Atomic layer deposition and properties of ZrO₂/Fe₂O₃ thin films

Kristjan Kalam^{*1}, Helina Seemen¹, Peeter Ritslaid¹, Mihkel Rähn¹, Aile Tamm¹, Kaupo Kukli^{1,2}, Aarne Kasikov¹, Joosep Link³, Raivo Stern³, Salvador Dueñas⁴, Helena Castán⁴ and Héctor García⁴

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Address:

¹Institute of Physics, University of Tartu, W. Ostwald 1, 50411 Tartu, Estonia, ²Department of Chemistry, University of Helsinki, P. O. Box 55, FI-00014 Helsinki, Finland, ³National Institute of Chemical Physics and Biophysics, Akadeemia tee 23, 12618 Tallinn, Estonia and ⁴Department of Electronics, University of Valladolid, Paseo Belén 15, 47011 Valladolid, Spain

Email:

Kristjan Kalam* - kristjan.kalam@ut.ee

* Corresponding author

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Abstract

Thin solid films consisting of ZrO_2 and Fe_2O_3 were grown by atomic layer deposition (ALD) at 400 °C. Metastable phases of ZrO_2 were stabilized by Fe_2O_3 doping. The number of alternating ZrO_2 and Fe_2O_3 deposition cycles were varied in order to achieve films with different cation ratios. The influence of annealing on the composition and structure of the thin films was investigated. Additionally, the influence of composition and structure on electrical and magnetic properties was studied. Several samples exhibited a measurable saturation magnetization and most of the samples exhibited a charge polarization. Both phenomena were observed in the sample with a Zr/Fe atomic ratio of 2.0.

Introduction

Doped ZrO_2 has been a subject of interest because of several potential applications, for example, in microelectronics as a memory material [1]. Also, doping a dielectric film with a magnetic material might provide the structural distortion required to stabilize the ferroelectric phase thus resulting in a multiferroic material, which would allow an additional degree of freedom in device design [2].

ZrO₂ doped with various chemical elements has been studied for several applications and different processes have been em-

ployed to prepare the samples. Ca- and Mg-stabilized cubic zirconia, prepared by pulsed laser deposition (PLD), has shown ferromagnetic properties [3]. Magnetic properties of PLD-synthesized ZrO₂, doped with Co, Fe, Mn or Ni, have been studied [4], showing that doping ZrO₂ with Mn results in a significantly higher saturation magnetization than doping ZrO₂ with the other transition metals studied. Fe–ZrO₂ nanocomposite thin films have been synthesized using a solid state reaction between the Zr and Fe₂O₃ layers, and their composition, structure, chemical stability and magnetic properties were characterized

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[5]. Upon annealing at 500 °C, the metastable cubic phase of ZrO2 was stabilized and ferromagnetic hysteresis of the nanocomposite film was confirmed. Saturation magnetization was measured to be ≈173 emu/g [5]. Undoped ZrO₂, prepared by pulsed electron beam deposition [6] or reactive DC magnetron sputtering [7], also exhibited ferromagnetic properties. The undoped ZrO₂ exhibited ferromagnetic properties mainly driven by oxygen vacancies. Monoclinic and tetragonal phases with similar amounts of oxygen vacancies were compared and ferromagnetism was only observed in the case of the tetragonal phase [7]. Microwave-assisted combustion synthesis of a powder and the subsequent sintering of samples were used to fabricate ZrO2 doped with Co, and ferromagnetism in such samples was confirmed [8]. Different nanostructures of undoped ZrO₂ were prepared by catalyst-assisted PLD and all structures were found to possess ferromagnetic behavior [9]. Ferromagnetism was also observed in annealed Co and Fe co-doped ZrO₂, prepared by the sol-gel method [10]. Samples of Mn- and Fe-stabilized cubic zirconia were obtained by a co-precipitation method and no ferromagnetism was observed in such samples [11]. Phase diagrams for the ZrO₂-FeO system were described [12] and the influence of thermal treatment on the phase development in ZrO₂-Fe₂O₃ and HfO₂-Fe₂O₃ systems was assessed [13].

ALD of ZrO₂ from ZrCl₄ and O₃ has been studied [14]. Reactions between Fe(acac)₃ adsorbing on zirconia surfaces [15,16] has been studied as well. Phase stabilization of ZrO2 by Fe doping was investigated by using ALD [17]. Ferromagnetism in ALD-grown Fe₃O₄/ZrO₂/Fe₃O₄ multilayer nanotubes has been demonstrated [18], while the precursors for distinct solid oxide layers constituting these samples were ferrocene/ozone and tetrakis(dimethylamido)zirconium(IV)/water for Fe₂O₃ and ZrO₂, respectively, and Fe₂O₃ was reduced to Fe₃O₄ after the growth. In another study [19], ZrO₂/Fe thin films were prepared by ALD from β-diketonate precursors and ozone. After annealing at 600 °C in N2 flux for 60 s, the films exhibited ferromagnetic properties. The ferroelectric properties of ALDgrown undoped zirconia have also been investigated [20]. ZrO₂ thin films were deposited by remote plasma ALD from tetrakis(dimethylamido)zirconium(IV) and oxygen plasma, and were found to exhibit ferroelectric behavior. ALD and physical vapor deposited ZrO2 were compared from the viewpoint of ferroelectric behavior [21], whereby ALD precursors were Zr-based metal organic precursors (TEMAZr) and H₂O. Both fabrication methods provided samples with antiferroelectric behavior.

Here, ZrO₂ films doped with Fe₂O₃ were grown by atomic layer deposition from zirconium chloride and ferrocene precursors. The purpose of this study was to investigate the effects of the

Zr/Fe cation ratio on the film structure as well as on the magnetic and electrical properties, and to examine whether alternately deposited iron and zirconium oxides form dilute solids or become segregated upon thermal deposition and processing. The goal of the magnetic and electrical measurements was to evaluate the ability of ZrO₂/Fe₂O₃ films to polarize in both electric and magnetic fields and thus clarify to what extent may such materials exhibit multiferroic behavior.

Experimental

The ZrO₂/Fe₂O₃ films were grown in a low-pressure (200-260 Pa) flow-type in-house built hot-wall ALD reactor [22] at 400 °C. Zirconium tetrachloride, ZrCl₄ (Aldrich, 99.99%), and ferrocene, Fe(C₅H₅)₂ (ABCR, 99%), were used as zirconium and iron precursors, respectively. Ozone, O3, was used as the oxidizer. Nitrogen, N2 (99.999% purity, AGA), was applied as the carrier and purging gas. At the temperature chosen (400 °C) the ZrO₂ grows efficiently from ZrCl₄ and O₃ [14], and this temperature is also sufficiently high to ensure efficient growth also for iron oxide from cyclopentadienyls and ozone [23]. ALD growth of Fe₂O₃ from ferrocene and ozone is even possible at 200 °C, but at such a low temperature, the required pulse duration is very long. A maximal growth rate was achieved when ferrocene pulses were 40 s long and ozone pulses were 200 s [24]. ZrCl₄ and Fe(C₅H₅)₂ were evaporated at 161-163 °C and 61-63 °C, respectively, from open boats inside the reactor and transported to the substrates by the carrier

Ozone was produced from O_2 (99.999% purity, AGA) using a BMT Messtechnik 802 N generator. The ozone concentration, measured using a BMT Messtechnik 964 analyzer, was 245–250 g/m³ in the experiments, which is 17.1–17.5%. The estimated ozone flow rate from the generator was about 67 sccm, while the carrier gas flow rate was kept at about 220 sccm.

The ZrO_2/Fe_2O_3 films were grown by alternately applying certain amounts of constituent binary oxide growth cycles. The ZrO_2/Fe_2O_3 cycle ratio was varied as 10:1, 10:3, 10:5, 10:10, and 5:5. The cycle times used for both ZrO_2 and Fe_2O_3 were 5–5–5–5 s for the sequence of metal precursor pulse–purge– O_3 pulse–purge. Each deposition started and ended with pulses of ZrO_2 in order to make the film symmetrical from electrode to electrode in terms of the chemical composition. Altogether, 220–235 ALD cycles were deposited to obtain each thin film.

The films were grown on various substrates: Si(100) and highly doped conductive Si substrates covered by a 10 nm TiN film grown by chemical vapor deposition. Before deposition, the Si(100) was rinsed in a mixture of H₂SO₄/H₂O₂ 5:2 and heated

on a hot plate at 80 °C. After that, the Si(100) was cleaned with distilled water in an ultrasound container. The next step was to clean the samples with a 7% solution of HF and again clean with distilled water under ultrasound. During the time after cleaning and before placing the samples in a reactor, a $\rm SiO_2$ layer of a few nanometers thick grew on the sample because of contact with surrounding air.

Selected samples on Si(100) were annealed at 850 °C and 1000 °C in air for 15 and 10 minutes, respectively. For electrical measurements, the films deposited on TiN substrates were supplied with platinum electrodes, with an area of 0.204 mm², that were electron-beam-evaporated on top of the films. In addition, for investigating deposition uniformity in terms of conformal growth, stacked three-dimensional structures of silicon substrates with an aspect ratio of 1:20 were used. Such three-dimensional trenched or stacked structures are similar to those used in micro- and nanoelectronics as capacitor bottom electodes, allowing the exploitation of capacitor arrays with enlarged memory density [25-27].

The crystal structure was evaluated by grazing incidence X-ray diffractometry (XRD), using an X-ray diffractometer (SmartLab, Rigaku) using Cu K α radiation, which corresponds to an X-ray wavelength of 0.15406 nm. A spectroscopic ellipsometer (model GES5-E) was used for the evaluation of the film thickness. The Tauc–Lorentz model was used for the modelling of samples in the range of 300–1000 nm. An X-ray fluorescence (XRF) spectrometer (Rigaku, ZSX 400) with the software program ZSX (version 5.55) was used to evaluate the elemental composition of the films. The surface morphology of the films and a cross-section of an ALD coated stack were evaluated by scanning electron microscopy (SEM) using a dual beam SEM (FEI, Helios NanoLab 600). An FEI Titan Themis 200 device was used for transmission electron microscopy (TEM) imaging of selected samples.

The electrical measurements were performed by means of an Agilent DXO-X 3104 digital oscilloscope with a built-in wave

generator. The standard Sawyer–Tower experiment was carried out by applying a periodic triangular-shaped stimulus and recording the voltage loop data from the oscilloscope. The charge values were obtained from the sensed voltage across a stated capacitance.

Magnetic measurements were performed using a vibrating sample magnetometer (VSM) option of the physical property measurement system (14T, Quantum Design) by scanning the magnetic field from -1 T to 1 T parallel to the film surface at room temperature.

Results and Discussion Film growth and composition

The film thickness varied between 15 and 40 nm. Regarding the composition, the cation ratio of mixed films in the as-deposited state, measured by X-ray fluorescence spectrometry, varied from Zr/Fe 0.15-50. Table 1 gives a list of the samples with a description of the growth cycle sequences as well as the observed range of composition in terms of the constituent metal ratio. Possibly, due to the different adsorption and nucleation rates of different metal precursors, the marked deviations in the chemical composition could not be avoided. One can also see that the oxide cycle ratio does not directly match with the cation ratio. This was actually expected because the nucleation of every oxide is slower with the growth rate retarded during the first few deposition cycles, and the content of metal deposited does not linearly correspond to the amount of additive cycles at the early stages of the growth. The nonlinearity of the thickness as a function of growth per cycle was also observed with HfO2 [28] and ZrO₂ [29].

The growth rate for a single deposition was different for each sample in the reactor because of the reactor type and positioning of the samples (Figure 1). The growth rate decreased with increasing distance between the sample and the valve, which releases precursors into the reactor. Pure Fe₂O₃ exhibited the highest growth rate of 0.18 nm/cycle at a position, which was closest to the precursor valve. The growth rate of

Table 1: List of the ZrO_2/Fe_2O_3 films, with constituent oxide cycle ratios and full growth cycle sequences indicated, deposited on Si substrates (subjected to magnetometry) and TiN substrates (subjected to electrical measurements). The range of cation ratios is due to the variation of the ratios measured at different locations on the substrate holder and are due to the film thickness growth rate profiles along the gas flow direction. In these experiments, lower cation ratios were measured on TiN substrates as compared to Si substrates.

Cycle ratio	Cycle sequence	Zr/Fe cation ratio	Thickness on Si	Thickness on TiN
10:3	17 × [10 × ZrO ₂ + 3 × Fe ₂ O ₃] + 10 × ZrO ₂	10	23 nm	21 nm
10:5	$15 \times [10 \times ZrO_2 + 5 \times Fe_2O_3] + 10 \times ZrO_2$	0.34-2.0	22 nm	15 nm
10:10	$11 \times [10 \times ZrO_2 + 10 \times Fe_2O_3] + 10 \times ZrO_2$	0.21-1.7	26 nm	19 nm
5:5	$22 \times [5 \times ZrO_2 + 5 \times Fe_2O_3] + 5 \times ZrO_2$	0.14-0.16	21 nm	36 nm

ZrO₂ was 0.12 nm/cycle at the same position. Mixtures of these oxides had lower growth rates than individual Fe₂O₃ or ZrO₂, varying around 0.1 nm/cycle (Figure 1).

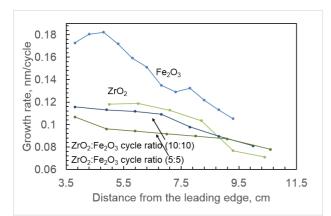


Figure 1: Growth rates of Fe_2O_3 , ZrO_2 and their mixtures with different cycle ratios as a function of sample position in the ALD reactor (distance between the sample and valve). All samples of mixed oxides exhibited lower growth rates than pure Fe_2O_3 or ZrO_2 . Mixtures with cycle ratios not shown in the image exhibited growth rates very similar to the shown mixtures.

The thickness (growth rate) profiles along the gas flow direction may be regarded as characteristic of cross-flow atomic layer deposition reactors, in which the substrate holder is designed in a way that leaves one edge of the substrate pointing towards the gas flow inlet and the other, trailing, edge closer to the pumping line. In the case of such arrangement, the gas flow in the reactor chamber is parallel to the surface plane from one fixed point towards another. Although the precursor pulses are separated by purge periods, the separation is not perfect during the actual operation. The adsorption waves may partially overlap and meet in the vicinity of the substrate surface, causing the chemical vapor deposition to be less controlled. At the leading edge, the film thickness is usually higher and gradually decreases towards the trailing edge. One can recognize this phenomenon on the basis of the curves in Figure 1. The extent of the profile depends on the reactivity of the precursors, process temperature, and pulse time parameters, and has been experimentally observed, for example, in an early study devoted to ALD of ZrO₂ from ZrCl₂ and H₂O [30] and later in the case of ALD of ZrO₂ from ZrCl₄ and O₃ [14]. This was generalized in a more theoretical study on the adsorption process during ALD [31]. Such profiles are likely and expected to be different for each constituent material.

Although the thickness and composition were not constant during a single deposition in the present study, all comparisons between different depositions have been made using samples that were at the same location in the reactor (i.e., that have the same distance from the leading edge). The analysis of the light residue in the host materials similar to those grown in the present study has been described earlier in terms of the most critical residue components, i.e., chlorine and hydrogen [14]. The time-of-flight elastic recoil detection analysis verified quite low residual chlorine content with an average amount of 0.14 \pm 0.01 atom % in the films grown at 300 °C. Similar, rather low amounts were found in the films grown at 400 °C. The amount of hydrogen was 0.30 \pm 0.03 atom % in the film grown at 300 °C, and decreased to 0.14 \pm 0.03 atom % in the film grown at 400 °C.

Film structure

It was shown that doping ZrO₂ with Fe₂O₃ stabilized the tetragonal/cubic phase (PDF Card 70-6628) of ZrO₂ for all substrates as long as the cation ratio of Zr/Fe exceed 2.0 (Figure 2 and Figure 3).

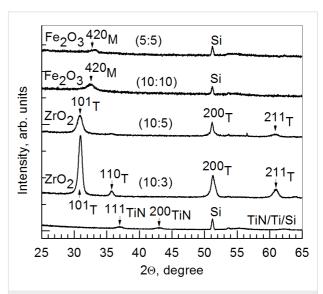


Figure 2: Grazing incidence X-ray diffraction (XRD) patterns for ZrO_2/Fe_2O_3 films deposited on TiN with ZrO_2/Fe_2O_3 cycle ratios and thickness indicated in the labels. The Miller indices are attributed to corresponding monoclinic (M) and tetragonal (T) phases of pure Fe_2O_3 and ZrO_2 , respectively. The cation ratio of Zr/Fe in the films deposited on TiN with cycle ratios of 5:5, 10:10, 10:5 and 10:3 were 0.15, 1.7, 2.0 and 10, respectively.

Lamperti et al. also determined the phase stabilization of ZrO₂ by Fe doping [17] or otherwise in an ozone-based ALD of ZrO₂ [32]. The stabilization of the tetragonal/cubic phase in Fe-doped ZrO₂ has been confirmed by de Souza et al. [33] for samples fabricated by freeze-drying process and Kuryliszyn-Kudelska et al. [34] for samples prepared via wet chemistry. For the samples grown on TiN substrates, films with a higher Fe₂O₃ concentration did not show peaks assignable to any zirconia phase, but demonstrated reflections from crystallized Fe₂O₃ in the monoclinic phase (PDF Card 016-0653) (Figure 2). Ternary iron–zirconium oxide phases were not recognized. It is worth

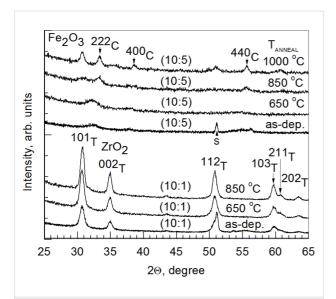


Figure 3: Grazing incidence X-ray diffraction (XRD) patterns for ZrO_2/Fe_2O_3 films deposited on Si(100) with ZrO_2/Fe_2O_3 cycle ratios and annealing temperatures indicated in the labels. The Miller indices are attributed to corresponding tetragonal (T) and cubic (C) phases of pure ZrO_2 and Fe_2O_3 , respectively.

noting that, despite few existing works indicating the target material with the stoichiometry of ZrFe₂O₅, the synthesis of ZrFe₂O₅ has not been convincingly completed. Instead, one can believe that the solubility of iron oxide in zirconia is rather low and, in addition, it requires rather aggressive heat treatments. No crystallographic traces of ternary ZrFe₂O₅ have been registered, and distinct phases of Fe₂O₃ and ZrO₂ have also earlier determined on the basis of the XRD patterns [13,35]. In our study, also alternative, zirconium-rich, ternary phases Zr₂FeO_x (PDF Card 019-0646), Zr₄Fe₂O_{0.6} (PDF Card 038-1168), and Zr₆Fe₃O (PDF Card 017-0559) were not recognized in XRD patterns. Samples on TiN and Si(100) substrates were different in their composition due to their different position in the reactor and the effect of a certain lateral composition profile. This is probably caused by differences in the adsorption rates of the precursors and concurrently different growth rate profiles (Figure 1) due to the appearance of the thickness profiles as described above.

Annealing the samples of lightly doped zirconia caused the grain growth in metastable tetragonal/cubic phase (Figure 3). Annealing samples with a ZrO_2/Fe_2O_3 cycle ratio (10:1) and Zr/Fe = 50 resulted in the increased crystallization of the tetragonal/cubic phase of ZrO_2 (Figure 3). This was also observed in the case of Zr/Fe = 10 (not shown). A similar result was also obtained by Štefanić et al. [13] who also showed that lightly Fe-doped ZrO_2 crystallized in tetragonal/cubic zirconia phase after annealing. In the case of Zr/Fe = 0.34 and a cycle ratio of 10:5, the film was almost amorphous in the as-deposited state,

but the crystallization of tetragonal/cubic ZrO₂ and cubic Fe₂O₃ (PDF Card 04-0755) was observed upon annealing at 1000 °C in air for 10 min (Figure 3).

Since the ultimate goal of any materials synthesis process are its application at the industrial scale, preliminary steps illustrating the potential ability to grow electrically or magnetically interesting films conformally on three-dimensional substrates should be considered. For that purpose, the films were deposited on a 3D stacked silicon surface with an aspect ratio of 1:20 (Figure 4). It can be seen in Figure 4 that the film is indeed uniform on the 3D surface and it is also important to demonstrate that atomic layer deposition provides a solution for how to deposit the film inside the pattern, without covering the top. The film surface scans on planar substrates indicated certain changes in the morphology, likely related to crystal growth and roughening of the surface already in as-deposited state (Figure 4).

Constituent layers of the sample can be seen in the transmission electron microscope image (Figure 5). The interface between the oxide film and TiN electrode layer is distinguishable; the film is dense and the crystal growth has started immediately together with the nucleation process. The TEM image proves the formation of a stacked layer structure.

Electrical and magnetic properties

Most samples exhibited charge polarization, as shown in Figure 6. Charge polarization was found to decrease as the amount of Fe₂O₃ in the films increased (Figure 6). The charge polarization–applied voltage loops, in principle, look similar to those observed earlier for materials grown by ALD and targeted as ferroelectrics, such as LiNbO₃ [36,37], Bi₄Ti₃O₁₂ [38], or Pb(Zr,Ti)O_x [39]. However, these loops are not yet to be attributed to the properties of a purely ferroelectric material, but an overwhelming contribution from the interfacial polarization should be taken into account. On the other hand, for comparison, in the case of ALD-grown non-centrosymmetric orthorhombic phase of HfO₂ stabilized by doping with foreign cations [40], well-defined ferroelectric hysteresis was recorded.

For stabilized hafnia [40] the electrical polarization charge clearly tended to saturate upon the incremental increase in the external electric field strength. Along with the backward sweep of the applied voltage towards zero field, the polarization in the material initially tended to retain a value close to the saturation level, then decreased more rapidly to a certain finite level at zero field, allowing one to consider that there exists a remnant polarization value, characteristic of classical ferroelectric behavior. Furthermore, the application of a minimum external field in the opposite direction was required to draw the internal

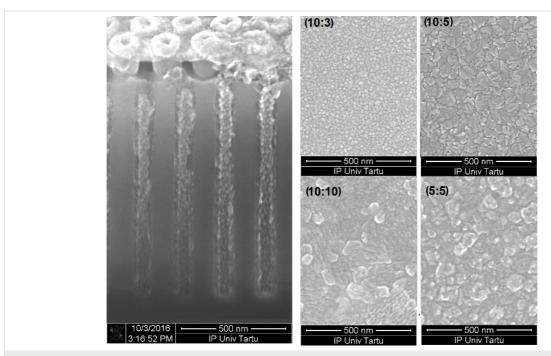


Figure 4: SEM image of a ZrO_2/Fe_2O_3 film grown in a 3D stacked structure (left panel) and images of ZrO_2/Fe_2O_3 film surfaces (right panel). The cycle ratios of the respective films are shown on each image of the right panel.

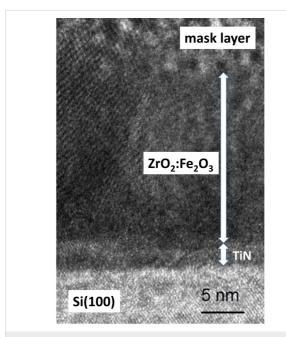


Figure 5: TEM image of a ZrO_2/Fe_2O_3 sample with cycle ratio 10:3 on TiN. 231 atomic layer deposition cycles were deposited.

polarization to zero, allowing one to account for a coercive force [40].

In the doped ZrO₂ films grown in the present study, any saturation level for polarization was actually not quite achieved at either polarity of the external field, as the polarized charge kept increasing with the voltage. Most probably, the charge (which is responsible for the polarization in the material deposited) was due to the electric field that drifted from an electrode to the counterelectrode. That charge, either electronic or ionic, may become trapped at the interface between the metal oxide layer and the electrode, giving rise to the interfacial polarization. It must be considered that in the vicinity of an interface a material is always most prone to defects and its lattice the most open. This is due to the usual mismatch between the crystal structure of the substrate and functional layer. The 10 nm thick TiN electrode layers were nanocrystalline, without preferred orientation, demonstrating weak and broad 111 and 200 reflections (Figure 2). ZrO₂ and Fe₂O₃ phases do not possess lattice structure, allowing commensurate growth on TiN. The epitaxial relationship between magnetite (Fe₃O₄) and titanium nitride might be considered [41], but magnetite was not recognized in this study in the XRD patterns and, most importantly, the very first layers deposited in contact with the electrodes were always zirconium oxide. The zirconium oxide component in the films is evidently, as revealed by the X-ray diffraction studies, crystallized in the form of a tetragonal/cubic polymorph, which, in turn, is indicative of the defective nature of the material. The majority of defects may arise from the oxygen deficiency, which may be considered as one of the most influential factor stabilizing the tetragonal/cubic polymorph [42] in addition to the impurities deforming the lattice. Consequently, the electrical charge becomes carried to and trapped at the interface layer under certain polarity, and an opposite polarity with in-

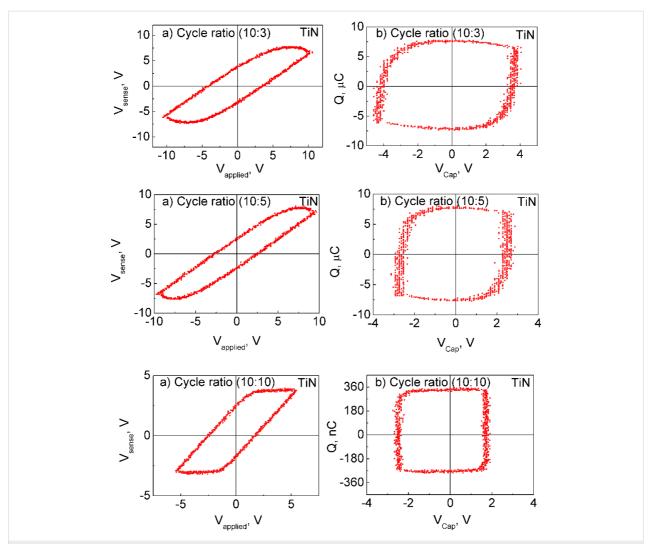


Figure 6: Each of the three panels show a sensed voltage—applied voltage curve on the left (a) and a polarization charge—applied voltage curve on the right (b) for Pt/ZrO₂/Fe₂O₃/TiN/Ti/Si(100)/Al samples. ZrO₂/Fe₂O₃ cycle ratios are shown on each image.

creasing, oppositely directed field is required to release the charge from the traps for the subsequent drift towards the counterelectrode. The current density to applied electric field curves are shown in Figure 7 to support the given explanation. The charge polarization values of the samples correlate with the respective leakage current through samples. The leakage currents in the film grown with ZrO₂/Fe₂O₃ cycle ratio 10:5 (not shown) were similar to those in the film grown with cycle ratio 10:3, which were considerably higher than the leakage currents through other samples. Also, the charge polarization values were the highest for the samples grown with ZrO₂/Fe₂O₃ cycle ratios 10:3 and 10:5 (Figure 6). Other mixtures possessed both lower leakage currents and considerably lower charge polarization values. Probably, in the mixture of iron and zirconium oxides, the amount of defects (in particular oxygen vacancies) is increased due to the substitutive exchange between metal ions of different valence, resulting in an increase also in

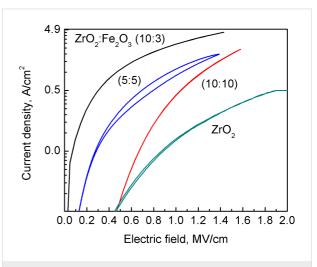


Figure 7: Current density–applied electric field curves for samples with ZrO_2/Fe_2O_3 cycle ratios shown in the image.

the leakage currents. Pure ZrO₂ exhibited the lowest leakage current (Figure 7) and did not show any significant charge polarization at all.

However, most samples in the present study exhibited ferromagnetic-like behavior at room temperature as shown in Figure 8. Saturation magnetization, reaching M/S = 5×10^{-6} A as the maximum value obtained per unit area, could be ob-

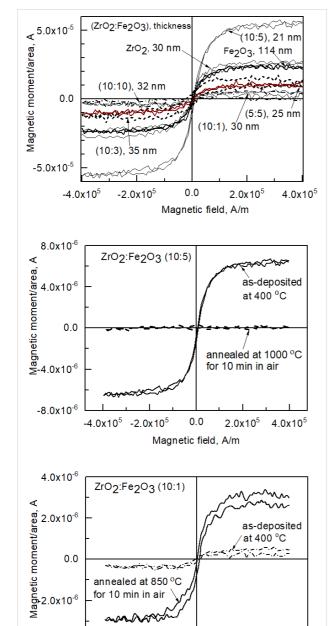


Figure 8: Selected room temperature magnetization–field curves for ${\sf ZrO_2/Fe_2O_3}$ films. Cycle ratios and thicknesses are indicated in the labels.

0

Magnetic field, A/m

2x10⁵

5x10⁵

-2x10⁵

-4.0x10⁻⁶

-5x10⁵

served in most samples, but no definite hysteresis was determined in any case. The ZrO_2/Fe_2O_3 sample with cycle ratio 10:5 (with cation ratio Zr/Fe=2.0) showed considerably higher saturation magnetization with the value very similar to that measured by Sangalli et al. [19]. It is also worth noting that pure ZrO_2 exhibited higher saturation magnetization than other mixed samples. In the case of a cycle ratio of 10:1 and Zr/Fe=50, annealing this sample at 850 °C resulted in the increase of the observed saturation magnetization. Annealing the sample with a cycle ratio 10:5 and Zr/Fe=0.34 at 1000 °C had an opposite result and removed the saturation magnetization of the film (Figure 8).

For comparison with the literature, Myagkov et al. [5] have synthesized Fe-Fe₃O₄-ZrO₂ nanocomposite films by thermally depositing and oxidizing Fe films, followed by the sequential deposition of Zr layers and annealing of the double metal oxide-metal stacks. In this study, saturative magnetization appeared in the samples heat treated above 250 °C and was recorded up to 500 °C, above which nonferromagnetic phases supposedly started to form. In other studies by de Souza et al. [33], Kuryliszyn-Kudelska et al. [34] and Okabayashi et al. [10], only paramagnetic behavior in Fe-doped ZrO₂ was observed. One can note that, for example, in the work by Okabayashi et al. [10], annealing was presumably required in order to initiate crystallization in sol-gel-synthesized (and therefore initially amorphous) films. Magnetization-field hysteresis loops were recorded in the ZrO2 films co-doped with both iron and cobalt, whereas in the films doped only with 1% iron, only paramagnetic behavior was observed [10]. In our films in the present study, the material layers were evidently crystallized already in the as-deposited state, which was sufficient for the appearance of magnetization loops, However, the annealing procedure may have resulted in the formation of nonferromagnetic phases, possibly due to the increased contribution from monoclinic ZrO2 due to the partial compensation of the oxygen deficiency in the films upon heat treatment in air. It is, however, worth noting that the recrystallization or transformation into monoclinic phase cannot be regarded as a fast process, because there appeared only very weak traces of the monoclinic polymorph, expressed by an almost insignificantly low additional reflection at 31.3°, otherwise being indicative of 111 reflection of monoclinic ZrO₂ (PDF Card 37-1484). It is also worth noting that magnetization even in ZrO2 that is not doped with ferromagnetic metals has been investigated and theoretically predicted [43]. In another study on electron beam evaporated undoped ZrO2 films, ferromagnetic-like magnetization hysteresis was recorded, whereby the saturation magnetization was positively correlated with the amount of metastable tetragonal/cubic zirconia in relation to the stable monoclinic phase in the films [6].

Conclusion

Zirconium oxide mixed with iron oxide thin films with various cycle ratios of constituent oxides were grown by ALD to thicknesses ranging from 15 to 40 nm from ZrCl₄, Fe(C₅H₅)₂ and O₃. Most of the films exhibited charge polarization and saturation magnetization. The highest numerical values of these physical quantities were achieved in samples with a relative low iron content (Zr/Fe > 2.0). The saturation magnetization per unit area in a film with ZrO₂/Fe₂O₃ cycle ratio of 10:5 was about 5×10^{-5} A and the charge polarization at 0 V applied voltage was about 8 μ C. These samples were crystallized in the tetragonal/cubic phase of ZrO₂. As-deposited samples, in which the relative cation content for zirconium decreased below 2.0, did not show any peaks assignable to zirconia.

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References

- Leskelä, M.; Niinistö, J.; Ritala, M. Atomic Layer Deposition. In Comprehensive Materials Processing; Cameron, D., Ed.; Elsevier: Amsterdam, Netherlands, 2014; Vol. 4, pp 101–123. doi:10.1016/B978-0-08-096532-1.00401-5
- Hill, N. A. J. Phys. Chem. B 2000, 104, 6694–6709. doi:10.1021/jp000114x
- Dimri, M. C.; Khanduri, H.; Kooskora, H.; Kodu, M.; Jaaniso, R.; Heinmaa, I.; Mere, A.; Krustok, J.; Stern, R. J. Phys. D: Appl. Phys. 2012, 45, 475003. doi:10.1088/0022-3727/45/47/475003
- Hong, N. H.; Kanoun, M. B.; Goumri-Said, S.; Song, J.-H.; Chikoidze, E.; Dumont, Y.; Ruyter, A.; Kurisu, M. J. Phys.: Condens. Matter 2013, 25, 436003. doi:10.1088/0953-8984/25/43/436003
- Myagkov, V. G.; Bykova, L. E.; Bayukov, O. A.; Zhigalov, V. S.; Tambasov, I. A.; Zharkov, S. M.; Matsynin, A. A.; Bondarenko, G. N. J. Alloys Compd. 2015, 636, 223–228. doi:10.1016/j.jallcom.2015.02.012
- Ning, S.; Zhan, P.; Xie, Q.; Li, Z.; Zhang, Z. J. Phys. D: Appl. Phys. 2013, 46, 445004. doi:10.1088/0022-3727/46/44/445004
- Ning, S.; Zhang, Z. RSC Adv. 2015, 5, 3636–3641. doi:10.1039/C4RA11924J
- Sahoo, T. R.; Manoharan, S. S.; Lim, S. H.; Salamanca-Riba, L. G. Synth. React. Inorg., Met.-Org., Nano-Met. Chem. 2008, 38, 280–283.
- Rahman, M. A.; Rout, S.; Thomas, J. P.; McGillivray, D.; Leung, K. T. J. Am. Chem. Soc. 2016, 138, 11896–11906. doi:10.1021/jacs.6b06949
- Okabayashi, J.; Kono, S.; Yamada, Y.; Nomura, K. AIP Adv. 2011, 1, 042138. doi:10.1063/1.3662044
- 11. Yu, J.; Duan, L. B.; Wang, Y. C.; Rao, G. H. *Physica B* **2008**, *403*, 4264–4268. doi:10.1016/j.physb.2008.09.015

- 12. Bechta, S. V.; Krushinov, E. V.; Almjashev, V. I.; Vitol, S. A.; Mezentseva, L. P.; Petrov, Yu. B.; Lopukh, D. B.; Khabensky, V. B.; Barrachin, M.; Hellmann, S.; Froment, K.; Fischer, M.; Tromm, W.; Bottomley, D.; Defoort, F.; Gusarov, V. V. J. Nucl. Mater. 2006, 348, 114–121. doi:10.1016/j.jnucmat.2005.09.009
- Štefanić, G.; Gržeta, B.; Nomura, K.; Trojko, R.; Musić, S. J. Alloys Compd. 2001, 327, 151–160. doi:10.1016/S0925-8388(01)01401-3
- Kukli, K.; Kemell, M.; Köykkä, J.; Mizohata, K.; Vehkamäki, M.; Ritala, M.; Leskelä, M. *Thin Solid Films* **2015**, *589*, 597–604. doi:10.1016/j.tsf.2015.06.033
- 15. Van Der Voort, P.; van Welzenis, R.; de Ridder, M.; Brongersma, H. H.; Baltes, M.; Mathieu, M.; van de Ven, P. C.; Vansant, E. F. Langmuir 2002, 18, 4420–4425. doi:10.1021/la025679u
- 16. de Ridder, M.; van de Ven, P. C.; van Welzenis, R. G.; Brongersma, H. H.; Helfensteyn, S.; Creemers, C.; Van Der Voort, P.; Baltes, M.; Mathieu, M.; Vansant, E. F. J. Phys. Chem. B 2002, 106, 13146–13153. doi:10.1021/jp0211640
- Lamperti, A.; Cianci, E.; Ciprian, R.; Sangalli, D.; Debernardi, A.
 Thin Solid Films 2013, 533, 83–87. doi:10.1016/j.tsf.2012.11.127
- Pitzschel, K.; Bachmann, J.; Montero-Moreno, J. M.; Escrig, J.; Görlitz, D.; Nielsch, K. Nanotechnology 2012, 23, 495718. doi:10.1088/0957-4484/23/49/495718
- Sangalli, D.; Cianci, E.; Lamperti, A.; Ciprian, R.; Albertini, F.;
 Casoli, F.; Lupo, P.; Nasi, L.; Campanini, M.; Debernardi, A.
 Eur. Phys. J. B 2013, 86, 211–216. doi:10.1140/epjb/e2013-30669-3
- Lin, B.-T.; Lu, Y.-W.; Shieh, J.; Chen, M.-J. J. Eur. Ceram. Soc. 2017, 37, 1135–1139. doi:10.1016/j.jeurceramsoc.2016.10.028
- Pešic, M.; Hoffmann, M.; Richter, C.; Mikolajick, T.; Schroeder, U.
 Adv. Funct. Mater. 2016, 26, 7486–7494. doi:10.1002/adfm.201603182
- Arroval, T.; Aarik, L.; Rammula, R.; Kruusla, V.; Aarik, J.
 Thin Solid Films 2016, 600, 119–125. doi:10.1016/j.tsf.2016.01.024
- Tamm, A.; Dimri, M. C.; Kozlova, J.; Aidla, A.; Tätte, T.; Arroval, T.;
 Mäeorg, U.; Mändar, H.; Stern, R.; Kukli, K. J. Cryst. Growth 2012, 343, 21–27. doi:10.1016/j.jcrysgro.2011.09.062
- 24. Martinson, A. B. F.; DeVries, M. J.; Libera, J. A.; Christensen, S. T.; Hupp, J. T.; Pellin, M. J.; Elam, J. W. J. Phys. Chem. C 2011, 115, 4333–4339. doi:10.1021/jp110203x
- Jacqueline, S.; Domengès, B.; Voiron, F.; Murray, H. Semicond. Sci. Technol. 2013, 28, 045018. doi:10.1088/0268-1242/28/4/045018
- Ritala, M.; Niinistö, J. ECS Trans. 2009, 25, 641–652. doi:10.1149/1.3207651
- 27. Bae, C.; Shin, H.; Nielsch, K. MRS Bull. 2011, 36, 887–897. doi:10.1557/mrs.2011.264
- Hughes, K. J.; Engstrom, J. R. J. Vac. Sci. Technol., A 2012, 30, 01A102. doi:10.1116/1.3625564
- Kukli, K.; Ritala, M.; Aarik, J.; Uustare, T.; Leskelä, M. J. Appl. Phys.
 2002, 92, 1833–1840. doi:10.1063/1.1493657
- Ritala, M.; Leskelä, M. Appl. Surf. Sci. 1994, 75, 333–340.
 doi:10.1016/0169-4332(94)90180-5
- Siimon, H.; Aarik, J. J. Phys. D: Appl. Phys. 1997, 30, 1725–1728. doi:10.1088/0022-3727/30/12/006
- Lamperti, A.; Lamagna, L.; Congedo, G.; Spiga, S.
 J. Electrochem. Soc. 2011, 158, G221–G226. doi:10.1149/1.3625254
- de Souza, A. O.; Ivashita, F. F.; Biondo, V.; Paesano, A., Jr.; Mosca, D. H. J. Alloys Compd. 2016, 680, 701–710. doi:10.1016/j.jallcom.2016.04.170

- 34. Kuryliszyn-Kudelska, I.; Arciszewska, M.; Małolepszy, A.; Mazurkiewicz, M.; Stobinski, L.; Grabias, A.; Kopcewicz, M.; Paszkowicz, W.; Minikaev, R.; Domukhovski, V.; Nedelko, N.; Dobrowolski, W. J. Alloys Compd. 2015, 632, 609–616. doi:10.1016/j.jallcom.2015.01.257
- Chakma, S.; Moholkar, V. S. RSC Adv. 2015, 5, 53529–53542. doi:10.1039/c5ra06148b
- Østreng, E.; Sønsteby, H. H.; Sajavaara, T.; Nilsen, O.; Fjellvåg, H.
 J. Mater. Chem. C 2013, 1, 4283–4290. doi:10.1039/C3TC30271G
- Sønsteby, H. H.; Fjellvåg, H.; Nilsen, O. Adv. Mater. Interfaces 2017, 4, 1600903. doi:10.1002/admi.201600903
- Vehkamäki, M.; Hatanpää, T.; Kemell, M.; Ritala, M.; Leskelä, M.
 Chem. Mater. 2006, 18, 3883–3888. doi:10.1021/cm060966v
- 39. Zhang, F.; Perng, Y.-C.; Choi, J. H.; Wu, T.; Chung, T.-K.; Carman, G. P.; Locke, C.; Thomas, S.; Saddow, S. E.; Chang, J. P. J. Appl. Phys. 2011, 109, 124109. doi:10.1063/1.3596574
- Martin, D.; Müller, J.; Schenk, T.; Arruda, T. M.; Kumar, A.; Strelcov, E.; Yurchuk, E.; Müller, S.; Pohl, D.; Schröder, U.; Kalinin, S. V.; Mikolajick, T. Adv. Mater. 2014, 26, 8198–8202. doi:10.1002/adma.201403115
- 41. Xiang, H.; Shi, F.; Rzchowski, M. S.; Voyles, P. M.; Chang, Y. A. Appl. Phys. Lett. 2010, 97, 092508. doi:10.1063/1.3484278
- 42. Guo, X. J. Phys. Chem. Solids **1999**, 60, 539–546. doi:10.1016/S0022-3697(98)00301-1
- Máca, F.; Kudrnovský, J.; Drchal, V.; Bouzerar, G. Appl. Phys. Lett. 2008, 92, 212503. doi:10.1063/1.2936858

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