The Alzheimer's β amyloid $(A\beta_{1-39})$ [monomer in an implicit solvent](http://dx.doi.org/10.1063/1.2907718)

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Results from replica-exchange and regular room temperature molecular dynamics simulations of the Alzheimer's β amyloid (A β_{1-39}) monomer in an implicit solvent are reported. Our data indicate that at room temperature, the monomer assumes random-coil and soluble conformations. No beta content is observed which therefore seems to be a product of oligomerization and aggregation of monomers. © *2008 American Institute of Physics*. DOI: [10.1063/1.2907718](http://dx.doi.org/10.1063/1.2907718)

I. INTRODUCTION

Alzheimer's disease (AD) is a neurological disorder that affects approximately 26×10^6 people worldwide. Its molecular hallmark is the presence of senile plaques made up of insoluble deposits of amyloid beta $(A\beta)$ peptides.¹ These peptides, $A\beta_{1-39}$ to $A\beta_{1-42}$, are released from amyloid precursor protein and have a chain length of 39–42 amino acid.^{2[,3](#page-4-2)} Mass spectrometry and solid state NMR spectroscopy indicates the formation of either parallel or antiparallel β strands (depending on the length of the fragment) $4\overline{-7}$ $4\overline{-7}$ that are placed perpendicular to the fibril axis. Less is known about the structure of $A\beta$ monomer as it easily aggregates and has a noncrystalline nature. The NMR study of unmodified fragments of $A\beta40$ and $A\beta42$ monomers in water indicated the regions of localized stability, but they do not resolve the peptide at atomic resolution.⁸ Circular dichroism (CD) experiments show that monomeric $A\beta$ is disordered, indicating 70% random, \sim 25% β sheet, and 5% α helix structure.^{9,[10](#page-4-7)} Whether these frequencies correspond to distinct clusters of conformations is not known.

Computational approaches provide an alternate means to study the fibrillogenesis process. Ma and Nussinov 11 suggested that at high temperature fragments $A\beta_{16-35}$ and $A\beta_{10-35}$ form a strand-loop-strand structure with parallel β sheet, with an interior salt bridge between Asp23 and Lys28 residues as a major element of the fibril structure. On the other hand, computational studies using coarse-grained as well as atomistic atom models suggest the formation of antiparallel β sheets by A β_{16-22} subpeptides.¹² The relation between this antiparallel structure and the nucleation and fibril forming properties of in-register parallel structures formed by the full length $A\beta_{1-40}$ and $A\beta_{1-42}$ peptides is not clear, and no simulation has so far reproduced the experimentally measured secondary structure contents for $A\beta$ peptides.

These discrepancies reflect the numerical difficulties of computational aggregation studies. Regular canonical simulations at low temperatures tend to be trapped in one of a huge number of local minimum energy states in a rough energy landscape. This difficulty is alleviated, in part, by the replica-exchange method, also known as parallel tempering[.13](#page-4-10)[,14](#page-4-11) Here, a number of replicas are simulated at different temperatures and periodically swapped between the neighboring temperatures. Over the course of simulation, each replica walks along a ladder of temperatures. In that way, a replica will not only find local minima but can also escape from them. As a consequence, sampling at low temperature is much faster than when all computational resources are set solely in that temperature.¹³

Following earlier work,¹⁵ we report in this article replica-exchange molecular dynamics (MD) simulations of $A\beta_{39}$ monomer. The prohibitive cost of simulations with explicit solvent molecules forces us to approximate the interaction between the peptide and the surrounding water by an implicit solvent model. We observe that at room temperature, solvated $A\beta_{39}$ monomer does not assume a unique folded conformation but exists as a mixture of rapidly interconverting conformations. Three distinct clusters dominate, characterized by a random-coil-like structure and only limited helicity is present. The conformation with the lowest energy is part of one of these clusters and is characterized by turn regions spanning Ala21-Asn27 residues. Slightly higher in energy are conformations with solely a turn around Val12- Leu17 residues and suchconformers that have a helical region in the N-terminal, i.e., similar to the one observed in NMR studies done in trifluoroethanol (TFE)/water solution.¹⁶

II. MATERIAL AND METHODS

All calculations were carried out with the AMBER 9.0 (assisted model building with energy refinement)^{[17](#page-4-4)} simulation package using the ff99 force field.¹⁸ The effect of solvation was approximated by a generalized Born solvent-accessible surface area¹⁹ implicit solvent model bond radii of 0.09 Å, solvent dielectric constant of 78.5, and surface tension of 0.005 kcal/mol A^{-2}). No attempt was made to modify the solvent parameters to capture their temperature dependence. The SHAKE algorithm was used to constrain all bond lengths to their equilibrium distances. 20

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The $A\beta_{39}$ peptide consists of 39 amino-acid residues: [Asp-Ala-Glu-Phe-Arg-His-Asp-Ser-Gly-Tyr-Glu-Val-His-His-Gln-Lys-Leu-Val-Phe-Phe-Ala-Glu-Asp-Val-Gly-Ser-Asn-Lys-Gly-Ala-Ile-Ile-Gly-Leu-Met-Val-Gly-Gly-Val. Canonical simulations were started both from NMR structures (PDB Id.: 1AML) and a completely extended conformation. 5000 steps of steepest decent minimization was followed by an initial equilibration run consisting of 50 000 steps of MD at 2 fs time steps with a cutoff radius of 10 Å.

The replica-exchange MD was implemented by using the Sander MD program. The 32 two replicas (fully extended configuration) were simulated over a range of temperatures from 200 to 640 K 200, 207, 215, 223, 232, 241, 250, 260, 270, 280, 291, 302, 313, 325, 338, 351, 364, 378, 392, 407, 423, 439, 456, 474, 492, 510, 530, 550, 571, 593, 616, 640). Exchange attempts are made after every 0.02 ps. Replica temperatures are maintained at the appropriate temperatures by a combination of velocity reassignment and a Langevin thermostat with a collision frequency²¹ of 5 ps⁻¹. Initial 5 ns are discarded and only the last 85 ns are analyzed. For comparison, we have also performed five regular canonical MD simulations of 90 ns length. The temperature is set by velocity reassignment from a Maxwell–Boltzmann distribution at 291 K and maintained at that temperature by using a Langevin thermostat with a collision frequency of 5 ps^{-1} .

Using multiscale modeling tools for structural biology $(MMTSB)$,^{[22](#page-4-18)} all structures collected in the replica-exchange MD production phase are grouped into clusters of similar geometry. Clustering is done according to the root mean square deviation (RMSD) over backbone atoms between configurations. Only clusters containing more than ten members are analyzed further. Final average structures are con-verted to the PDB format with program PYMOL.^{[23](#page-4-19)} For calculations of molecular properties, all snapshot structures were used and analyzed with the program DSSP.^{[24](#page-4-20)} Atoms within helical signatures (classes H alpha helix, G 3-helix $(3/10)$ helix), I 5 helix $(\pi$ helix)) are counted and compared to the total number of atoms resulting in values for the relative helical content. For the calculation of the specific heat, 25 we compute the heat capacities from finite differences, i.e., $C_p = [E(T_{n+1}) - E(T_{n-1})]/(T_{n+1} - T_{n-1})$, with $E(T_n)$ as the average total energy at some temperature T_n . The radius of gyration and the fraction of helical content are calculated by using the AMBER ptraj module and DSSP.²⁴

III. RESULTS AND DISCUSSION

In order to probe the structure of $A\beta_{39}$ monomer, we have performed multiple replica-exchange and regular canonical MD (at $T = 291$ K) simulations starting with different initial configurations. Figure [1](#page-1-0) shows the random walk for one replica in temperatures. Associated with it is a walk in energy, allowing the replica to escape local minima. As a consequence, reliable physical quantities can be calculated over the whole range of temperatures. An example is the specific heat plotted in Fig. [1,](#page-1-0) giving an approximate melting temperature of 330–360 K, i.e., above our temperature of interest $(T=291 \text{ K})$. With the exception of a single canonical MD run, that got stuck in a local minimum with higher

FIG. 1. (a) Time series of temperature exchange for one of the replicas in the replica-exchange MD simulation. (b) Specific heat as a function of temperature. The line connecting the data points is guide for the eyes.

energy than found in all other runs, all simulations set at *T*=291 K to similar U-shape structures, however, the energy is \approx 48 kcal/mol lower in the replica-exchange simulation. The improved sampling at room temperature is also seen in Fig. [2,](#page-1-1) where we compare the distribution of dihedral angles

FIG. 2. (Color online) Distribution of dihedral angles (ϕ, ψ) of Gly25 for $T=291$ K obtained (a) from a replica-exchange MD simulation and (b) from regular canonical MD simulation.

FIG. 3. (Color) The three dominant clusters of $(A\beta_{39})$ configurations seen in the simulations.

 (ϕ, ψ) of Gly25. At room temperature, the replica-exchange simulation samples a much wider part of configuration space as the conventional canonical MD run. 26

At room temperature, the $A\beta$ monomer does not have a unique folded conformation in water but adopts one of the

several low energy structures as seen, for instance, in Fig. [3.](#page-2-0) The root mean square displacement based clustering analysis²² shows that the ensemble at 291 K consists of 37 clusters. Of these, 26 can be further grouped by visual inspection into three super clusters. These three clusters differ little in their average potential energies [cluster A, 946.09(4) kcal/mol; cluster B, 962.59(2) kcal/mol; cluster C , 949.10(2) kcal/mol]. All three clusters share as a common theme the U-shape configurations that are strongly amphipathic. The longer arm is made mostly of hydrophilic and charged amino-acid residues, while the shorter arm is formed exclusively by hydrophobic branched residues. This "U-shape" is detrimental to the structure adopted by the sequence in the full protein. On the other hand, NMR results show conformations varying from α -helical structure in a nonpolar solution⁸ to disordered N-terminal and C-terminal tails with consistent turn regions as determined by electrospray mass spectroscopy and implicit solvent MD. 26,27 26,27 26,27 The side chain of most hydrophobic residues found in the shorter arm point outward, providing a hydrophobic platform for an association between $A\beta$ monomer. This is consistent with the simulation by Tarus *et al.*^{[28](#page-4-24)} of dimerization of $A\beta_{10-35}$ using umbrella sampling and MD. They show also that hydrophobic interactions between monomers are the stabilizing forces and that the dimerization process involves substantial structure reorganization of both the C and N termini.

A more thorough analysis reveals differences between the clusters. The dominant one, cluster A, is a random coil structure with turn around Ala21-Asn27 and appears with a frequency of 58%. It agrees well with the NMR studies of the $A\beta$ fibril.⁵ The turn around of these residues are not unexpected as the Gly25, Ser26, and Asp227 residues have a high probability to be within a beta turn.²⁹ Only marginal secondary content $(3-10$ helix, α helix and turns as determined by DSSP program, 24 24 24 but no β structures) are observed for configurations of this cluster. The region around Leu17- Ala21 residues is well defined and its structure fluctuates considerably less than rest of the peptide. This is consistent with NMR studies, $8,30,31$ $8,30,31$ $8,30,31$ which suggest a central hydrophobic region, Leu17-Ala21, with well defined structure in the $A\beta$ peptide, and this region is found to be critical for fibril formation. 32 However, cluster A deviates from the ones found in NMR experiments of $A\beta$ in TFE/water.¹⁶ The average main chain RMSD between cluster A structure and the 20 NMR models of the 1AML.PDB is 9.89 Å, where structural superposition is done by using sequence alignment on 1–39 residues.

Configurations of cluster B are also random coil but show a turn around Val12-Leu17. This cluster appears at *T*=291 K with a frequency of 29%. Only in these configurations, we observed the salt bridges between Asp23- Lys $28/16$ residues (Fig. [3](#page-2-0)) that have been reported in various experiments as controlling the aggregation rate of $A\beta$.^{[33](#page-4-30)} In agreement with Baumketner and Shea,³⁴ we find that this salt bridge is easily formed and stabilizes the U-type structure. Other polar contacts that stabilize the structure are the ones between carboxyl group of the side chain of Asp7 and the amino group in the side chain of Lys28, between the side chain C= σ of Glu1 and the side chain NH₂ group of Arg5, and between the side chain amidegroup of Lys16 and the side chain $C = 0$ of Asp7. The optimized structures have a short right handed α helix region with dihedral angle $-60 \lt \phi, \psi \lt -45$ extending around Gly15-Val18 residues. The β type1 turns are present around Arg5-Ser8 residues

with (ϕ, ψ) values of the $(i+1)$ th residue $(-3.2, -23.1)$ and $(i+2)$ th residue $(-59.2, 1.0)$ in degrees; and around Leu17-Phe20 residues with ϕ and ψ values of $(i+1)$ th residue $(-55.7, -23.6)$ and $(i+2)$ th residue $(-97.0, 2.2)$. γ turns regions are also observed, stabilized by an extensive although not a regular array of hydrogen bond between the $C=O$ group of the *i*th residue and the NH group of the $(i+2)$ th residue. Carbonyl-carbonyl interactions are also observed with various distances such as $dO_i - C_{i+1}$, dC_i — O_{i+1} , etc., and these distances lie in the range of $3.5-4.2$ Å.

Finally, 13% of configurations are rich in helicity and form cluster C. Configurations in this cluster are similar to the structure that was determined in a 40% TFE/water solution. The average main chain RMSD between cluster C configurations and the 20 NMR models of the 1AML entry in the protein data bank is 2.2 Å. The structural comparison is done by using sequence alignment on 1–39 residues. Cluster C exhibits a helix around Gln11-Glu22 residues, another around Ile32-Met35 residues, and a coiled region around Val24-Ala30 residues. All three elements are in agreement with the NMR experiment¹⁶ in 40% (v/v) TFE/water solution.

Our results indicate that the experimentally measured secondary structure contents are thus the consequences of elemental contributions from three clusters. About 90% of configurations belong to either cluster A or B, and show characteristics found also in the fibrillar conformations. Cluster A exhibits the turn between residues around Asp23 and Lys28 found in NMR studies of the $A\beta$ fibrillar structure.⁵ However, only the smaller cluster B has the salt bridges between Asp23-Lys28/16 residues that have been reported in various experiments 33 as controlling the aggregation rate of $A\beta$. Following the comparison with the experimental results, we conjecture that configurations of this cluster are especially prone for oligomerization and aggregation. The remaining $\approx 10\%$ of configurations include the helix rich structures of cluster C. These are similar to the tertiary structure of $1AML.PDB.¹⁶$ While that structure was obtained by NMR for the monomer in a TFE/water mixture, our results show that this structure appears also in the more physiologically relevant case of the monomer in water.

The percentage of ordered residues in aqueous $A\beta_{39}$ in our simulations agrees reasonably with the NMR data $35,36$ $35,36$ that estimate the amount of disordered structure to be 60%– 80%, and α helices to be 10%. However, our simulations did not reveal the origin of the 25% β content found in CD studies for disordered $A\beta_{39}$ in water. We conjecture that the higher percentage of ordered structure (70% random, 25% β , and 5% helical), and especially the β content, in the CD experiment results from the fact that these experiments probe not only the monomer (as our work) but also a dynamic equilibrium between monomers and oligomers (i.e., dimer and trimer). If this assumption is correct, it would indicate that β -structures are only formed in the processes of oligomerization, and are not a feature of the monomer. The brain microenvironmental influences are required to convert these configurations in β sheets during the process of β -amyloidosis in AD.

Since the $A\beta_{39}$ molecule does not contain any regular secondary structure element (except a very short small region), most of the potential hydrogen bonding sites are free, and it is the hydration of these moieties and other hydrophilic groups that makes them soluble. Solubility is also supported by the free hydrogen bonding sites at the ends of both the arms.

IV. CONCLUSIONS

Our study demonstrates that increased computational power and sophisticated sampling techniques allow now a detailed analysis of the folding and interactions of $A\beta$ peptides. Replica-exchange MD simulations illustrate that $A\beta_{39}$ monomer at room temperature does not have a unique structure in water but adopt one of the several possible lowenergy conformations. Three distinct families of mostly random coil-like structures dominate, with one cluster having a turn around residues of 21–27, another with few helical elements present, and a third having a turn around residues of 12–17. The exposed free hydrogen bonding sites indicate that the monomer is soluble. However, no stable β -structure has been observed. This indicates that the sheets observed in fibrillar $A\beta$ form only later in the process of oligomerization and aggregation while at the onset of the process are monomers that exist as soluble random coil configurations. We hope that further analysis of the interactions stabilizing the "U-like" structure in the monomer, with a longer arm rich in hydrophilic and charged residues and the shorter arm dominated by hydrophobic residues, will help to understand the process of dimerization and oligomerization of $A\beta$.

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