia, engineering of NPs for stimuli-responsive MRI can be realized for the precise imaging of tumors and detection of their pathological conditions with MRI.

Despite the recent significant progress and promising results in preclinical animal models, the investigation, evaluation, and optimization of NP-based MRI CAs for clinical use are still in their infancy. To facilitate and enable the successful translation of these CAs from laboratories to clinics, we outline several considerations for future research on the interdisciplinary cooperation among engineers, material scientists, biologists, clinicians, and others toward future (molecular) MRI (Fig. 24).

1. Although some factors/strategies have been explored by regulating the key parameters that influence the perturbation of protons by the dipolar field of NPs,

we believe that more novel factors/strategies for the design of high-relaxivity CAs are on the horizon. In addition, because the biodistribution, excretion, and pharmacokinetics of NPs are dependent on their size, shape, and surface status, rigorous and extensive attention should be paid to deep understanding the correlations among nanoparticle characteristics (*e.g.*, size, shape, and surface status), *in vivo* metabolism, and MRI contrast enhancement efficiency. Through a basic understanding of how these factors work together, future clinical applications may be realized. Synergistic integration of these factors may offer promising approaches in fabricating low-toxicity and high-performance CAs that are amenable to clinical translation which will eventually benefit patients.<sup>18</sup>

2. Efficient systemic delivery of CAs to the targeted tissues is of utmost importance for accurate and sensitive detection by MRI, which can include passive accumulation of elaborately synthesized NPs to maximize enhanced permeability and retention (EPR) effect, as well as active targeting of NPs *via* various ligands such as antibodies, antibody fragments, peptides, small molecules, and proteins. The impacts of nanoparticle size, surface modification, and shape on proton relaxation have been well studied, while their relationship with the EPR effect remains to be further elucidated. Generally, ultra-small sized renal clearable NPs exhibit poor EPR effect.<sup>215</sup> Future studies should focus on the specific interaction between MRI contrast performance of CAs and EPR effect to develop a clinically relevant and useful strategy for obtaining high-relaxivity NPs with an optimal balance between renal clearance and strong EPR effects over a suitable time scale (*e.g.*, less than a few hours).<sup>216</sup>
3. The hypoxic microenvironment in various tumors has received tremendous attention from researchers because it is closely related to the tumor's aggressiveness and resistance to therapy,<sup>217,218</sup> and a number of NPs have been used to effectively eradicate hypoxic tumors in preclinical animal models.<sup>219-222</sup> Although in this review article we have discussed various strategies based on NPs for efficient MRI, there has been no literature report on hypoxia-responsive NPs as MRI CAs except on several chemical complexes for hypoxia-related MRI such as  $T_1$  or  $T_2$  CAs,<sup>223</sup>  $^{19}\text{F}$ -based probes,<sup>224</sup> and CEST probes.<sup>86</sup> We believe that research on this subject is likely already in progress, and the NPs that respond to hypoxia *via*, for example, valency change, size/morphology variation, assembly, and disassembly for enhanced MRI, will likely play multiple roles in accurate and noninvasive delineation of hypoxic areas, more precise and effective treatment, and real-time and noninvasive monitoring of therapeutic efficacy.
4. The next generation of MRI CAs will be developed not only based on design principles of NPs described in this review article, but also on the concurrent development of advanced nanotechnology, novel MRI sequences and scanners, and new biological applications. The synthesis of these NPs needs further optimization, and NPs that are synthesized in an easy, low-cost, and large-scale way is prerequisite to future clinical translation, a challenge which may be overcome by seeking more facile, reproducible, and economic one-pot syntheses without harmful reagents. With an in-depth understanding of the involved

processes, we firmly believe that good biocompatibility and biodegradability, high relaxivity, and novel MRI contrast mechanisms can all be achieved with the future generation of novel inorganic NPs to usher in the new era of 21st century precision medicine (Table 3).

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## Biographies



### Dalong Ni

Dalong Ni received his PhD degree in 2016 from the Shanghai Institute of Ceramics, Chinese Academy of Sciences (SICCAS), under the supervision of Prof. Wenbo Bu. He then joined the Department of Radiology at the University of Wisconsin-Madison as a postdoctoral fellow under the supervision of Prof. Weibo Cai. His research interests focus on

design, synthesis, and biomedical applications of multifunctional nanoplateforms for imaging and therapy applications, such as MRI and MRI-guided cancer therapy.



### **Wenbo Bu**

Wenbo Bu received his PhD degree from Nanjing Tech University (China) in 2002. He has been a full-time professor in SICCAS since 2008. In 2016, he moved from SICCAS to East China Normal University. His research mainly focuses on design and synthesis of multifunctional rareearth- based nanomaterials for cancer imaging and therapeutic applications.



### **Emily B. Ehlerding**

Emily B. Ehlerding received her MS degree in Medical Physics from the University of Wisconsin- Madison in 2016 and her BS in Chemistry and Physics from Manchester University in 2014. She is currently a PhD candidate in Medical Physics at UW-Madison, working with antibody-based platforms for multimodality imaging of immunotherapy-related targets.



### **Weibo Cai**

Weibo Cai received his PhD degree in Chemistry from University of California, San Diego, in 2004 and is now an associate professor at the University of Wisconsin- Madison (<http://mi.wisc.edu>). He has published more than 230 peerreviewed articles (H-index 4 60) and has

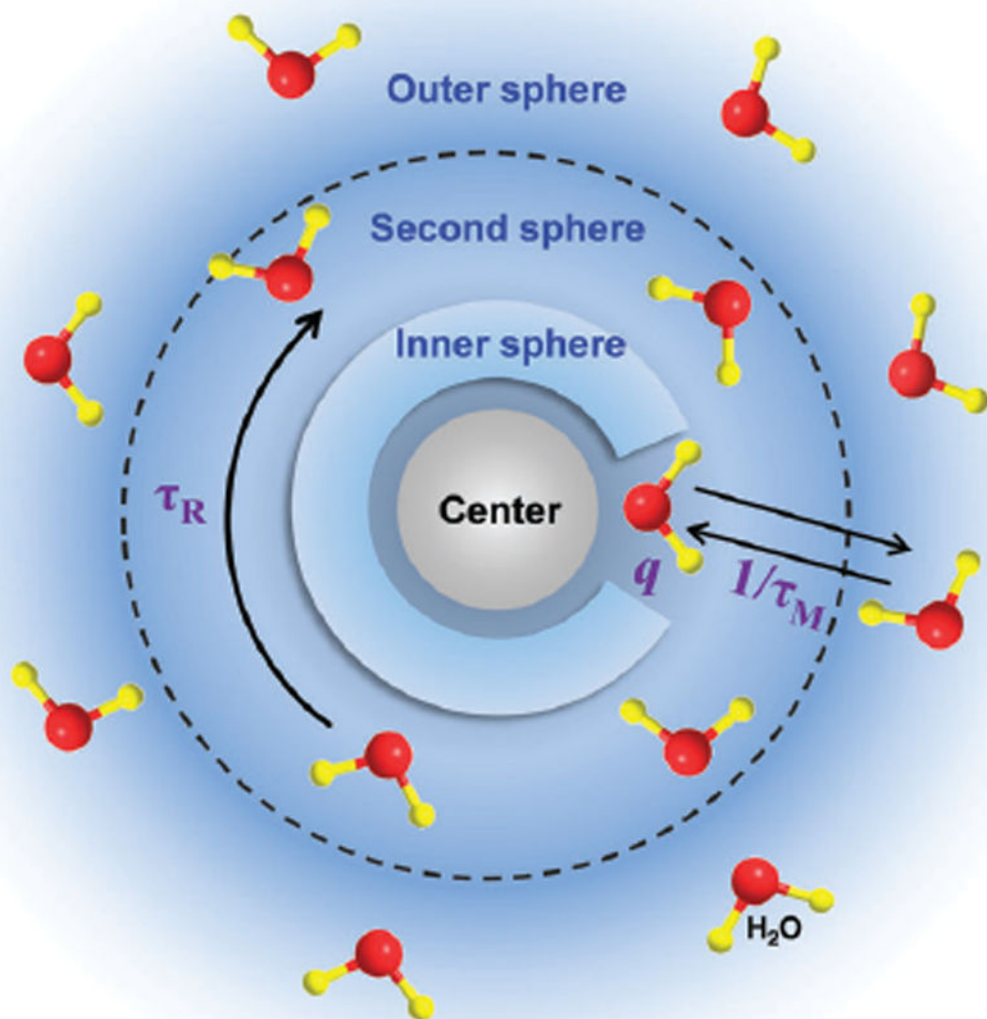


received many awards such as the Society of Nuclear Medicine and Molecular Imaging (SNMMI) Young Professionals Committee Best Basic Science Award (2007), the European Association of Nuclear Medicine Springer Prize (2011 & 2013), American Cancer Society Research Scholar (2013–2017), and UW-Madison Vilas Distinguished Achievement Professor Award (2017). His trainees have received more than 80 awards to date, such as the 2012 Berson-Yalow Award from SNMMI.



### **Jianlin Shi**

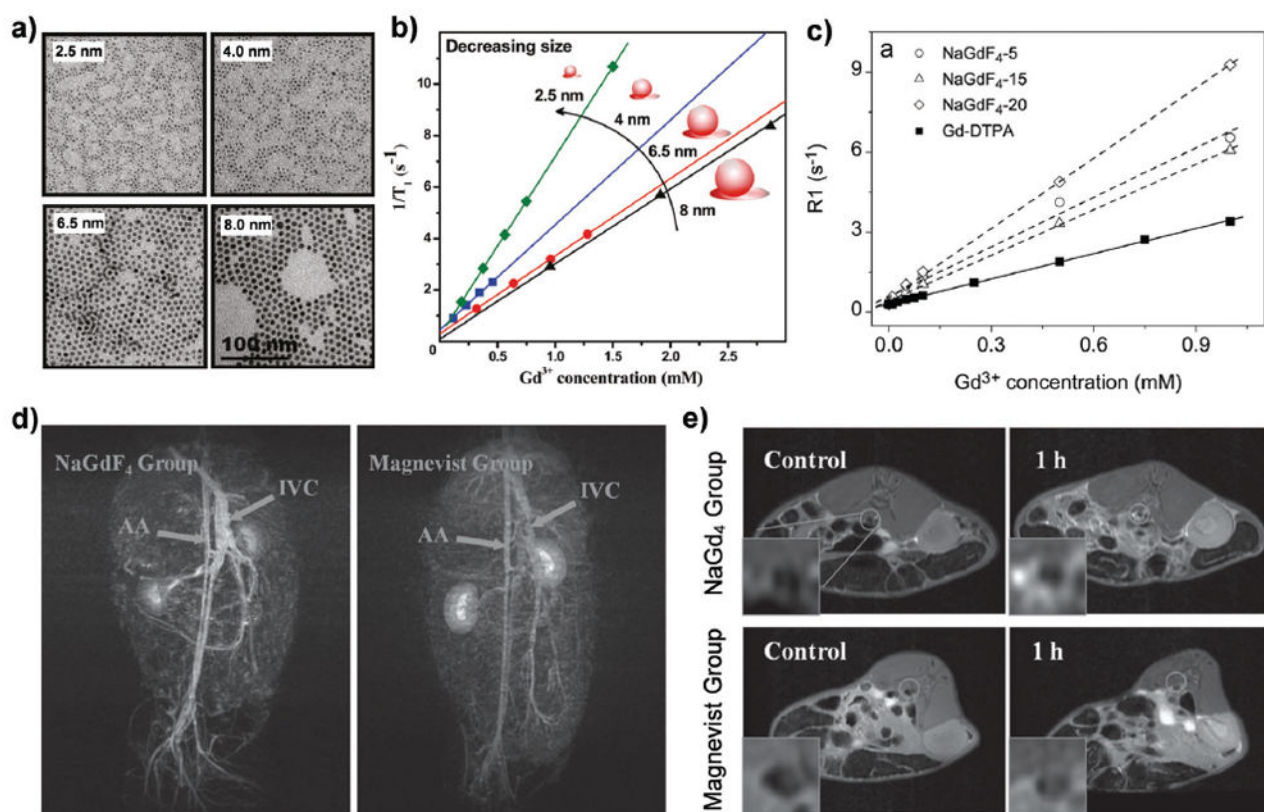
Jianlin Shi received his PhD degree from SICCAS. He has been a full-time professor at SICCAS since 1994. His research interests include the synthesis of mesoporous silica nanoparticles and mesoporous silica-based nanocomposites for biomedical applications. He has published over 400 scientific papers, which have been cited more than 22 000 times by other scientists, and has an H-index of 79.



**Fig. 1.** Inner sphere, secondary sphere, and outer sphere can influence relaxation rates of MRI CAs.

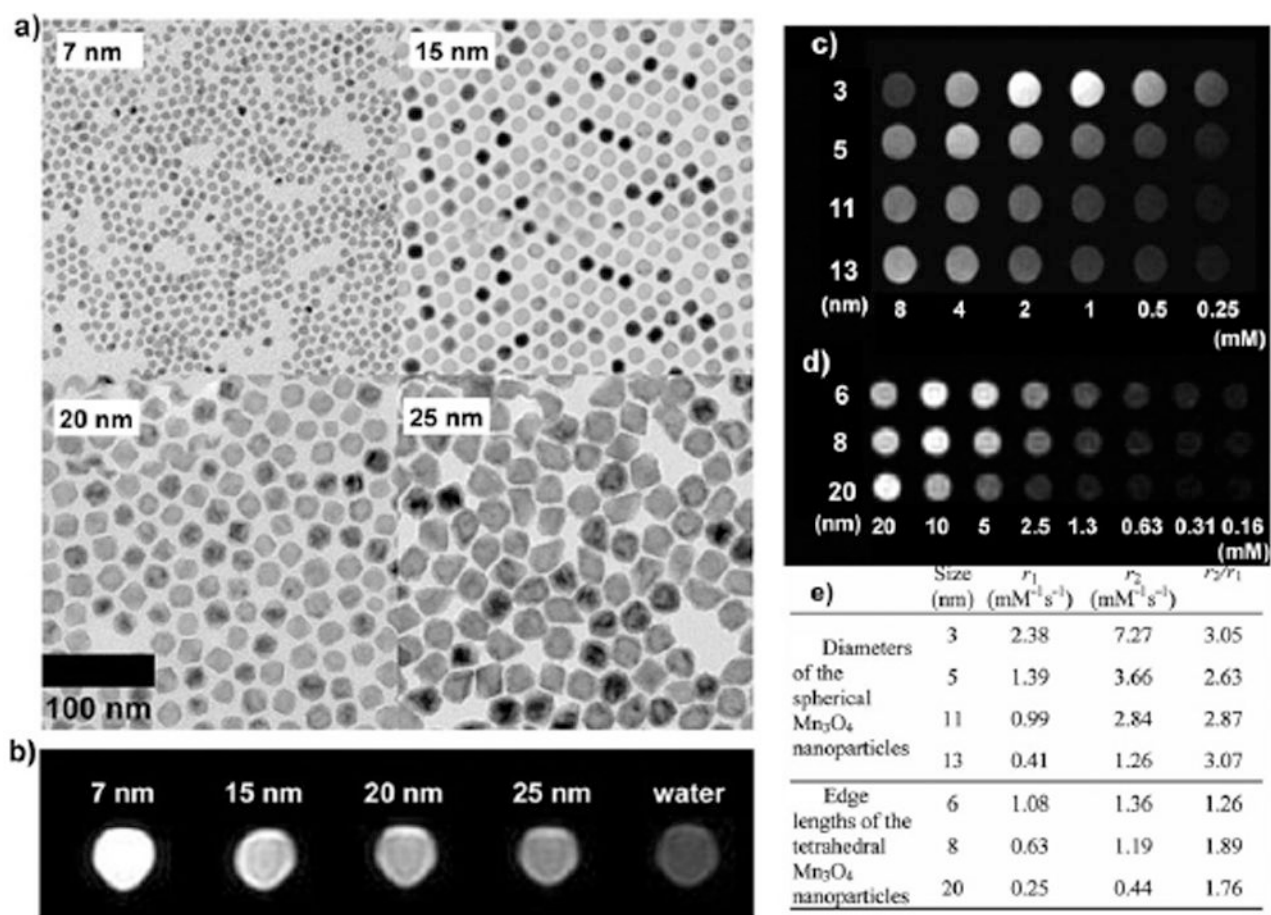


**Fig. 2.** Factors influencing relaxation of MRI CAs: size, surface, and shape effects at the nanoscale.



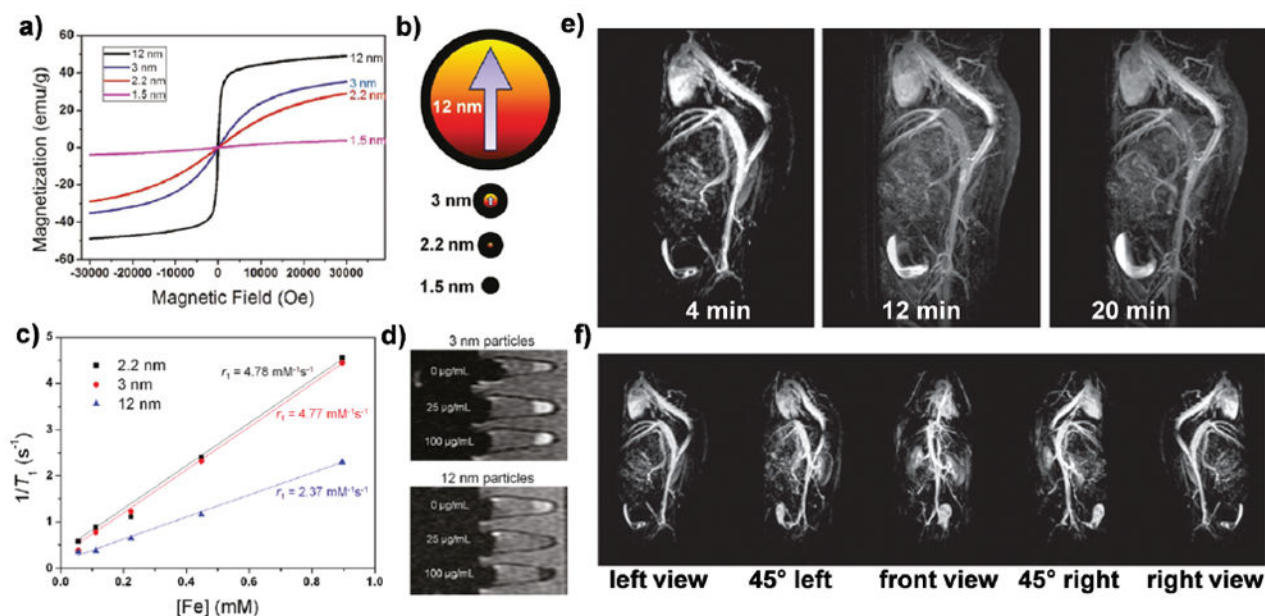
**Fig. 3.**

(a) TEM images of oleate-stabilized NaGdF<sub>4</sub> NPs of varied particle sizes (scale bar is the same for all images). (b) Plots of  $1/T_1$  versus Gd<sup>3+</sup> concentration for varied NaGdF<sub>4</sub> NPs in water (1.5 T). Reproduced with permission from ref. 49. Copyright 2011, American Chemical Society. (c)  $R_1$  relaxivity of aqueous solutions containing NaGdF<sub>4</sub> NPs of different sizes or Gd–DTPA with different concentrations of Gd<sup>3+</sup>. Reproduced with permission from ref. 41. Copyright 2013, American Chemical Society. (d) Magnetic resonance angiography of rabbits within 3 min after injection of ultra-small NaGdF<sub>4</sub> nanodots or Magnevist. AA: abdominal aorta; IVC: inferior vena cava. (e) Transverse cross-sectional images of rabbit atherosclerotic plaques before and after injection of ultra-small NaGdF<sub>4</sub> nanodots or Magnevist at the same dosage (13 mg Gd per kg). Reproduced with permission from ref. 40. Copyright 2014, Wiley-VCH.

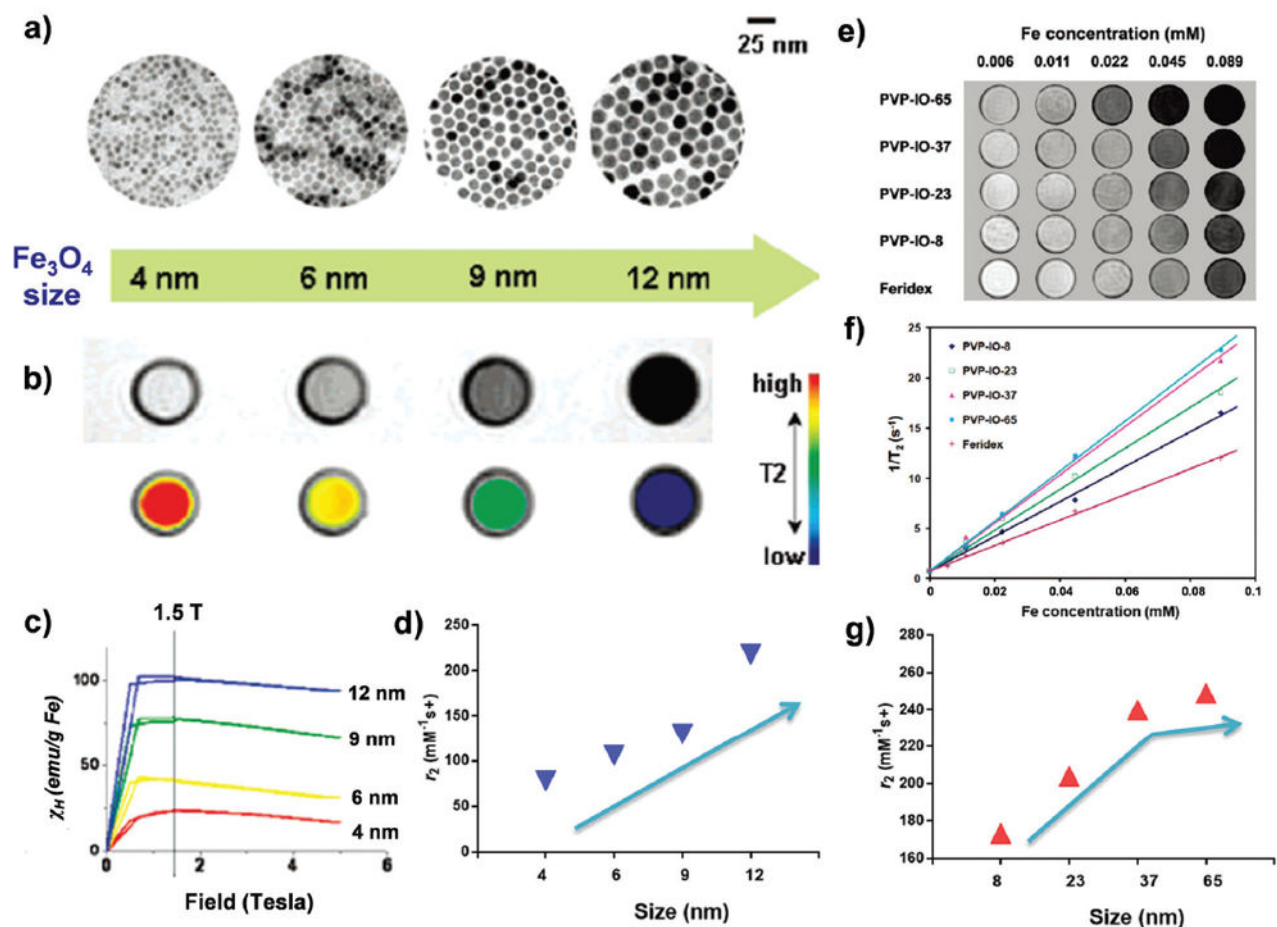


**Fig. 4.**

(a) TEM images of MnO NPs of varied sizes (7, 15, 20, and 25 nm) in water. (b)  $T_1$ -weighted MR image of MnO NPs in a 3.0 T clinical MRI system. Reproduced with permission from ref. 69. Copyright 2007, Wiley-VCH. Size-dependent  $T_1$ -weighted MR images and relaxivities of water-dispersible spherical (c) and tetrahedral (d)  $Mn_3O_4$  NPs at varied  $Mn^{2+}$  concentrations. (e) The  $r_1$  and  $r_2$  values and  $r_2/r_1$  ratios of  $Mn_3O_4$  MPs from c and d. Reproduced with permission from ref. 132. Copyright 2012, Wiley-VCH.



**Fig. 5.** (a) Field-dependent magnetization curves at 300 K for iron oxide NPs of varied sizes. (b) Description of spin canting effect (canting layer = 0.9 nm) in iron oxide NPs of varied sizes. Red and black colors represent magnetic cores and magnetically disordered shells. (c) Plots of  $1/T_1$  over concentration of iron oxide NPs of 3 nm and 12 nm diameter. (d)  $T_1$ -weighted MR images of MCF-7 cell pellets after 24 h incubation with iron oxide NPs of 3 nm and 12 nm in diameters. Reproduced with permission from ref. 135. Copyright 2011, American Chemical Society. (e)  $T_1$ -weighted magnetic resonance angiography (MRA) of a mouse injected with ultra-small iron oxide NPs at 7 T. (f) Five different perspectives of MRA images, which were extracted from a 3D scan, at 4 minutes post-injection. Reproduced with permission from ref. 142. Copyright 2017, National Academy of Sciences.



**Fig. 6.**

(a) TEM images of iron oxide NPs of 4, 6, 9, and 12 nm. (b) Upper panel: Size-dependent  $T_2$ -weighted MR images of iron oxide NPs in aqueous solution at 1.5 T; lower panel: size-dependent changes from red to blue in color-coded MR images based on  $T_2$  values. (c) The magnetization of iron oxide NPs measured by a SQUID magnetometer. (d) Size-dependent  $r_2$  values of iron oxide NPs. Reproduced with permission from ref. 79. Copyright 2005, American Chemical Society. (e) Size-dependent  $T_2$ -weighted MR images of PVP-coated iron oxide NPs in aqueous solution with various concentrations at 7 T. (f) Plots of  $1/T_2$  against Fe concentration of PVP-coated iron oxide NPs. (g) Size-dependent  $r_2$  values of iron oxide NPs. Reproduced with permission from ref. 138. Copyright 2010, American Chemical Society.

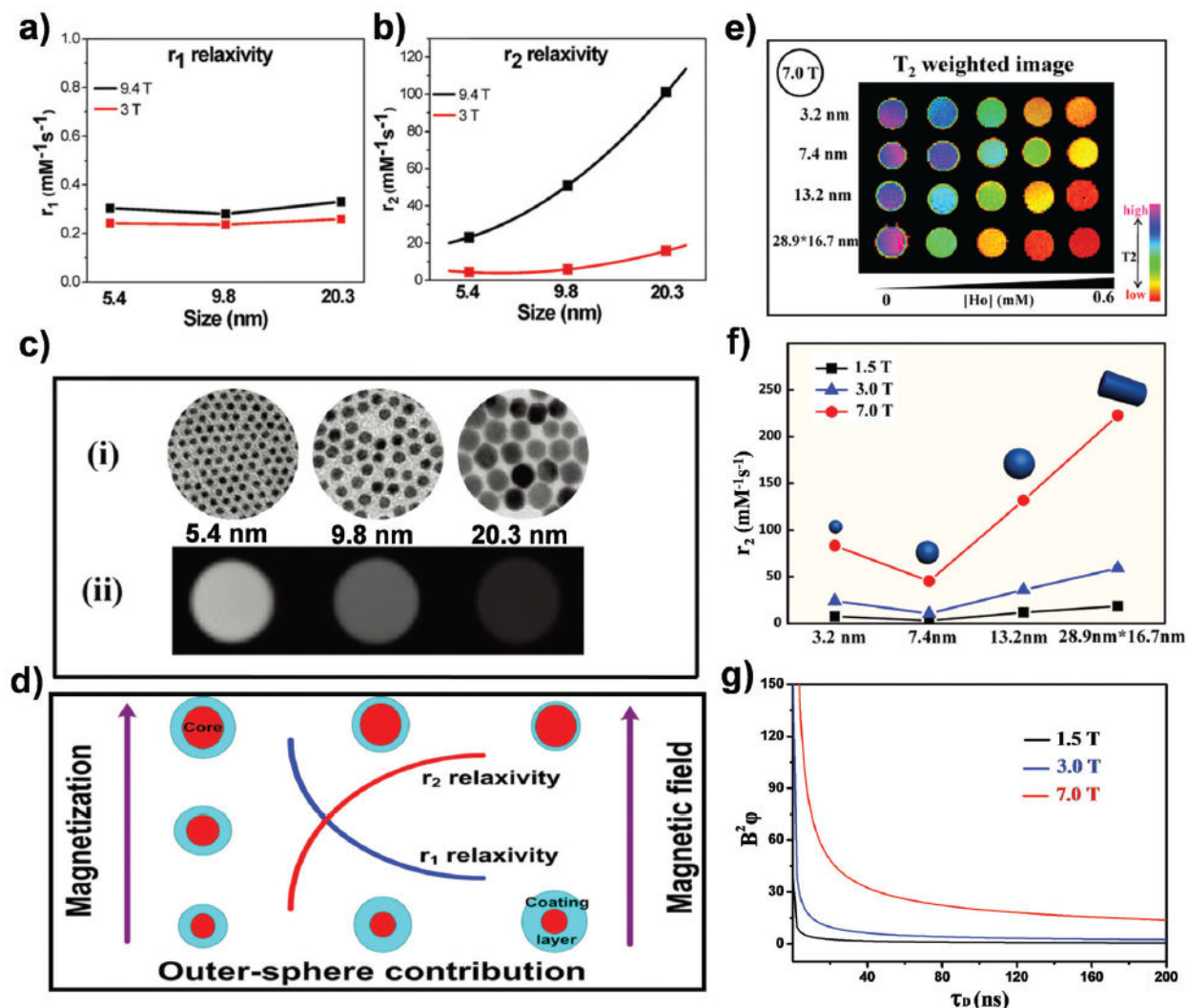
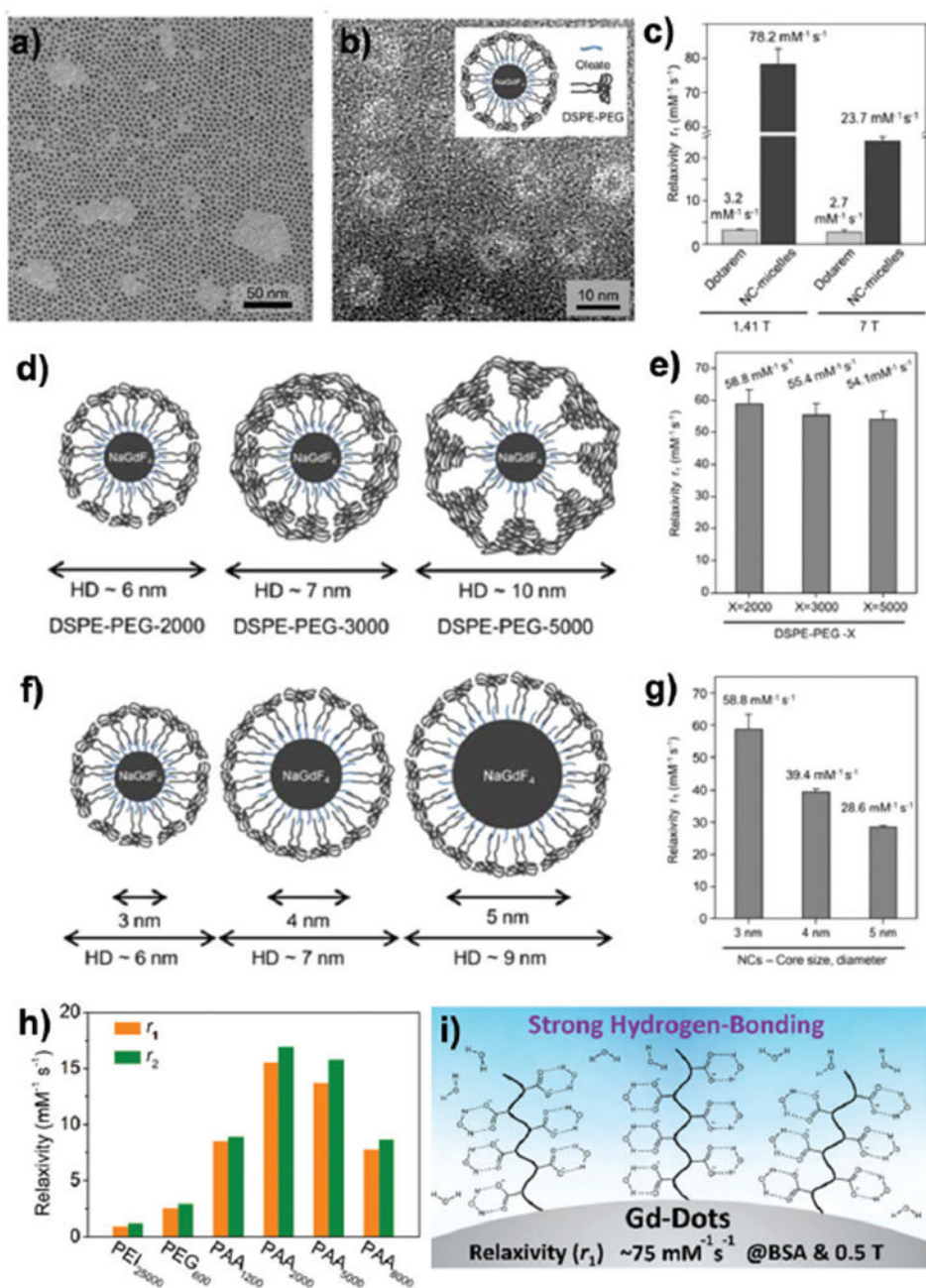


Fig. 7.

(a)  $r_1$  and (b)  $r_2$  values of NaDyF<sub>4</sub> NPs (three different sizes) at 3 and 9.4 T. (c) (i) TEM images of 5.4, 9.8, and 20.3 nm NaDyF<sub>4</sub> NPs, (ii) phantom MR images of NPs at 1.0 mM Dy<sup>3+</sup> concentration at 9.4 T. Reproduced with permission from ref. 19. Copyright 2012, American Chemical Society. (d) Schematic of the effect of size and a coating layer on the relaxation of NaDyF<sub>4</sub> and NaHoF<sub>4</sub> NPs. Reproduced with permission from ref. 20. Copyright 2016, American Chemical Society. (e)  $T_2$ -weighted MR images of NaHoF<sub>4</sub> NPs of varied sizes in aqueous solutions at 7.0 T. (f)  $r_2$  values obtained for varied sized NaHoF<sub>4</sub> NPs at 1.5 T, 3.0 T, and 7.0 T. (g) Simulated curves of  $B^2\phi$  from Curie contribution versus diffusion correlation time ( $\tau_D$ ), which is related to NP size. Reproduced with permission from ref. 21. Copyright 2016, Elsevier Ltd.





**Fig. 8.** (a) A TEM image of ultra-small oleate-stabilized NaGdF<sub>4</sub> NPs of 3 nm diameter. (b) A representative negative-stained TEM image of DSPE–PEG-coated NP micelles. Inset: Schematic of a NaGdF<sub>4</sub> NP confined within a DSPE–PEG micelle. (c)  $r_1$  values of a clinical MRI contrast agent (Dotarem) and the compact NP micelles at low (1.41 T) and high (7 T) fields. (d) Schematic and (e)  $r_1$  values of NaGdF<sub>4</sub> NPs (3 nm) confined within DSPE–PEG micelles with varied PEG chain length. (f) Schematic and (g)  $r_1$  values of NaGdF<sub>4</sub> NPs confined within DSPE–PEG micelles with variable core NP size (3–5 nm). Reproduced with permission from ref. 39. Copyright 2016, American Chemical Society. (h)  $r_1$  and  $r_2$  values of

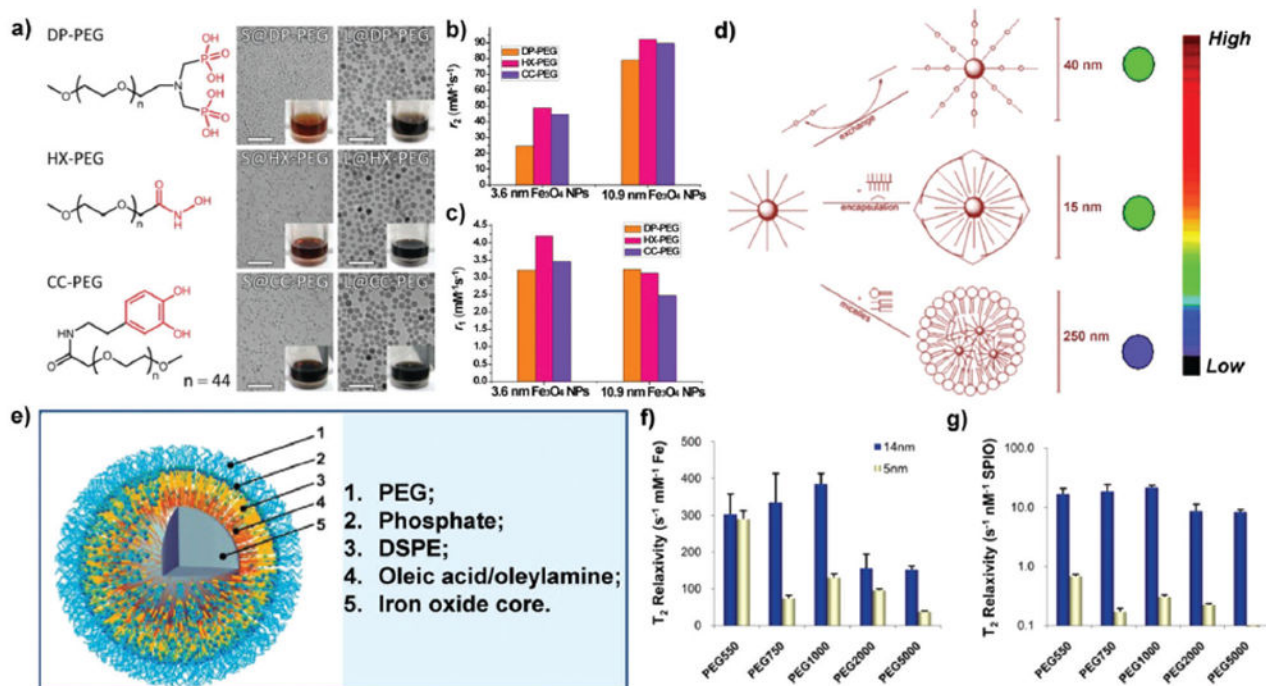
NaGdF<sub>4</sub> NPs coated with different ligands measured at 0.5 T. (i) Schematic of the strong hydrogen-bonding capacity of PAA to water molecules to improve  $r_1$ . Reproduced with permission from ref. 36. Copyright 2017, American Chemical Society.

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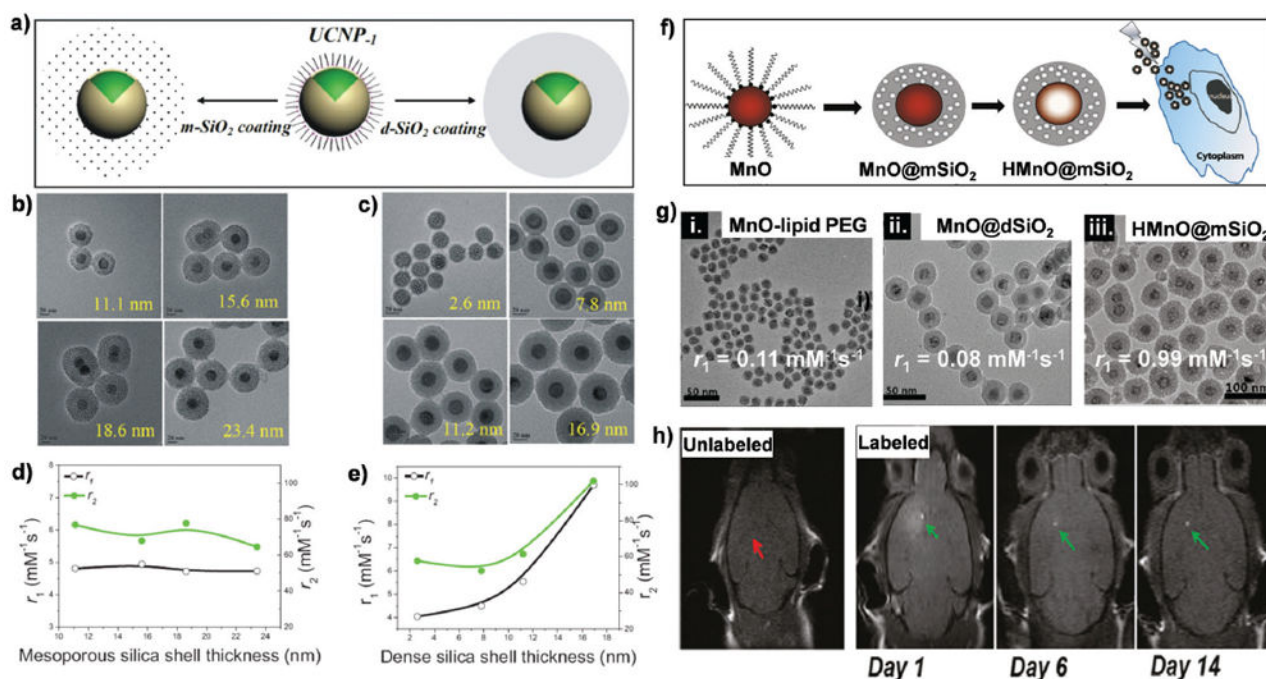
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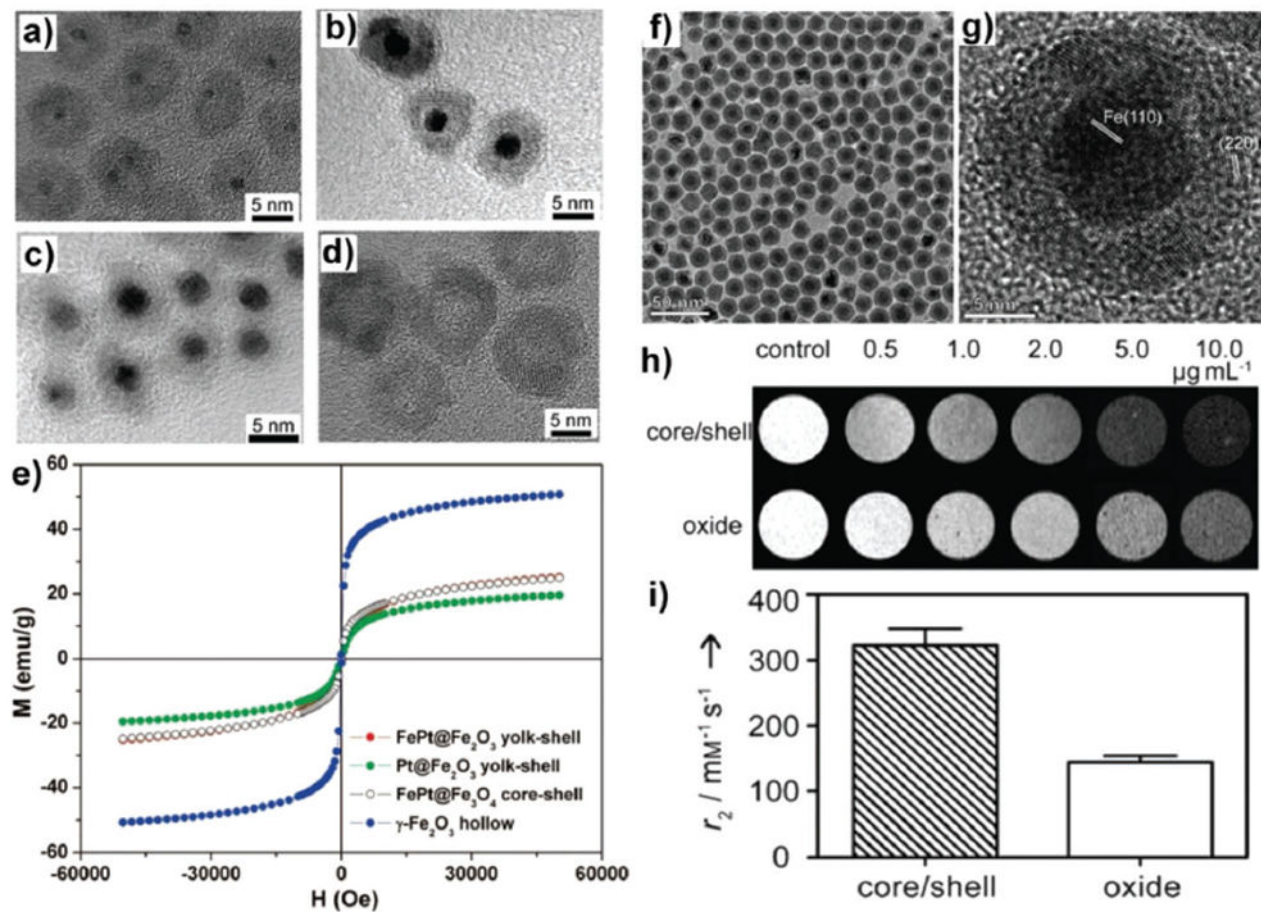
**Fig. 9.**

(a) Left panel: Chemical structures of PEGs used for exchanging the hydrophobic ligands of 3.6 nm(S: small) and 10.9 nm(L: large) iron oxide NPs; right panel: TEM images of resulting PEGylated NPs (scale bars: 50 nm); insets: photographs of aqueous solutions of PEGylated NPs at an equal Fe concentration of 20 mM. Comparison of (b)  $r_2$  and (c)  $r_1$  values of different kinds of PEGylated NPs. Reproduced with permission from ref. 149. Copyright 2014, Wiley-VCH. (d) Left panel: Schematic of three transfer approaches: ligand exchange of oleic acid with a water-soluble polymer (top), coating of individual NPs with the amphiphilic polymer (middle), and embedding into lipid micelles (bottom). Reproduced with permission from ref. 147. Copyright 2007, American Chemical Society. (e) Schematic of a nanoparticle with 4.8 nm iron oxide core and DSPE-PEG1000 coating. (f)  $T_2$  relaxivity of iron oxide NPs at a constant iron concentration. (g)  $T_2$  relaxivity of NPs on a per-particle basis. Iron oxide NPs with two core sizes, 5 and 14 nm, and five PEG sizes, 550, 750, 1000, 2000, and 5000 Da, were evaluated. Reproduced with permission from ref. 105. Copyright 2010, American Chemical Society.

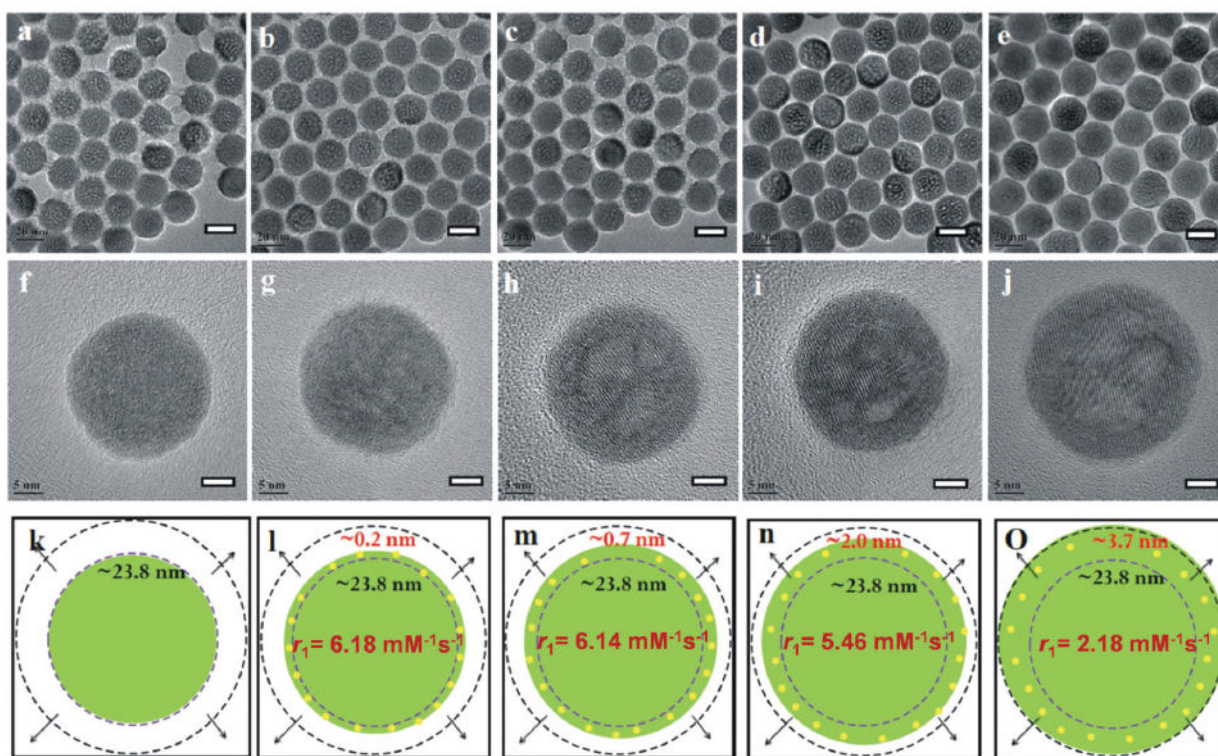


**Fig. 10.**

(a) Schematic illustration of coating UCNPs with dSiO<sub>2</sub> and mSiO<sub>2</sub> shells. (b) TEM images of (b) UCNP@mSiO<sub>2</sub> and (c) UCNP@dSiO<sub>2</sub> with varied shell thicknesses. Scale bar: 20 nm. (d) The plot of  $r_1$  and  $r_2$  versus mSiO<sub>2</sub> shell thickness. (e) The plot of  $r_1$  and  $r_2$  versus dSiO<sub>2</sub> shell thickness. Reproduced with permission from ref. 159. Copyright 2014, Wiley-VCH. (f) Schematic of the synthesis of HMnO@mSiO<sub>2</sub> NPs and labeling of mesenchymal stem cells. HMnO denotes hollow structure manganese oxide NPs. (g) TEM images and  $r_1$  value of MnO–lipid PEG, MnO@dSiO<sub>2</sub>, HMnO@mSiO<sub>2</sub> NPs. (h) No MRI contrast enhancement (red arrow) was detected in mice transplanted with unlabeled mesenchymal stem cells into the brain, whereas hyperintense signals (green arrows) were detected in mice transplanted with HMnO@mSiO<sub>2</sub>-labeled mesenchymal stem cells, which were still visible 14 days after injection in the brain. The scheme and figures (with minor modifications) are licensed under the ACS AuthorChoice license. Reproduced with permission from ref. 62. Copyright 2011, American Chemical Society.

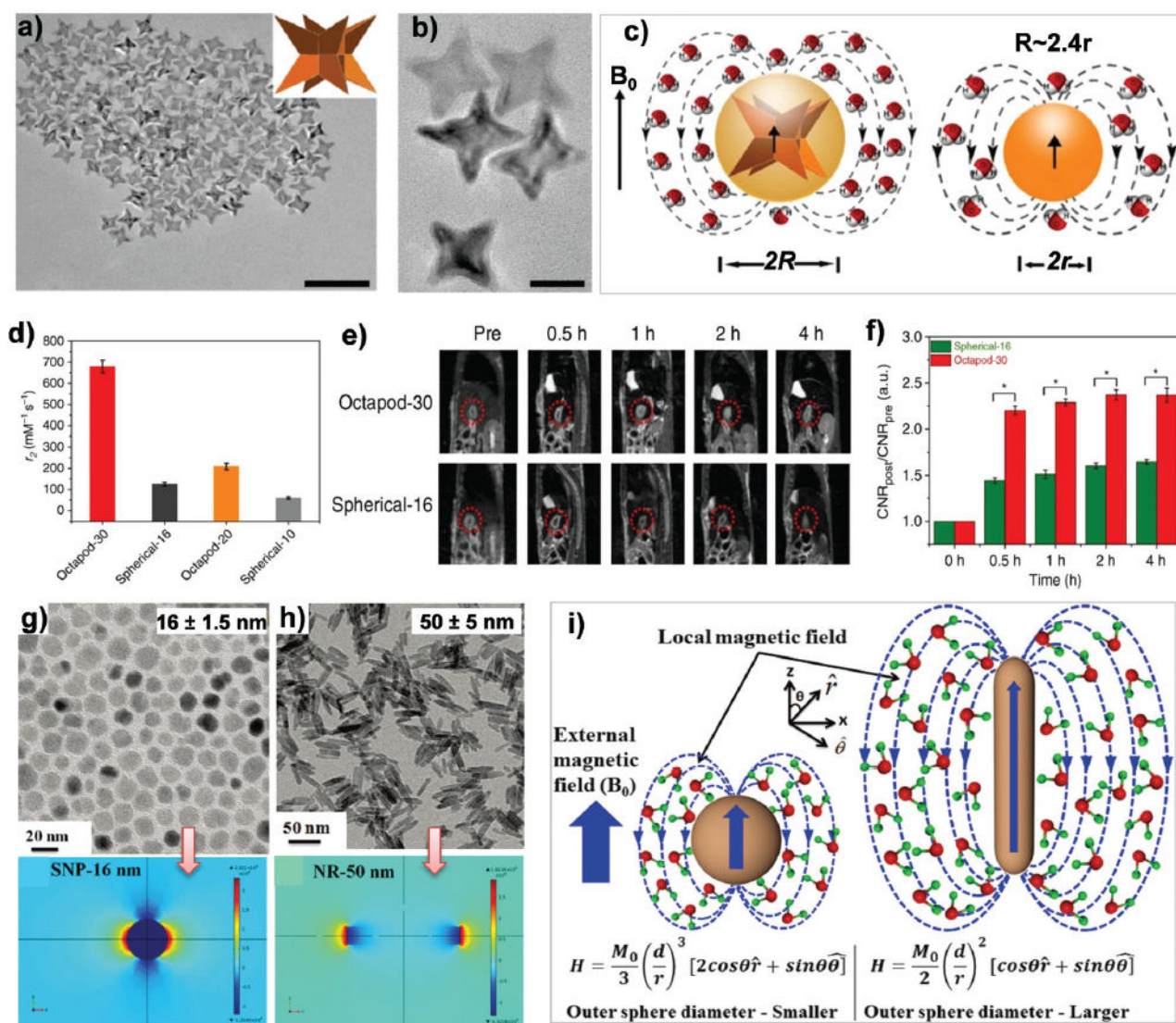


**Fig. 11.** HRTEM images of (a) FePt@Fe<sub>2</sub>O<sub>3</sub> yolk-shell NPs, (b) Pt@Fe<sub>2</sub>O<sub>3</sub> yolk-shell NPs, (c) FePt@Fe<sub>3</sub>O<sub>4</sub> core-shell NPs obtained by the seed-growth method, and (d)  $\gamma$ -Fe<sub>2</sub>O<sub>3</sub> hollow NPs. (e) Room-temperature field-dependent magnetization measurements of different NPs. Reproduced with permission from ref. 77. Copyright 2008, American Chemical Society. (f) TEM image of 16 nm iron/iron oxide core/shell NPs. (g) HRTEM image showing a core of single-crystal a-Fe and a shell consisting of multiple domains of iron oxide. (h) T<sub>2</sub>-weighted MR images at 9.4 T of iron/iron oxide core/shell NPs and iron oxide NPs. (i)  $r_2$  values of the core/shell and oxide NPs determined from the same samples as in (h). Reproduced with permission from ref. 163. Copyright 2011, Wiley-VCH.



**Fig. 12.**

TEM images and schematic of Gd-free core (a, f and k) and core@NaGdF<sub>4</sub> with varied NaGdF<sub>4</sub> shell thicknesses of 0.2 nm (b, g, and l), 0.7 nm (c, h and m), 2.0 nm (d, i, and n), and 3.7 nm (e, j, and o). Yellow dots in l, m, n, and o represent Gd<sup>3+</sup> ions. The  $r_1$  values of silica-coated water-soluble core@NaGdF<sub>4</sub> NPs of various sizes are shown in the middle of the schematic illustrations. Reproduced with permission from ref. 164. Copyright 2011, Wiley-VCH.



**Fig. 13.** Shape effects on  $T_2$  contrast. (a) TEM image (scale bar, 100 nm) and (b) higher magnification TEM image (scale bar, 20 nm) of octapod-30 NPs of uniform four-armed iron oxide NPs. (c) A cartoon showing octapod and spherical iron oxide NPs. With the same geometric core volumes, the octapod NPs have much larger effective volumes (radius,  $R$ ) than the spherical NPs (radius,  $r$ ). (d) Comparison of  $T_2$  values of different iron oxide NPs. (e) *In vivo* sagittal MR images and (f) quantification of signal-to-background ratio ( $*P = 0.01$ ) of orthotopic liver tumor models at 0, 0.5, 1, 2 and 4 h after intravenous injection of octapod-30 and spherical-16 iron oxide NPs. Reproduced with permission from ref. 137. Copyright 2013, Nature Publishing Group. (g and h) Top: TEM images of 16 nm iron oxide NPs (g) and 50 nm NRs (h). Bottom: Local magnetic field generated by the NRs and spherical  $\text{Fe}_3\text{O}_4$  NPs of equivalent material volumes under an applied magnetic field of 3 T. (i) Schematic of the quantum mechanical outer sphere model of  $\text{Fe}_3\text{O}_4$  NPs and NRs of the

same material volume. Reproduced with permission from ref. 80. Copyright 2015, Royal Society of Chemistry.

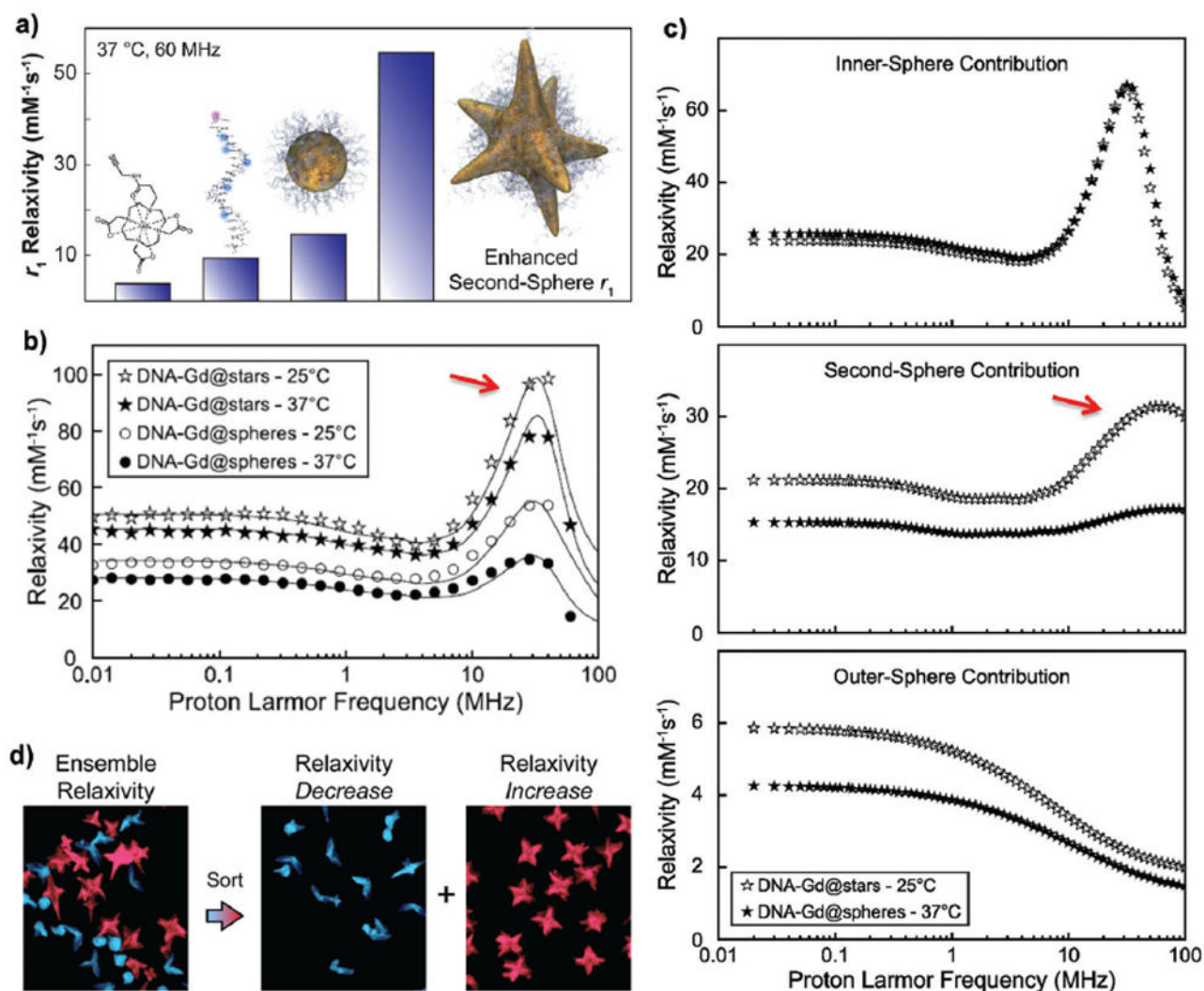
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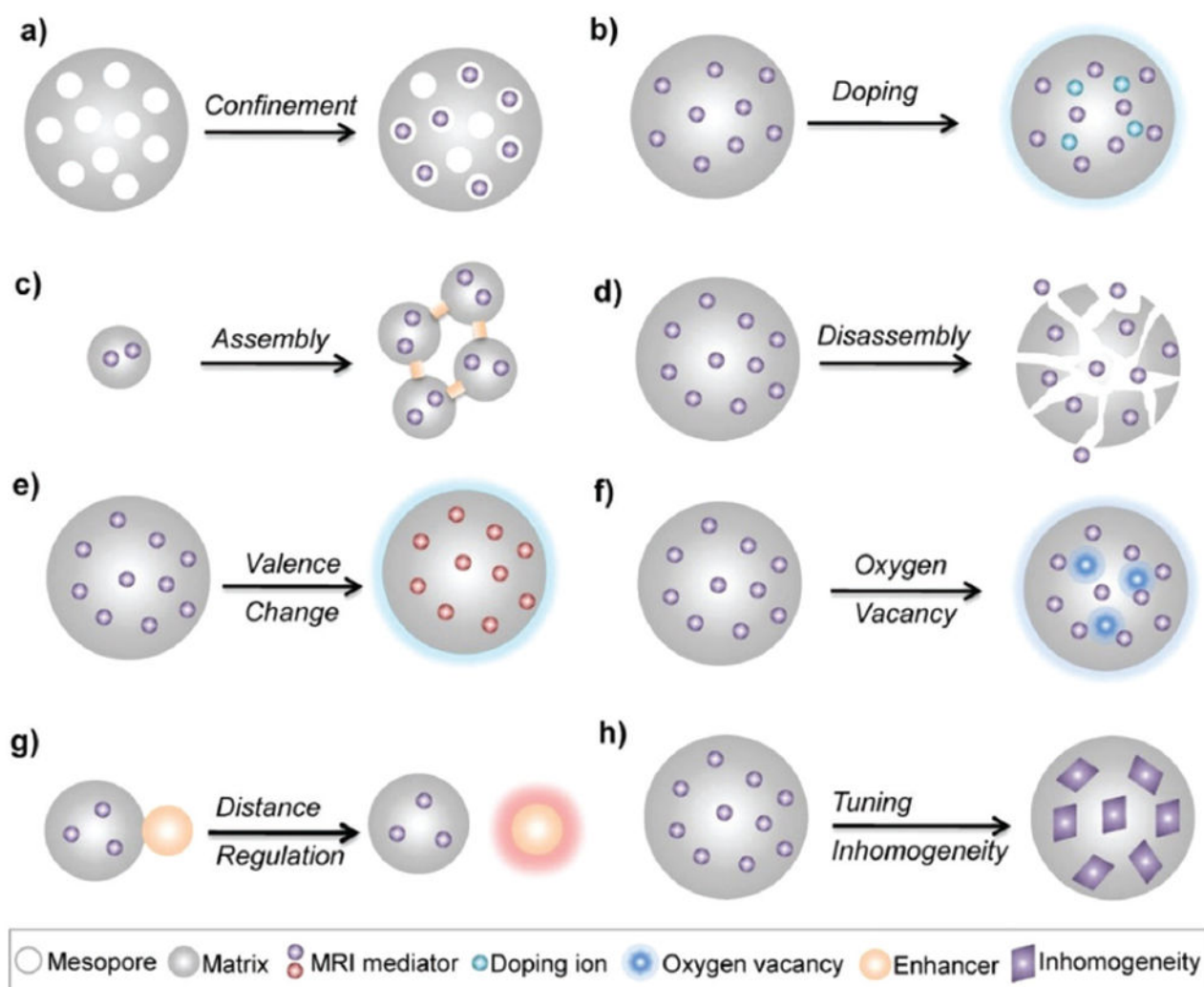
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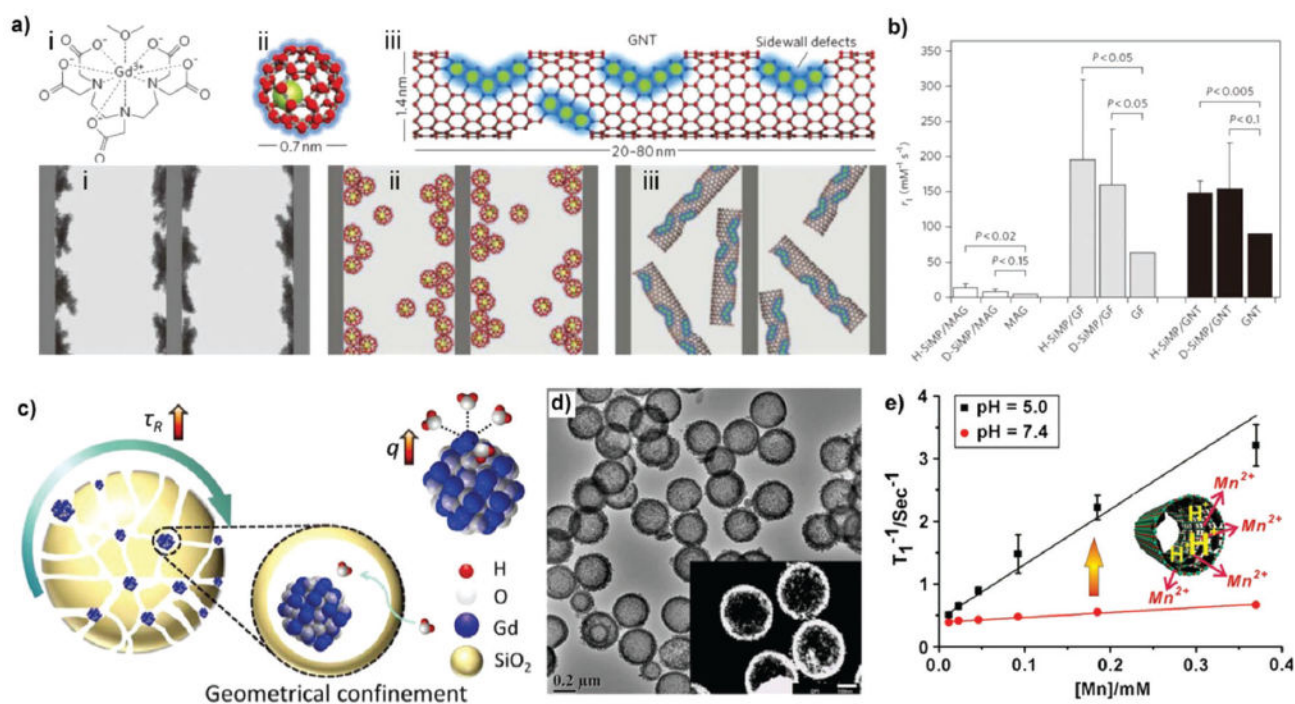




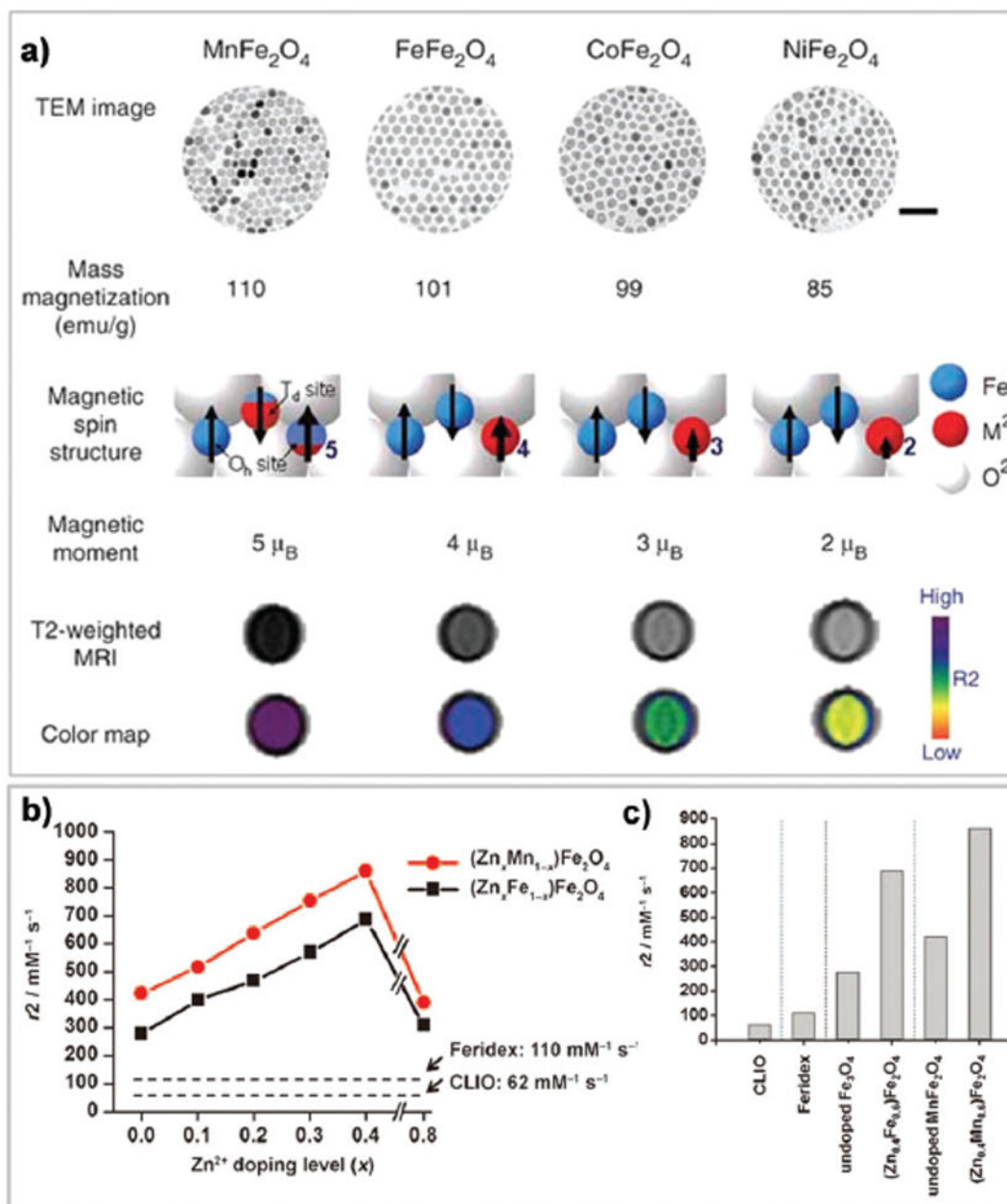
**Fig. 14.** Effects of NP shape on  $T_1$  relaxivity. (a) Schematic showing  $r_1$  values of different kinds of MRI CAs including  $\text{Gd}^{3+}$  chelates, DNA-Gd@spheres, and DNA-Gd@stars. (b) NMRD profiles for water solutions of DNA-Gd@stars and DNA-Gd@spheres. (c) Simulated deconvolution of DNA-Gd@stars NMRD profiles into their inner, secondary, and outer sphere contributions. Reproduced with permission from ref. 95. Copyright 2015, American Chemical Society. (d) Schematic showing the dependence of relaxation on branch numbers of DNA-Gd@stars. Reproduced with permission from ref.94. Copyright 2016, American Chemical Society.



**Fig. 15.** Schematic of engineering strategies to obtain high-relaxivity CAs for efficient MRI contrast enhancement.

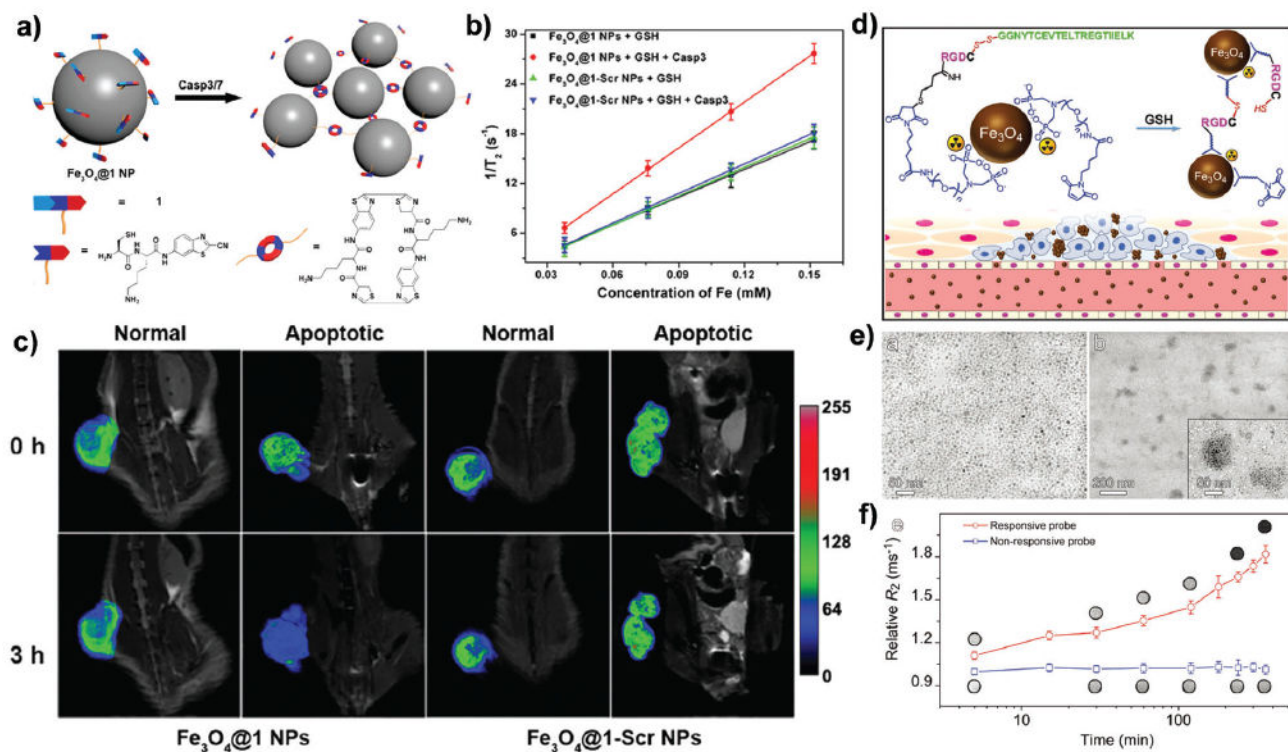
**Fig. 16.**

(a) Schematics showing (i) Magnevist (MAG), (ii) gadofullerenes (GF), and (iii) gadonanotubes (GNT) and cartoons showing Magnevist, GFs and GNTs entrapped within the porous structure of SiMPs. (b)  $r_1$  values of six MRI CA nanostructures in comparison with corresponding Gd-based CAs at 1.41 T and 37 °C. Reproduced with permission from ref. 186. Copyright 2010, Nature Publishing Group. (c) Schematic of  $Gd_2O_3$  confined in MSNs. Reproduced with permission from ref. 187. Copyright 2016, Royal Society of Chemistry. (d) TEM image of hybrid mesoporous composite nanocapsules (HMCNs). Inset: STEM image with scale bar = 100 nm. (e) Relaxivity of an aqueous suspension of HMCNs after 4 h of soaking in buffer solutions at pH 5.0 and 7.4 at 37 °C. Reproduced with permission from ref. 61. Copyright 2012, Elsevier Ltd.

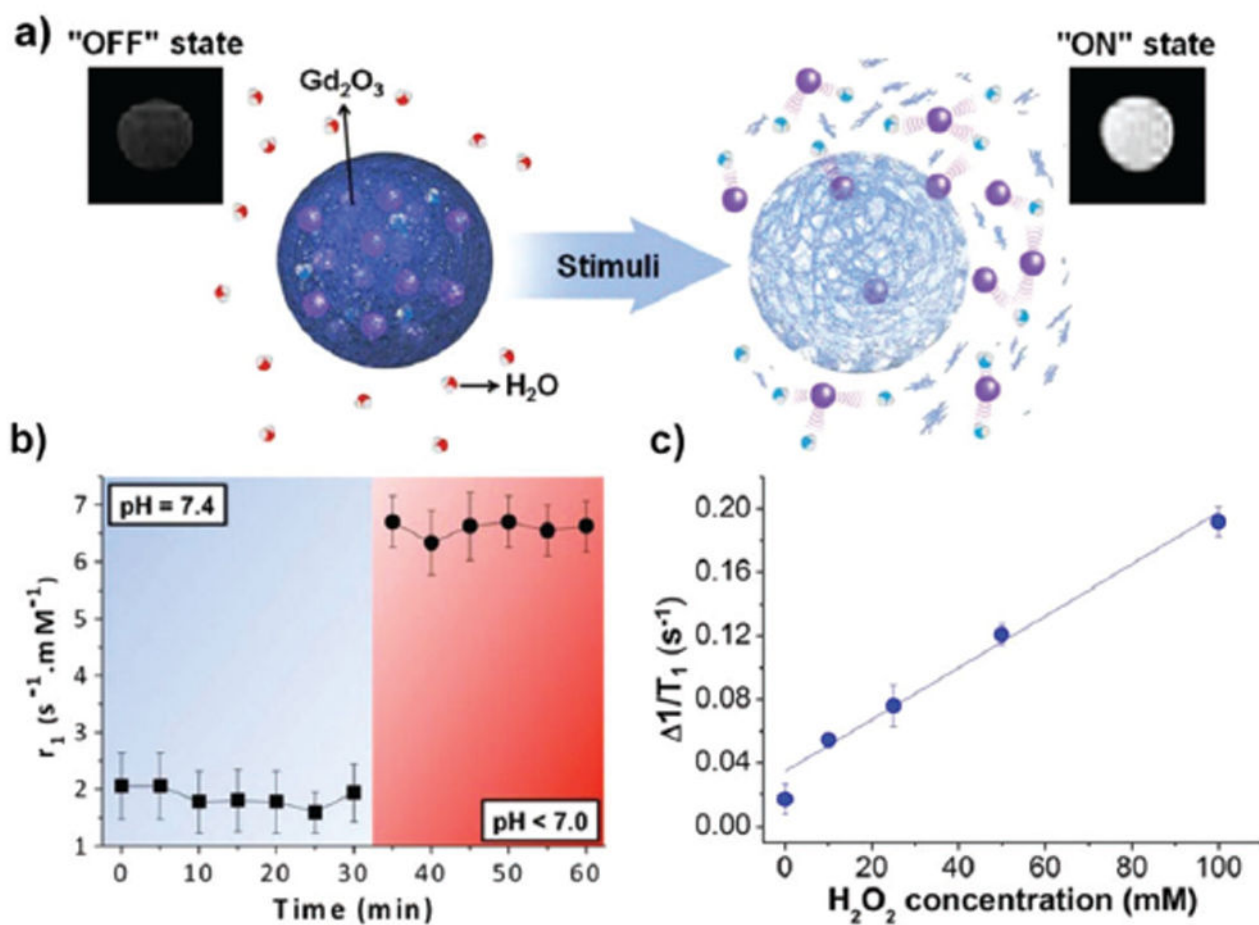


**Fig. 17.**

(a) TEM images, mass magnetization values, magnetic spin structures, magnetic moments,  $T_2$ -weighted images and colormaps of MnFe<sub>2</sub>O<sub>4</sub>, FeFe<sub>2</sub>O<sub>4</sub>, CoFe<sub>2</sub>O<sub>4</sub>, and NiFe<sub>2</sub>O<sub>4</sub> NPs. Reproduced with permission from ref. 84. Copyright 2007, Nature Publishing Group. (b) Plots of  $r_2$  versus Zn<sup>2+</sup> doping level in (Zn<sub>x</sub>M<sub>1-x</sub>)Fe<sub>2</sub>O<sub>4</sub> (M = Mn<sup>2+</sup>, Fe<sup>2+</sup>) NPs at 4.5 T. (c) Comparison of  $r_2$  values of NPs, showing that Zn<sup>2+</sup> doped NPs have significantly enhanced MRI contrast when compared to conventional iron oxide NPs. Reproduced with permission from ref. 83. Copyright 2009, Wiley-VCH.

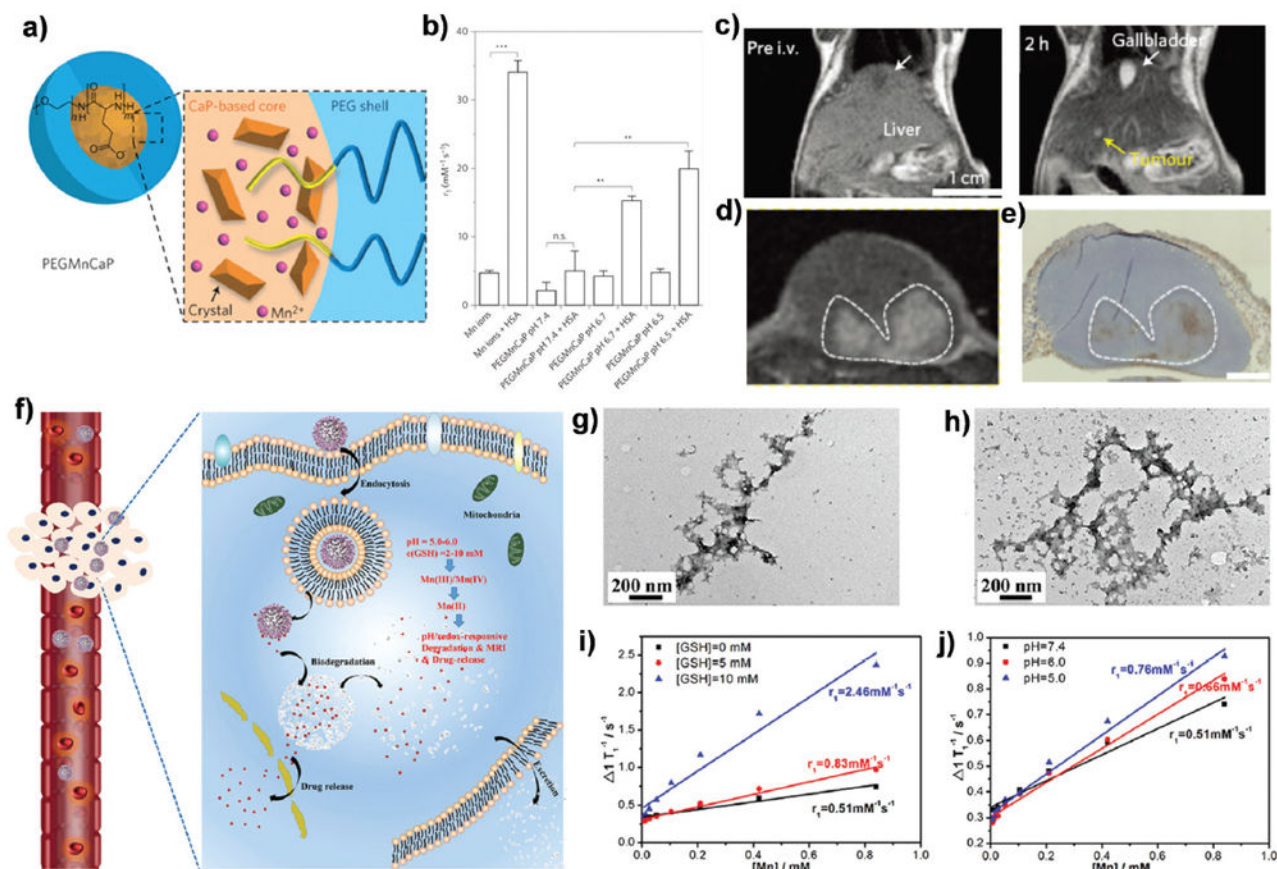
**Fig. 18.**

(a) Schematic of intracellular caspase-3/7-triggered aggregation of  $\text{Fe}_3\text{O}_4@1$  NPs. (b) Plots of  $1/T_2$  versus metal concentration in the presence or absence of caspase-3/7. (c) *In vivo*  $T_2$ -weighted coronal images of  $\text{Fe}_3\text{O}_4@1$  NPs or  $\text{Fe}_3\text{O}_4@1$ -Scr NPs in saline or DOX-treated (i.e., apoptotic) mice in 0 h (top) or 3 h (bottom) post injection. To make tumor apoptotic, about  $8 \text{ mg kg}^{-1}$  of doxorubicin (DOX) was injected intravenously, once every 4 days three times. Reproduced with permission from ref. 197. Copyright 2016, American Chemical Society. (d) Schematic of  $^{99\text{m}}\text{Tc}$ -labeled  $\text{Fe}_3\text{O}_4$  NPs and their responsiveness to GSH-triggering within the tumor microenvironment to form aggregates through inter-particle crosslinking reactions. (e) TEM images of the nonresponsive probe (left panel) and responsive probe (right panel) after being treated with GSH. (f) Temporal evolution of transverse relaxation rate  $R_2$  for both the responsive probe and nonresponsive control recorded on a 3.0 T MRI scanner during incubation with GSH (inset:  $T_2$ -weighted images of probe solutions acquired at different incubation time points). Reproduced with permission from ref. 198. Copyright 2017, Wiley-VCH.

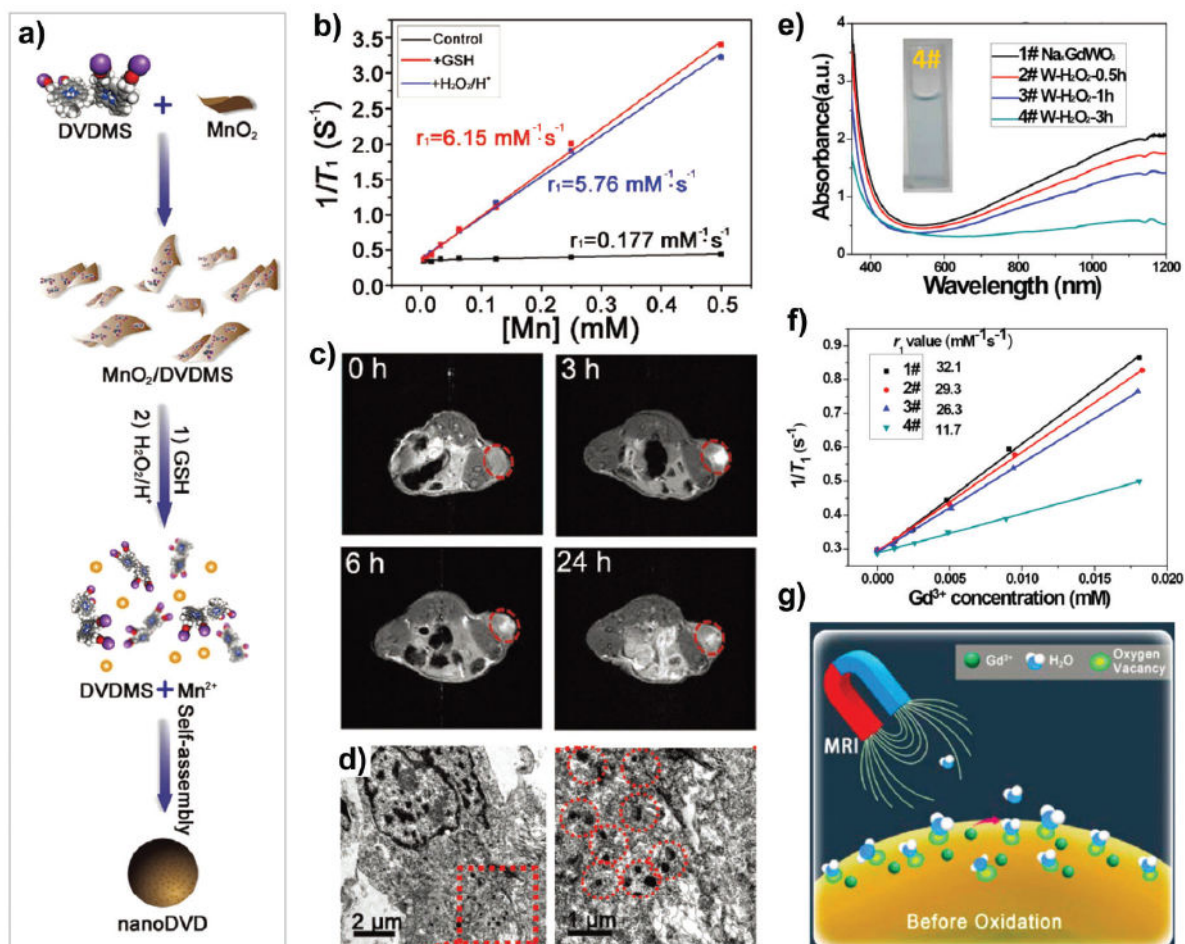


**Fig. 19.**

(a) Schematic of how a degradable polymermatrix is able to control the interaction of water molecules with  $Gd_2O_3$  NPs (purple spheres). (b) Magnetic relaxivity of  $Gd_2O_3$  NPs encapsulated in pH-responsive materials demonstrates a jump from neutral pH to mild acidity. (c) Increasing concentrations of  $H_2O_2$  results in corresponding increases in the  $T_1$  relaxation rates of  $Gd_2O_3$  NPs encapsulated in an  $H_2O_2$ -responsive polymer. Reproduced with permission from ref. 131. Copyright 2013, American Chemical Society.

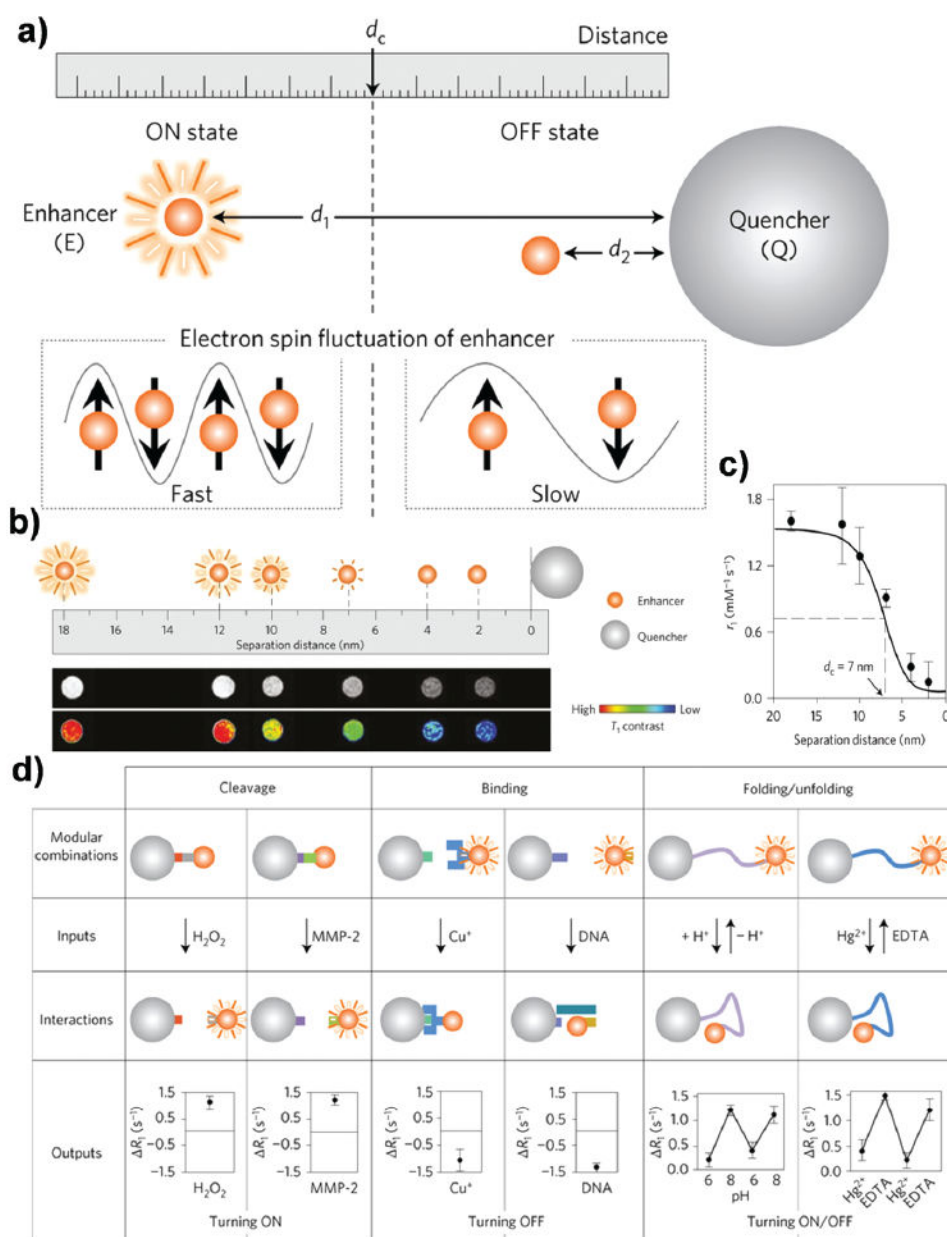
**Fig. 20.**

(a) Schematic of the hybrid structure of PEGMnCaP. PEGMnCaP consists of a CaP-based core and a PEG shell. The  $\text{Mn}^{2+}$  ions are trapped in the CaP core. (b)  $r_1$  value of PEGMnCaP in physiological environments at different pH levels, with and without proteins (for example, HSA). (c) MR images of liver metastasis using 1 T MRI scanner after i.v. injection of PEGMnCaP NPs. Scale bar, 1 cm. (d) MR images of a hypoxic region within a C26 tumor at 1 T, 4 h after the i.v. injection of PEGMnCaP NPs. (e) Staining of tumor tissues with pimonidazole confirmed that the hypoxic regions (brown) are at the same location as the tumor regions with higher MRI contrast enhancement. Scale bar, 1 mm. Reproduced with permission from ref. 52. Copyright 2016, Nature Publishing Group. (f) Schematic of the disassembly of PEG/Mn-HMSNs through “manganese extraction” and release of  $\text{Mn}^{2+}$  component intracellularly. TEM images showing structural evolution of Mn-HMSNs after biodegradation at a GSH concentration of (g) 5.0 and (h) 10 mM at pH 5.0 for 48 h. (i) Plot of  $1/T_1$  versus Mn concentration for PEG/Mn-HMSN solution at varied GSH concentrations and (j) at varied pH. Reproduced with permission from ref. 208. Copyright 2016, American Chemical Society.

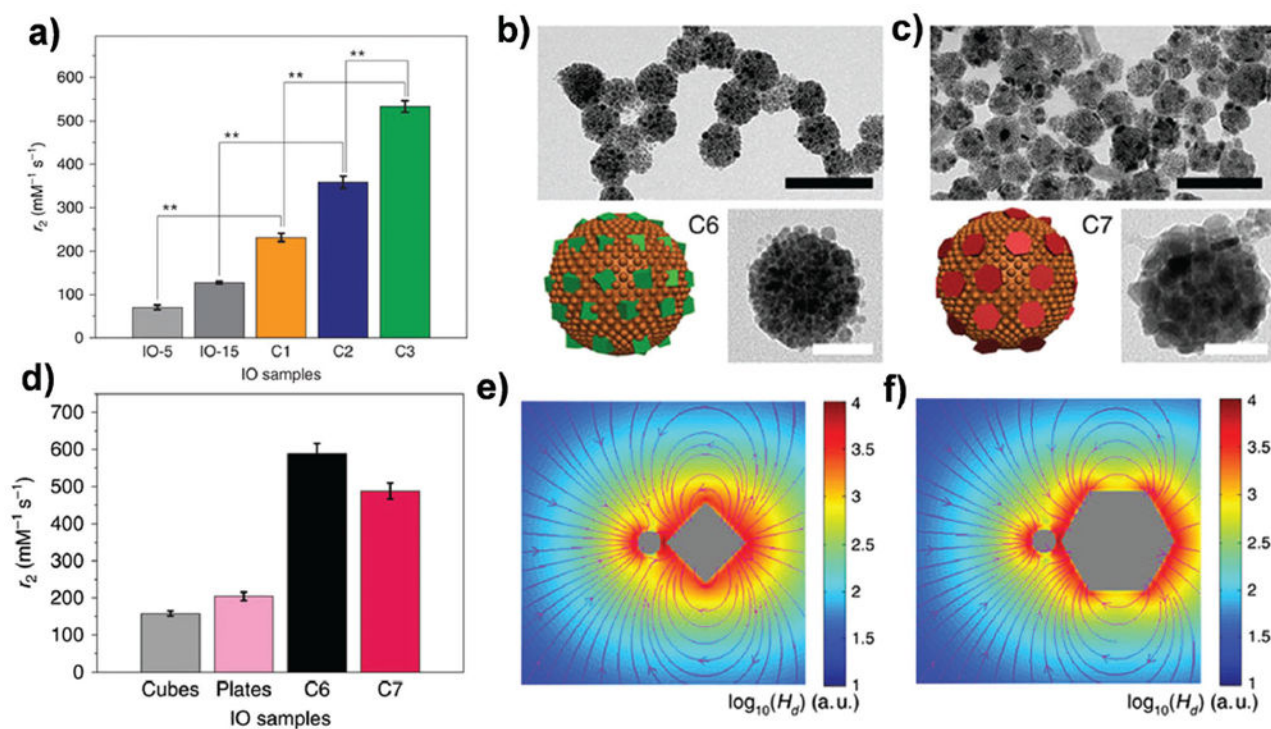
**Fig. 21.**

(a) Schematic illustration showing the fabrication process of  $\text{MnO}_2/\text{DVDMS}$  and the reaction in GSH solution or with  $\text{H}_2\text{O}_2$  (pH 5.5). (b) Plot of  $1/T_1$  versus  $[\text{Mn}]$  for  $\text{MnO}_2/\text{DVDMS}$  (black line),  $\text{MnO}_2/\text{DVDMS} + \text{GSH}$  (red line), and  $\text{MnO}_2/\text{DVDMS} + \text{H}_2\text{O}_2/\text{H}^+$  (blue line) solutions. (c)  $T_1$ -weighted MR images before and after injection of  $\text{MnO}_2/\text{DVDMS}$ . (d) TEM images of MCF-7 tumor sections at 24 h after injection of  $\text{MnO}_2/\text{DVDMS}$ . Reproduced with permission from ref. 211. Copyright 2017, Wiley-VCH. (e) UV-vis-NIR absorption spectra and (f) plot of  $1/T_1$  versus  $\text{Gd}^{3+}$  concentration of PEG- $\text{Na}_x\text{GdWO}_3$  nanorods after oxidation with  $\text{H}_2\text{O}_2$  for varied time periods: 1#: 0 h; 2#: 0.5 h; 3#: 1 h; 4#: 3 h. (g) Schematic of the affinity of oxygen vacancies for oxygen atoms and its impact on the interaction between water molecules and  $\text{Gd}^{3+}$  ions. Reproduced with permission from ref. 37. Copyright 2017, American Chemical Society.



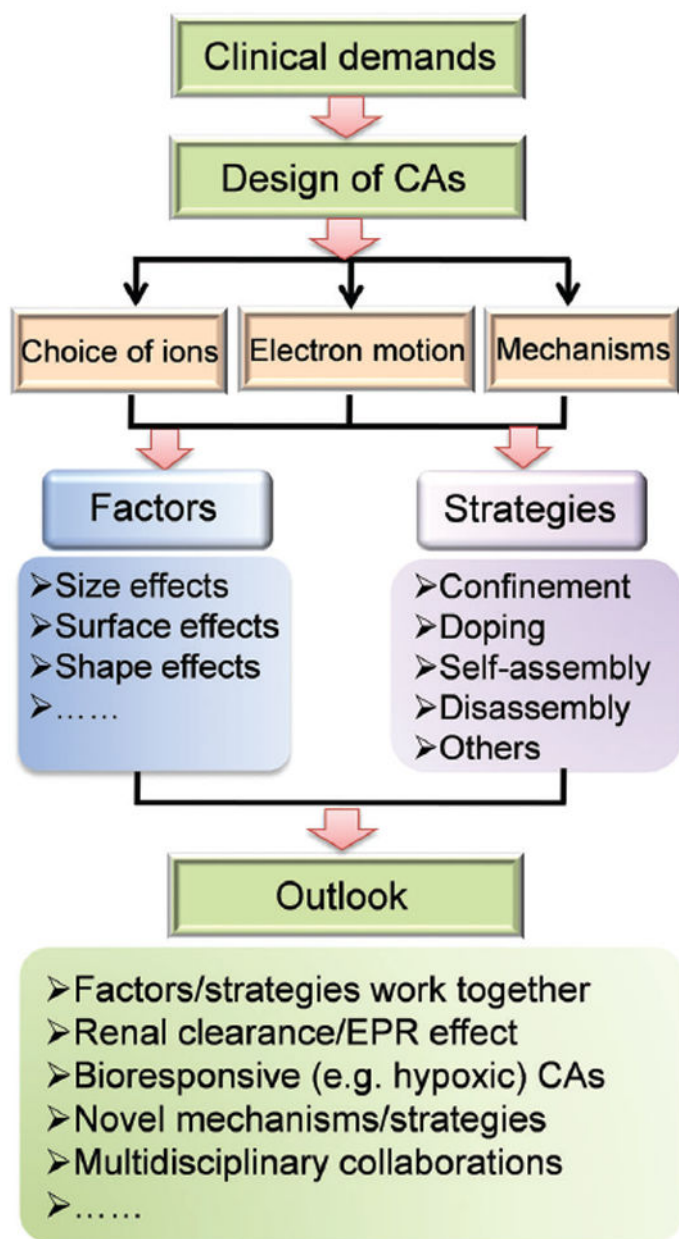


**Fig. 22.** (a) Schematic representation of distance-dependent magnetic resonance tuning. (b) Schematic of a nanoscale ruler that shows a variable  $T_1$  MRI signal, dependent on the separation distance between the paramagnetic enhancer (Gd-DOTA) and superparamagnetic quencher (*i.e.*, 12 nm  $\text{Zn}_{0.4}\text{Fe}_{2.6}\text{O}_4$  particle). Lower panel:  $T_1$ -weighted MR image and color map image of a solution containing the enhancer and quencher with varied separation distances. (c) A plot of  $r_1$  values *versus* separation distance. (d) Schematic of modular combinations for the preparation of distance-dependent magnetic resonance sensors, operated using three different modes of interactions (cleavage, binding, and conformational changes) and subsequent  $T_1$  MRI signal outputs. Reproduced with permission from ref. 213. Copyright 2017, Nature Publishing Group.



**Fig. 23.**

(a) Columns showing  $r_2$  values of iron oxide (IO) clusters, as well as the single IO-5 and IO-15 NPs. C1: clusters of 5 nm NPs only; C2: clusters of 15 nm NPs only; C3: clusters of mixed 5 nm NPs and 15 nm NPs. (b and c) TEM and high-resolution TEM images, as well as cartoons of clusters C6 (IO cubes) and C7 (IO nanoplates), respectively. (d) Columns showing  $r_2$  values of IO clusters C6 and C7, as well as the single IO cubes and plates. (e and f) Simulation models and calculated stray fields for clusters C6 and C7, respectively. Color bars represent  $\log_{10}(H_d)$ , where  $H_d$  is the calculated stray field. Reproduced with permission from ref. 214. Copyright 2017, Nature Publishing Group.



**Fig. 24.** Summary and outlooks of designing inorganic NPs for MRI.

Table 1

Outer orbital and quantum parameters of metal ions

Ion	Orbital	S	L	J	$(\mu_{\text{eff}}/\mu_B)$
Mn <sup>2+</sup>	3d <sup>5</sup>	5/2	0	5/2	5.92
Mn <sup>4+</sup>	3d <sup>3</sup>	3/2	3	9/2	3.87
Fe <sup>2+</sup>	3d <sup>6</sup>	2	2	4	4.90
Fe <sup>3+</sup>	3d <sup>5</sup>	5/2	0	5/2	5.92
Co <sup>2+</sup>	3d <sup>7</sup>	3/2	3	9/2	3.87
Co <sup>3+</sup>	3d <sup>6</sup>	2	2	4	4.90
Cu <sup>2+</sup>	3d <sup>9</sup>	1/2	2	5/2	1.73
Eu <sup>2+</sup>	4f <sup>7</sup>	7/2	0	7/2	7.94
Eu <sup>3+</sup>	4f <sup>6</sup>	3	3	0	—
Gd <sup>3+</sup>	4f <sup>7</sup>	7/2	0	7/2	7.94
Dy <sup>3+</sup>	4f <sup>9</sup>	5/2	5	15/2	10.65
Ho <sup>3+</sup>	4f <sup>10</sup>	2	6	8	10.61

Table 2

Representative inorganic nanoparticle size effect on relaxivity

NPs	Shape	Approximate size (nm)	Surface coating	Relaxivity (mM <sup>-1</sup> s <sup>-1</sup> )	Field (T)	Ref.
NaGdF <sub>4</sub>	Sphere	2.5/4.0/6.5/8.0	PVP	$r_1 = 7.2/4.5/3.3/3.0$	1.5	49
NaGdF <sub>4</sub>	Sphere	5/15/20	PEG	$r_1 = 6.2/5.7/8.78$	3.0	41
NaGdF <sub>4</sub>	Sphere	2.1/3.2/8.6	PAA	$r_1 = 47.3/27.5/15.3$	0.5	36
NaGdF <sub>4</sub>	Sphere	3/4/5	PEG	$r_1 = 58.8/39.4/28.6$	1.4	39
Gd <sub>3</sub> O <sub>2</sub>	Sphere	2.2/3.8/4.6	SiPEG	$r_1 = 8.8/8.8/4.4$	7.0	45
Gd <sub>3</sub> O <sub>2</sub>	Sphere	4.4/5.4/10.1	Albumin	$r_1 = 18.49/16.22/12.26$	1.5	129
MnO	Irregular	7/15/20/25	PEG	$r_1 = 0.37/0.18/0.13/0.12$	3.0	69
Mn <sub>2</sub> O <sub>3</sub>	Sphere	3/5/11/13	PEG	$r_1 = 2.38/1.39/0.99/0.41$	1.5	132
Mn <sub>2</sub> O <sub>3</sub>	Tetrahedron	6/8/20	PEG	$r_1 = 1.08/0.63/0.25$	1.5	132
Fe <sub>2</sub> O <sub>3</sub>	Sphere	2.2/3/12	PEG	$r_1 = 4.78/4.77/2.37$	1.5	135
MnFe <sub>2</sub> O <sub>4</sub>	Sphere	2/3/3.9	PEG	$r_1 = 8.43/8.23/6.98$	3.0	136
Fe <sub>3</sub> O <sub>4</sub>	Sphere	4/6/9/12	DMSA	$r_2 = 78/106/130/218$	1.5	79
Fe <sub>3</sub> O <sub>4</sub>	Octapod	20/30	HDA-G2	$r_2 = 679/209$	7.0	137
Fe <sub>3</sub> O <sub>4</sub>	Sphere	8/23/37/65	PVP	$r_2 = 173/204/240/249$	7.0	138
Fe <sub>3</sub> O <sub>4</sub>	Rod	30 <sup>a</sup> /40/50/60/70	PEI	$r_2 = 312/381/427/545/608$	3.0	80
MnFe <sub>2</sub> O <sub>4</sub>	Sphere	6/9/12	DMSA	$r_2 = 208/265/358$	1.5	84
NaDyF <sub>4</sub>	Sphere	5.4/9.8/20.3	PEG	$r_2 = 32/51/101$	9.4	19
NaDyF <sub>4</sub>	Rod	19 <sup>a</sup> /20/25/35	PEG	$r_2 = 75/65.2/91.4/204.4$	9.4	20
NaHoF <sub>4</sub>	Sphere	10.6/14.1/16.4	PEG	$r_2 = 31.9/51.5/70.4$	9.4	20
NaHoF <sub>4</sub>	Sphere	3.2/7.4/13.2	PEG	$r_2 = 83.4/45.4/131.7$	7.0	21

<sup>a</sup>Length of nanorods.

**Table 3**

Full names and their corresponding abbreviations

Full name	Abbreviation	Full name	Abbreviation
5-Aminolevulinic acid	5-ALA	Mesoporous silica nanoparticles	MSNs
Catechol	CC	Mercaptosuccinic acid	MSA
Chemical exchange saturation transfer	CEST	Motional averaging regime	MAR
Computed X-ray tomography	CT	Nanoparticles	NPs
Dense silica	dSiO <sub>2</sub>	Nanorods	NRs
Diethylenetriaminepentaacetic acid	DTPA	Nuclear magnetic relaxation dispersion	NMRD
Diffusion correlation time	$\tau_D$	Number of bound water molecules	$q$
Diphosphate	DP	Poly(acrylic acid)	PAA
Contrast agents	CAs	Photoacoustic	PA
Distearyl phosphatidylethanolamine	DSPE	Polyethylene glycol	PEG
2,3-Dimercaptosuccinic acid	DMSA	Polyethylenimine	PEI
Effective magnetic moment	$\mu_{\text{eff}}$	Polyvinylpyrrolidone	PVP
Electron spin angular momentum	$S$	Poly(maleic anhydride- <i>alt</i> -1-octadecene)	PMAO
Electron orbital angular momentum	$L$	Poly(maleic acid)-octadecene	PMO
Enhanced permeation and retention	EPR	Reticuloendothelial system	RES
Glutathione	GSH	Silicon microparticles	SiMPs
1-Hexadecylamine	HDA	Rotational correlation time	$\tau_R$
High-resolution	TEM	HRTEM Saturation magnetization	$M_s$
Hollow mesoporous silica nanoparticles	HMSNs	Sinoporphyrin sodium	DVDMS
Hydroxamate	HX	Solomon–Bloembergen–Morgan	SBM
Lanthanide	Ln	Static dephasing regime	SDR
Magnetic resonance angiography	MRA	Transmission electron microscopy	TEM
Magnetic resonance imaging	MRI	Total electron angular momentum	$J$
Magnevist	Gd–DTPA	Upconversion nanoparticles	UCNPs
Mesoporous silica	mSiO <sub>2</sub>	Water residence time	$\tau_M$