

# Floating and Tether-Coupled Adhesion of Bacteria to Hydrophobic and Hydrophilic Surfaces

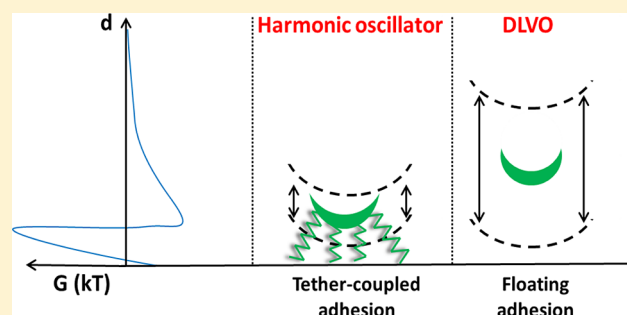
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## S Supporting Information

**ABSTRACT:** Models for bacterial adhesion to substratum surfaces all include uncertainty with respect to the (ir)-reversibility of adhesion. In a model, based on vibrations exhibited by adhering bacteria parallel to a surface, adhesion was described as a result of reversible binding of multiple bacterial tethers that detach from and successively reattach to a surface, eventually making bacterial adhesion irreversible. Here, we use total internal reflection microscopy to determine whether adhering bacteria also exhibit variations over time in their perpendicular distance above surfaces. Streptococci with fibrillar surface tethers showed perpendicular vibrations with amplitudes of around 5 nm, regardless of surface hydrophobicity. Adhering, nonfibrillated streptococci vibrated with amplitudes around 20 nm above a hydrophobic surface. Amplitudes did not depend on ionic strength for either strain. Calculations of bacterial energies from their distances above the surfaces using the Boltzmann equation showed that bacteria with fibrillar tethers vibrated as a harmonic oscillator. The energy of bacteria without fibrillar tethers varied with distance in a comparable fashion as the DLVO (Derjaguin, Landau, Verwey, and Overbeek)-interaction energy. Distance variations above the surface over time of bacteria with fibrillar tethers are suggested to be governed by the harmonic oscillations, allowed by elasticity of the tethers, piercing through the potential energy barrier. Bacteria without fibrillar tethers “float” above a surface in the secondary energy minimum, with their perpendicular displacement restricted by their thermal energy and the width of the secondary minimum. The distinction between “tether-coupled” and “floating” adhesion is new, and may have implications for bacterial detachment strategies.



## INTRODUCTION

Bacterial adhesion occurs to many different surfaces in a wide variety of environments, and is either desirable such as in many bioreactor systems, soil remediation or to intestinal surfaces in the human body or is, among others, the cause of severe infections, food spoilage or microbially induced corrosion.<sup>1</sup> Accordingly, there is an ongoing quest to model bacterial adhesion. This quest is on the one hand led by biochemists trying to discover more and more specific ligand–receptor systems facilitating bacterial adhesion to surfaces, while on the other hand physico-chemists attempt to design generally valid models that explain and predict adhesion of bacteria to surfaces by treating living organisms as (bio)colloidal particles.<sup>2,3</sup>

The two most common physicochemical approaches used to model bacterial adhesion are surface thermodynamic<sup>4,5</sup> and (extended) DLVO (Derjaguin, Landau, Verwey and Overbeek)-types of analyses.<sup>6,7</sup> Application of surface thermodynamics involves the measurement of contact angles with different liquids, followed by calculation and comparison of the free energies of the substratum and bacterial cell surfaces to

yield an interfacial free energy of adhesion. Negative values for the interfacial free energy of adhesion are assumed to be predictive for bacterial adhesion to occur, but this has never become a generally valid observation, possibly because the surface thermodynamic requirement of reversibility is seldom or never met and the interface between a bacterium and a substratum surface is a dynamic one, changing over time.<sup>8,9</sup> DLVO-types of analyses calculate the interaction energy between a (bio)colloidal particle and a substratum surface as a function of distance between the particle and the substratum surface. Under most relevant conditions, electrostatic interactions in bacterial adhesion are repulsive, that together with attractive Lifshitz–van der Waals forces yield a secondary interaction minimum at a distance of around 20–50 nm<sup>10</sup> from the substratum surface, a potential energy barrier that impedes close approach and a deep primary minimum close to the

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surface that can only be reached once a particle has overcome the potential energy barrier. Also DLVO-type approaches have never acquired a general validity across different bacterial strains and species. An important reason for this is that bacterial cell surfaces can possess a wide variety of surface appendages of different length, width, composition and surface density that have been suggested to be able to pierce the potential energy barrier and tether a bacterium to a surface. Moreover, possession of tethers with lengths that may range up to several micrometers<sup>9</sup> makes it impossible to adequately define the interaction distance in DLVO-type analyses of bacterial adhesion as it creates a multiscale roughness on the bacterial cell surface.<sup>11</sup> Several of the troublesome issues involved in the application of surface thermodynamic and DLVO-types of analyses of bacterial adhesion have been addressed in a new model of bacterial adhesion describing irreversible adhesion of bacteria as a result of the reversible binding of multiple tethers that detach from and successively reattach to a surface, resulting in irreversible adhesion of an organism as a whole.<sup>9</sup> The model puts bacterial adhesion a par with the irreversible adsorption of high-molecular-weight proteins to surfaces, mediated by multiple, reversibly binding molecular segments and was confirmed by in silico modeling of the key-observations underlying the model: (1) force–distance curves in single probe bacterial atomic force microscopy showing detachment events indicative of multiple binding tethers, (2) nanoscopic displacements of bacteria with relatively long autocorrelation times up to several seconds, in the absence of macroscopic displacement, (3) nanoscopic vibrational amplitudes of adhering bacteria parallel to a surface decreasing with increasing adhesion-forces, and (4) increases in mean-squared-displacements over prolonged time periods according to  $t^\alpha$  with  $0 < \alpha \ll 1$ , indicative of confined displacement.<sup>9</sup>

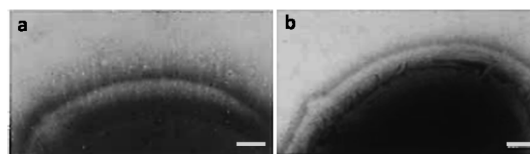
The role of adhesion forces acting perpendicular to a substratum surface may seem puzzling in a model that is based on parallel displacements of adhering bacteria over the surface. However, in silico modeling suggested that tether adhesion forces merely dictate the frequency with which individual tethers detach. This leaves the question open as to whether the distance of an adhering bacterium above a substratum surface also varies over time, similar as its position on a substratum surface, and whether tether-binding plays a role here too.

In order to answer this question, this paper aims to determine whether adhering bacteria exhibit variations over time in their distance perpendicular to substratum surfaces and, if so, whether these variations differ for two streptococcal strains with and without 91 nm long fibrillar, tethers. Total internal reflection microscopy (TIRM)<sup>12–14</sup> will be used to determine the variations in distance above different substratum surfaces over time, employing hydrophobic and hydrophilic glass surfaces as substrata. Distance variations will be related with the shape and width of the secondary interaction minimum and its distance from the substratum surface in a DLVO-type of analysis.

## ■ EXPERIMENTAL SECTION

### Bacterial Strains, Culture Conditions, and Harvesting.

*Streptococcus salivarius* HB7, possessing 91 nm long fibrillar tethers and its isogenic mutant HBC12, considered bald without demonstrable surface tethers,<sup>15</sup> were employed in this study (see Figure 1). Both *S. salivarius* strains are negatively charged, yet differ in surface hydrophobicity with *S. salivarius* HBC12, being slightly more hydrophilic than *S. salivarius* HB7.<sup>16</sup> Both *S. salivarius* strains were



**Figure 1.** Transmission electron micrographs of negatively stained (1% methylamine tungstate) sections of *S. salivarius* HB7 (a) and *S. salivarius* HBC12 (b). The bar denotes 100 nm. Adapted from Van der Mei et al. with permission from the publisher, Springer Nature.<sup>15</sup>

precultured in 10 mL of Todd Hewitt Broth (THB, OXOID, Basingstoke, UK), under static conditions. Precultures were grown for 24 h at 37 °C. After 24 h, precultures were inoculated into 200 mL of THB, and maintained under identical conditions for another 16 h. Streptococci were harvested by centrifugation at 5000g for 5 min at 10 °C, subsequently washed three times in 100 mL of adhesion buffer having an ionic strength of 57 mM (50 mM potassium chloride, 2 mM potassium phosphate, and 1 mM calcium chloride, pH 6.8) or ionic strength of 0.57 mM (10 times diluted adhesion buffer). Following this, bacteria were sonicated on ice three times for 10 s at 30 W (Vibra Cell model 375; Sonics and Materials Inc., Danbury, CT) to break bacterial chains and obtain single bacteria. Finally, bacteria were resuspended in adhesion buffer to a final concentration of  $3 \times 10^8$  bacteria/mL, as determined using a Bürker–Türk counting chamber.

**Preparation of Substratum Surfaces.** Glass microscope slides (15 mm × 15 mm; Thermo Scientific, 38116 Braunschweig, Germany) were used as substrata. Prior to each experiment, glass surfaces were cleaned by 10 min sonication at 100 W (Bransonic 2510E, Danbury, USA) in 2% Hellmanex (Hellma GmbH & CO., Müllheim, Germany), 99% ethanol and finally in ultrapure water (specific resistance >18 MΩ cm). Next, glass surfaces were treated with UV/ozone, yielding a hydrophilic surface while for the preparation of hydrophobic surfaces, glass slides were thoroughly dried after water-washing in an oven at 80 °C, followed by silanization in 0.05% (w/v) dimethyldichlorosilane (DDS Sigma-Aldrich) in 99% ethanol for approximately 15 min.

**Total Internal Reflection Microscopy.** Adhesion of streptococci onto uncoated and DDS-coated glass surfaces was established from adhesion buffer (0.57 mM or 57 mM) at room temperature. To this end, a streptococcal suspension was introduced into the circularly shaped flow chamber (14 mm diameter and 0.35 mm in height) of the instrument using a peristaltic pump at a flow rate of 300 μL/min during 60 min. Next the chamber was perfused with buffer, after which TIRM light scattering was measured. TIRM was performed on an objective-based total internal reflection fluorescence (TIRF) microscope (Nikon, Eclipse Ti with TIRF module, Tokyo, Japan), equipped with a high numerical aperture objective (Olympus, PLANO-APO 100×, 1.45, Tokyo, Japan) illuminated by a 488 nm laser (Melles Griot, Dynamic Laser, Salt Lake City, UT) laterally focused on the back focal plane. To avoid overexposure by the reflected laser beam, a spatial filter was employed to block the reflected beam in the back focal plane without image interference. The filter cube contained only a 488 dichroic mirror. Scattering light was captured on an electron multiplying, charge-coupled device camera (Andor, ixon DU-885BV Andor, Dublin, Ireland). Image size was cropped to 512 × 512 pixel resolutions to achieve a frame rate of 33 frames/s over 2000 frames. Prior to each experiment, the TIRM angle was verified with an out-coupling prism to an external reference. Light scattered by adhering bacteria was observed as two diffraction limited spots, separated by the bacterial diameter, recognized as a comet-orbit shape.<sup>17</sup>

For each adhering single bacterium, the relative  $z$ -displacement,  $z_t$  at time  $t$  with respect to the closest distance encountered, was calculated according to

$$z_t = -\ln\left(\frac{I_t}{I_{\max}}\right)d_p \quad (1)$$

with  $I_t$  the intensity of the scattered light in the evanescent field at time  $t$ ,  $I_{\max}$  the maximum intensity belonging to the distance of closest

bacterial approach and  $d_p$ , the penetration depth of the evanescent wave (185 nm). Absolute distances could not be obtained, because the evanescent wave intensity was not constant over the entire field of view. As a result the maximum intensity  $I_{\max}$  was different for each bacterium over the field of view. The vibrational amplitude  $\Delta z$  was calculated as the standard deviation of all  $z_i$ -displacements over the experimental time. All TIRM experiments were carried out in triplicates on uncoated and DDS-coated glass surfaces using different bacterial cultures for each experiment.

**Contact Angle Measurements.** Contact angles were measured on the uncoated and DDS-coated glass surfaces with water, formamide, and methyleneiodide. Three 0.5  $\mu\text{L}$  droplets of each liquid were randomly positioned on one microscope slide, employing three slides for one series of measurements with each of the above three liquids. Images of the droplets were recorded by a camera about 5 s after placing a droplet on a coverslips surface and the droplet contour digitized after gray-value thresholding, after which contact angles were calculated from the digitized contours using homemade software. Contact angles on bacteria were measured by preparing macroscopic bacterial lawns on membrane filters. Bacterial lawns were made by suspending bacteria to a concentration of  $3 \times 10^8 \text{ mL}^{-1}$  in water, followed by deposition on a cellulose acetate membrane filter (pore diameter 0.45  $\mu\text{m}$ ) placed on a fritted glass support by filtration of the suspension. At least three separate filters, from three different cultures were used for each bacterial strain tested. Strains deposited similarly in a smooth and even layer. The filters with their deposited bacterial lawn, were placed on a metal sample disc and allowed to air-dry for 30–90 min,<sup>18</sup> in order to obtain relatively stable, so-called “plateau” contact angles, indicative of bacteria in a hydrated state but without free water in between. Contact angles were subsequently measured as described above. The contact angles presented, represent the averages from three experiments with separate prepared surfaces as well as bacterial cultures.

Next, contact angles on each surface were converted to a Lifshitz–van der Waals ( $\gamma^{\text{LW}}$ ) and acid–base ( $\gamma^{\text{AB}}$ ) surface free energy component, while the acid–base component was split up into an electron-donating ( $\gamma^-$ ) and an electron-accepting ( $\gamma^+$ ) parameter according to

$$\begin{bmatrix} \sqrt{\gamma_{\text{water}}^{\text{LW}}} & \sqrt{\gamma_{\text{water}}^+} & \sqrt{\gamma_{\text{water}}^-} \\ \sqrt{\gamma_{\text{formamide}}^{\text{LW}}} & \sqrt{\gamma_{\text{formamide}}^+} & \sqrt{\gamma_{\text{formamide}}^-} \\ \sqrt{\gamma_{\text{methyleneiodide}}^{\text{LW}}} & \sqrt{\gamma_{\text{methyleneiodide}}^+} & \sqrt{\gamma_{\text{methyleneiodide}}^-} \end{bmatrix} \begin{bmatrix} \sqrt{\gamma^{\text{LW}}} \\ \sqrt{\gamma^-} \\ \sqrt{\gamma^+} \end{bmatrix} = \begin{bmatrix} (1 + \cos \theta_{\text{water}})\gamma_{\text{water}}/2 \\ (1 + \cos \theta_{\text{formamide}})\gamma_{\text{formamide}}/2 \\ (1 + \cos \theta_{\text{methyleneiodide}})\gamma_{\text{methyleneiodide}}/2 \end{bmatrix} \quad (2)$$

in which  $\gamma$  denotes the surface free energy and/or its components and parameters of the various liquids used or the solid surface to be analyzed, while  $\theta$  represents the contact angle of the different liquids.<sup>19</sup> Surface free energy components and parameters of the liquids used can be found in Supporting Information Table S1.

**Bacterial and Substratum Zeta Potentials.** To determine the zeta potentials of the two bacterial strains, particulate microelectrophoresis (Zetasizer nano-ZS, Malvern Instruments, Worcester-shire, UK) was carried out at low and high ionic strength (0.57 and 57 mM, respectively) at pH 6.8.<sup>20</sup> Streaming potential measurements were employed to determine the zeta potentials of uncoated and DDS-coated glass surfaces. To this end, glass slides with and without DDS-coating were mounted in a homemade parallel plate flow chamber, separated by a 100  $\mu\text{m}$  Teflon spacer.<sup>21</sup> A platinum electrode was located on each side of the chamber. The streaming potentials were measured at pressures ranging from 50 to 400 mbar, and each pressure was applied for 10 s in both directions. Following this, the zeta potentials were calculated by linear regression, i.e., linear least-squares fitting of the streaming potentials.

**DLVO Theory.** The DLVO theory, describes (bio)colloidal particle adhesion as a result of attractive Lifshitz–van der Waals and attractive or repulsive electrostatic forces. Accordingly, the interaction energy between a colloidal particle and a substratum can be expressed as a function of their separation distance ( $d$ ) as

$$G^{\text{TOT}}(d) = G^{\text{LW}}(d) + G^{\text{EL}}(d) \quad (3)$$

in which  $G^{\text{TOT}}$ ,  $G^{\text{LW}}$ , and  $G^{\text{EL}}$  represent the total, Lifshitz–van der Waals, and electrostatic interaction energies, respectively. The Lifshitz–van der Waals interaction between a spherical colloidal particle and a planar surface is given by

$$G^{\text{LW}}(d) = -\frac{A}{6} \left[ \frac{a}{d} + \frac{a}{d+2a} + \ln \left( \frac{d}{d+2a} \right) \right] \quad (4)$$

in which  $A$  is the Hamaker constant and  $a$  is the radius of the colloidal particle.<sup>19</sup> The Hamaker constant was derived from the Lifshitz–van der Waals component of interfacial free energy of adhesion according to

$$A = -12\pi d_0^2 \Delta G_{\text{Adh}}^{\text{LW}} \quad (5)$$

in which  $d_0$  is the distance of closest approach (0.157 nm).<sup>22–24</sup> Analogously, the electrostatic interaction can be calculated using measured zeta potentials according to

$$G^{\text{EL}}(d) = \pi \epsilon a (\zeta_b^2 + \zeta_s^2) \left\{ \frac{2\zeta_b + \zeta_s}{\zeta_b^2 + \zeta_s^2} \ln \frac{1 + \exp(-\kappa d)}{1 - \exp(-\kappa d)} + \ln[1 - \exp(-\kappa d)] \right\} \quad (6)$$

in which  $\epsilon$  refers to the dielectric permittivity of the medium, and  $\zeta_b$  and  $\zeta_s$  refer to the zeta potentials of the bacterium and substratum surface, respectively.  $1/\kappa$  is the Debye–Hückel length, given by

$$\kappa = \left[ \frac{e^2}{\epsilon k_B T} \sum_i z_i n_i \right]^{1/2} \quad (7)$$

in which  $e$  corresponds to the electron charge,  $k_B$  is the Boltzmann constant,  $T$  is the absolute temperature,  $z_i$  is the valence of ions present, and  $n_i$  is the number of ions per unit volume.

**Statistical Analysis.** All experiments were carried out in triplicates with separately prepared bacterial cultures as well as different surfaces, and all data are presented as means  $\pm$  standard deviations. Results were compared pairwise for the two different strains of bacteria for the effects of ionic strength by using a Student's  $t$  test.  $p < 0.05$  was considered to be statistically significant.

## RESULTS

Table 1 summarizes the contact angles measured with different liquids on both substratum surfaces and bacterial strains. On the basis of the water contact angles, uncoated and DDS-coated glass surfaces can be classified as hydrophilic and hydrophobic, respectively. Although *S. salivarius* HB7 was more hydrophobic

**Table 1. Contact Angles with Different Liquids for Uncoated and DDS-Coated Glass Surfaces as Well as for *S. salivarius* HB7 and *S. salivarius* HBC12<sup>a</sup>**

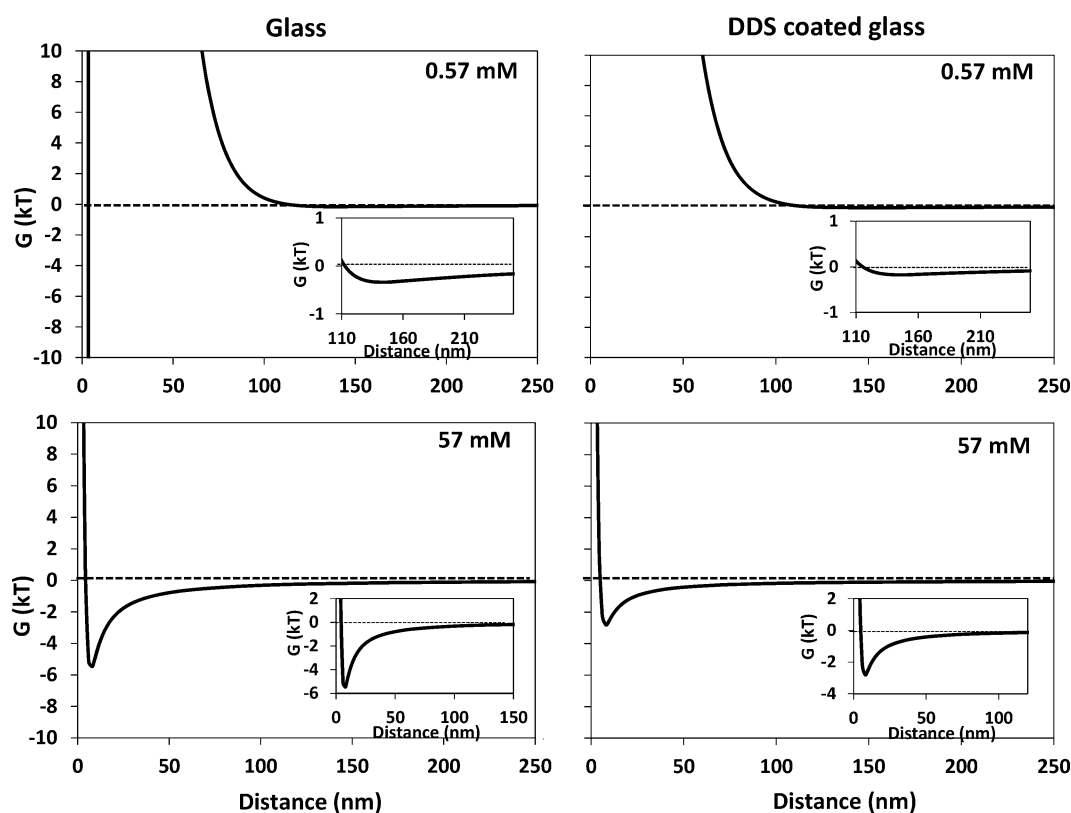
surface/bacterial strain	$\theta_{\text{water}}$ (deg)	$\theta_{\text{formamide}}$ (deg)	$\theta_{\text{methyleneiodide}}$ (deg)
glass	23.3 $\pm$ 1.5	19.3 $\pm$ 2.3	52.3 $\pm$ 6.4
DDS-coated glass	97.0 $\pm$ 1.7	74.3 $\pm$ 6.7	63.3 $\pm$ 3.8
<i>S. salivarius</i> HB7	34.3 $\pm$ 4.6	12.3 $\pm$ 2.5	24.0 $\pm$ 6.0
<i>S. salivarius</i> HBC12	21.6 $\pm$ 3.6	24.7 $\pm$ 4.9	38.0 $\pm$ 16.2

<sup>a</sup>Data represent averages with standard deviations over three droplets on three different glass surfaces of each type and bacterial lawns prepared from three different bacterial cultures.

**Table 2. Surface Free Energy Components and Parameters Together with Zeta Potentials for Uncoated and DDS-Coated Glass Surfaces as Well as for *S. salivarius* HBC12 and *S. salivarius* HB7, Respectively<sup>a</sup>**

	surface free energy components and parameters ( $\text{mJ m}^{-2}$ )			
	glass	DDS-coated glass	<i>S. salivarius</i> HBC12	<i>S. salivarius</i> HB7
$\gamma$	$54.9 \pm 1.1$	$27.7 \pm 1.0$	$51.5 \pm 2.6$	$57.0 \pm 1.2$
$\gamma^{\text{LW}}$	$34.0 \pm 3.0$	$26.8 \pm 1.3$	$40.1 \pm 7.9$	$46.3 \pm 2.1$
$\gamma^{\text{AB}}$	$21.1 \pm 3.8$	$0.4 \pm 0.4$	$11.5 \pm 8.8$	$10.6 \pm 3.2$
$\gamma^-$	$48.0 \pm 0.8$	$2.1 \pm 1.8$	$51.8 \pm 4.9$	$35.1 \pm 10.2$
$\gamma^+$	$2.3 \pm 0.9$	$0.4 \pm 0.7$	$0.9 \pm 0.9$	$1.0 \pm 0.8$
$\gamma^-/\gamma^+$	$20.6 \pm 0.9$	$4.1 \pm 3.2$	$59.7 \pm 5.5$	$35.1 \pm 13.5$
	zeta potentials (mV)			
	glass	DDS-coated glass	<i>S. salivarius</i> HBC12	<i>S. salivarius</i> HB7
0.57 mM	$-84.9 \pm 0.4$	$-55.2 \pm 1.4$	$-20.0 \pm 0.4$	$-16.9 \pm 2.4$
57 mM	$-39.5 \pm 0.6$	$-27.8 \pm 0.0$	$-8.2 \pm 0.2$	$-9.1 \pm 2.5$

<sup>a</sup>Data regarding surface energetics represent averages with standard deviations over three measurements on three different glass surfaces of each type and bacterial lawns prepared from three different bacterial cultures. Bacterial zeta potentials are averages with standard deviations over three experiments with different bacterial cultures, while substratum zeta potentials are averages with standard deviations over three streaming potential measurements with different uncoated and DDS-coated surfaces.



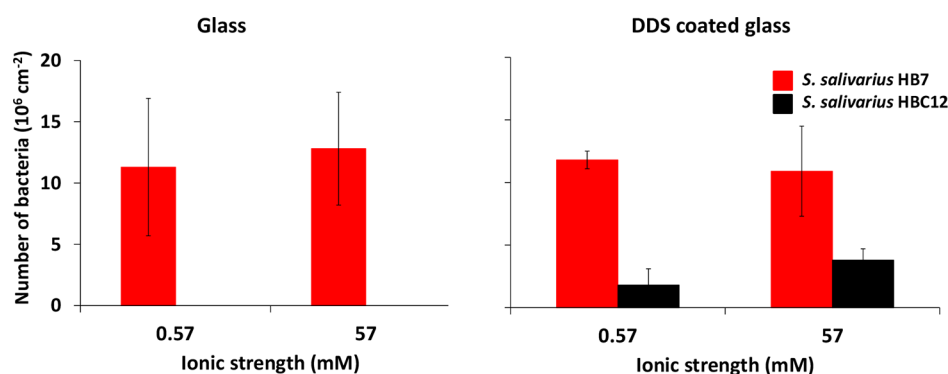
**Figure 2.** Interaction energies between *S. salivarius* HBC12 and uncoated or DDS-coated glass surfaces in different ionic strength suspensions, as a function of their surface-to-surface separation distance. Insets represent part of the interaction energy curve at a different scale to better visualize the secondary minimum.

than *S. salivarius* HBC12, both bacterial strains can be classified as hydrophilic.

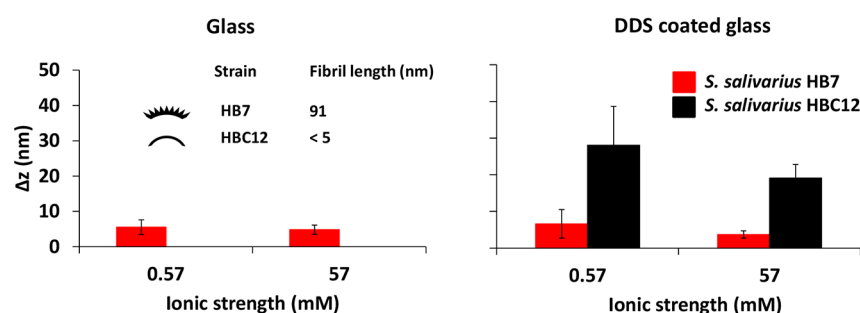
Surface free energy components and parameters, calculated from contact angles with the three different liquids were subsequently compiled in Table 2, most notably showing a small acid–base component for hydrophobic DDS-coated glass, due to both small electron-donating and accepting parameters. In addition, the ratio of electron-donating over electron-accepting parameters varied between the two substratum surfaces, indicative of different structuring of water molecules nearby the surface.<sup>19,25</sup> Both streptococcal strains had similar

acid–base components, but *S. salivarius* HBC12 had a much higher electron-donating surface free energy parameter than *S. salivarius* HB7, resulting in different ratios between their electron-donating and electron-accepting parameters. Zeta potentials of the different surfaces, also compiled in Table 2, were negative for all surfaces at both low and high ionic strength, while being significantly more negative in the low ionic than in the high ionic strength buffer.

The surface free energy components from Table 2 can be used together with their counterparts for water (see Supporting Information Table S1) to calculate the Lifshitz–van der Waals



**Figure 3.** Numbers of adhering *S. salivarius* HB7 and *S. salivarius* HBC12 on glass and DDS-coated glass surfaces at ionic strengths of 0.57 and 57 mM. Note that, on uncoated glass, the number of adhering *S. salivarius* HBC12 was too low for microscopic enumeration.



**Figure 4.** Vibrational amplitudes  $\Delta z$  above glass or DDS-coated glass surfaces for *S. salivarius* HB7 and *S. salivarius* HBC12 in low and high ionic strength buffers, obtained using TIRM. Note that on hydrophilic glass, the number of adhering *S. salivarius* HBC12 was too low for TIRM measurements. Data represent averages over three separate experiments with error bars indicating the standard deviations over three different bacterial cultures.

interfacial free energy of adhesion (eq 4), and subsequently using eq 5 to calculate the Hamaker constant for bacterial interaction with uncoated or DDS-coated glass in an aqueous suspension. Since the concept of interaction distance in the DLVO approach loses its meaning when the bacterial cell surface possesses a multiscale roughness,<sup>11</sup> such as due to fibrillar surface tethers in *S. salivarius* HB7, these calculations were only made for *S. salivarius* HBC12, yielding Hamaker constants of  $3.8 \times 10^{-21}$  and  $1.6 \times 10^{-21}$  J against glass and DDS-coated glass, respectively. DLVO interaction energies versus distance for *S. salivarius* HBC12 for glass and DDS-coated glass were subsequently calculated inserting these Hamaker constants and the zeta potentials from Table 2, into eqs 3, 4, and 6, assuming a bacterial radius of 500 nm,<sup>26</sup> and a Debye–Hückel length  $1/\kappa$  for the two ionic strengths (0.57 and 57 mM) of  $1.3 \times 10^{-8}$  and  $1.3 \times 10^{-9}$  m, respectively.

At low ionic strength (Figure 2), the secondary minimum was extremely shallow with a depth of  $0.5kT$  and  $0.25kT$  for uncoated and DDS-coated glass, respectively and located approximately 140 to 150 nm away from the substratum surface, respectively. Due to the decrease of electrostatic repulsion, the secondary minimum at 57 mM ionic strength was much deeper than that in 0.57 mM suspensions (Figure 2) and amounted around  $5kT$  and  $3kT$  for glass and DDS-coated glass, respectively, while being located approximately 15 nm from the surface. The potential energy barrier at close approach may be considered too high for a bacterium to overcome and adhere in the primary minimum as a whole in all cases.<sup>27</sup> Due to the relatively strong electron-donating and electron-accepting parameters of glass as compared with DDS-coated glass, both strains will also experience large monopolar repulsion at close

approach, that will be far less or absent on hydrophobic, DDS-coated glass than on hydrophilic glass (see also Table 2). For *S. salivarius* HBC12 on glass monopolar repulsion  $\Delta G^{\text{AB}}(d_0)$  amounts  $+30.4 \text{ mJ m}^{-2}$ , turning into attraction ( $-10.6 \text{ mJ m}^{-2}$ ) on DDS-coated glass, but since this is at the distance of closest approach  $d_0$ , it is not reflected in the interaction energies presented in Figure 2 according to the DLVO theory as applied.

As a first step in the TIRM measurements, the chamber was perfused with a bacterial suspension and the number of adhering streptococci enumerated, as summarized in Figure 3.

*S. salivarius* HB7 with its fibrillar tethers adhered in approximately equal numbers to uncoated and DDS-coated glass, regardless of ionic strength ( $11 \times 10^6$  and  $13 \times 10^6 \text{ cm}^{-2}$ ). *S. salivarius* HBC12 demonstrated no microscopically enumerable numbers to the uncoated glass surface, but on DDS-coated glass enumerable numbers were observed, amounting  $2 \times 10^6$  and  $4 \times 10^6 \text{ cm}^{-2}$  for the low and high ionic strength suspension, respectively.

Vibrational amplitudes,  $\Delta z$ , of *S. salivarius* HB7 adhering on a hydrophilic, uncoated glass surface (Figure 4) were relatively small, around 5 nm irrespective of ionic strength, while *S. salivarius* HBC12 adhered in too low numbers for TIRM experiments. *S. salivarius* HB7 also exhibited a relatively small vibrational amplitude  $\Delta z$  of around 5 nm on hydrophobic, DDS-coated glass, albeit here the vibrational amplitude was slightly higher at low ionic strength (not statistically significant;  $p > 0.05$ , Student's  $t$  test) than at high ionic strength. Strikingly, *S. salivarius* HBC12 demonstrated much higher vibrational amplitudes  $\Delta z$  on hydrophobic, DDS-coated glass than as observed for *S. salivarius* HBC7. In addition, these vibrational

amplitudes decreased slightly toward high ionic strength (not statistically significant;  $p > 0.05$ , Student's  $t$  test).

## DISCUSSION

Using TIRM, the variations in distance over time from a substratum surface to which they adhered, were measured for two strains of *S. salivarius* with and without fibrillar surface tethers. The strain with fibrillar tethers showed vibrational amplitudes of around 5 nm, regardless of ionic strength or substratum hydrophobicity. The strain without fibrillar tethers did not adhere in sufficient numbers to derive vibrational amplitudes on hydrophilic, uncoated glass, due to unfavorable thermodynamic conditions (interfacial free energy of adhesion<sup>27,28</sup> calculated from the data in Table 2: +26.5 mJ m<sup>-2</sup> due to strong monopolar repulsion). Oppositely, on hydrophobic, DDS-coated glass (interfacial free energy of adhesion: -12.3 mJ m<sup>-2</sup> in absence of strong monopolar repulsion), the non-fibrillated strain adhered reasonably well and vibrated perpendicularly above the surface with a 5-fold-higher amplitude around 25 nm, regardless of ionic strength. Previously, TIRM has been used to analyze the change in separation distance of these streptococcal strains adhering for only 5 min to a substratum surface upon reducing ionic strength.<sup>29</sup> When the ionic strength was reduced from 57 mM to 5.7 mM, the distance between the bacterial cell of *S. salivarius* HB7 and the substratum increased from 45 to 90 nm, which suggests that fibrils change from a compressed, side-on conformation to a fully stretched state. This conclusion was later confirmed by QCM-D (quartz crystal microbalance with dissipation) experiments<sup>30</sup> suggesting collapse of streptococcal fibrillar tethers within several minutes after contact with a substratum surface. The vibrational amplitude of the fibrillated strain observed here (5 nm) is not only much smaller than the distance at which a bacterium adheres above the surface and unaffected by ionic strength, but also much smaller than the fibrillar length or the distance above the surface measured before, probably because in our measurements 60 min of adhesion were allowed before measurements, causing tether collapse over time under influence of the adhesion forces arising from the substratum surface.

The perpendicular, vibrational amplitudes of adhering *S. salivarius* HBC12 without fibrillar surface tethers can be related with the DLVO interaction free energy curves but due to low numbers of adhering bacteria resulting from combined monopolar and low ionic strength electrostatic repulsions only at high ionic strength. Under high ionic strength conditions, there is a clear secondary minimum (Figure 2). Accounting for a thermal energy of  $1.5kT$  for a bacterium,<sup>31</sup> this allows a bacterium adhering in the secondary minimum to float and move away from and toward the substratum surface. This floating behavior is constrained by the width of the secondary minimum, while bacteria remain to adhere at an average distance above the surface dictated by the absolute secondary interaction minimum. The width of the secondary interaction minimum at  $1.5kT$  above the absolute minimum amounts around 15 to 20 nm, which coincides with the variations in distance ( $\Delta z$ ) above the surface observed using TIRM (Figure 4). Perpendicular, vibrational amplitudes of *S. salivarius* HBC12, adhering in a "floating" mode are much larger than of strain *S. salivarius* HB7, possessing fibrillar tethers (Figure 4). Tether coupling to the surface clearly restricts the vibrational amplitudes.

In order to obtain further evidence for a floating or tether-coupled mode of bacterial adhesion, the energy of adhering bacteria as a function of their distance above a substratum surface can be compared with the distant-dependent DLVO interaction energy (see Figure 2). The probability of a bacterium being located at a certain position  $z$  above a surface follows from the frequency histogram of bacterial positions around its equilibrium,  $P(z_t - \langle z_t \rangle)$ , which was related to the Boltzmann distribution<sup>32</sup> according to

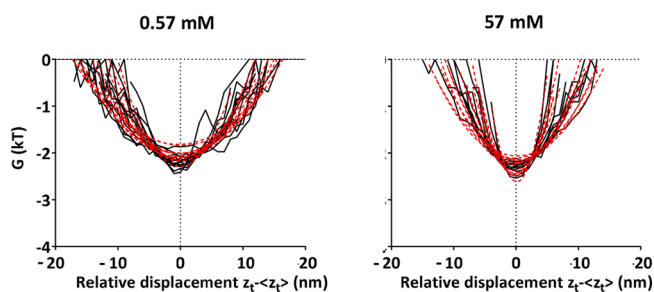
$$P(z_t - \langle z_t \rangle) = A \exp\left(-\frac{G(z_t - \langle z_t \rangle)}{k_B T}\right) \quad (8)$$

in which  $A$  is a normalization constant,  $G(z_t - \langle z_t \rangle)$  is the interaction energy at a position relative to the equilibrium  $z$ -position,  $\langle z_t \rangle$ , of a bacterium.

The interaction energy can now readily be calculated expressed in  $kT$  units according to

$$\frac{G(z_t - \langle z_t \rangle)}{k_B T} = -\log(P(z_t - \langle z_t \rangle)) + \log(A) \quad (9)$$

Neglecting  $\log(A)$  as a constant that merely defines the absolute energy level, the distance dependence of the interaction energy follows directly from the vibration amplitudes and associated probabilities that a position above the surface is occupied. Shapewise, the interaction energies of *S. salivarius* HB7 (Figure 5) are highly symmetrical and parabolic



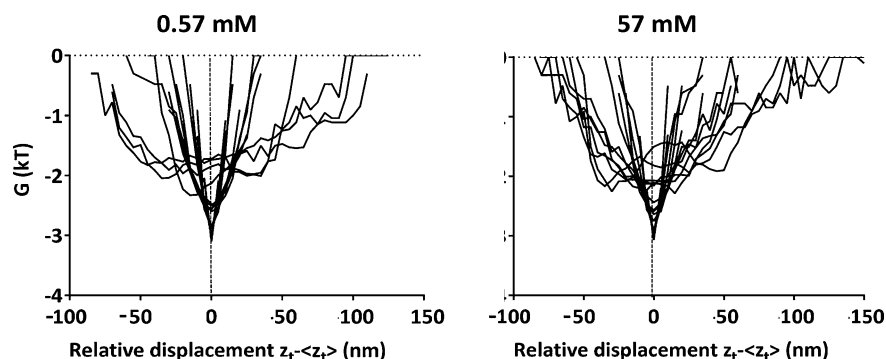
**Figure 5.** Comparison of the distance dependence of the interaction energy (calculated from perpendicular vibration amplitudes above the surface) between adhering *S. salivarius* HB7 with fibrillar surface tethers and hydrophobic, DDS-coated glass surface at two ionic strengths (0.57 and 57 mM) with the distance dependence calculated according to a harmonic oscillator model. Black lines represent the calculated interaction energy as a function of the relative displacement ( $z_t - \langle z_t \rangle$ ), and the red dotted lines represent their fitting to a harmonic oscillator model. The figure refers to 15 individual bacteria, each represented by one pair of black and red dotted lines, i.e., fits.

at both ionic strengths and can be well fitted to a harmonic oscillator model<sup>32</sup> according to

$$G(z_t - \langle z_t \rangle) = \frac{1}{2}k(z_t - \langle z_t \rangle)^2 \quad (10)$$

in which  $k$  is the spring constant of the tether, that can on average be calculated to be  $2.5 \times 10^{-5} \text{ N m}^{-1}$  regardless of ionic strengths. Thus, it can be concluded that streptococci with fibrillar surface tethers couple directly to a substratum surface, which requires piercing of the DLVO potential energy barrier by the tethers, which has been suggested before<sup>33</sup> but never backed-up with experimental evidence.

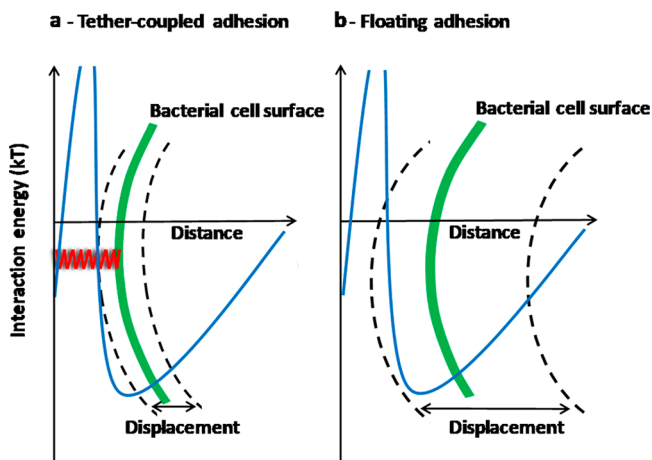
The distance dependence of the interaction energy calculated from vibration amplitudes for *S. salivarius* HBC12 is completely



**Figure 6.** Distance dependence of the interaction energy (calculated from perpendicular vibration amplitudes above the surface) between adhering *S. salivarius* HBC12 without fibrillar surface tethers and hydrophobic, DDS-coated glass surface at two ionic strengths (0.57 and 57 mM), showing an asymmetrical distribution around their equilibrium position. Black lines represent the calculated interaction energy as a function of the relative displacement ( $z_t - \langle z_t \rangle$ ). The figure refers to 15 individual bacteria, each represented by one line.

different than that for *S. salivarius* HB7 (compare Figures 5 and 6). Most notably, its distance dependence is not symmetrical around an equilibrium distance and therewith not according to a harmonic oscillator model but resembling the asymmetry of the DLVO secondary potential energy minimum (compare Figures 2 and 6). This confirms absence of tether-coupling and a mode of adhesion that we propose to call “floating-adhesion” above the surface.

In summary, we have provided experimental evidence for the existence of two modes of bacterial adhesion, as schematically summarized in Figure 7. Bacteria with fibrillar surface tethers



**Figure 7.** Schematics of tether-coupled adhesion of a bacterium (a) and floating adhesion (b). Note that tether-coupled adhesion requires piercing of the elastic tether (indicated as a red spring) through the DLVO-potential energy barrier.

adhere to a substratum surface in an irreversible fashion by tether-coupling to the surface, i.e. piercing of the tether through the potential energy barrier (Figure 7a). Distance variations above the surface over time are governed by the harmonic oscillations allowed by the spring. Bacteria without fibrillar surface tethers adhere in the secondary energy minimum with their perpendicular displacement over time restricted by the width of the secondary minimum at  $1.5kT$  above the minimum itself (Figure 7b). The distinction between “tether-coupled” and “floating” adhesion is new, and may have implications for bacterial detachment strategies, since detachment of tether-coupled bacteria may involve disrupting the bond of multiple

tethers with a surface while “floating” adhesion will be disrupted by decreasing the depth of the secondary interaction minimum, which is relatively easy, e.g., by changing prevailing ionic strength conditions or slight-rinsing of the substrata with adhering bacteria (data not shown). (For details on detachment forces involved in slight-rinsing see Gomez-Suarez, et al.<sup>34</sup>)

## ■ ASSOCIATED CONTENT

### 📄 Supporting Information

The Supporting Information is available free of charge on the ACS Publications website at DOI: 10.1021/acs.langmuir.7b04331.

Surface free energy parameters and components of water, formamide, and methyleneiodide used for contact angle measurements (PDF)

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### Notes

The authors declare the following competing financial interest(s): H.J.B. is also director of a consulting company SASA BV. The authors declare no potential conflicts of interest with respect to authorship and/or publication of this article.

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