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Dynamic membrane interactions of antibacterial and antifungal biomolecules, and amyloid peptides, revealed by solid-state NMR spectroscopy*

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Abstract

A variety of biomolecules acting on the cell membrane folds into a biologically active structure in the membrane environment, it is, therefore, important to determine the structures and dynamics of such biomolecules in a membrane environment. While several biophysical techniques are used to obtain low-resolution information, solid-state NMR spectroscopy is one of the most powerful means for determining the structure and dynamics of membrane bound biomolecules such as antibacterial biomolecules and amyloidogenic proteins; unlike X-ray crystallography and solution NMR spectroscopy, applications of solid-state NMR spectroscopy are not limited by noncrystalline, non-soluble nature or molecular size of membrane-associated biomolecules. This review article focuses on the applications of solid-state NMR techniques to study a few selected antibacterial and amyloid peptides. Solid-state NMR studies revealing the membrane inserted bent a-helical structure associated with the hemolytic activity of bee venom melittin and the chemical shift oscillation analysis used to determine the transmembrane structure (with α -helix and 3_{10} helix in the N- and C-termini, respectively) of antibiotic peptide alamethicin are discussed in detail. Oligomerization of an amyloidogenic islet amyloid polypeptide (IAPP, or also known as amylin) resulting from its aggregation in a membrane environment, molecular interactions of the antifungal natural product amphotericin B with ergosterol in lipid bilayers, and the mechanism of lipid raft formation by sphingomyelin studied using solid state NMR methods are also discussed in this review article. This article is part of a Special Issue entitled "Biophysical Exploration of Dynamical Ordering of Biomolecular Systems" edited by Dr. Koichi Kato.

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Keywords

Antibacterial peptide; Antifungal natural product; Amyloidogenic peptide; IAPP; Membrane environment; Lipid raft; Solid-state NMR

1. Introduction

A variety of biomolecules located and functioning in the membrane that surrounds living cells fold into a biologically active structure in a membrane environment. The cell membrane is in the liquid crystalline state which is an orientationally ordered state and some biomolecules form biologically active structures by interacting with the cell membrane. For example, antimicrobial peptides do not form structures in solution, while they form a structure after they interact with a membrane. Several models (the aggregate, the barrelstave, the toroidal pore and the carpet models [1]) have been proposed to explain the mechanism by which interaction with a membrane results in peptide ordering. The initial step in all these models of membrane perturbation is the adsorption of the peptide at the ordered membrane surface. In the aggregate model, the peptide acts similar to a detergent and results in the formation of micelles. In the barrel-stave and toroidal pore models, the peptides adopt a transmembrane orientation. In the barrel-stave model the peptides form a structure similar to that of an ion channel and this mechanism of action is favored for long and hydrophobic peptides. For shorter and positively charged peptides, the toroidal pore mechanism is favored since each peptide is surrounded by phospholipids, which limits electrostatic repulsion between charged peptides [2,3]. Finally, the carpet model is similar to the toroidal pore model with the difference that the peptide sits at the surface of the membrane [4]. In both the toroidal and carpet models, the peptides induce a positive curvature of the membrane. In addition, some of the antimicrobial peptides have been shown to translocate across the lipid bi-layers without inducing significant disruption to the cell membrane [5-7].

The structures and orientations of membrane-associated peptides and proteins in membrane environments can be elucidated using structural and orientational constraints derived from magnetically anisotropic interactions as studied by solid-state NMR spectroscopy [8–14]. One such type of interaction is chemical-shift anisotropy (CSA). Usually, this type of interaction appears as a dynamically averaged chemical-shift pattern from which the orientation of peptides and proteins in a lipid-bilayer environment can be derived [15]. Another interaction affording structural constraint is dipolar interaction between coupled nuclei, which can be combined with chemical-shift constraints to provide detailed structural information regarding membrane-associated peptides and proteins [16]. These constraints provide the dynamic structures and orientations of membrane-associated peptides and proteins [17]. Although complete structures have been determined for only a limited number of compounds, these constraints provide detailed local structure and topology as well as orientation.

Solid-state NMR analysis of biomolecules is normally performed using ¹³C and ¹⁵N nuclei, which exhibit spin-1/2. Solid-state NMR spectra yield enormously broadened signals, with

line widths on the order of 20 kHz. Magic angle spinning (MAS) experiments combined with cross-polarization (CP) and high-power decoupling (CP-MAS) provide high-resolution ¹³C and ¹⁵N NMR signals for selectively or uniformly labeled membrane-bound peptides and proteins. In solid-state NMR, homomiclear and heteronuclear dipolar interactions can be recoupled using various pulse sequences combined with multidimensional NMR techniques under MAS conditions. These experiments enable the correlation of ¹³C and ¹⁵N signals for assignment to amino acid residues, and chemical shifts and distance constraints. These data can then be used to determine the high-resolution structures of membrane-associated peptides and proteins [18,19].

In this review, we describe solid-state NMR methods used to determine the membrane bound structures, orientations and dynamics of antibacterial peptides, antifungal natural products, and amyloid forming peptides. We have focused on the systems that are investigated in the authors' laboratories as examples to illustrate the power of various solidstate NMR techniques to study membrane-bound biomolecules.

2. Solid-state NMR spectroscopy for elucidating the structures and orientations of membrane bound biomolecules

2.1. Chemical shift anisotropy

The orientation of peptides bound to a magnetically aligned lipid bilayer can be determined by analyzing the ¹³C or ¹⁵N chemical shift anisotropy (CSA) tensors as reported in the literature. The use of ¹³C CSA tensors of carbonyl carbons in the peptide chain is illustrated in Fig. 1. When the peptide is completely rigid, 150 ppm of CSA span will be observed in a low-temperature experiment {Fig. 1(a) and (f)). In contrast, when the temperature is increased above the gel-to-liquid crystalline phase transition temperature (Tc), the molecules exhibit a large-amplitude rotational motion about the bilayer normal. In this case, an axially symmetric powder pattern is observed as shown in lecithin-melittin bilayer systems (Fig. 1(b) and (g)) [15,20,21].

Dimyristoyl-phosphatidylcholine (DMPC)-melittin lipid bilayers under excess hydration and static conditions have been shown to spontaneously align in the presence of an external magnetic field [12]. The entire melittin molecule forms an α -helix in the membrane-bound state, and the helical axis rapidly rotates about the bilayer normal with a tilt angle ζ and phase angle γ (Fig. 1A). In that case, the ¹³C NMR signal of the carbonyl carbon in the magnetically oriented lipid bilayer appears at the δ_{\perp} position of the axially symmetric powder pattern (Fig. 1e). However, it must be stressed that the α -helix does not rotate about the helical axis, but rather about the bilayer normal. In this dynamic state, the CSA, $\delta = \delta_{//} - \delta_{\perp}$, can be expressed as [15]

$$\Delta \delta = \frac{3}{2} \sin^2 \zeta \left(\delta_{11} \cos^2 \gamma + \delta_{33} \sin^2 \gamma - \delta_{22} \right) + \left(\delta_{22} - \frac{\delta_{11} + \delta_{33}}{2} \right), \quad (1)$$

where δ_{11} , δ_{22} , and δ_{33} are the principal values of the ¹³C chemical-shift tensor of the carbonyl carbon nucleus in the rigid state. In Eq. (1), δ values oscillate as a function of γ , with an oscillation amplitude of (3/2) sin² ζ , referred to as the chemical-shift oscillation [15]. When an α -helical peptide has a large tilt angle ζ , the δ value changes significantly, as shown in Fig. 2A. Using this property of δ , the tilt angle ζ of the α -helical axis can be determined with respect to the bilayer normal by analyzing the anisotropic ¹³C chemical-shifts of the carbonyl carbons of the consecutive amino acid residues in the α -helícal region (Fig. 2). When a peptide forms an ideal α -helical structure, it can be assumed that the interpeptide plane angle for consecutive peptide planes is 100°. This tilt angle can be accurately obtained by analyzing the root-mean-square deviation (RMSD) for the observed and calculated CSAs using Eq. (2).

$$\text{RMSD} = \left[\sum_{1}^{N} \left\{ \left(\Delta \delta_{obs}\right)_{i} - \left(\Delta \delta_{calc}\right)_{i} \right\}^{2} / N \right]^{\frac{1}{2}} \quad (2)$$

where N is the number of ${}^{13}C$ labeled residues in the α -helical region.

In chemical-shift oscillation analyses, lipid bilayers are not necessarily aligned with the magnetic field.

2.2. Dipolar interaction

Similar information can be obtained by observing ¹⁵N—¹H dipolar interactions in the peptide backbone. Polarization inversion spin exchange at the magic angle (PISEMA), a type of separated-local-field 2D NMR spectroscopy, provides excellent resolution in the dipolar dimension in correlation spectra between 15N chemical-shift values and ¹⁵N–¹H dipolar interaction [22,23]. PISEMA pulse sequence was developed based on the flip-flop Lee-Goldberg (using phase- and frequency-switched) pulse sequence, which is used to spinlock ¹H spins along the magic angle to suppress homonuclear dipolar interactions among the highly abundant proton spins. Sample heating is a concern in PISEMA experiments, particularly when analyzing hydrated samples containing membrane proteins. This problem can be overcome by drastically reducing the radio frequency (RF) power [24–26] or using a resonance coil designed to minimize sample heating [27]. Offset problems encountered in the PISEMA pulse sequence are also overcome by using the broadband PISEMA technique [28].

A characteristic circular pattern wheel, known as a polarity index slant angle (PISA), is observed in PISEMA (correlation of ¹⁵N chemical shift and ¹H–¹⁵N dipolar interactions) 2D NMR spectra for a helical structure of a protein embedded in membrane. The shape of the wheel is sensitive to the tilting angle of the α -helix with respect to the lipid bilayer normal [29–31] (Fig. 3). By examining the PISA wheel, one can evaluate the amino acid residues involved in the α -helix, even if the amino acid sequence is unknown. When the helical axis is parallel to the bilayer normal, all the amide sites have a very similar orientation relative to the direction of the applied magnetic field and therefore, all backbone NH resonances overlap with the same ¹H-¹⁵N dipolar couplings and ¹⁵N chemical-shift frequencies.

Whereas tilting the helix away from the membrane normal breaks the symmetry about the magnetic field, introducing a variation in the orientations of the amide N–H bond vectors relative to the field direction (Fig. 3B). This is manifested in the observed PISEMA spectra as dispersions of both the $^{1}H^{-15}N$ dipolar couplings and the ^{15}N chemical-shift frequencies, as shown in Fig. 3B. Thus, the tilt angle of the helical axis can be accurately determined. PISEMA and related multidimensional solid-state NMR techniques have enabled high-resolution structural studies of membrane proteins from fluid lamellar phase lipid bilayers without the need for crystallization or freezing of the sample. Many studies have reported the dynamic structures of a variety of membrane associated peptides and proteins [9,16,23,29–31], Since the interpretation of helical wheels to the structures of transmembrane domains, in particular, is relatively easy, the use of aligned samples in solid-state NMR spectroscopy is attractive and has the potential to be a high throughput approach with the advent of higher magnetic fields and non-uniform sampling approaches.

2.3. PDSD and DARR 2D MAS NMR spectroscopy

MAS [32] technique seeks to suppress the anisotropic interactions in order to achieve highresolution "solution-like" spectra by a rapid mechanical rotation of the sample at the magic angle (54.74°). As the anisotropic interactions that give rise to the broadening observation in NMR spectra of solid materials exhibit a $(3\cos^2\theta - 1)/2$ dependence, MAS induces a scaling in the anisotropic interaction. When the sample rotation frequency is slower than the anisotropic interaction, an incomplete averaging occurs, which is manifested in the solidstate NMR spectrum as resonances separated (from the isotropic resonance) by multiples of the rotation speed of the sample, which are known as spinning side bands.

Both ¹³C and ¹⁵N nuclei are frequently observed in membrane-associated peptides and proteins. In this case, CP techniques [33,34] can be used to enhance the intensities of signals from less abundant nuclei by transferring the polarization from abundant protons in the sample. Furthermore, combining CP, high-power proton decoupling, and MAS (CP-MAS) [35] enables the acquisition of high-resolution signals for biological molecules as well demonstrated for numerous systems.

The significant improvements in resolution afforded by MAS comes at the expense of suppressing the valuable anisotropic interactions that contain the rich structural and dynamical information about the system under investigation. To address this, a number of methods have been developed that permit the selective reintroduction of these anisotropic interactions to facilitate analysis. Some of the details about these methods and their applications to biological systems can be found in the literature [14,15].

Single amino acid can be identified by, for example, their characteristic chemical shift and connectivity pattern among carbon resonance. As mentioned above, a number of methods are available for reintroducing homonuclear dipolar ¹³C–¹³C couplings. A popular class of experiments relies on the proton-mediated reintroduction of carbon homonuclear dipolar couplings. This approach is usually referred to as longitudinal magnetization transfer by spin diffusion and facilitate zero-quantum flip-flop transition among ¹³C nuclei. In its basic form, this process is known as proton-driven spin diffusion (PDSD) [36]. In PDSD, the decoupling field on the proton channel is switched off (while ¹³C magnetization is stored along the z-

axis) during the spin diffusion mixing time to reintroduce heteronuclear dipolar couplings, thus increasing the efficiency of flip-flop transitions. This reintroduction occurs, however, only to higher order and scales inversely with the spinning frequency. Other methods, such as dipolar-assisted rotational resonance (DARR) [37,38], the phase-alternated recoupling irradiation scheme [39], or mixed rotational and rotary resonance [40] outperform PDSD with increased spinning speeds of the sample, because the heteronuclear dipolar couplings are reintroduced independent of the MAS frequency. However, because the flip-flop transition appears in the correlation terms of the effective Hamiltonian, the spin diffusion efficiency invariably decreases with increasing spinning frequency, although the transfer can still be efficient at very high MAS frequencies [41].

2.4. REDOR

Solid-state NMR allows the determination of inter-nuclear distances quite accurately by measuring recoupled dipolar interactions using re-coupling experiments that are otherwise suppressed by MAS. The experimentally measured interatomic distances can be used to elucidate the three-dimensional (3D) structure of membrane bound biomolecules. The rotational echo double resonance (REDOR) [42] is one such re-coupling MAS technique that has been used to recouple the relatively weak heteronuclear dipolar interactions under MAS by applying a π pulse synchronously with the MAS rotor period. Consequently, the transverse magnetization cannot be refocused completely at the end of the rotor cycle, leading to a reduction in the echo amplitude (Fig. 4a). The extent of the observed reduction in the echo amplitude as a function of the number of rotor periods depends on the strength of the heteronuclear dipolar interaction. This method has extensively been used to determine relatively remote interatomic distances in the range of 2 to 8 Å.

Transverse magnetization which precesses about the static magnetic field due to dipolar interaction under the MAS condition moves back to the same direction at every rotor period because the integration of dipolar frequency, ω_D , over a rotor period is zero. Consequently, the rotational echo signals are refocused at the end of each rotor period (Fig. 4b). When a π pulse is applied to the S nucleus, which is coupled with the I nucleus, it inverts the precession direction of the magnetization of the observed I nucleus within a rotor period. Consequently, the magnetization vector of the I nucleus cannot move back to the same direction after one rotor period, therefore, the amplitude of the echo intensity decreases (Fig. 4a). The extent of the reduction of the rotational echo amplitude yields the interatomic distances. To evaluate the REDOR echo amplitude theoretically, one must consider the averaging precession frequency, $\overline{\omega_D}$, in the presence of a π pulse at the center of the rotor period over one rotor cycle as follows:

$$\overline{\omega_D(\alpha,\beta,t)} = \pm \frac{1}{T_r} \left(\int_0^{\frac{T_r}{2}} \omega_D dt - \int_{\frac{T_r}{2}}^{T_r} \omega_D dt \right) = \pm \frac{D}{T_r} \sqrt{2} \sin 2\beta \sin \alpha, \quad (3)$$

where a is the azimuthal angle and β is the polar angle defined by the IS intemuclear vector with respect to the rotor axis. Therefore, the phase angle, $\Phi(a,\beta)$, for the N_C rotor cycles can be given by

$$\Delta \Phi(\alpha, \beta) = \omega_D(\alpha, \beta) N_C T_r, \quad (4)$$

where T_r is one MAS rotor period and N_C is the number of rotor cycles. Finally, the echo amplitude can be obtained by averaging over orientations as follows:

$$S_f = \frac{1}{2\pi} \int_{\alpha} \int_{\beta} \cos[\Delta \Phi(\alpha, \beta)] d\alpha \sin\beta d\beta \quad (5)$$

Therefore, the normalized echo difference, S/S_0 can be given by:

$$\frac{\Delta S}{S_0} = \frac{S_0 - S_f}{S_0} = 1 - S_f \quad (6)$$

The REDOR echo amplitude can be evaluated more rigorously by using the density matrix approach, and by considering the pulse width in the calculation of the analysis of REDOR [43]. The REDOR based experimental approaches made it possible to elucidate the 3D structure of N-acetyl-Pro-Gly-Phe [43]. It is also important to consider the case where the observed nucleus (I_1) can be coupled to two other heteronuclei $(S_1 \text{ and } S_2)$ [44]. The experimentally measured S/S_0 for such a three-spin system strongly depends not only on I_1 - S_1 and I_1 - S_2 dipolar couplings but also on the S_1 - I_1 - S_2 angle. REDOR dephasing patterns due to the recoupled dipolar couplings between the selectively labeled carbonyl ¹³C in bovine lactpferrampin (LFampinB) and ³¹P in phospholipids were measured by ¹³C-³¹P₃ multi-spin REDOR analysis, and the results revealed that LFampinB is located in the interfacial region of the membrane [45]. A detailed analysis of the dipolar couplings among multiple spins and Boltzmann statistics analysis of REDOR dephasing curves can be found in the literature [46-48]. Overall, the REDOR based experiments are pretty robust and the results are relatively easy to interpret. Since there are very few steps to be optimized, REDOR is an excellent choice for structural studies on peptides or proteins that are selectively labeled with the required heteronuclear isotopes. The most often used isotopes are a combination of ¹³C, ¹⁵N, ¹⁹F and ³¹P.

3. Dynamic structure of antimicrobial peptides in membrane environments

3.1. Structure determination of membrane-bound antimicrobial peptides

There is considerable interest in understanding the mechanism of action of antimicrobial peptides (AMPs) that function by interacting with the lipid bilayer of the cell membrane of bacteria. The bacterial selectivity and potency of an AMP depends on its structural interactions with the lipid bilayer of the bacteria or mammalian cell. Since the physicochemical properties of the lipid bilayer varies with the lipid composition, and the

lipid composition of a bacterial cell membrane is different from that of a mammalian cell membrane, the structural interactions of an AMP can be very different in bacterial vs mammalian membranes. Therefore, while it is important to investigate the high-resolution structures of an AMP in a fluid lipid bilayer that mimics the native membrane, it is highly challenging and such samples are not amenable for most biophysical techniques. Fortunately, solid-state NMR can be used to study the fluid lipid bilayers and the high-resolution structures of AMP.

Structures, topologies, and orientations of a number of membrane-bound peptides in lipidbilayer environments have been investigated recently using solid-state NMR in order to fully understand their biological functions. Solid-state NMR is particularly well suited for elucidating the dynamics, topologies, orientations, and high resolution structures of peptides in the lipid bilayer environment using membrane mimetics such as lipid vesicles, bicelles or lipid nanodiscs/macrodiscs. While several studies have reported valuable structural and functional insights into a variety of AMPs (including magainins, MSI-78, LL-37, pardaxin, protegrin and other designed peptides), the use of chemical shift oscillation analysis in elucidating the detailed structure and orientation of the membrane bound state of antibacterial peptides such as melittin [15,20,21], bonbolitin II [49], lactoferrampin [45] and alamethicin [50] reported in the literature are highlighted below.

3.1.1. Melittin—Melittin, a hexacosapeptide with the primary structure GIGAVLKVLTTGLPALISTIKRKRQQNH₂, is the major component of bee venom [51]. Melittin exhibits potent hemolytic activity [52] in addition to indication of voltagedependent ion conductance across a planar lipid bilayer at low concentration [53]. Melittin also causes selective micellization of lipid bilayers as well as membrane fusion at high concentration [54]. As the temperature is decreased and approaches the gel phase, the membranes break down into small particles [55]. Upon raising the temperature above the Tc, the small particles reform into unilamellar vesicles and ultimately multilamellar vesicles.

Giant vesicles with diameters of ~ 20 μ m were observed upon optical microscopic analysis of melittin-DMPC bilayers at 27.9 °C (Fig. 5B) [20]. When the temperature was lowered to 24.9 °C (Tc = 23 °C for neat DMPC bilayers), the surface of the vesicles became blurred, and dynamic pore formation was visible in microscopic images collected at different exposure times [56]. These vesicles disappeared completely at 22.9 °C. It was thus found that melittin-lecithin bilayers undergo reversible fusion and disruption near the respective Tc.

Static ³¹P dipolar-decoupled (DD)-MAS spectra of melittin-DMPC bilayers hydrated with Tris buffer were recorded at various temperatures (Fig. 5A) [56]. Immediately after the sample was placed in the magnetic field, the ³¹P NMR spectrum of an axially symmetric powder pattern characteristic of the liquid-crystalline phase was recorded at 40 °C. The upper field edge (δ_{\perp}) was more intense than the lower field edge ($\delta_{//}$) as compared with a normal axially symmetric powder pattern. This finding indicates that the DMPC bilayer is partially aligned with the applied magnetic field, with the bilayer plane assuming a parallel orientation by forming elongated vesicles, which are referred to as magnetically oriented vesicle systems (MOVS) [56], When the temperature was lowered to 30 °C, the intensity of

the upper field edge of the powder pattern increased further, leading to a spectrum exhibiting almost complete alignment with the magnetic field. At 25 °C, the axially symmetric powder pattern appeared again. At 20 °C, this spectrum changed to a broad envelope of the powder pattern, with round edges due to the presence of a large-amplitude motion in addition to a rotational motion about the molecular axis associating with lateral diffusion of the lipid molecules. At 10 °C, the isotropic ³¹P NMR signal dominated near 0 ppm because of the isotropic rapid tumbling motion of small particles caused by melittin-induced lysis of the larger vesicles. The same axially symmetric powder patterns reappeared as a result of fusion when the temperature was raised from 10 °C to 25 °C, leading to the formation of larger spherical vesicles (Fig. 5A).

It was shown that melittin forms a pseudo-transmembrane amphophilic α -helix [15]. However, melittin can remain in a homogeneous monomeric form in hydrophobic environments when the lipid bilayer enters the lipid-crystalline phase above the Tc (Fig. 6 right). At temperatures close to the Tc, the hydrophilic sides of melittin molecules associate with one another such that the hydrophobic sides face the lipid molecules, thereby causing greater phase separation and partial disordering of the lipid bilayer (Fig. 6 middle). At temperatures below the Tc, a large number of melittin molecules associate with each other. Consequently, small lipid-bilayer particles became surrounded by a belt of melittin molecules and are released from the vesicle, resulting in membrane disruption. Subsequently, entire vesicles dissolve in the buffer solution (Fig. 6 left).

The diether lipid, ditetradecylphosphatidylcholine (DTPC) incorporated by ¹³C labeled melittin, was mechanically aligned between stacked glass plates, then ¹³C NMR spectra were recorded [57,58] as a function of angle between the bilayer planes and the magnetic field at temperatures above and below Tc. An analysis of the ¹³C chemical shifts of the carbonyl resonances of each residue, clearly indicated that the peptide was well aligned in the bilayers and reoriented about the bi-layer normal. The reduced CSA and the chemical shift were consistent with melittin adopting a helical conformation with a trans-bilayer orientation in the lipid membranes, with a greater than 120° angle between the helices in the crystal form [59,60].

Fully hydrated melittin-lipid bilayer also spontaneously aligns along a static magnetic field by forming elongated vesicles, with the long axis parallel to the magnetic field. Using this magnetic orientation property of the membrane, the melittin has been determined to be aligned in a transmembrane fashion (Fig. 1C and (h)) [20].

The secondary structure of melittin bound to a DMPC bilayer can be determined based on conformation-dependent ¹³C chemical shifts referenced to those of model systems [61–63]. The isotropic ¹³C chemical shifts of the [l-¹³C]Gly3, [l-¹³C]Val5, [l-¹³C]Gly12, [l-¹³C]Leul6, and [l-¹³C]Ile20 residues in melittin were determined to be 172.7, 175.2, 171.6, 175.6, and 174.8 ppm, respectively, indicating that all of these residues are involved in the α -helix [15,20,21].

The dynamic structures of melittin bound to dipalmitoylphosphatidylcholine (DPPC), dilauroylphosphatidylcholine (DLPC) and dímyristoylphosphatidylgrycerol (DMPG) were

investigated by analyzing the ¹³C anisotropic and isotropic chemical shifts of selectively ¹³C-labeled carbonyl carbons of melittin under static and MAS conditions [15,21]. Axially symmetric chemical-shift powder patterns were observed for [1-13C]Gly3, [1-13C]Ala4, [1-13C]Val5, [1-13C]Glyl2, [1-13C]Leul6, and [1-13C]Ile20. Analysis of the chemical-shift oscillations (Fig. 7) indicated that melittin molecules form two differently oriented α -helices and diffuse laterally to rotate rapidly around the lipid bilayer normal, with tilt angles of -33° , -36° and -32° for the N-terminal helix, and 21° , 25° and 30° for the C-terminal helix in DLPC, DPPC and DMPG vesicles, respectively. Although the RMSD values could be determined, the tilt angles ζ and $-\zeta$ could not be distinguished because of the symmetry relation associated with Eq. (2). As a result, the REDOR method [42] was used to measure the inter-atomic distance between [1-13C]Val8 and [15N]Leul3 to further identify the bending a-helical structure of melittin. The result provided interhelical angles of 126°, 119° and 118° in DLPC, DPPC and DMPG lipid bilayers, respectively [15,21]. The interhelical angle in the range of 139° to 145° was also reported in DTPC lipid bilayers [58]. These analyses led to the conclusion that the a-helices of melittin molecules penetrate the hydrophobic core of the bilayers incompletely as a pseudo-transmembrane structure, which induces fusion and disruption of the vesicles.

MD simulation for the melittin-DMPG system indicated that the basic residue Lys7 is located slightly within the lipid core region and interacts with one DMPG molecule to disturb the lower surface of the lipid bilayers, and this interaction is associated with melittin's membrane disruption activity (Fig. 7C) [21].

3.1.2. Alamethicin—Alamethicin is a 20-residue antibiotic peptide from *Trichoderma viride* [64]. One of the major amino acid sequences is Ac-Aib-Pro-Aib-Ala-Aib-Ala-Gln-Aib-Val-Aib-Gly-Leu-Aib-Pro-Val-Aib-Gln-Gln-Phol and contains 8 α-amino isobutyric acid (Aib) residues. In addition, the N-terminus is acetylated, and the C-terminus is comprised of an L-phenylalaninol residue [65].

Alamethicin exhibits voltage-dependent ion channel activity in membranes [66]. Alamethicin has a high affinity for lipid bilayers, and binds to the surface of lipid bilayers and inserts into the membrane. Various channel models have been proposed to determine its ion channel activity, such as the barrel stave model [67]. Alamethicin channels are formed by parallel bundles of 3–12 transmembrane helical monomers surrounding a central waterfilled pore [68]. The ion channel activity of alamethicin makes it a suitable model for investigating voltage-dependent ion channel proteins [69].

X-ray crystallographic data indicate that alamethicin assumes an overall helical structure with a kink at Prol4, and the N- and C-termini assume α - and 3₁₀-helical structures, respectively [67]. In solution, NMR studies have shown that the N-terminus forms an α -helix and that the C-terminus can assume a variety of helical structures depending on the conditions [70,71]. Another study revealed that poly-Aib forms a 3₁₀-helix [72].

Two-dimensional separated local-field ¹⁵N NMR spectra were obtained using a motion model of alamethic in which the peptide rotates about the bilayer normal. The spectra revealed an N–H dipolar coupling of 17 kHz, indicating an N–H bond angle of 24° with

respect to the magnetic field (Bo). These data are consistent with an α -helical conformation inserted along the bilayer normal [73].

The conformation of alamethicin in mechanically oriented phospholipid bilayers has been studied using ¹⁵N solid-state NMR in combination with molecular modeling and MD simulations. Alamethicin variants labeled with ¹⁵N at different positions, in combination with substitution of three Aib residues with Ala residues, were examined. Alamethicin with ¹⁵N-labeled Ala6, Val9, and Val15 was incorporated into a phospholipid bilayer at a peptideto-lipid molar ratio of 1:8. Anisotropic 15N chemical shifts and ${}^{1}H{}^{-15}N$ dipolar couplings obtained using this system determined that alamethic n assumes a largely linear α -helical structure that spans the membrane with the molecular axis tilted by $10-20^{\circ}$ relative to the bilayer normal. In comparison, a straight α -helix tilted by 17° and a slightly kinked MD structure tilted by 11° relative to the bilayer normal showed a good compatibility with the structure and orientation of membrane bound alamethicin obtained from solid-state NMR [74]. Measurement of the orientation-dependent ¹H-¹⁵N dipole-dipole coupling, ¹⁵N anisotropic chemical shift, and ²H quadrupole coupling parameters for a single residue, combined with analysis of the anisotropic interaction for the Aib8 residue, provided detailed information regarding the helix-tilt angle, wobble, and oscillatory rotation around the helical axis in the membrane-bound state of alamethicin [75].

Proton-decoupled ¹⁵N solid-state spectra of a sample of alamethicin uniformly labeled with ¹⁵N and reconstituted into oriented palmitoyloleoylphosphatidylcholine (POPC) and DMPC membranes showed that alamethicin adopts a transmembrane orientation [76], Two-dimensional ¹⁵N chemical-shift ¹H–¹⁵N dipolar-coupling solid-state NMR correlation spectroscopy (i.e. PISEMA) analysis suggested that in the transmembrane configuration, alamethicin adapts a mixed $\alpha/3_{10}$ -helical structure with a tilt angle of 8.9° with respect to the bilayer normal [77].

Further detailed structural and orientational analyses of membrane-bound alamethicin have been carried out using solid-state NMR spectroscopy by analyzing the chemical shift oscillation patterns [12,50]. ¹³C chemical-shift interactions were observed in [1-¹³C]-labeled alamethicin. The isotropic chemical-shift values indicated that the entire length of alamethicin forms a helical structure. The CSA of the carbonyl carbon of isotopically labeled alamethicin was also analyzed under the assumption that alamethicin molecules rotate rapidly about the bilayer normal. It was concluded that the adjacent peptide plane forms an angle of 100° or 120° upon assumption of an α - or 3₁₀-helix, respectively. Anisotropic data were acquired for four and seven different residues at the N- and C-termini, respectively. The observed chemical-shift oscillation patterns for the ¹³C CSAs of the carbonyl carbons are shown in Fig. 8A. This chemical-shift oscillation pattern clearly indicated that there are two helices with different tilt angles. RMSD analysis indicated that the N- and C-terminal helical axes tilt 17° and 32°, respectively, relative to the lipid bilayer normal (Fig. 8B). The dihedral angles of adjacent peptide planes of the N- and C-termini were found to be 100° and 120°, respectively. It was, therefore, determined that the N- and C-termini form α - and 3₁₀-helices, respectively (Fig. 8) [12,50].

Hexagonal alamethic in ion channel model is illustrated in Fig. 8C and D, based on the previously obtained structure by MD simulation [74]. The membrane bound structure of alamethic in determined by solid-state NMR analysis is used as components of a hexagonal alamethic oligomer [50]. It is of interest to note that Gln7 and Gln18 residues face to the inside of the ion channel, and thus the inside of the helical bundle shows a hydrophilic character suggesting the high ion conductivity inside the channel.

4. Protein aggregation and amyloid formation in a membrane environment

4.1. Protein aggregation is implicated in devastating neurodegenerative diseases

Protein aggregation is in general an exciting area of research and has implications in many disciplines including biomedicine and bioengineering. Particularly, misfolding and aggregation of a variety of amyloid proteins are implicated in aging-related diseases such as Alzheimer's disease, Parkinson's disease, Huntington's disease, and type-2 diabetes [78-80], Most of these amyloid proteins are typically unstructured as monomers, and are known to aggregate to form oligomeric intermediates, protoflbers and Anally they end up as amyloid fibers. The process of formation of the species is intriguing and varies with experimental conditions and the nature of the peptide or protein [81]. Physiological studies have shown that the process of aggregation of amyloid-beta (typically containing 39, 40 or 42 residues) causes neuronal cell death and memory loss in Alzheimer's disease [82], while the aggregation of islet amyloid polypeptide (IAPP, or also known as amylin) kills insulinproducing islet cells only in the case of type-2 diabetics [83]. Therefore, there is considerable interest in understanding the kinetics of aggregation, physicochemical properties and structures of intermediates formed in the aggregation pathway, and polymorphic structures of amyloid fibers. Numerous low-resolution microscopic and diffraction studies and high-resolution solid-state NMR investigations have provided a wealth of information on amyloid fibers [84], Though most of the amyloid fibers reported are found to contain beta-sheet structures, studies have shown polymorphism in the structure of amyloid fibers. Amyloid fibers have also inspired the development of nano-fibril materials for various potential chemical, biomedical and biotechnological applications. However, for many of these amyloid peptides/proteins, functional studies have demonstrated that oligometric intermediates are the most toxic species and matured fibers are typically not toxic. To better understand the toxic nature of oligomeric intermediates and to develop chemical compounds to avoid their formation, many research groups are currently investigating the preparation and physicochemical characterization, thermodynamic stability and kinetics of formation, high-resolution structure, and heterogeneity of amyloid oligomers.

4.2. Membrane interaction of an amyloid peptide/protein is similar to that of an antimicrobial peptide

Most importantly, the cellular toxicity induced by the amyloid aggregation is mediated or amplified by the lipid components of the cell membrane [85,86]. In fact, the membrane interaction, oligomerization and cell toxicity of an amyloid protein are somewhat similar to that of the well-studied antimicrobial peptides. In addition to the membrane interaction and aggregation similarities between antimicrobial peptides and amyloid proteins/peptides, studies have reported the amyloid fiber formation by antimicrobial peptides and antibacterial

properties of some amyloid proteins. Therefore, the biophysical and biochemical approaches successfully developed to study the mechanism of antimicrobial peptides can be applied to understand the structure and function of amyloid proteins, though amyloid proteins are more challenging.

4.3. Cell membrane catalyzes amyloid aggregation that results in cell toxicity

Because of the link between the cellular toxicity and amyloid aggregation in various aging related amyloid diseases, the role of a lipid membrane on amyloid aggregation must be investigated to fully understand the cellular toxicity of an amyloid peptide or protein; it is remarkable that lipid membrane components dramatically alter most aspects of amyloid aggregation [86-88]. However, these are very difficult tasks to be accomplished, as there are numerous factors that control the formation of oligomers in solution, which include temperature, pH, ionic concentration, metals (such as copper, zinc, aluminum, calcium, and iron), and concentration of protein/peptide [89,90]. In addition to these difficulties, the need for investigating amyloid aggregation in a membrane environment is further complicated by additional variables like the nature of lipids (for example, charge of head group and hydrophobic length of acyl chains), membrane fluidity and curvature, lipid bilayer structure, lipid domains (for example, raft domains), and the presence of other membrane components (for example, cholesterol) [86,87]. Nevertheless, recent studies have made significant contributions in disentangling the many roles of the cell membrane on amyloid aggregation and the formation of toxic species by effectively combining in-vitro and in-vivo studies. The in-vitro studies utilized a combination of various biophysical techniques and membrane mimetics to investigate membrane-protein interactions. Solid-state NMR spectroscopy has elucidated the membrane-induced conformation of $A\beta$, in which the disordered N-terminal segment is followed by the stable C-terminal β -sheet [91,92]. Membrane binding affinity, low and high resolution structures and topologies of membrane-bound amyloid species, and changes in membrane order/disorder and structure are several important aspects reported to date in the literature [86,87,93-100]. NMR experiments revealed the helical structural domains of $A\beta$ upon interaction with lipid bilayers. Zwitterionic lipid bilayers were used in this study to lower the binding affinity and to slower the aggregation in order to trap the membrane-bound non-fibril structure of A β -1–40 [100]. Interestingly, the helical structured membrane-bound region of the peptide is similar to the 3_{10} helical structure determined from a solution sample of A β -1–40 [101]. A recent study demonstrated the effect of membrane thinning by probing the interaction of AB with lipid bilayers varying in hydrophobic thickness [102]. The thinner membrane enabled easy membrane fragmentation induced by the A β peptide; this finding has correlation with the lipid degradation and the formation of fatty acids due to the cell aging of neuronal cells and also due to the physiology of Alzheimer's disease. We refer the readers to recent review articles reported in the literature that comprehensively cover these aspects [86–88] and to several recent publications that reported novel findings [93-100]. The ability to measure amyloid peptide induced membrane disruption by solid-state NMR experiments is briefly discussed below by using IAPP as an example.

4.4. Solid-state NMR reveals the detergent-type membrane disruption by IAPP

Islet amyloid polypeptide (IAPP, also known as amylin) is a 37-residue peptide hormone whose aggregation to form amyloid fibers is implicated in the loss of insulin producing pancreatic beta-cells [83,103,104]. Thioflavin-based fluorescence experiments, circular dichroism (CD) and NMR experiments have been used to discern the aggregation kinetics, structure and the effects of temperature, pH, zinc and insulin [105–116]. Human-IAPP forms "micelle-like" mesoscopic aggregates whereas the rat-IAPP does not [90]. It is also reported that micelle formation correlated with amyloidogenicity of the peptide: the naturallyoccurring C-amidated human-IAPP exhibited about 10 times less CMC than the nonamidated while the human-IAPP-1-19 fragments of both human and rat did not form micelles [109]. Studies have shown that the presence of lipid influences the aggregation of IAPP [86,93–98,117–125]. Both 1–19 and full-length human-IAPP were shown to be toxic to islet cells whereas the 1–19 and full-length rat-IAPP did not exhibit toxicity [116]. A twostep mechanism of membrane disruption has been reported for IAPP: non-fiber structures of IAPP form pores in the first step and fiber-formation dependent "detergent-type" membrane disruption in the second-step [93]. The pore-formation can be suppressed by acidic pH and the fiber formation can be suppressed by insulin, zinc or polyphenolic compounds such as curcumin and EGCG [86,89,90,110]. These studies explain the roles that acidic pH, insulin and zinc play in keeping human-IAPP in a nontoxic mode in non-diabetics. NMR structural studies reported high-resolution three-dimensional structures of membrane-bound IAPP [126–128]. To avoid membrane disruption, detergent micelles were used to trap the helical species of IAPP. To further probe the fiber-formation dependent "detergent-type" membrane disruption by human-IAPP, solid-state NMR experiments were performed on aligned lipid bilayers [128,129]. These NMR studies demonstrated that the C-terminal domain of human-IAPP plays a vital role in fiber formation [108], and thus human-IAPP-20-29 peptide fragment (SNNFGAILSS) which forms fibers [130], was used to demonstrate the membrane disruption process. ³¹P NMR experiments on anode aluminum oxide (AAO) nanotubes revealed the fragmentation of aligned lipid bilayers in the presence of human IAPP-20-29 (Fig. 9) [129]. The appearance of an isotropic ³¹P peak illustrates the formation of small "micelle-like" fragments caused by human-IAPP-20-29. At the same time, these results also demonstrated that a non-fiber forming peptide, rat-IAPP-20-29 (SNNLGPVLPP), does not disrupt membrane [131]. These results demonstrate the unique ability of solid-state NMR experiments to probe "detergent-type" membrane disruption by amyloid peptides (Fig. 10). Studies have shown that the presence of anionic lipids enhance the binding affinity [86], negative curvature strain inducing PE lipids enhanced membrane fragmentation [96], and cholesterol to modulate the membrane mediated aggregation of human-IAPP [97]. The presence of calcium ions has been shown to promote human-IAPP aggregation and membrane fragmentation via a carpet-like mechanism [94]. For more details on the biological properties of IAPP, we refer readers to studies reported in the literature [132– 138].

5. Dynamics of amphotericin B and sphingomyelin in membrane bilayers

Isotope-labeled peptides and proteins suitable for solid-state NMR measurements are readily available because general methods for isotope labeling are well established. On the other

hand, isotope labeling of small molecules such as natural products remains challenging because typically they must be chemically synthesized, and the synthetic route is different for each compound. This makes it more difficult and challenging to apply solid-state NMR to small molecules than to peptides and proteins, and consequently applications of solidstate NMR to membrane-associated natural products remain limited. Thus, studies combining organic synthesis and solid-state NMR are necessary to reveal the dynamic structures of natural products in membrane environments. As examples of such studies, we first present a molecular interaction analysis of amphotericin B with sterol in a hydrated bilaycr, followed by solid-state NMR applications to sphingomyelin, a major component of lipid rafts.

5.1. Amphotericin B

Amphotericin B (AmB) is a clinically important antifungal drug used for over 50 years to treat deep-seated systemic fungal infections [139,140]. Despite its serious side effects, until recently the clinical importance of AmB remained unchanged due to the lack of better alternatives and the rarity of resistant strains. It is generally accepted that AmB forms ion channels with sterols and/or extracts sterols from cell membranes [141,142], leading to cell damage or cell death. Its selective toxicity against fungi over animals is believed to result from its stronger interaction with ergosterol (Erg), an abundant sterol in fungal membranes, than with cholesterol, the major sterol in animal cell membranes [143].

Direct interaction between AmB and sterol, particularly Erg, was verified by ²H NMR measurements of deuterated Erg [144]: the fast axial rotation of deuterated Erg in lipid bilayers was inhibited by the presence of AmB. Insight into the mode of interaction of AmB with Erg were obtained by measuring the intermolecular REDOR between ¹³C-labeled Erg and F-labeled AmB (Fig. 11) [145]. An F nucleus was chemically introduced at the C-14 position of AmB (14-F-AmB) [146], and Erg was ¹³C-labeled by culturing yeast cells in the presence of 2-¹³C-acetate, furnishing skipped ¹³C-labeled Erg (Fig. 11). The REDOR experiments unexpectedly suggested that Amp-Erg complexes interacted not only in parallel but also in anti-parallel with each other in lipid bilayers: the REDOR effect from the F nucleus of AmB to the C-19 in Erg indicated a parallel interaction, while that to the C-26/C-27 or C-21 in the Erg side chain suggested an antiparallel alignment in the lipid bilayer (Fig. 11). Previously proposed AmB-Erg assembly models [141,142] suggested that AmB adopts a parallel interaction, but the REDOR results raise the possibility of a different mode of bimolecular recognition, namely, antiparallel interaction, in the AmB-Erg complex.

To obtain a clearer evidence for the presence of the antiparallel orientation, we chemically synthesized 26,27- $^{13}C_2$ -Erg, in which only the C-26 and C-27 positions were selectively ^{13}C -labeled with 100% enrichment [145]. ^{13}C - ^{19}F REDOR experiments with 14-F-AmB and 26,27- $^{13}C_2$ -Erg demonstrated that the ^{13}C atoms in the Erg side chain closely approach the ^{19}F atom of the AmB head group (Fig. 12), thus further supporting an antiparallel binary interaction between AmB and Erg in lipid bilayers. The REDOR effect is saturated at ca. 30%, which probably indicates that the antiparallel alignment occurs at ca. 30%. The interatomic distance between the F of AmB and C-26/C-27 of Erg was estimated to be 5.5

and 7.0 Å, where the values are interchangeable, by fitting the REDOR plots to theoretical curves (Fig. 12).

Next, to obtain distance information for the parallel orientation, 32-F-AmB was synthesized and intermolecular REDOR measurements between 32-F-AmB and $26,27^{-13}C_2$ -Erg were performed [147]. Interestingly, the REDOR effect plateaued at 70% (Fig. 13 top), consistent with the aforementioned 30% saturation of the REDOR effect for the antiparallel complex. The results indicate the coexistence of parallel and antiparallel orientations for AmB and Erg at a ratio of 7:3. Theoretical curve fitting provided estimates of the F-32/C-26 and F-32/C-27 distances of 5.2 and 6.8 Å, or vice versa (Fig. 13 top).

We further synthesized 4-¹³C-Erg and conducted intermolecular REDOR measurements of the other two combinations (Fig. 13 bottom) [147,148]: 14-F-AmB and 4-¹³C-Erg for the parallel orientation, and 32-F-AmB and 4-¹³C-Erg for the antiparallel orientation. The REDOR data not only supported the coexistence of parallel and antiparallel orientations at a ratio of 7:3, but also provided additional distance information: 6.4 Å for F-14/C-4 and 6.0 Å for F-32/C-4.

Based on these distance data as well as our previous structure-activity relation study [148], a geometry search for AmB and Erg complexes was performed and the most likely geometries for the AmB-Erg complex were obtained for the parallel and antiparallel orientations (Fig. 14] [147]. The structures show that two molecules are in face-to-face van der Waals (VDW) close contact with each other both in the parallel and antiparallel orientations. The face-to-face contact is less effective with cholesterol, which would explain the weaker activity of AmB in cholesterol-containing animal cell membranes.

5.2. Sphingomyelin

The lipid rafts hypothesis was proposed in the late 1990s [149], and its implications for membrane processes such as signal transduction, protein sorting, and viral infection have been revealed [150]. In general, cholesterol (Cliol) and sphingolipids such as sphingomyelin (SM), are considered essential for the formation of lipid rafts [151–153]. In addition, SM is known to be the target of some proteinous toxins such as equinatoxins and lysenin [154,155], and it is also reported that the raft formation affects membrane disruptive activities of AMPs [156]. Thus, the dynamics of SM in lipid bilayers is important for understanding not only the mechanism of raft formation but also its influence on the function of membrane-active proteins and peptides including AMPs.

The structure of SM is characterized by the presence of amide and hydroxy groups, which presumably are responsible for the formation of collateral intra- and intermolecular hydrogen bonds in membranes, making SM a preferable raft constituent compared with phosphatidylcholine. In particular, hydrogen bonds between SM and Choi have drawn the most attention as a key factor in raft formation [157–161]. However, the molecular basis of raft formation including inter-molecular hydrogen bonds, has not been fully elucidated, in part, because few NMR experiments have been conducted due to the difficulty of stable isotopical labeling of SM.

We used three approaches to reveal the structure and dynamics of SM molecules in lipid rafts. The first was conformational analysis of SM using bicelles [162]. This study showed that SM forms bicelles in the presence of a detergent, and the conformation of SM in the small bicelles was revealed using liquid state NMR. The second approach was dynamics analysis of the SM alkyl chains using ²H NMR [163–165]. Seventeen different siteselectively deuterated SMs were prepared and site-specific dynamics of SM was revealed in raft and non-raft membranes. In the third approach, the orientation of the SM amide plane in lipid bilayers was analyzed using dipolar couplings and CSAs [166] as described below.

The orientation of the SM amide plane can be defined by the rotation axis direction and the order parameter *S* (Fig. 15). *S* represents the degree of wobbling of the rotation axis, and varies from 1 to 0. When S = 1, no molecular wobbling with respect to the rotation axis is present, and when S = 0, every orientation is allowed. To determine the rotation axis direction and order parameter, we prepared the ¹³C, ¹⁵N doubly-labeled SMs, l'-¹³C-2-¹⁵N-SM and 2'-¹³C-2-¹⁵N-SM (Fig. 15). The former sample provided l'-¹³C and ¹⁵N CSAs, and the intramolecular l'-¹³C-2-¹⁵N dipole coupling constant in hydrated bilayers with and without Choi, and the latter gave the intramolecular 2'-¹³C-2-¹⁵N dipole coupling. The CSAs and intramolecular dipole couplings were obtained through slow MAS and REDOR experiments, respectively [164].

The rotation axis was defined by the polar coordinates α (azimuthal angle) and β (polar angle) with respect to the ¹³C CSA tensor axes (Fig. 16A). The CSA tensor axis directions were determined by theoretical calculation based on the gauge including atomic orbital (GIAO) method. The CSAs and intramolecular dipole couplings can be calculated using the α , β , and S values, and therefore it is possible to find the combination of α , β , and S that best reproduces the experimental data by minimizing the RMSD. Fig. 16B shows the 2D RMSD plots for SM and SM-Chol membranes at S values of 0.67 and 0.84, respectively, which provide the minimum RMSD values. Using these data, the rotation axis direction and order parameter S of the SM amide in hydrated bilayers were determined, as shown in Fig. 16C. Interestingly, this demonstrated that the SM amide orientation is not affected significantly by the presence or absence of Choi, which suggests that direct hydrogen bonding between the SM amide and Chol may not be as strong as previously proposed. This is also in line with our previous reports that demonstrated that Choi is more deeply positioned in SM than in the PC membrane [163–165], because the formation of a direct hydrogen bond between the SM amide and the OH group of deeply situated Choi should induce some changes in the position and orientation of the amide group. We also noticed that the orientation and conformation around the SM amide is suitable for forming intermolecular hydrogen bonds with neighboring SM amides, as shown in Fig. 16C. Given that the amide group in SM can act as both a hydrogen-bond donor and an acceptor, while the ester groups in phosphatidylcholine function as hydrogen-bond acceptors, it can be assumed that the intermolecular hydrogen-bonds between the SM amide groups, and subsequent formation of a hydrogen-bond network or cluster among SM molecules, is an important molecular basis for raft formation.

What, then, is the role of Chol in raft formation? As demonstrated in this study, Choi significantly enhanced the order parameter S of the SM amide, from 0.67 to 0.84, which

means that Chol restricts rotation axis wobbling of the SM amide. On the other hand, as described above, Chol does not change the amide orientation (Fig. 16C), which suggests that Chol is less likely to directly interact with the SM amide group. It was previously demonstrated that Chol is positioned deep in the SM membrane and significantly restricts the motion of the middle region of the SM alkyl chains [163–165]. Therefore, overall, Chol restricts the dynamic motion of the middle region of SM alkyl chains, which indirectly promotes the formation of hydrogen bonds between SM amide groups. These intermolecular hydrogen bonds would restrict the wobbling of the amide group, and lead to the formation of stable lipid rafts.

6. Conclusion

In this review article, the ability of solid-state NMR spectroscopy to provide biologically meaningful information on the structure and function of membrane bound biomolecules arc highlighted. Although there are several reports in the literature on the structure of melittin, the most important dynamic structure of melittin in a membrane environment is obtained using solid-state NMR experiments that reported a pseudo-transmcmbrane bending a-helical structure and the ability of melittin to disrupt the lipid bilayer. The solid-state determined structure of alamethicin was found to be a transmembrane structure with the N- and Ctermini forming an α -helix and a 3₁₀-helix, respectively, and these helices are associated with the ion channel formation. Solid-state NMR studies demonstrated that the C-terminal domain of human-IAPP plays a vital role in fiber formation, and a human-IAPP-20-29 peptide fragment was used to demonstrate the membrane disruption process, while non-fiber forming rat-IAPP-20-29 does not disrupt the membrane. By using distances measured from solid-state NMR, geometry for the amphotericin B ergosterol complex was found to take both parallel and anti-parallel orientations. The geometry further explained the selective activity of amphotericin B toward an ergosterol-containing fungal membrane over a cholesterol-containing animal cell membrane. The formation of a hydrogen-bond network among sphingomyelin molecules plays an important role in the raft or membrane domain formation. It was also revealed that cholesterol is located deep in the sphingomyelin membrane and significantly restricts the motion of the middle region of sphingomyelin alkyl chains, which lead to the formation of stable lipid rafts. While these are a few selected examples to demonstrate the applications of solid-state NMR techniques, there are many studies that have reported high-resolution structures of membrane proteins that are beyond the scope of this review article. We refer the readers to some of the recent important solidstate NMR approaches for structural and dynamical studies of membrane-associated proteins [167–177]. The advancements in the development of higher field magnets, faster spinning MAS probes, and lipid nanodiscs would significantly overcome the remaining challenges in the membrane arena.

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Fig. 1.

A: Direction of the principal axis of the ¹³C chemical shift tensor of the C=0 group. B: The helical axis rotates about the magnetic field (H₀). C: the helical axis rotates about the bilayer normal. ¹³C NMR spectral patterns of the C=0 carbons corresponding to the orientation of the *a*-helix with respect to the surface of the magnetically oriented lipid bilayers. Simulated spectra were calculated using $\delta_{11} = 241$, $\delta_{22} = 189$, and $\delta_{33} = 96$ ppm for the rigid case (a), rotation about the unique axis without orienting to the magnetic field (slow MAS condition) (b), MAS (c), rotation about the magnetic field (d), and rotation about the bilayer normal (e). ¹³C NMR of [1-¹³C]Ile4-melittin-DMPG bilayer system in the rigid state at – 60 °C (f), in the slow MAS condition at 40 °C (g), in the oriented condition at 40 °C (h), and in the fast MAS condition at 40 °C (i) [12].



Fig. 2.

(A): Chemical-shift oscillation curves of chemical-shift anisotropy, $\delta = \delta_{//} - \delta_{\perp}$ against the phase angle, γ , of the peptide plane for the amino acid residues as a function of the tilt angles, ζ of the *a*-helical axis relative to the bilayer normal (Z' axis) [12]. (B): ¹³C chemical shift powder patterns of the [l-¹³C]Gly3-, [l-¹³C]Ala4-, [l-¹³C]Val5, [1-¹³C] Leu16, [l-¹³C]Ile17, and [l-¹³C]Ile20-melittins-DMPG bilayer systems at 40 °C [21].



Fig. 3.

The generation of "PISA wheels". (A) Definitions of τ and ρ for an α -helix. $\tau = 0^{\circ}$ occurs when the helix axis, h_3 , is parallel to $B_0 \cdot \rho = 0^{\circ}$ occurs when the projection of B_0 onto a plane perpendicular to, h_3 , makes an angle of 0° with h_1 , the radial axis of the helix that passes through the C α carbon of Leu26 of the M2-TMP helix [31]. (B) "Circles" drawn for one of the dipolar transitions rising average values for the ¹⁵N chemical shift tensor elements ($\sigma_{11} = 31.3$, $\sigma_{22} = 55.2$, $\sigma_{33} = 201.8$ ppm) and the relative orientations of the ¹⁵N–¹H dipolar and ¹⁵N chemical shift tensor, given by $\theta = 17^{\circ}$, the angle of the peptide plane between σ_{22} and υ (parallel to the N–H bond). Note that the center of the PISA wheel falls on a line that passes through the ¹⁵N isotropic chemical shift (96 ppm) at 0 kHz on the dipolar scale [31].



Fig. 4.

Pulse sequence to observe rotational echo. (a) Transverse magnetizations do not refocus at the rotor period T_r in the REDOR pulse sequence, (b) Transverse magnetizations refocus at the rotor period T_r as the rotational echo.



Fig. 5.

A. Effect of temperature variation on the ³¹P NMR spectra of melittin-DMPC bilayer systems. Arrows indicate the direction of the temperature variation. H_0 indicates the direction of the static magnetic field. The shapes of the vesicles are depicted to the right of the corresponding spectra [56]. B. Optical microscopic pictures of melittin-DMPC vesicle systems at 22.9, 24.9 and 27.9 °C [56].



Fig. 6.

Schematic representation of the process of morphological changes in melittinlecithin bilayer systems in the absence (A) and presence (B) of an applied magnetic field (H_0) [56].



Fig. 7.

Schematic representation of the structures and orientations of melittin. (A): Bound to DPPC and (B): DLPC lipid bilayers. (Top): Chemical shift oscillation curves for [1-¹³C]Gly3, Ala4, Val5, Leu16, Ile17, and Ile20 of melittin in DPPC (A) and DLPC (B) lipid bilayers. (Bottom): The N- and C-terminal helices of melittin bound to DPPC and DLPC lipid bilayers are inserted into the bilayer with tilt angles of 36° and 25° in DPPC and 33° and 21° in DLPC, respectively [15]. (C): Structure of melittin bound to DMPG bilayers obtained from MD simulation [21].



Fig. 8.

(A): Chemical shift oscillation patterns of alamethicin. Chemical shift oscillation curves were obtained from the chemical shift anisotropies of the N-terminus (AIa6, Gln7, Val9 and Aib10) and the C-terminus (Val15, Aib16, Aib17, and Gln18). The tilt angles of the N- and C-termini were determined to be 17° and 32° , respectively. The dihedral angles of the peptide planes between the n– and n + 1-residues of the N- and C-termini were determined to be α - and 3_{10} -helices, respectively. (B): Structure and topology of alamethicin bound to a DMPC bilayer, as determined from chemical shift oscillation data [50]. Side view (C) and top view (D) of the hexameric oligomer of alamethicin in the membrane environment.



Fig. 9.

Phosphorus-31 chemical shift spectra of aligned DMPC bilayers containing various mol% of IAPP in AAO nanotubes: (A) 0%, (B) 0.5% human-IAPP20–29, (C) 0.75% human-IAPP20–29, (D) 1% human-IAPP20–29, (E) 1% human-IAPP20–29 after incubation for 48 h, (F) 1.25% human-IAPP20–29, (G) 0.5% rat-IAPP20–29, (H) 1% rat-IAPP20–29.



Fig. 10.

A schematic model for membrane fragmentation by human IAPP20–29. (A) At lower concentrations, the peptide aggregates on the lipid bilayer surface to form an intermediate state that can extract phospholipid molecules from the bilayer, leading to peptide–lipid vesicles (B). At higher concentrations, fiber formation follows an alternate pathway bypassing the formation of this intermediate (C).





Structures of 14-F AmB and skipped ¹³C-ergosteroL (top) and intermolecular ¹³C–¹⁹F REDOR spectra of the 1:1 mixture in hydrated bilayers. *S* is a difference spectrum obtained by subtracting a ¹⁹F-irradiated spectrum *S* from a non-irradiated spectrum, S_0 (full echo). Numbers in tile spectra denote the carbon atoms of Erg.



Fig. 12.

Intermolecular ${}^{13}C{-}^{19}F$ REDOR observation between 14-F-AmB and 26,27- ${}^{13}C_2$ -Erg for the antiparallel orientation. The internuclear distances are derived by fitting the REDOR plots to theoretical curves (bottom right).



Fig. 13.

Intermolecular ${}^{13}C{-}^{19}F$ REDOR observation between 32-F-AmB and 26,27- ${}^{13}C_2$ -Erg for the parallel orientation (top), and the other two combinations of REDOR measurements using $4{}^{13}C$ -Erg (bottom).



Fig. 14.

Structures of the most stable complexes of AmB and Erg for the parallel and antiparallel orientations.



Fig. 15.

Schematic summary of orientation analysis of the SM amide plane in lipid bi-layers. The orientation of the SM amide, which is defined by the rotation axis direction and the order parameter S, was analyzed using CSAs and dipole coupling constants, both of which were derived from synthetic ¹³C and ¹⁵N doubly labeled sphingomyelins.

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Fig. 16.

(A) The rotation-axis direction is defined by the polar coordinates *a* (azimuthal angle) and β (polar angle) with respect to the ¹³C CSA tensor axes. The oxygen atom on C-l' was omitted for clarity. (B) 2D RMSD contour plots for SM at S = 0.67 and SM-Chol at S = 0.84. (C) The rotation-axis direction of SM in hydrated membranes. The orange axes indicate the rotation axes.