ORIGINAL ARTICLE

Extremely high-efficiency and ultrasimplified hybrid white organic light-emitting diodes exploiting double multifunctional blue emitting layers

Bai-Quan Liu[1], Lei Wang[1], Dong-Yu Gao[2], Jian-Hua Zou[1], Hong-Long Ning[1], Jun-Biao Peng[1] and Yong Cao[1]

Numerous hybrid white organic light-emitting diodes (WOLEDs) have recently been developed. However, their efficiency is not comparable to that of their best all-phosphorescent WOLED counterparts, and the structures are usually complicated, restricting their further development. Herein, a novel concept is used to achieve a hybrid WOLED, whose crucial feature is the exploitation of double multifunctional blue emitting layers. The three-organic-layer WOLED exhibits a total efficiency of 89.3 and 65.1 lm W^{-1} at 100 and 1000 cd m^{-2}, respectively, making it the most efficient hybrid WOLED reported in the literature so far. Significantly, the efficiencies of hybrid WOLEDs have, for the first time, been demonstrated to be comparable to those of the best all-phosphorescent WOLEDs. In addition, the device exhibits the lowest voltages among hybrid WOLEDs (i.e., 2.4, 2.7 and 3.1 V for 1, 100 and 1000 cd m^{-2}, respectively). Such remarkable performance achieved from such an ultrasimplified structure opens a new path toward low-cost commercialization.

INTRODUCTION

White organic light-emitting diodes (WOLEDs) have aroused both academic and industrial interest due to their promising applications in displays and lighting[1–3]. Compared with all-fluorophor or all-phosphor WOLEDs, hybrid WOLEDs—a term that generally refers to devices utilizing blue fluorescent emitters and orange/green-red phosphorescent emitters—are supposed to be more promising for commercialization owing to their extraordinary characteristics, such as good color stability, excellent lifetime, very-high color rendering index (CRI) as well as high efficiency[4–20]. Because all excitons (i.e., singlets and triplets) can be well managed by dint of efficient structures, hybrid WOLEDs with the theoretical 100% internal quantum efficiency can be realized, leading to high efficiency[4–6]. In addition, because hybrid WOLEDs avoid using blue phosphors, which are detrimental to color stability, stable colors can be obtained[7–9]. Moreover, various blue fluorophors exhibit deep blue colors, making it easy to construct very-high-CRI (≥90) hybrid WOLEDs[10–12]. Finally, there is no perfect blue phosphor in terms of lifetime[13], whereas hybrid WOLEDs have been demonstrated to show long lifetimes ($>3 \times 10^5$ h at 1000 cd m^{-2}) because of the introduction of stable blue fluorophors[14].

Forrest *et al.*[4] built a hybrid WOLED by inserting interlayers between fluorescent and phosphorescent emitting layers (EMLs), which can prohibit the unfavorable exciton quenching as well as energy transfer, achieving a maximum forward-viewing power efficiency (PE) of 22.1 lm W^{-1}. Later, Leo *et al.*[15] reported the

triplet-harvesting concept to fabricate hybrid WOLEDs, in which blue fluorophors possess high triplet energies. In their device, the interlayer is unnecessary and the fluorescent emitters 4P-NPD (*N,N'*-di-1-naphthalenyl-*N,N'*-diphenyl-[1,1':4',1'':4'',1'''-quaterphenyl]-4,4'''-diamine) can fulfill the requirements, attaining a PE of 33.9 lm W^{-1} at 100 cd m^{-2}. To further simplify the structure, Zhang *et al.*[16] realized a single-EML hybrid WOLED using a fluorescent material as both the host of the orange guest and the sky-blue emitter. With a low-doping guest concentration (0.1%), the generated singlets and triplets in the device were separated to be well harnessed by the host and guest, respectively, achieving the efficiency of 39.5 lm W^{-1}. More recently, Ma *et al.*[17] organized the devcie using 4P-NPD as a blue guest, suppressing the mutual quenching between the fluorophor and phosphors due to the bipolar mixed-host. Thus, the interlayer was averted and a maximum efficiency of 41.7 lm W^{-1} was recorded. According to the above mentioned findings, the preparation of hybrid structures has been demonstrated to be an effective approach to realize WOLEDs.

From the point of view of real commercialization, structures should be as simplified as possible, which can shorten fabrication procedures, reduce structural heterogeneities, decrease voltages and lower cost. However, the structures of these state-of-the-art hybrid WOLEDs are complicated[4,15–17]. Even for the two-color, single-EML WOLED, four organic layers are needed[16]. Another drawback of single-EML hybrid WOLEDs is the too low guest concentration (0.1%), making it very

[1]Institute of Polymer Optoelectronic Materials and Devices, State Key Laboratory of Luminescent Materials and Devices, South China University of Technology, Guangzhou 510640, China and [2]New Vision Opto-Electronic Technology Co., Ltd, Guangzhou 510530, China
Correspondence: JH Zou, Email: zou1007@gmail.com; JB Peng, Email: psjbpeng@scut.edu.cn

hard to the controlling as well as the reproduction of coevaporation procedures[18,19]. In addition, the efficiency of these state-of-the-art hybrid WOLEDs is not sufficient. Despite the p-i-n method[15] or n-doping technology with an additional electron injection layer[17], the maximum efficiency is only ~ 40 lm W^{-1}, indicating that no hybrid WOLEDs can show comparable efficiencies to their best all-phosphorescent WOLEDs counterparts[21–27]. Moreover, the efficiency roll-off problem restricts the evolution of hybrid WOLEDs. For example, although bipolar hosts have been utilized, the PE at 1000 cd m^{-2} is only 19.7 lm W^{-1}, less than half of the maximum efficiency[16]. Therefore, effort is needed to alleviate these difficulties.

Herein, we propose a novel concept to organize hybrid WOLEDs. The crucial feature is the exploitation of double multifunctional blue EMLs, which can simplify the structure, guarantee the blue emission, boost the efficiency and reduce the efficiency roll-off. One of the blue EMLs—N,N′-di(naphthalene-1-yl)-N,N′-diphenyl-benzidine (NPB)—functions as a bridge, allowing a portion of electrons to flexibly pass through to reach the orange EML and then form excitons with holes at the orange guest molecules, broadening both the charges and excitons distribution. In addition, the other blue EML—bis[2-(2-hydroxyphenyl)-pyridine] beryllium (Bepp$_2$)—also functions as an electron transporting layer (ETL), reducing the heterojunction between ETL and EML. Using this novel concept, all generated excitons can be effectively harvested.

MATERIALS AND METHODS
Device structure of the hybrid WOLED
As displayed in Figure 1, the optimized WOLED (W1) has the structure: indium tin oxide (ITO)/1,4,5,8,9,11-hexaazatriphenylene hexacarbonitrile (HAT-CN, 100 nm)/NPB (20 nm)/NPB: bis(2-phenyl-4,5-dimethylpyridinato)[2-(biphenyl-3-yl)pyridinato] iridium(III) [Ir(dmppy)$_2$(dpp), 35 nm, 15%]/NPB (4.5 nm)/Bepp$_2$ (35 nm)/LiF

(1 nm)/Al (200 nm), in which ITO acts as the anode, HAT-CN acts as the hole injection layer, Ir(dmppy)$_2$(dpp) acts as the orange phosphor, B-E-I and B-E-II denote the 4.5 nm NPB blue EML and Bepp$_2$ blue EML, respectively. All materials were commercially purchased, and no further purification was performed. Without breaking the vacuum at a base pressure of 2.5×10^{-6} Torr, each layer was thermally deposited. Utilizing correspondingly independent quartz crystal oscillators, both hosts and guests deposition rates were recorded for the doping process. Epoxy glue and glass lids were used to encapsulate the fabricated devices immediately under a nitrogen atmosphere. All emission areas were 3×3 mm^2. The Commission International de I'Eclairage (CIE) coordinates, electroluminescent (EL) spectra, CRIs and color-correlated temperatures (CCTs) of the prepared WOLEDs were measured via the Konica Minolta CS2000 spectra system (Konica Minolta Company, Osaka, Japan). The current density (J)-voltage-luminance characteristics were measured via the computer-controlled source meter (Keithley 2400, Tektronix Company, Cleveland, OH, USA) as well as multimeter (Keithley 2000) with a calibrated silicon photodiode. The total efficiencies obtained in this work were calculated according to a previously described method[4,15,16].

RESULTS AND DISCUSSION
As depicted in Figure 2a, W1 possesses the peak forward-viewing current efficiency (CE) of 49.8 cd A^{-1} and the highest PE of 62.5 lm W^{-1}. Because the total emitted power is typically used to characterize illumination sources[4,15,16], W1 has the highest total PE of 106.3 lm W^{-1}. At 100 and 1000 cd m^{-2}, the forward-viewing PE are 52.5 and 38.3 lm W^{-1}, corresponding to the total efficiency of 89.3 and 65.1 lm W^{-1}, respectively, indicating that the efficiency at high luminance has already overtaken fluorescent lamps (with typical PE of 40–70 lm W^{-1}). To the best of our knowledge, this is the most
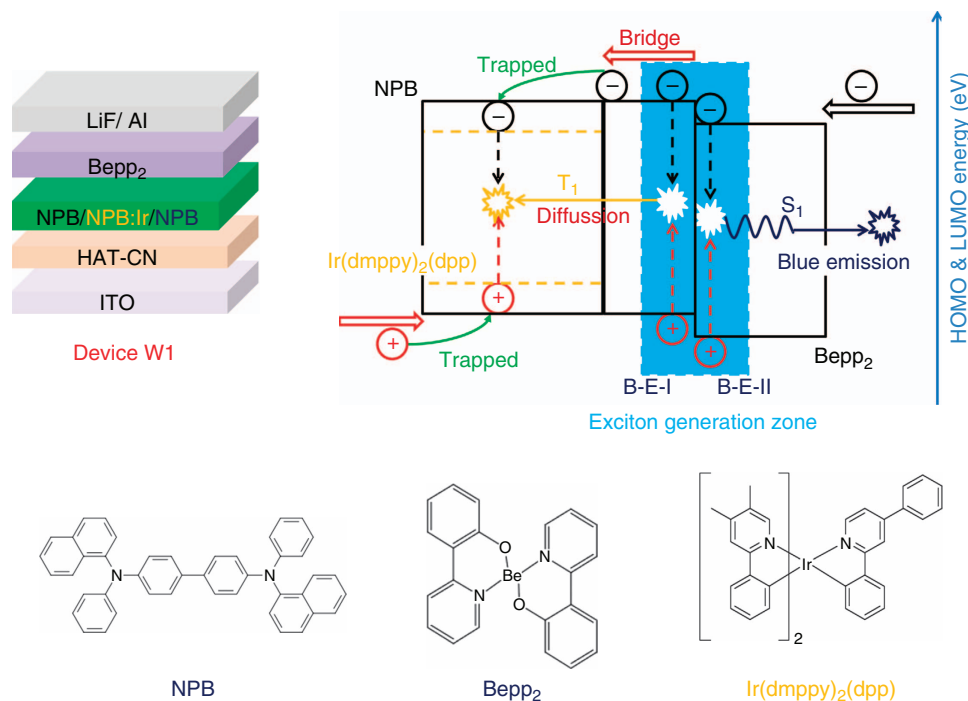

Figure 1 Top: The structure of W1 and schematic diagrams of the working mechanisms. The sky-blue filled region is the main exciton generation zone. Owing to the bridge-like function of B-E-I, a portion of excitons can also be generated at the orange EML via the charge trapping process. Bottom: chemical structures of emitters.

efficient hybrid WOLED yet reported in the literature. Even at the display-relevant luminance of 100 cd m^{-2}, the efficiency (52.5 lm W^{-1}) is much higher than the maximum efficiency of state-of-the-art hybrid WOLEDs[4,15–17], a significant step in hybrid WOLEDs. High efficiencies can be achieved at high luminances, which implies that the efficiency roll-off is greatly reduced[4,15–17]. Significantly, the efficiencies of hybrid WOLEDs, for the first time, have been demonstrated to be comparable to those of best all-phosphorescent WOLEDs[20–26]. However, even the best WOLEDs[21–28] are not as simplified as our device, indicating a significant breakthrough in the simplified WOLEDs.

The CIE coordinates are (0.45, 0.49) and (0.43, 0.48) at 100 and 1000 cd m^{-2}, respectively, and the color-correlated temperatures are 3447 and 3554 K at 100 and 1000 cd m^{-2}, respectively. Similar to previous hybrid WOLEDs, the colors are located within the acceptable warm-white range[10,15–17,29]. Because of the two-color system, the CRI is not high enough (i.e., 47 at 1000 cd m^{-2}). However, they are still higher than that of sodium lamps (CRI = 20) and some two-color WOLEDs[5,18,19] and are comparable to those of mercury lamps. Because recent works reported that two-color devices with CRI ⩾ 80 can meet the call for high-quality lighting[19,30], we believe that high-efficiency, high-CRI and simplified WOLEDs can be realized if more advanced materials are utilized. Jou *et al.*[31] reported that candle-emission OLEDs with yellowish orange light are more physiologically friendly at night due to the low blue intensity. Clearly, W1 can fulfill this requirement, further demonstrating the advantage of this novel WOLED.

As displayed in Figure 2b, W1 exhibits very low driving voltages. The turn-on voltage is 2.4 V (for 1 cdm^{-2}). The voltages are 2.7 and 3.1 V at 100 and 1000 cd m^{-2}, respectively. To our knowledge, these voltages are the lowest among hybrid WOLEDs[4–20,32]. It can be expected that lower voltage as well as higher efficiency may be obtained if p-i-n structures are utilized[15].

To gain insight into the working mechanism of W1, it is necessary to understand the effect of double blue multifunctional EMLs. Because NPB and Bepp$_2$ are *p*-type and *n*-type materials, respectively, coupled with the existing energy barriers between these two materials, the main exciton generation zone can be inferred to be located at the B-E-I/B-E-II

interface, leading to singlets and triplets (with a ratio of 1:3) being simultaneously generated at the interface[5], as shown in Figure 1. Singlets are harvested by NPB and Bepp$_2$, whereas the unused triplets are diffused to the orange region and then harvested by the orange guest because the diffusion lengths of singlets and triplets are ~ 3 and 100 nm, respectively[4]. Thus, all generated excitons at the main exciton generation zone are consumed for light. For these emission processes, the triplet energies (T_1) of NPB and Bepp$_2$ are 2.3 and 2.6 eV, respectively[32,33], which are higher to meet the call for the emitter Ir (dmppy)$_2$(dpp) ($T_1 < 2.25$ eV)[34], ensuring that triplets can be effectively used. B-E-I confines the majority of electrons injected from the cathode to meet holes at the B-E-I/B-E-II interface, ensuring excitons are mainly generated at this interface; otherwise no blue emission can be furnished. Then, B-E-I functions as a bridge because it can allow the unconfined electrons to flexibly pass through B-E-I to reach the orange EML and then form excitons with holes injected from the anode at the orange guest molecules, resulting in high efficiency and reduced efficiency roll-off. By contrast, B-E-II functions as an ETL because of the high electron mobility of Bepp$_2$ (10^{-4} cm^2/V s)[32], which can simplify the structures and reduce the heterojunction between the EML and ETL, enhancing the performance[3]. Moreover, by dint of double multifunctional blue EMLs, the exciton generation zone is greatly widened, which is beneficial to the high performance[8]. Apart from the above factors, since the hole mobility and T_1 of NPB are high enough[33], there is no need to adopt other materials as the hole transport layer (HTL) and host of the orange emitter. Thus, the structure is dramatically simplified (i.e., only three organic layers) and the heterojunctions (between HTL and orange EML or between the orange EML and B-E-I) are eliminated, further enhancing the performance. Therefore, although <25% of the singlets are consumed for the blue emission due to the bridge-like NPB blue EML, all excitons can be harvested via the smart device engineering mentioned above, leading to an extremely high-efficiency hybrid WOLED with ultrasimplified structure.

To further understand the significance of B-E-I, we developed a device (W21) without B-E-I, with the other layers similar to those of W1. For W21, no blue emission was observed, as shown in Figure 3a.
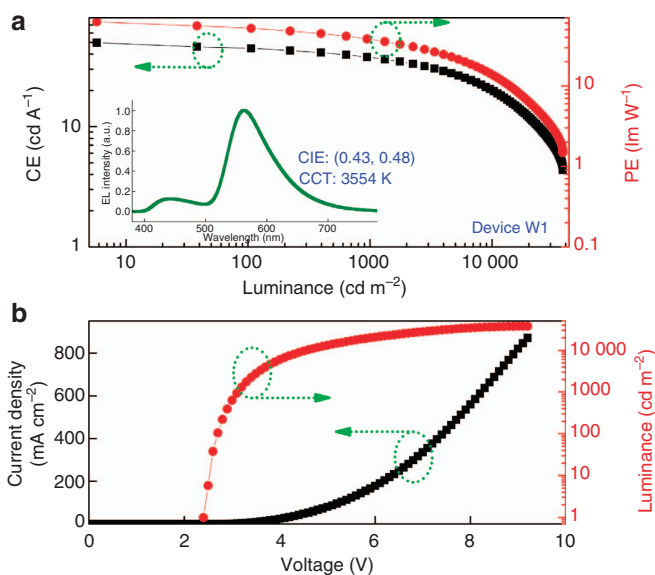

Figure 2 (**a**) Forward-viewing CE as well as PE of W1. Inset: the spectrum of W1 (1000 cd m^{-2}). (**b**) J-voltage-luminance curves.
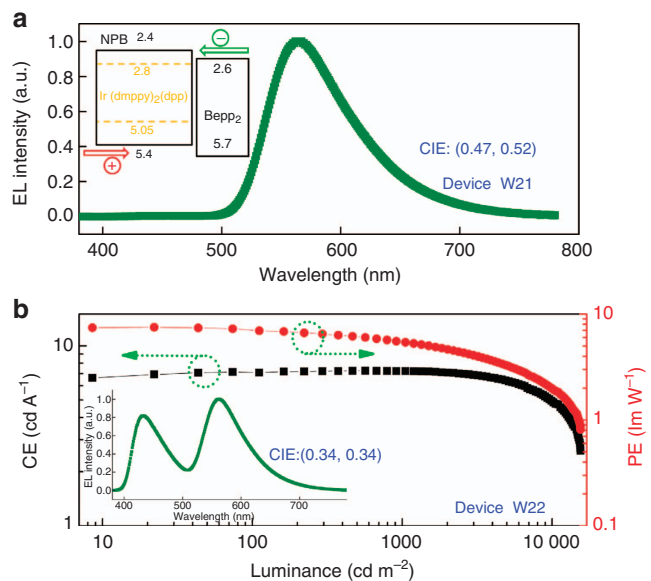

Figure 3 (**a**) the EL spectrum of W21 (1000 cd m^{-2}). (**b**) CE and PE for W22. Inset: the EL spectrum of W22 (1000 cd m^{-2}).

Without B-E-I, the main exciton generation zone is located at the orange EML/Bepp$_2$ interface. Because both singlet energy (S_1) and T_1 of Ir(dmppy)$_2$(dpp) are much lower compared with Bepp$_2$, both singlet and triplet excitons are easily utilized by this orange emitter, leading to only orange emission[3].

NPB and Bepp$_2$ exhibit similar peak emissions (~440 nm) due to the similar S_1[32,33], which indicates that the blue emission is produced by them. To prove that the blue emission originates from B-E-I and to illustrate the significance of Bepp$_2$ functioning as the B-E-II/ETL, we developed a device (W22) by replacing Bepp$_2$ with 35 nm 2,2′,2″-(1,3,5-benzinetriyl)-tris(1-phenyl-1-H-benzimidazole) (TPBi); the other layers were similar to those of W1. TPBi is a well-known electron transport material with a high T_1 (2.74 eV) and almost the same electron mobility as Bepp$_2$, whose lowest unoccupied molecular orbital (LUMO) and highest occupied molecular orbital (HOMO) are 2.7 and 6.2 eV, respectively[35]. By using TPBi, the blue emission originating from Bepp$_2$ was excluded. As expected, white emission was still achieved by combining the blue emission from NPB (~440 nm) with the orange emission from Ir(dmppy)$_2$(dpp), as shown in Figure 3b (inset). However, the maximum PE is only 7.5 lm W^{-1}, which is 8.3 times lower than that of W1, indicating the significance of Bepp$_2$, as shown in Figure 3b.

To prove that Bepp$_2$ contributes blue emission in W1, we developed a device (W3) by replacing B-E-I with 4.5 nm 1-bis[4-[N,N-di(4-tolyl)amino]phenyl]-cyclohexane (TAPC) to exclude the NPB emission; the other layers were similar to those of W1. Because TAPC is a typical p-type material with a T_1 of 2.9 eV and a hole mobility of 10^{-2} cm^2 V^{-1} s whose LUMO and HOMO are 1.8 and 5.3 eV, respectively[32], the exciton generation zone is located at the TPAC/Bepp$_2$ interface. For device W3, it is difficult for electrons to pass through TAPC due to the large LUMO barrier between TAPC and Bepp$_2$ (0.8 eV) as well as the negligible electron mobility of TAPC, and it is difficult for triplets to diffuse to the orange EML due to the high T_1 [31]. Therefore, TAPC acts as an electron- and exciton-blocking layer, which can be used to observe the Bepp$_2$ emission. As expected, no white emission was observed, because only blue emission (~440 nm) originating from Bepp$_2$ and very weak orange emission were produced, as shown in Figure 4 (inset). Because most triplets decay nonradiatively and singlets are harvested only by Bepp$_2$, the maximum efficiency of W3 is only 2.2 lm W^{-1}, as shown in Figure 4. These results also demonstrate the importance of B-E-I.

Finally, to illustrate the bridge-like function of B-E-I, we developed a device (W4) by replacing B-E-I with 4.5 nm 4P-NPD; the other layers were similar to those of W1. Because 4P-NPD exhibits blue emission, high hole mobility and high T_1 (2.3 eV), the exciton generation zone is located at the 4P-NPD/Bepp$_2$ interface[5]. Thus, double blue EMLs are also formed in W4. However, the electron mobility of 4P-NPD is negligible, indicating that it is difficult for electrons to pass through the 4.5 nm 4P-NPD layer to reach the orange region[5]. Therefore, 4P-NPD cannot function as a bridge. Compared with 4P-NPD, NPB exhibits somewhat bipolar characteristics because it can even be an ETL[36]. This bipolar characteristic makes NPB act as a bridge, rendering it possible for electrons to pass through B-E-I to reach the orange EML. Then, Ir(dmppy)$_2$(dpp) can effectively trap these electrons because NPB (2.4 V)[33] possesses a much higher LUMO than Ir(dmppy)$_2$(dpp) (2.8 eV)[34]. Moreover, holes can be effectively trapped by Ir(dmppy)$_2$(dpp) because NPB (5.4 eV)[33] possesses a lower HOMO than Ir(dmppy)$_2$(dpp) (5.05 eV)[34]. Therefore, as shown in Figure 1, the direct exciton formation can also produce the orange emission in the host-guest system [NPB-Ir(dmppy)$_2$(dpp)]. This emission mechanism can effectively boost
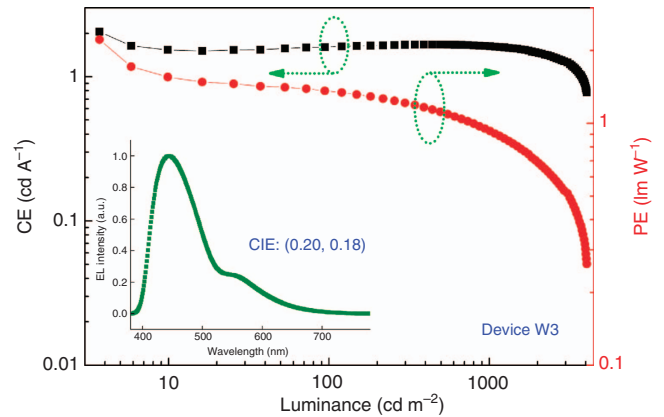
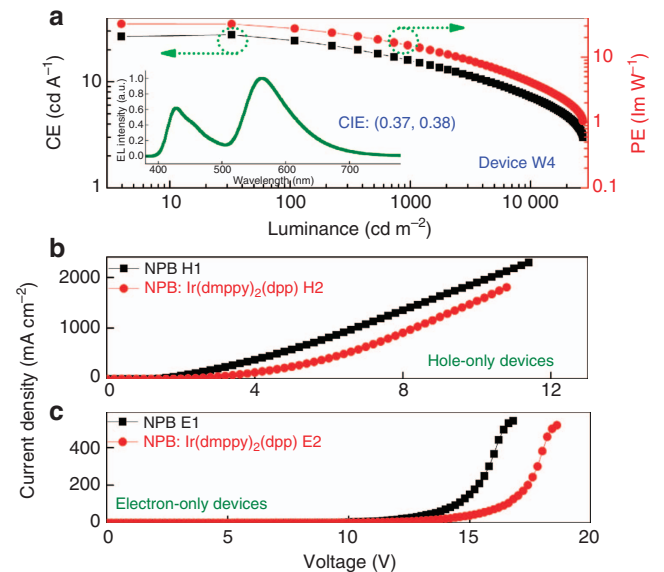**Figure 4** Forward-viewing CE and PE for W3. Inset: the EL spectrum (1000 cd m^{-2}).

Figure 5 (**a**) CE as well as PE for W4. Inset: the spectrum of W4 (1000 cd m^{-2}). (**b**) J-voltage characteristic for hole-only devices. (**c**) J-voltage characteristic for electron-only devices.

efficiency because there is no need to excite the host and the charge leakage is decreased[3]. Furthermore, the charge trapping procedure can reduce efficiency roll-off because trapped charges do not suffer the diffusive transport loss, and the reduced charge/exciton accumulation at the EML edges decreases the triplet-triplet annihilation and triplet-polaron quenching[12]. Thus, both the charges and excitons distribution are effectively broadened by the bridge-like function of B-E-I, enhancing the performance.

To clarify the above analyses, we measured the performance of W4, as shown in Figure 5a. W4 shows the peak efficiency of 32.3 lm W^{-1} as well as 15.0 lm W^{-1} at 1000 cd m^{-2}, which are 93.5% and 152.3% lower than those of W1, respectively. The performances indicate that W4 exhibits a much more serious efficiency roll-off than W1, demonstrating the importance of the bridge-like function of B-E-I. For device W4, because all electrons are confined at the 4P-NPD/Bepp$_2$ interface, a strong exciton density is formed at this interface, which is detrimental to the efficiency and efficiency roll-off[3].

Therefore, although 4P-NPD has a high photoluminescence quantum efficiency of 92%[5], W4 exhibits an unsatisfactory performance.

To illustrate the charge trapping process in the NPB-Ir(dmppy)$_2$ (dpp) system, we have developed hole-only as well as electron-only devices comprising configurations of ITO/HAT-CN (100 nm)/NPB (20 nm)/NPB: guest (35 nm)/NPB (20 nm)/Al (200 nm) and ITO/LiF (1 nm)/Bepp$_2$ (50 nm)/NPB: guest (35 nm)/NPB (4.5 nm)/Bepp$_2$ (35 nm)/LiF (1 nm)/Al (200 nm), respectively, in which guest signifies nothing for devices H1 and E1 and 15% Ir(dmppy)$_2$(dpp) for devices H2 and E2. As displayed in Figure 5b, because of the hole trapping effect of Ir(dmppy)$_2$(dpp), it is reasonable that the J of H2 is much lower than that of H1. In addition, as displayed in Figure 5c, the J of E2 is lower than that of E1 because of the electron trapping effect of Ir (dmppy)$_2$(dpp). Thus, the main exciton generation and emission mechanism in the NPB-Ir(dmppy)$_2$(dpp) system is the direct charge trapping on this guest[3]. In other words, orange emission is obtained from the combined mechanisms of triplets diffusion and direct exciton formation on Ir(dmppy)$_2$(dpp).

Then, to further demonstrate the importance of the novel concept, even higher efficiency can be expected because the photoluminescence efficiency (q) of blue emitters is not high enough[37], and p-i-n methods or outcoupling technologies have not been adopted in this work[12,22]. Supposing the Lambertian type is the emission pattern, a connection of the external quantum efficiency (η_{ext}) and PE can be described as:[3]

$$PE \propto \frac{\eta_{ext}}{U} \quad (1)$$

where U is the operational voltage. η_{ext} can be described as:

$$\eta_{ext} = \eta_{out} \cdot r \cdot q \cdot \gamma \quad (2)$$

where η_{out}, r and γ represent a outcoupling factor, a ratio of excitons which potentially decay radiatively and a charge balance, respectively. As a result, PE can be described as:

$$PE \propto \frac{\eta_{out} \cdot r \cdot q \cdot \gamma}{U} \quad (3)$$

To further boost the efficiency, U should be low while η_{out}, r, q and γ high. Hence, much higher efficiency can be expected if more advanced blue emitters can be obtained. Besides, the efficiency can be improved if p-i-n methods or outcoupling technologies are used[12,22].

CONCLUSIONS

In conclusion, by smartly exploiting double multifunctional blue EMLs, the novel three-organic-layer device shows the highest efficiencies and lowest voltages among hybrid WOLEDs, despite not using p-i-n technology. Previously, hybrid WOLEDs have demonstrated stable colors, very-high CRIs and long lifetimes. Here, we have demonstrated that extremely high efficiency and low voltage can be achieved via ultrasimplified structures. The efficiencies of the hybrid WOLEDs have, for the first time, been demonstrated to be comparable to those of the best all-phosphorescent WOLEDs. Remarkably, the structure is the most simplified among the best WOLEDs, indicating a significant breakthrough in the simplified WOLEDs. Such exceptionally high performance is obtained via an ultrasimplified structure, which opens a new path toward the low-cost commercialization of WOLEDs in the display and lighting markets.

CONFLICT OF INTEREST

The authors declare no conflict of interest.

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