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Exploiting delayed transitions to sustain semiarid ecosystems after catastrophic shifts

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Semiarid ecosystems (including arid, semiarid and dry-subhumid ecosystems) span more than 40% of extant habitats and contain a similar percentage of the human population. Theoretical models and palaeoclimatic data predict a grim future, with rapid shifts towards a desert state, with accelerated diversity losses and ecological collapses. These shifts are a consequence of the special nonlinearities resulting from ecological facilitation. Here, we investigate a simple model of semiarid ecosystems identifying the so-called ghost, which appears after a catastrophic transition from a vegetated to a desert state once a critical rate of soil degradation is overcome. The ghost involves a slowdown of transients towards the desert state, making the ecosystem seem stable even though vegetation extinction is inevitable. We use this model to show how to exploit the ecological ghosts to avoid collapse. Doing so involves the restoration of small fractions of desert areas with vegetation capable of maintaining a stable community once the catastrophic shift condition has been achieved. This intervention method is successfully tested under the presence of demographic stochastic fluctuations.

1. Introduction

A major consequence of the habitat-degradation processes associated with the Anthropocene is the accelerated loss of habitats and worldwide species extinctions [1–4]. Moreover, the nature of this decay is likely to be catastrophic, and evidence from field, experimental and palaeoclimatic data along with modelling efforts strongly indicate that such decay will end up in collapse [5–7]. Such ecological shifts are a consequence of the multistable nature of ecosystems, which is especially relevant in semiarid ecosystems. Indeed, recent field studies of vegetation spatial patterns in drylands have identified multiple ecosystem states [8].

Because of the global relevance of drylands [9,10], dedicated efforts have been made towards forecasting the so-called green-desert transition from a vegetated ecosystem to a desert area [6,11–13]. These warning signals are associated with the presence of increasing fluctuations that are known to diverge close to critical points [7] as well as with changes in spatial patterns [14]. The definition and quantification of these warning signals have been a very active research area over the last decade [15]. Less attention has been given to another type of non-linear phenomenon associated with transient behaviour when a parameter μ surpasses a critical value μ_c . It involves the presence of a special class of very long transients, also called delayed transition or ghost [16–20].

The importance of transients in ecological systems has been recently discussed [21,22]. Importantly, transients have been characterized in experimental

systems with *Tribolium* sp. [23] as well as in field data for the crab *Cancer magister* [24]. In a ghost-related transient, the system will appear stationary, thus creating the false impression that the system is stable, despite its assured collapse. In particular, ghosts follow a universal scaling law [17,25] and the length of transients, τ , scales following an inverse square-root law:

$$\tau \sim (\mu_c - \mu)^{-1/2}, \quad (1.1)$$

with μ and μ_c being the control parameter and the bifurcation value, respectively. This scaling law has been identified in physical systems [16,26], in mathematical metapopulation models for autocatalytic species [27], low-dimension [28–30] and high-dimensional replicator models exhibiting cooperation [31,32], and single-species discrete models with Allee effects [33]. In this sense, the presence of very long transients can force us to redefine the concept of steady state [21,22,34] and provide essential clues to the persistence and organization of populations over long timescales.

If ghosts appear in green-desert transitions, the take-home message is simple but leads to a grim picture: a given ecosystem might be on its way towards collapse despite its apparent stability. In this paper, we show that ghosts are indeed expected in the dynamics of semiarid ecosystem models. However, a positive message is that the ghost can be exploited to ensure ecosystem's persistence. We introduce a simple intervention method based on the recovery of small fractions of desert areas with vegetation as a way to keep dry-land dynamics near the phase space regions where delays occur, thus allowing the ecosystem to remain stable once the bifurcation towards desertification has occurred.

2. Mathematical model

The starting point of our study will be the model developed by Kéfi and co-workers as a coarse-grained approximation of semiarid ecosystems [11]. This model captures the nonlinear interactions due to facilitation and the relevant alternative states found in semi-arid ecosystems (vegetated area, fertile soil with the potential of hosting vegetation and desert areas). In this context, we will not explicitly take into account spatial structure [35–37], an explicit coupling with soil moisture [38,39], or some sources of heterogeneity [40], which can play a relevant role. Here, we aim to understand the properties and management of long transients close to the green-desert shift found in a minimal model capturing the main mechanisms of facilitation.

The original model [11] considers a spatial lattice model in which each site can be in one of three alternative states: vegetated (V), non-vegetated but fertile (S) and destroyed (D). Each state can transition to another with some probability, e.g. $\omega(V \rightarrow S)$. In this model, these transition probabilities are given by

$$\begin{aligned} \omega(D \rightarrow S) &= r + f\rho_V, \\ \omega(S \rightarrow D) &= d, \\ \omega(S \rightarrow V) &= (\delta\Delta_V + (1 - \delta)\rho_V)(b - c\Delta_V) \end{aligned}$$

and $\omega(V \rightarrow S) = m$.

Because of the nature of semiarid ecosystems, the transition probabilities are driven not only by the nearest vegetated neighbours (ρ_V) but also by non-local processes that depend on the fraction of global vegetation (Δ_V). Here, r is the

spontaneous regeneration rate of a site due to, e.g., the arrival of seeds by the wind. The parameters f and d stand for the facilitation and the soil degradation rates, respectively. Regarding the spread of the vegetation, the dispersion of the seeds is modelled by the constant δ (ratio of vegetation due to the whole vegetated sites) and by the development (b) and death (c) of the seeds. Finally, the rate of vegetation decay is given by a probability m .

The microscopic transition rules provide the rates of growth and death that can be incorporated into a time-continuous (mean-field) modelling framework. The resulting equations are expressed in terms of balances between different processes:

$$\begin{aligned} \frac{dD}{dt} &= \omega(S \rightarrow D)S - \omega(D \rightarrow S)D, \\ \frac{dV}{dt} &= \omega(S \rightarrow V)S - \omega(V \rightarrow S)V \\ \text{and } \frac{dS}{dt} &= \omega(V \rightarrow S)V - \omega(S \rightarrow V)S \\ &\quad + \omega(D \rightarrow S)D - \omega(S \rightarrow D)S. \end{aligned}$$

Ignoring spatial correlations, we have $\Delta_V = \rho_V = V$. This assumption leads us to the following system of differential equations:

$$\begin{aligned} \frac{dD}{dt} &= dS - rD - fVD, \\ \frac{dV}{dt} &= (b - cV)SV - mV \\ \text{and } \frac{dS}{dt} &= V(m - (b - cV)S + fD) + rD - dS. \end{aligned}$$

The state variables of the system (D, V, S) correspond to the ratio of the area that is occupied by each different state, with $D + V + S = 1$. Thus, we can reduce the three-variable system to a two-variable one by means of the linear relation $D = 1 - (S + V)$. The two-variable system is now

$$\frac{dV}{dt} = V((b - cV)S - m) \quad (2.1)$$

and

$$\frac{dS}{dt} = (r + fV)(1 - S - V) - dS - V[(b - cV)S - m]. \quad (2.2)$$

Given that we consider the whole semiarid ecosystems area, no income terms are included. In our case, the laws and conditions that affect the ecosystem depend on the variables in the ecosystem. For this reason, there is no spontaneous recovery of the degraded state to the non-occupied state (i.e. $r = 0$).

3. Results

3.1. Catastrophic shifts in semiarid ecosystems

The model given by equations (2.1) and (2.2) displays bistability: for a low soil degradation rate, d , two possible exclusive stable states are present: an ecosystem with vegetation and a completely desert ecosystem. Actually, this system has four fixed points. The first fixed point is ($V^* = 0, S^* = 0$), which corresponds, in the case of being stable, to the extinction of vegetation and the destruction of the fertile soil, with a dominance of the desert state. The extinction state is locally stable (see electronic supplementary material, Section S1), and the system displays an abrupt transition. This catastrophic shift leads to a collapse of the vegetated

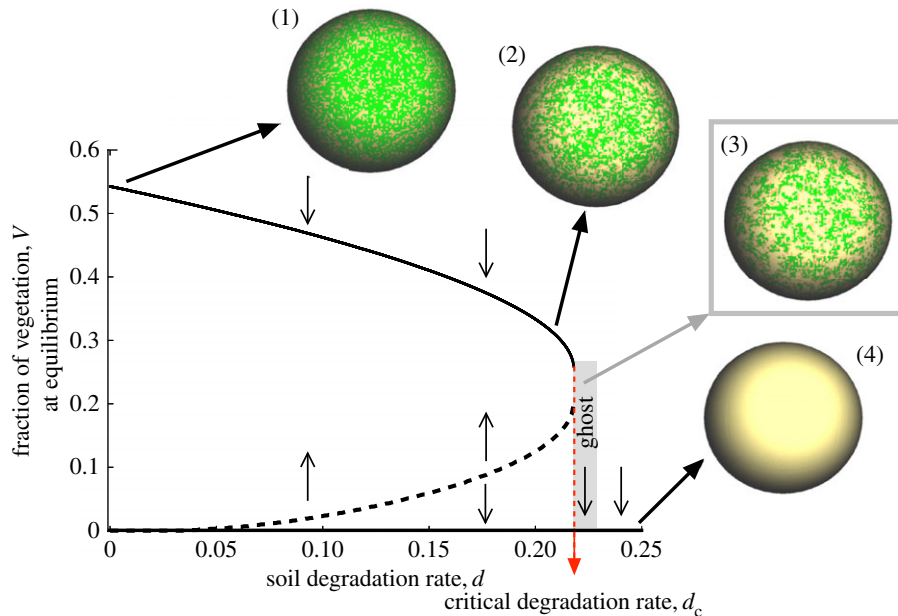


Figure 1. Catastrophic transition occurring at the critical value of the soil degradation rate, d_c (indicated with a dashed red line), found in a mean field version of the spatial model by Kéfi *et al.* [11]. The model considers three variables: vegetated area (V), non-vegetated area with fertile soil (S) and desert area (D). The diagram displays the equilibrium of V against soil degradation (d) using $b = 0.3$, $c = 0.15$, $m = 0.1$ and $f = 0.9$. The upper (solid line) and lower (dashed line) branches are the stable (node) and unstable (saddle) fixed points, respectively. Typical of saddle–node bifurcations, the unstable branch displays a kind of symmetry and increases monotonically until the two fixed points collide. After the bifurcation occurring at $d = d_c$, a ghost causing a delayed transition appears. The ghost involves extremely long times towards the desert state (note that $V^* = 0$ is a stable state given by the fixed point P_0^* ; see electronic supplementary material, Section S1), indicating an apparently stationary state of the vegetation (the small, thin arrows indicate the stability of the branches and of the origin; figure 2). The spheres are spatial snapshots at $t = 100$ using $d = 0.001$ (1); $d = 0.21$ (2); $d = 0.22 \gtrsim d_c$ (3) and $d = 0.25$ (4).

ecosystem once $d > d_c$, and the system falls into the desert state. These results are summarized in figure 1, where the fraction of vegetated habitat is plotted against d . The increase in d involves a monotonous decrease of vegetation. As mentioned, at a given critical value $d = d_c$, an abrupt extinction takes place due to a saddle-node bifurcation. In the bifurcation diagram, note that the two branches of the fixed points (one being stable and the other unstable) approach each other as d increases. Specifically, under the parameter values used in our analyses (i.e. $b = 0.3$, $c = 0.15$, $m = 0.1$ and $f = 0.9$), the bifurcation value would be $d_c \approx 0.22$ (see Section S2 for an accurate computation of d_c).

The phenomenon that we are investigating here takes place right after the bifurcation. When $d \gtrsim d_c$, the system has a single steady state: the desert state. However, for a value of d close enough to d_c , the time to extinction T_e rapidly diverges. This time divergence always occurs when bifurcation values are approached [17] and leads to a paradoxical situation: the system appears to be in a steady state, but it suddenly collapses after a long time period. This phenomenon is also depicted in figure 1 by means of the spatial representation of the ecosystem. The snapshots in figure 1 display the vegetation pattern using four values of d . For case (1), the ecosystem is fully vegetated, while for cases (2) and (3), the fraction of vegetation is intermediate. However, case (3) is obtained with a value of d past the bifurcation point, but the system remains vegetated for a long time, eventually decaying into a completely desert state, represented by the sphere in case (4).

3.2. Dynamics near the green-desert transition

The detailed dynamics under different d values used in figure 1 are displayed in figure 2 by means of trajectories in

the (V, S) phase space, confined to a simplex, as the sum of all of the variables is constant (i.e. $V + S + D = 1$). When $d = 0.001$, the attractor involving a fraction of vegetated V and of fertile S soil is close to the boundary of the simplex, and the basin of attraction of the fixed point ($V = 0, S = 0$) is extremely small (see the time series in figure 2(1) and snapshot 1 in figure 1). Typically, for a given value of d and for a given initial fraction of vegetated area $V(0)$, trajectories in phase space go to the slow manifold (hereafter SM) and then approach the stable fixed point. The SM corresponds to the set of points in the phase space from which the attractor is reached. These dynamics can be clearly observed in the simplexes as well as in the time series of figure 2. As d increases, the basin of attraction of the fixed point $(0, 0)$ (shown in red) becomes larger. For this case, the system is less vegetated (see sphere (2) in figure 1). Under this scenario, the dynamics within the basin of attraction of the fixed point involving a vegetated ecosystem makes the orbits to rapidly flow to the SM of this equilibrium point. Once this manifold is reached, there is a slow approach towards this equilibrium (see the thick trajectory in the zoom of figure 2(2)).

The two previous scenarios involved a stable green state, provided that $V(0)$ and $S(0)$ were large enough. Once d is slightly above its critical value (figure 2(3)), all of the trajectories end up in the desert state attractor $(0, 0)$. However, the ghost (or saddle-node remnant) appears close to the region where the stable node and the saddle collide, leading to a saddle-node bifurcation [17]. Now, the flows undergo a very slow passage through a bottleneck before reaching the desert state. This is confirmed by the slow decay displayed by the $V(t)$ time series. The extinction dynamics becomes very fast at high d (see time series in figure 2(4)). Indeed, for a large d , trajectories usually miss the ‘gap’ on the slow

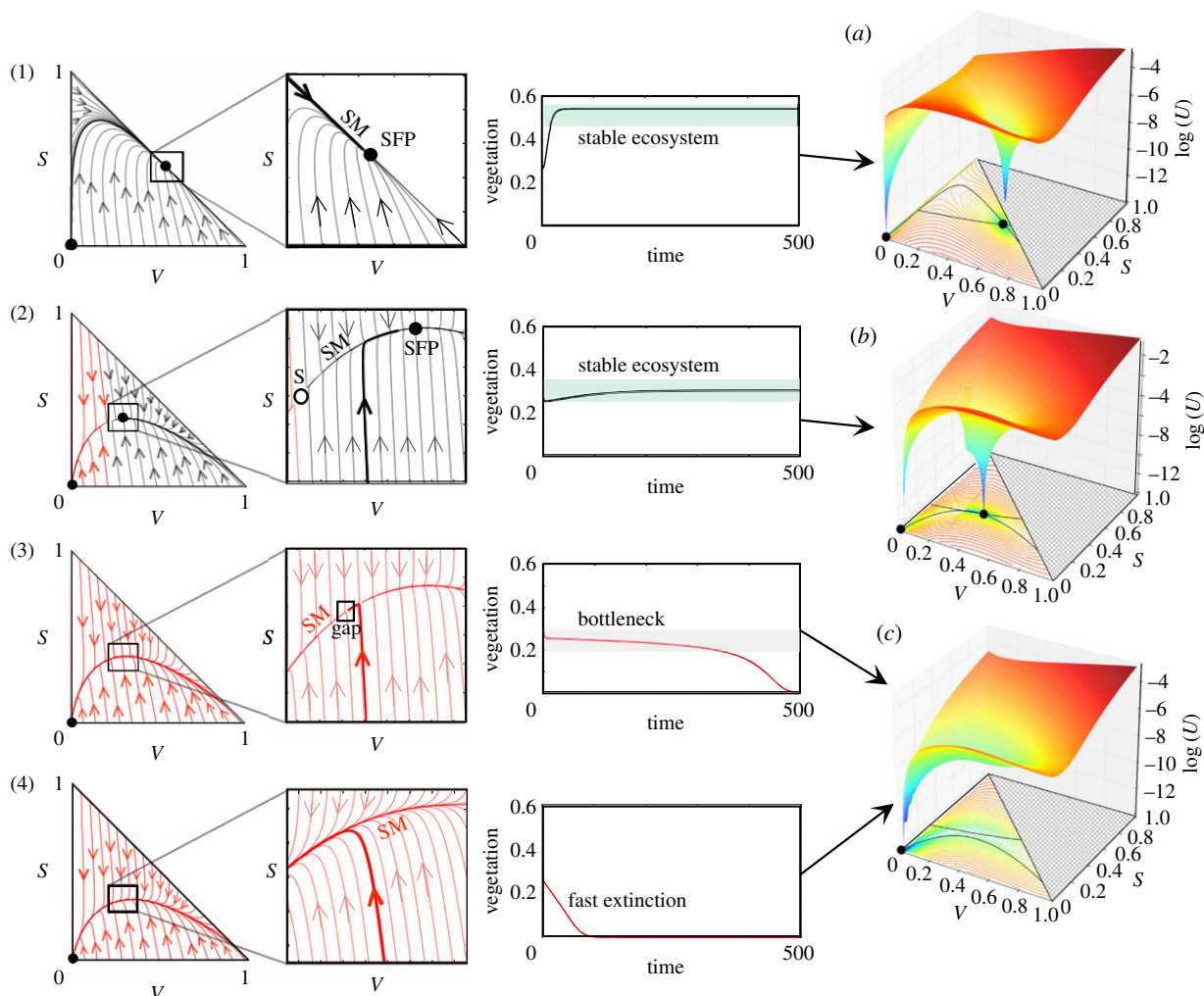


Figure 2. Dynamics as the degradation rate of fertile soil, d , increases, with: (1) $d = 0.001$, (2) $d = 0.21$, (3) $d = 0.22$ and (4) $d = 0.3$. Examples (1) and (2) are computed for $d < d_c$, where the system is bistable (SFP: stable fixed point (black circles); S: saddle (white circle)). As d increases (2), the basin of attraction for the desert state grows (illustrated with red orbits). The time series for both cases show the persistence of the vegetation. In panels (3) and (4), the critical point d_c has been overcome. While close to d_c (3), the trajectories undergo a slow passage once they reach the slow manifold (SM). The time series display extinction after a long time period (note the small gap at the left of the trajectory on the manifold, which indicates that the orbit is delayed). When the value of d is further increased, extinction occurs rapidly (4). The thick trajectories in the zooms and in the time series in all panels are obtained with the initial condition $V(0) = 0.28$ and $S(0) = 0.2$. Panels (a–c) in the right column: quasi-potential landscapes, $\log(U)$ (see electronic supplementary material, Section S4), before the saddle-node bifurcation (a); at the bifurcation value (b); and once the bifurcation has taken place (c). Specifically, we use (a) $d = 0.001$, (b) $d = 0.2$ and (c) $d = 0.3$; the other parameters are the same used in figure 1. The nullclines (electronic supplementary material, Section S2) for each panel are displayed on the V – S plane. The colour lines on the V – S plane are the isodines indicating the value of the quasi-potential function. The solid circles are stable fixed points.

manifold, resulting in shorter transient times towards full desertification.

Another way to see the qualitative changes associated with the emergence of a ghost is provided by the so-called potential functions or quasi-potentials [1]. These surfaces provide information on the stability of the fixed points. Specifically, stable and unstable fixed points appear as valleys and peaks, respectively. The quasi-potential for equations (2.1) and (2.2) in its general form reads $U(V, S) = -\int [(dV/dt)^2 + (dS/dt)^2] dt$ (see electronic supplementary material, Section S4 for further details). The different surfaces associated with the previous flows are displayed in the right-hand side of figure 2. Two minima are clearly visible for $d < d_c$, whereas a drastic change occurs for $d \approx d_c$, where the quasi-potential function has a single minimum (figure 2c). The previous two-well shape has now been replaced by an elongated, almost flat channel associated with the ghost and the observed bottleneck of the trajectories.

3.3. Delayed green-desert transitions

The dependence on d of the times to extinction, T_{ev} , is displayed in figure 3a, where we can appreciate the almost stationary behaviour of the green state within the bottleneck, as compared with the steady ecosystem before the critical point. For $d < d_c$ but close to criticality, the green state is stable, with a vegetation cover of $V \approx 27\%$ (dashed line in figure 3a). If d is slightly increased, the system will achieve a seemingly stable fraction of vegetated habitat (of about 25%), which actually still holds at $t \approx 3.5 \times 10^5$. However, as the bifurcation has already occurred, the system rapidly collapses after this extremely long delay at $t \approx 350\,000$. Several examples are represented for slightly larger values of d . The same behaviour is found, although here, the times to extinction are shorter. Indeed, the dependence of the extinction times on the distance of the bifurcation parameter (d) from the bifurcation value (d_c) is shown to follow the inverse square-root scaling law (figure 3b), in agreement with

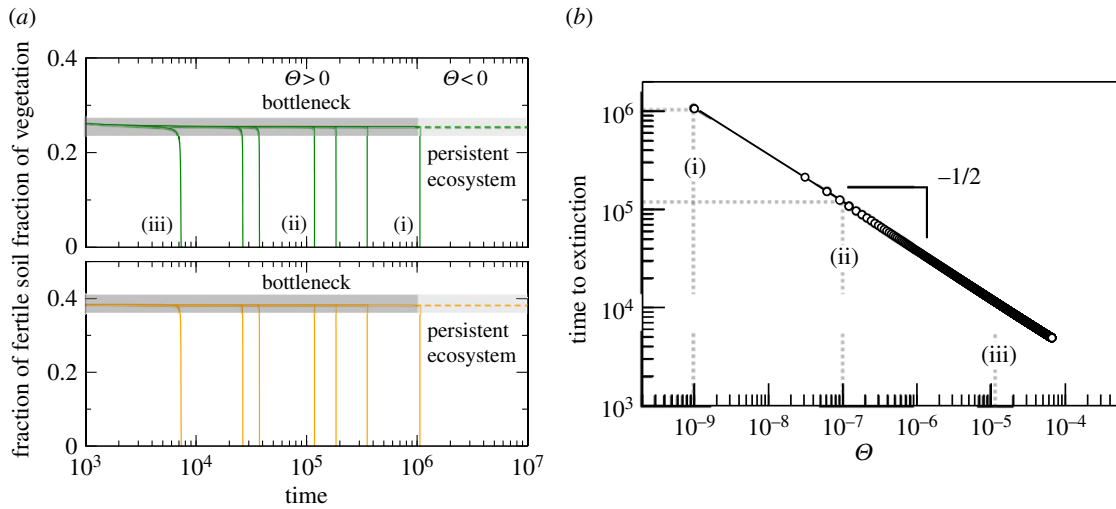


Figure 3. Time to extinction against $\Theta = d - d_c$, with d and d_c being the degradation rate of fertile soil and the bifurcation value, respectively. (a) Ecosystem persistence below the bifurcation threshold (dashed lines, $d = 0.127$, with $\Theta < 0$) and dynamics towards desertification above the bifurcation threshold (solid lines with $\Theta > 0$; time series (i)–(iii) illustrate three examples of transients to the desert state, which are also displayed in panel (b)). Note that in both panels, we used the same parameter values used in figure 1, except for d . (b) Inverse square-root scaling law dependence on extinction transients as parameter d increases above d_c (shown in log–log scale).

previous works on delayed transitions tied to saddle–node bifurcations [16–18,28,33]. Here, T_e is plotted against the distance to the bifurcation value using $\Theta = d - d_c$, and the power-law dependence $T_e \sim 1/\sqrt{\Theta}$ is found numerically.

The delaying capacities of the ghost are known to depend on the initial conditions [27]. Usually, for large enough initial conditions, the ghost captures the orbits and the delays appear. This phenomenon is illustrated in electronic supplementary material, figure S3, which displays extinction times for values of d close to the bifurcation (by using $10^3 \times 10^3$ different initial conditions). Here, two well-differentiated regions are found. One region where times are extremely long ($T_e \approx 5 \times 10^5$ in electronic supplementary material, figure S3a) exists at the right-hand side of the place where the saddle and the node collided (white area in the simplex of electronic supplementary material, figure S3). The second one, indicated in black (electronic supplementary material, figure S3a), shows the initial conditions for which the orbits go to the desert state very rapidly. Indeed, extremely close to the bifurcation value, these two regions indicate initial conditions with extremely long transients or initial conditions for which extinction is almost immediate. Once the bifurcation value is further increased, the times decrease substantially. In both cases, long transients are about 1.5 orders of magnitude below the times of electronic supplementary material, figure S3a. Note that in electronic supplementary material, figure S3, we display time series for values at the right side of the simplex (upper panels with long transients) as well as to the left, where the transients are much shorter.

3.4. Controlling catastrophic shifts

The nature of the ghost behaviour allows us to enormously expand the extinction time by using small external interventions with suitable frequencies and magnitudes. In our case study, the interventions would involve restoring vegetation in small desert patches. The restoration parameters are the fraction of replanted area (expressed as the amount of

vegetation increased, ΔV) and the restoration frequency (hereafter ‘freq’). Note that the time lapse between interventions is assumed to be much shorter than the timescale of the ecosystem dynamics. In this context, our restoration events are assumed to be effectively instantaneous. The intervention strategies can range from very rare (low-frequency) interventions affecting large areas to small but very frequent events. The effect of different values for both parameters ΔV and freq leads to different dynamical scenarios (see below).

The response of the ecosystem to a given intervention strategy depends on the actual ecosystem conditions and in particular on the increase of the soil degradation rate beyond the tipping point. To evaluate the response of the ecosystem, we have simulated the external interventions incorporating the external increase of vegetated area, performed at a given frequency (see electronic supplementary material, section S5). Figure 4a–c summarizes the nature of the intervention and its effects. Periodic perturbations create a cyclic behaviour (figure 4a), in which the trajectories are constantly re-injected close to the ghost. The effect of the intervention is to preserve (or return) the system within the region where the ghost has a strong delaying effect (see also electronic supplementary material, figure S4). When the pulses fail in re-injecting the trajectory towards the regions influenced by the ghost—i.e. to the right of the ‘gap’ on the slow manifold—the ecosystem decays to the desert state. On the other hand, when trajectories are efficiently sent to the ghost, the ecosystem remains in an stable green state.

To quantify how optimal the interventions are, we have to take into account that they depend on two parameters: freq and ΔV . For this reason, we obtain a surface that optimizes the intervention parameters for a given Θ , given by a Pareto front (electronic supplementary material, figure 4c). However, to find an optimal intervention strategy for the persistence of a given semiarid ecosystem, it is necessary to design a strategy that optimizes both parameters simultaneously. In our case, the optimality of an intervention is evaluated with a simple cost function, defined as the fraction

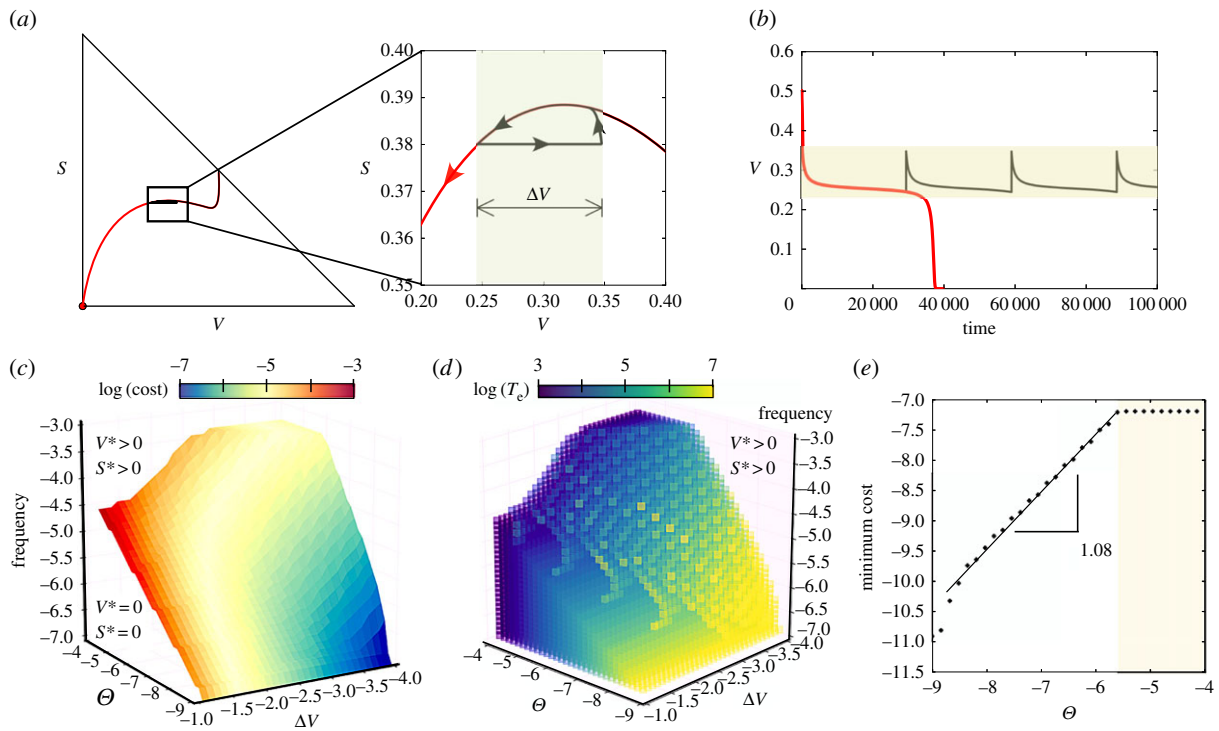


Figure 4. (a,b) Ghost-related restoration dynamics. (a) Phase portrait with a trajectory (red) that achieves extinction after suffering a delay of $t \approx 4 \times 10^4$. Overlapped, we display the dynamics using the same initial condition but applying the intervention (black orbit), here using $\Delta V = 0.1$ and $\text{freq} = 3 \times 10^{-4}$. The enlarged view displays the trajectories during the interventions, which return the orbits near the ghost, giving way to an externally forced oscillatory-like behaviour (black orbit). The red orbit displays the ecosystem's extinction without the intervention. (b) Time series for both orbits displayed in (a). The parameters for these calculations are equal to those used in the previous figures. (c) Limit between extinction and persistence of the vegetation as a function of the frequency and ΔV when d increases above the bifurcation threshold $\Theta = d - d_c$. (d) Time to ecosystem extinction towards desert in the same space as in (c), i.e. extinction times below the surface displayed in (c). (e) Minimum cost to preserve vegetation at increasing Θ using $\Delta V = 10^{-4}$. The straight line shows a power-law fit with exponent 1.08 (with correlation coefficient 0.9965; see also electronic supplementary material, figure S5 for further examples). The saturation of the minimum cost in panel (e), indicated with a yellow transparency, is due to the fixed maximum frequencies and values of ΔV used to perform the analyses. Note that all of the axes in panels (c–e) are in logarithmic scale.

of desert areas replanted per unit time ((plants/area)·time⁻¹) that is needed to sustain the system avoiding desertification, with the following equation:

$$\text{cost} = \Delta V \times \text{freq.}$$

In this study, the cost is assumed to increase if the fraction of replanted area and/or the frequency of planting increases. The changes in the cost as the degradation rate increases are shown in figure 4e. Note that the relation between the increase of degradation (Θ) and the optimal cost of intervention is almost close to a hyperbolic function (linear in log–log scale with exponent 1.08). The slope of this cost changes depending on the frequency of the interventions and on the area replanted. In the case of fixing the desired area to be replanted at each intervention, the optimal frequency will correspond to the Pareto front surface (electronic supplementary material, figure S5a). For larger amounts of added vegetation, ΔV , the increase in the cost function and the optimal intervention frequency are lower. The slope goes from 0.5 for large interventions ($\Delta V = 0.1$) to 1.08 for very small replantings ($\Delta V = 10^{-4}$); see electronic supplementary material, figure S5b.

As previously mentioned, another possible scenario may be found depending on the magnitude and frequency of the interventions, which may not be enough to preserve the ecosystem but would increase the extinction times (see electronic supplementary material, figure S6). Electronic supplementary

material, figure S6a displays the increase in the extinction times with respect to the non-interventions case in the space (ΔV , freq) as a function of Θ . A particular example of this delayed transition is shown in electronic supplementary material, figure S6b. Here, the orbits can be re-injected multiple times until they escape from the influence of the ghost towards ecosystem collapse. Electronic supplementary material, figure S6c displays several cuts done in the three-dimensional space, as displayed in electronic supplementary material, figure S6a. Generically, these kinds of interventions are less costly than those ensuring full ecosystem preservation, and they could be enough to save time, as a larger intervention or ecosystem change can be performed. In the case of semiarid ecosystems, the livestock policy and an active change in the ecosystem (e.g. increase of the humidity) are changes that may return the system to a lower degradation soil condition.

3.5. Interventions under stochasticity

The previous results have focused on the catastrophic transition towards the desert state found in the mathematical model, which is deterministic. Here, we extend our analyses by considering stochastic fluctuations due to the ecosystem's finite size, which are not gathered by the deterministic model. To consider stochasticity, we performed stochastic simulations using the *Gillespie method* [41,42] and implemented

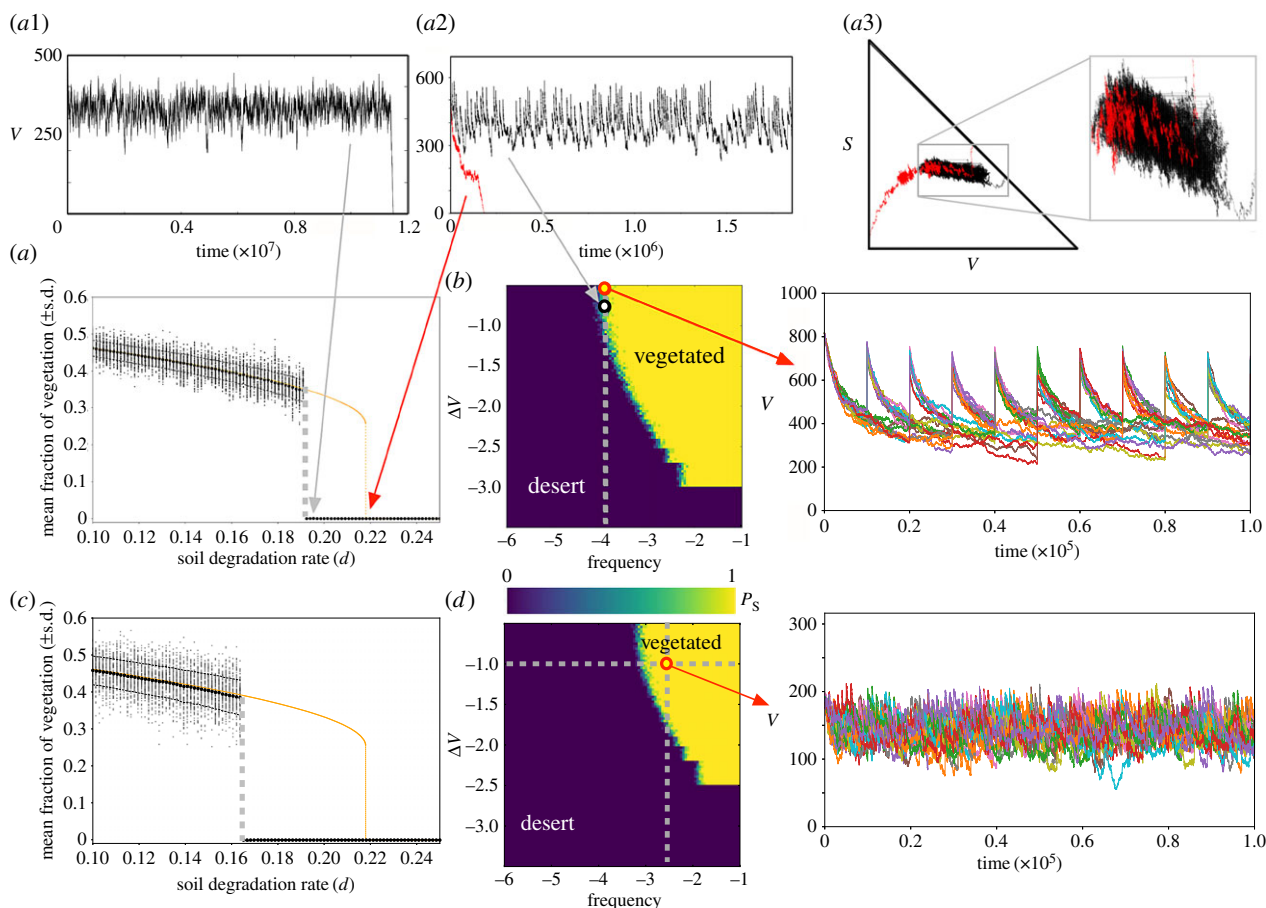


Figure 5. (a,c) Bifurcation diagrams with the deterministic (orange) and stochastic (black, overlapping the values of V at equilibrium for 500 runs) dynamics with the degradation rate of the fertile soil, d , using two system sizes: $\mathbb{S} = 10^3$ (a) and $\mathbb{S} = 10^{2.5}$ (b). Stochasticity involves a lower critical value for d , labelled d_c^s . Panel (a1) shows the delayed transition near to the critical point ($d_c^s < d = 0.196$). An example of the interventions is shown in panel (a2) for $d = 0.22$: the red curve is the non-intervened simulation (vegetation extinction), while the black time series persists using $\Delta V = 10^{0.8}$ and $\text{freq} = 10^{-4}$. The same trajectories are shown in the simplex in panel (a3). (b,d) Survival probability, P_s , in the space of the intervention parameters (freq, ΔV). For each one of these parameter spaces, we show (with red arrows) the sustained dynamics plotting 25 time series. The interventions analyses have been performed for $d = 0.22$, with the other parameters having the same values used in previous figures. In all of the panels, except for the lower three (using $\mathbb{S} = 10^{2.5}$), we have used $\mathbb{S} = 10^3$ (see also electronic supplementary material, figures S7 and S8).

the transition probabilities introduced in §2. Despite our main goal being to test the robustness of the interventions considering demographic fluctuations, we note that the impact of stochasticity on delayed transitions remains poorly explored. Also, we must note that, as our model has an absorbing boundary (i.e. the desert state), the system will always achieve this absorbing state under stochasticity. However, this mean time to extinction is extremely long, typically being of order $\exp(\mathbb{S})$, \mathbb{S} being the system's size or carrying capacity [43], compared to the timescales we are looking at.

Firstly, we computed how the stationary states of vegetation V change as the degradation rate of fertile soil increases under stochasticity. Hereafter, sizes of two different systems will be used: $\mathbb{S} = 10^{2.5}$ and $\mathbb{S} = 10^3$ total sites. The simulations for both sizes indicate that the catastrophic extinction of V will take place before the predictions of the deterministic model (figure 5a,c). As expected, the stochastic critical value of d , labelled d_c^s , becomes lower for smaller values of \mathbb{S} . This is the natural outcome of the so-called noise-induced transitions. In noise-induced transitions, the systems' size (which determines the intensity of the noise) can act as a bifurcation parameter [44]. In our model, noise is decreasing the threshold of the degradation rate of fertile

soil causing the catastrophic transition. Interestingly, the dynamics for $d \gtrsim d_c^s$ also display the characteristic long plateau typical of delayed transitions (figure 5a.1), which now is found just after the stochastic bifurcation. The delay for the stochastic system was quantified by computing the mean extinction times for V for values of d above d_c^s . The results, displayed in electronic supplementary material, figure S7 for both values of \mathbb{S} , indicate that extinction times increase by about two orders of magnitude when comparing $d = 0.3$ to $d \gtrsim d_c^s$. This delaying behaviour is clearly seen in the time series of electronic supplementary material, figure S7. For values of d close to d_c^s (see panels (1) and (4) in electronic supplementary material, figure S7), the stochastic trajectories undergo the plateau, with some of them being extremely long. The delaying behaviour is lost when d increases further away from d_c^s (panels (2, 3) and (5, 6) in electronic supplementary material, figure S7).

The same intervention strategy performed in the previous section was tested with stochastic simulations. Figure 5 indicates that the intervention method is able to sustain vegetation after the catastrophic extinction (yellow areas in panels (b) and (d)). Specifically, we display the probability of survival of the vegetated state, P_s , for different values of the intervention parameters (freq, ΔV), setting $d = 0.22 >$

d_c^s . P_S is computed as the number of replicates surviving at $t = 10^7$ from a total of 50 independent runs. The region where vegetation is sustained becomes larger for larger values of S . The dynamics tied to the interventions are displayed in figure 5 by means of sustained stochastic trajectories (displayed up to $t = 10^5$ but sustained until $t = 10^7$). An example of the stochastic delayed dynamics is displayed in the simplex (V, S) , similarly to the analysis shown in figure 4a for the deterministic model. Finally, several other examples of successful and failed interventions are displayed in electronic supplementary material, figure S8. We must note that strong stochasticity (small S) and being far away from the transition involve a much larger frequency of interventions, increasing the intervention costs for sustaining the vegetation. An example of the former can be seen in panels (b) and (d) of figure 5, as the boundary separating the vegetated state from the desert state is displaced to the upper-right corner of the space (freq, ΔV), meaning that the cost (as defined in the previous section) increases. That is, for larger S , the vegetation can be sustained by using lower frequencies and lower values of ΔV to keep the ecosystem stable at a lower cost.

4. Conclusion

In this study, we have introduced a dynamical restoration approach aimed at protecting drylands facing tipping points. This approach is based on a simple but well-established model for semiarid ecosystems [11]. The model exhibits bistable behaviour, specifically a green state (vegetated) and a desert state, separated by a sharp breakpoint involving a saddle–node bifurcation (first-order phase transition). We have shown that an intrinsic property of this model, namely the presence of a ghost, can be exploited to restore endangered semiarid vegetation by means of small, periodic replantings.

Our model considers three variables: vegetated soil, fertile soil and desert. The presence of a ghost has been identified as a function of soil degradation levels. Because of the apparently stationary nature of transients close to the tipping point, a given system might have already crossed the bifurcation but display vegetation cover until a rapid decline occurs. The shift is irreversible: a decrease in degradation d will not restore previous vegetated states once the desert state has been achieved. As a consequence, prevention strategies should help to avoid collapse, even when the tipping point has already been crossed. Such a goal has been shown to be achievable. We have quantified how these interventions (in terms of newly vegetated areas and frequency of application) might be useful depending on how far the system is placed beyond the bifurcation. Moreover, this intervention strategy has been tested successfully under demographic stochastic fluctuations, under which delayed transitions are also found. As expected, stochasticity involves a different bifurcation value, which results in lower critical degradation rates of fertile soil at decreasing systems' sizes. The lower soil degradation rates make the ecosystem more fragile under stochasticity. However, the intervention methodology also has been shown to work under demographic noise. Here, it is important to highlight two aspects. The first is that delayed transitions are extremely local phenomena. That is, the ghost effect rapidly disappears as d

increases beyond d_c . Another important point is that our approach neglects spatial correlations, which are known to be crucial in ecological systems [45–47]. It is known that space usually enlarges transients [21,22]. However, the impact of space in delayed transitions remains unknown for ecological systems as well as for other spatially extended systems exhibiting saddle-node bifurcations.

Another important point is that the success of the interventions might depend on the spatial patterns involved in each habitat. Depending on the local or regional climatic conditions (e.g., moisture, slope and degradation) the vegetation can be homogeneous or organized in gaps, stripes and spots [48–50]. That is, for the same fraction of vegetated area introduced during the interventions, differences in the dynamics could arise depending on where vegetation is replanted and thus on the specific spatial patterns of the existing vegetation at the moment of applying the interventions. These points should be addressed in future research. Finally, a practical intervention will require the right choices, such as regarding the specific biological components used to restore local patches. Restoration ecology has already considered different strategies to speed the recovery of degraded land [51], including explicit consideration of alternative states [52,53]. Drylands have also been studied under the restoration perspective [54], and several strategies have been presented, involving different restoration approaches [55,56]. In our suggested scenario, plants play a central role as islands of fertility [57] along with the rich communities inhabiting soil crusts [57–59]. Both plants and soil crust interact and could be used as potential restoration agents, but we should not discard other, perhaps complementary strategies, such as synthetic biology approaches [60,61] or resource augmentation [55,62] with a local enrichment of soil patches in nutrients and/or moisture. The different methods outlined above should also be considered in potential practical alternatives to our basic intervention approach.

Data accessibility. All the needed methods and results to replicate this article are part of the manuscript and the electronic supplementary material.

Author's contributions. R.S. and B.V. designed the mean field model. B.V. and J.S. analysed the mathematical model and processed the data. B.V. performed the stochastic simulations. All authors wrote the manuscript and gave final approval for publication.

Competing interests. We have no competing interests.

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