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## Evaluating Separation Selectivity and Collision Cross Section Measurement Reproducibility in Helium, Nitrogen, Argon, and Carbon Dioxide Drift Gases for Drift Tube Ion Mobility-Mass Spectrometry

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## Abstract

Previous ion mobility (IM) studies have demonstrated that varying the drift gas composition can be used to enhance chemical selectivity and resolution, yet there are few drift gas studies aimed at achieving quantitatively reproducible mobility measurements.. Here, we critically evaluate the conditions necessary to achieve reproducible collision cross section (CCS) measurements in pure drift gases (helium, nitrogen, argon, and carbon dioxide) using a commercial uniform field drift tube instrument. Optimal experimental parameters are assessed based on convergence of CCS measurements to reproducible values which compare with literature values. A suite of calibration standards with diverse masses, biological classes, and charge states are examined to assess chemical selectivity and resolution achievable in each drift gas. Results indicate nitrogen and argon perform similarly and are sufficient for most applications where high resolving power and high peak capacity are desired. Carbon dioxide exhibits more selectivity for resolving structurallyheterogeneous compounds, which may be preferable in specific analyte pair separations. Helium demonstrated modest separation capabilities but has utility for comparison to theoretical values and previously published work. In drift gases other than nitrogen, pressure differentials up to 230 mTorr between the drift tube and upstream chamber were optimal for improving correlation to literature values, while in nitrogen, the recommended pressure differential of 150 mTorr was found appropriate. We present recommended experimental parameters as well as gas-specific CCS measurements for structurally-homogeneous sets of analytes which are suitable for use by other laboratories as standards for purposes of instrument calibration and overall assessment of IM separation performance.

## **Graphical Abstract**



#### Keywords

Uniform Field Ion Mobility; Alternative Drift Gases; Gas Polarization Effects; Single Peak Resolving Power; 2-Peak Resolution; Peak Capacity; Phosphazenes; Quaternary Ammoniums; Carbohydrates; Polyalanines

## Introduction

Ion mobility spectrometry (IM) is an important analytical technique for the rapid separation and identification of a wide variety of chemical compounds [1-6]. In a conventional drift tube, the separation of chemical species results from numerous, near-thermal collision between analyte ions and a chemically-inert drift gas. Total IM separation times occur on the order of milliseconds, and thus are both high throughput and sufficiently fast to allow for integration with other analytical separation techniques [7, 8]. Whereas in conventional gas and liquid phase chromatography, it is well-established that the chemical selectivity of the separation can be enhanced by choosing different stationary phases, IM is more appropriately described as a gas-phase electrophoretic technique, and as such does not incorporate a stationary phase. Instead, varying the drift gas composition has been demonstrated to affect separation efficiency in a similar manner as changing solvent conditions in capillary electrophoresis to tune separation selectivity. Most notably in highfield IM techniques such as high-field asymmetric waveform ion mobility spectrometry (FAIMS) and differential mobility spectrometry (DMS), the carrier gas composition is commonly altered to increase the resolution for closely spaced analytes [9-15]. For conventional low field IM techniques such as drift tube (DTIMS) and traveling wave (TWIMS), this practice is less common, although several notable cases exist [16-18]. For example, seminal work from Hill and coworkers demonstrated enhanced analyte selectivity on an ambient pressure drift tube when varying the drift gas from low to high polarizability (helium, nitrogen, argon, and carbon dioxide) [19-22]. Eberlin and coworkers reported improved separation of isomeric haloanilines, carbohydrates, and petroleum constituents when operating a traveling wave instrument in CO<sub>2</sub> as opposed to conventional N<sub>2</sub>gas [23-26]. Recently, Yost and coworkers demonstrated increased resolving power for several isobaric steroids, analyzed in a reduced-pressure drift tube instrument, using CO<sub>2</sub> as the drift gas [27]. In many cases, however, the more polarizable drift gases (e.g., Ar, CO<sub>2</sub>, and N<sub>2</sub>O) are not reported to dramatically improve the overall IM peak capacity and resolution for chemically-similar analytes, including peptides, amino acids, structural isomers, and protein

conformers, in comparison to He or N<sub>2</sub> drift gases [28-31], A large survey of various drift gases (He, No, Ar, CO<sub>2</sub>, N<sub>2</sub>O, SF<sub>6</sub>) recently conducted on a commercial drift tube instrument by Fjeldsted and coworkers investigated the separation of isomeric carbohydrates, fluoroalkyl phosphazenes, and various small molecule pesticides. Higher resolution was generally observed for drift gases with access to the highest resolving powers in the instrumentation, namely conventional He and N<sub>2</sub>. Some enhanced selectivity and resolution was observed for the more polarizable drift gases (CO<sub>2</sub>, N<sub>2</sub>O, and SF<sub>6</sub>), although results were specific to the analyte pairs being investigated [32]. The literature examining the roles various drift gases play on IM separation capabilities are complicated and must be taken case-by-case. Given the potential analytical power that the drift gas composition can play in enhancing the IM resolution, there is a need for exploring IM separations in alternate drift gases across a broader range of analytes, and reporting quantitatively-comparable metrics of separation capabilities obtained from various laboratories. To accomplish this, a common framework for acquiring and comparing ion mobility data obtained in different drift gases needs to be established.

In addition to providing analyte structural information, collision cross section (CCS) values are useful for comparing IM measurements obtained using different IM instrumentation and techniques [1, 33-35] and recently have been used as an additional molecular description combined with accurate mass and retention time information for high confidence identification of unknowns arising from complex samples [34, 36-40]. While some fundamental work has been done using matrix-assisted laser desorption ionization (MALDI) [41, 42], the majority of CCS measurements has been conducted using electrospray ionization (ESI), including the work here. For alternate drift gas work, separations conducted across different instrument platforms, such as drift tube and traveling wave, cannot be directly compared without utilizing a normalized measurement such as CCS [33]. Comparisons of IM separation performance are hindered by the fact that the majority (95%) of CCS measurements have been reported only in conventional He and N<sub>2</sub> drift gas [1]. In addition to the general lack of reporting CCS measurements in alternate drift gas work, a large-scale study of the effects of drift gas on a wide mass range, multiple biological classes, and varying charge states has yet to be accomplished. In this work, we investigate four of the most commonly used drift gases, helium (He), nitrogen (N<sub>2</sub>), argon (Ar), and carbon dioxide (CO<sub>2</sub>), and establish instrument conditions necessary to achieve high reproducibility and enhanced separation efficiency for diverse sets of structurally-homogeneous analytes representing various charge states and chemical classes. We assess the separation efficiency in each gas using a combination of single-peak resolving power, two-peak resolution, and peak capacity, and provide detailed reporting of the CCS measured in each drift gas for low field DTIMS.

## Methods

#### Instrumentation

A commercial drift tube ion mobility-mass spectrometer (6560, Agilent Technologies) equipped with a thermally-assisted electrospray ionization source (Jet Stream, Agilent) was used for all measurements described herein [43]. The drift gases used were high purity N<sub>2</sub>,

He, Ar, and CO<sub>2</sub> (Air Liquide). N<sub>2</sub>, He, and Ar were UHP grade supplied at 99.999% purity, whereas CO<sub>2</sub> was Coleman grade, at 99.99% purity. The high-pressure funnel (HPF), trap funnel (TF), and drift tube (DT) regions of the instrument, shown in the schematic in Figure 1, were supplied with either He, N2, Ar, or CO2 using a commercially-upgraded drift gas manifold (Alternate Gas Kit, Agilent). Briefly, the gas kit consists of a precision closed-loop pressure controller (640B, MKS Instruments) which makes real-time adjustments to the pressures based on readings from a capacitance manometer (CDG-500, Agilent) mounted on the DT chamber. Prior to being directed into the instrument, each drift gas was passed through a passive purifier specific for He, Ar (RMSH-4, Agilent), N<sub>2</sub> (RMSN-4, Agilent), or CO<sub>2</sub> (P600-2, VICI Metronics). Instrument tuning utilized the vendor autotune function in standard mass range (m/z 3200 mode). For He, tuning was performed in N<sub>2</sub> before switching to He to mitigate issues related to gas breakdown at high voltage, manifesting as uncorrelated low m/z noise in the 2D ion mobility-mass spectrometry (IM-MS) spectrum. For  $N_2$ , Ar, and  $CO_2$ , tuning was performed in the chosen drift gas. During gas switching, the DT was allowed to equilibrate for at least 60 minutes, after which, pressures were confirmed to be stable and minor adjustments made as necessary. The HPF was operated at 4.80 Torr for N2 and Ar, 4.35 Torr for He, and 4.30 Torr for CO2. The TF was operated at 3.80 Torr for N<sub>2</sub>, 3.74 Torr for Ar, and 3.72 Torr for both CO<sub>2</sub> and He. CCS measurements utilize a stepped-field procedure in which several drift fields are surveyed in order to determine the ion transit time within the mobility drift region. For N<sub>2</sub>, Ar, and CO<sub>2</sub> the specific voltages used for the DT and associated ion transfer optics (TF exit, rear ion funnel, and post-IM hexapole) are identical to settings used in a standardized CCS method described in a recent interlaboratory study [34]. Specifically, this standardized, stepped-field method incorporates seven electric field strengths (E/N, where E is electric field and N is the molecular number density) between 8 and 15 Td with a static pressure of 3.95 Torr in the DT. For He, the same DT pressure (3.95 Torr) is utilized, however, a different range of electric fields are surveyed between 6 and 13 Td to optimize resolving power across a broad mass range while mitigating electrical breakdown of the gas. Detailed experimental settings for each drift gas are provided in Table S1 of the supporting information. The trap release time (referred to as ion gating time in other IM instrumentation) was set to 100 µs for all experiments to maximize instrument resolving power [44], and the TF exit voltage (defining the ion "injection" voltage into the DT) was kept at 10 V, minimizing end effects which otherwise shift drift times to lower values [45]. Ion source conditions were as follows: 325°C gas temperature, 13 L/min drying gas flow rate, 275°C and 12 L/min sheath gas temperature and flow rate, and nebulizer pressure at 20 psig. Additional experimental conditions for each drift gas are summarized in Table S1. The vendor-supplied software (LC/MS Data Acquisition B.07.00, Agilent) was used for all data acquisition.

#### Variable Pressure Experiments

In prior work with He drift gas, it was found that the operational pressures of the HPF and TF had an effect on the measured CCS, presumably due to changes in gas purity within the DT [35, 46]. In order to better characterize these pressure dependencies, the gas pressures were varied in increments of 0.10 Torr and 0.02 Torr for the HPF and TF regions, respectively, while keeping the DT at 3.95 Torr. This allowed CCS measurements to be obtained at multiple pressure combinations for both regions. The upper pressure limits

surveyed for N<sub>2</sub>, Ar, and CO<sub>2</sub> were based upon pressures utilized for normal N<sub>2</sub> operation (i.e., 4.80 Torr in the HPF and 3.80 Torr in the TF) [32]. Operational limits of the flow controller (calibrated for N<sub>2</sub>) restricted helium experiments to a maximum of 4.63 Torr in the HPF with 3.80 Torr in the TF. To obtain CCS values comparable to those reported in the literature, measurements were taken above the minimum operational pressure bounds of the HPF and TF; however, despite the proximity to operational pressure limits, CCS measurements were reproducible and subsequently represent the values reported in this work.

#### Effective Length

All CCS values observed in all drift gases are listed in Table S2, with CCS values not observed in all drift gases listed in Table S3, and calculated using a drift tube effective length of 78.06 cm. The effective length differs slightly from the geometric length of the instrument (ca. 78.12 cm), and is determined using a previously-reported protocol whereby CCS values for the MS tuning mixture (hexakis(fluoroalkoxy)phosphazenes, HFAP, *m/z* 322-*m/z* 2722) are scaled to reference measurements obtained using a gridded DT instrument [34, 47]. The effective length determined from this procedure is instrument-specific and was used to calculate CCS values for all drift gases.

#### **Chemical Standards**

Isopropanol, acetonitrile, water, and formic acid (Optima LC-MS grade) were obtained from Fisher Scientific. The chemical standards were purchased from several vendors, summarized in Table S4. Tetraalkylammonium salts (TAA) were received as dry powder and reconstituted in isopropanol. All chemical standards except the HFAP tuning mixture were prepared at 10  $\mu$ g/mL and directly infused into the ESI source at a flow rate of 10  $\mu$ L/min. Poly-DL-alanine and maltose standards were prepared in methanol:water (50:50% v:v) with 0.1% formic acid. The MS tuning mixture containing HFAP (ESI-L Low Concentration Tuning Mixture, Agilent) is supplied from the vendor as a mixture of components dissolved in acetonitrile:water (95:5% v:v), and was prepared as per the vendor instructions by diluting the solution by a factor of ten using acetonitrile:water (98:2% v:v).

## **Results and Discussion**

#### Analyte Selection

The analytes chosen for this study were selected as belonging to sets of structurallyhomogeneous molecules spanning a wide mass coverage (Figure 2a). Specifically, HFAP increases by symmetric additions of fluoroalkyl (CF<sub>2</sub>) subunits to the six terminal ethers, while TAA cations increase symmetrically by alkyl (CH<sub>2</sub>) subunits, and the carbohydrates represent oligosaccharides with repeating  $\beta$ -D-glucose units linked by an  $\alpha(1 \rightarrow 4)$ glycosidic bond. Although the MS tuning mixture also contains m/z 118 (betaine), this molecule is not a HFAP, and thus was not included in this study. Poly-DL-alanine has been used in previous IM studies for CCS calibration [48, 49] and, in addition to possessing repeating alanine subunits, forms multiple charge states from ESI, allowing for the investigation of charge state effects in different drift gases.

#### **Pressure Effects**

Pressure conditions in the ion transfer funnels prior to the drift tube were systematically investigated for each drift gas to determine the effect on the measured CCS. The cesium cation ( $Cs^+$ ) was used for quantitative comparisons of the CCS measurement due to ease of ionization, lack of multiple conformations, and CCS measurements have been previously documented in the literature for the drift gases investigated in this present study. The Cs<sup>+</sup> CCS results for select combinations of pressure conditions are presented in Figure 3, representing the upper (red boxes) and lower limits of pressures (blue boxes) surveyed for each funnel region. In all cases, the drift tube pressure was maintained at 3.95 Torr. The CCS value and corresponding variance reported in the literature for Cs<sup>+</sup> is also shown at the bottom of each panel (unfilled bars), as reported for He, N2, Ar, and CO2 [50]. While there was no significant change in the N2 CCS (DTCCSN2) values for Cs<sup>+</sup> (0.08% difference or less), CCS measurements in the other drift gases were sensitive to HPF and TF pressure changes, most notably with changes to the TF pressure, where as high as 47.96%, 7.99%, 0.81% differences in CCS were observed for He, CO<sub>2</sub>, and Ar respectively. Previously, it has been suggested that for this particular instrument configuration, the pressure difference between the TF and the DT affects the gas purity within the IM region, and our results support this idea, with lower pressure differences shifting the Cs<sup>+ DT</sup>CCS<sub>He</sub> to higher values, and lower CCS values for the other gases, implicating contamination of the drift tube with gas from the ion source, the latter of which is N2 sourced from the boil-off of a cryogenic tank [32, 46]. Though not reported previously, it was found in this study that lower pressures in the upstream HPF could also improve gas purity in the DT, and this effect is shown in Figure 3 for He, Ar, and CO<sub>2</sub>. As with the TF pressure, lowering the HPF pressure affected the Cs<sup>+</sup> CCS measurement by shifting it to lower values for He, and higher values for Ar and CO<sub>2</sub>. In He, the CCS continues to decrease as the HPF and TF pressures are reduced, with the lowest pressures surveyed exhibiting the closest CCS correlation to the literature. The CCS response to HPF and TF pressures indicate improved drift gas purity at lower HPF and TF pressures for drift gases other than nitrogen. At some pressures, the CCS response is minimal as the pressure is further lowered. Since the interpretation is that lower HPF and TF pressures show increased gas purity, as evidenced by better CCS agreement with literature values, the lowest pressures that could be achieved in He, Ar, and CO<sub>2</sub>, without signal degradation of the reported analytes, are recommended. This CCS shift exhibits some massdependence, as evidenced by comparing the CCS of the HFAP tuning mixture ions across various pressure conditions (cf., Figure S1). Specifically, the lower mass HFAP ions (m/z322 and m/z 622) experience a greater shift in CCS in response to changes in HPF and TF operational pressures, which we interpret as due to a greater magnitude of the so-called "end effect" documented for DTIMS, where ions subject to either higher fields or less collisional dampening (i.e., He) at the entrance to the IM region experience deeper penetration into the drift region [45, 51]. This delayed thermalization of ions results in a shorter effective drift length, and thus a shift in the apparent measured CCS to lower values as the TF operational pressure is decreased. For the lower limits of pressure, operating the TF and HPF regions at pressures below those recommended here was found to significantly degrade ion transmission in He, Ar, and CO<sub>2</sub>. Additionally for Ar, lower pressures resulted in non-linear behavior in the drift times obtained across multiple drift fields, while in He, gas discharge

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occurs, which manifests as uncorrelated spectral noise in the lower m/z range. The lower

pressure limits for  $N_2$  were chosen to overlap with the limits observed for the other drift gases. The high pressure limits for operating the TF and HPF regions for each drift gas were determined based on the recommended values previously established for  $N_2$  [32]. Recommended pressure settings for all drift gases evaluated are summarized in Table S1.

#### Ion Source Temperature Effects

CCS changes were also observed when adjusting the ion source gas temperatures (i.e., nebulizer sheath gas and source entrance drying gas). Specifically, when varying ion source gas temperatures, as much as a 2.7% change in the  $^{DT}CCS_{He}$  was observed, while the  $^{DT}CCS_{N2}$  shifted by a more modest amount, 0.3% or less. This result is interpreted as being due to changes in the gas flow dynamics between the ion source and the DT, which results in gas impurities being introduced into the drift region despite efforts to maintain constant operational pressures in the HPF, TF, and DT regions. Thus, the CCS results presented in this study are obtained under constant ion source temperature conditions which were selected in order to achieve close correlation of Cs<sup>+</sup> CCS values to those described in the literature.

#### Ion Injection Potential

The TF exit voltage establishes the potential bias between the TF and the DT. This potential affects the ion "injection" into the DT and plays a role in ion transmission, resolving power, ion internal energy, and the measured drift time.[45, 52, 53] In order to quantitatively evaluate these effects, the potential difference between the TF exit and DT was varied from 0 to 45V (35V in He) using the HFAP tuning mixture ions as a test system. Ion transmission, CCS, and resolving power were measured for each injection voltage (Figures S2-S4). It was found that in all drift gases, ion injection voltages between 5 and 15V yielded the highest ion transmissions based on peak area analysis (Figure S2), whereas resolving powers were found to be the highest for 10 and 15V. In this range, no strong correlation was observed for any particular HFAP ion, suggesting that between 10 and 15V the mass and mobilitydependent effects were minimal. Importantly, the ion injection voltage was found to shift the measured CCS for all gases surveyed, most dramatically for the low (0 and 5V) and high (15V) potential settings. For example,  $^{\text{DT}}\text{CCS}_{N2}$  differences relative to 10V for m/z 1522 was found to be between 2.13% and 1.24%, respectively. Specifically, low ion introduction voltages (5V) yielded high CCS values, whereas the high ion introduction voltages (15V) resulted in low CCS values regardless of the identity of the drift gas. These results can be interpreted as a consequence of different ion injection energies, specifically, for low voltages, ions spend more time at the DT entrance prior to being entrained in the drift field, whereas for high voltages, ions penetrate deeper into the drift region before reaching a steady-state drift velocity. It is somewhat surprising that, without an ion introduction potential (0V), ions still transfer against a pressure gradient into the DT, although this observation is not investigated further in this work. Also of note is that, despite what is typically observed for drift tube instruments, the higher ion injection potentials in this DTIMS configuration did not yield more ion signal, but rather decreased ion transmission and resolving power. The default setting for the injection potential (designated as the TF exit voltages in the software) is 10V, which yielded good results in terms of high ion

transmission and high resolving power for all gases and HFAP ions evaluated. As such, the TF exit voltage used for all subsequent experiments was 10V.

#### **CCS Measurements**

The CCS measurement reproducibility (precision) was found to depend on the drift gas. For example, 95% of the <sup>DT</sup>CCS<sub>He</sub> values exhibit an RSD of less than 1.4%, whereas in Ar, CO<sub>2</sub>, and N<sub>2</sub>, the precision was found to be much better, with the majority (95%) of CCS values exhibiting less than 0.6%, 0.5%, and 0.3% RSD, respectively (c.f., Figure S5). The cause of the lower precision in He is likely due to pressure fluctuations during CCS measurements as a result of the pressure controller valve orifice used in this work being optimized for N<sub>2</sub>, although it should be noted here that the vendor offers high flow valve configurations which should improve the pressure control when operating in He drift gas. In order to keep the recommended parameters in this work accessible to a wider network of researchers, an alternate pressure controller optimized for He was not used. Instead, the pressure controller shipped with the alternate gas kit was used for all four drift gases. As noted previously, both the HPF and TF pressures affect the CCS and operating the instrument at the lower HPF and TF pressures yield measurements which are closely aligned with canonical CCS values. Note that Cs<sup>+</sup> in He never matched with the literature value, even accounting for error. As only a single literature source could be found which reported <sup>DT</sup>CCS<sub>He</sub> values for the cesium cation, this observation could not be further evaluated. While HFAP ion CCS values have not been previously published for gases other than N2, HFAP ions were subsequently compared to previously published <sup>DT</sup>CCS<sub>N2</sub> values (Figure S6) [34, 54], demonstrating excellent agreements of 0.56% or less, with a relative standard deviation of measurements taken for this study of 0.31% or less. For poly-DL-alanine, CCS values obtained in this work showed less agreement with the literature, with a maximum bias of 1.7% observed for  $N_2$  and a 4.8% bias for He (Figure S7). This agreement is still seen as reasonable, given these measurements were from disparate instrument designs and the error reported for the literature values was estimated at 3% [34, 54]. Previously published He and N2 values for TAA cations [34, 54] were also compared to the current measurements (Figure S8), and it was found that, in general, the literature values were typically higher, with differences ranging from -2.96% to 3.67% observed among the different reports. Overall, the settings recommended in this study (Table S1) provide good CCS measurement reproducibility while corresponding reasonably well to prior measurements. Ongoing international efforts are aimed at further improvements in CCS measurement precision and accuracy.

#### **Chemical Class Behavior**

The gas-specific conformational space plots for all analytes investigated is shown in Figure 4. The low mass region is highlighted in each of the insets and depicts a region of IM-MS space where multiple chemical classes reside in close proximity. In general, all of the class-specific mobility-mass trends are qualitatively the same with the higher charge state ions exhibiting larger CCS values than the lower charge state ions, while for singly-charged ions, the CCS increases in the following order: HFAP<carbohydrates<peptides<TAA cations. As discussed in a number of previous reports, this conformational ordering is indicative of the relative gas-phase packing efficiency of each chemical class, with ions exhibiting less-

restrictive degrees of freedom (peptides and carbohydrates) are able to arrange themselves into more compact gas-phase structures [55, 56]. For singly-charged ions, the relative CCS values in  $N_2$  and Ar are similar, whereas  ${}^{DT}CCS_{Ar}$  values are all slightly lower than those measured in N<sub>2</sub>. Note that for multiply charged species, the CCS differences are greater between the various gases than singly charged species. This is apparent in Figure 4 when comparing the relative spacing between the poly-DL-alanine +2 and +3 ion trends (dark blue triangles) to the +1 trend (light blue triangles), as well as plotting CCS as a function of gas polarizability (Figure S9), where linear fits to the higher charge states exhibit larger slopes. In addition to charge state effects, the carbohydrates as a whole are strongly affected by the identity of the drift gas compared to the other classes investigated. This effect is shown in Figure 5, where the carbohydrates shift in drift time to a greater extent than other classes, resulting in an inversion in the IM elution orders. For example, in CO<sub>2</sub>, polyalanine 13-mer (blue dashed trace, m/2 942.5008) exhibits a lower measured drift time than the carbohydrate, maltohexaose (orange dashed trace, m/z 1013.3173). In all other drift gases, however, this relative elution order is reversed. This effect has been noted previously for small molecules with masses below 500 Da [19, 57] and this present work demonstrates that CCS inversion can occur at higher masses as well. This drift gas selectivity can be useful depending upon the ion species being separated. For instance, in the peak-to-peak resolution (Rpp) example shown in Figure 5, utilizing N2 or Ar would be ideal for separation of the lower mass ions shown, however, the best resolution for certain ion pairs is observed in CO<sub>2</sub> specifically between HFAP m/z 322 and maltose (labeled as analytes A and B). It is interesting to note that, with respect to the relative elution orders and resolutions observed, IM separations in Ar (39.96 Da, 1.64  $Å^3$ ) are most similar to those observed in N<sub>2</sub> (28.01 Da, 1.74 Å<sup>3</sup>), rather than CO<sub>2</sub> (43.99 Da, 2.91 Å<sup>3</sup>), [58] indicating CCS is more dependent on the polarizability of the drift gas than the neutral drift gas mass for the species measured here. While the experimental CCS is a function of both ion and drift gas size, the relative size of the drift gas as experienced by the ion is dominated by its polarizability, and thus plots of CCS versus polarization exhibit near linear trends [22]. This is shown in Figure 6a, with the select ions investigated plotted as a function of the drift gas polarizability. Of note is that, despite the relative differences in CCS observed across the different drift gases surveyed, ions of similar chemical class exhibit similar slopes in the CCS vs. polarizability plot.

#### **Drift Gas Selection Criteria**

In choosing the appropriate drift gas, analytical figures-of-merit such as single-peak resolving power ( $R_p$ ), two-peak resolution ( $R_{pp}$ ), and peak capacity are important considerations. Equations associated with these analytical figures-of-merit can be found in the supporting information. Peak capacity, a measure of the number of peaks that can occupy a given analytical separation space at half height, was found to be the highest in Ar, followed by N<sub>2</sub> at the lower mass range (Figure 6b, top panel), whereas N<sub>2</sub> outper-formed all of the other gases for higher mass analytes due to the ability for N<sub>2</sub> to access higher resolving powers in this mass range (Figure 6b, bottom panel). CO<sub>2</sub> exhibited the lowest peak capacities for the two mass ranges investigated, but certain systems may benefit from the increased analytical selectivity as previously described. Overall, N<sub>2</sub> would be recommended in most cases due to low %RSD, excellent peak capacity, and a consistently high resolving

power over a wide mass range. Helium may be chosen when comparison to computational or literature values is needed. Argon performed similarly to nitrogen but may be chosen over nitrogen if a slight increase in resolving power is needed for lower mass analytes.  $CO_2$  may be chosen to improve separation for specific analyte pairs.

## Conclusion

This work evaluates detailed experimental conditions necessary for operating a commercial drift tube instrument (Agilent 6560) with a variety of drift gases that include He, N<sub>2</sub>, Ar, and CO<sub>2</sub>. Criteria used to evaluate the optimal instrument settings needed for each gas include those parameters which provided the highest measurement repeatability of the CCS and the closest correspondence of CCS (e.g., cesium, tuning mixture ions, and poly-DL-alanines) to previously published values. Using these optimal operational settings, a large and highlyconsistent set of CCS measurements (N=280, 56 unique ions with 5 replicates each) for a variety of structurally-heterogeneous compounds (TAA cations, poly-DL-alanines, fluoroalkyl phosphazenes, and  $\alpha(1\rightarrow 4)$ -linked glucose oligosaccharides) are subsequently compiled, which are suitable as reference values for future studies. Comparison between the CCS response of the various classes, charge states, and masses surveyed in each drift gas is provided to aid future investigators in selecting the appropriate drift gas for their particular applications. The specific CCS trends observed were found to correlate strongly with the drift gas polarizability, thus providing a basis for predicting relative elution orders for gases not yet investigated. Overall, it was found that N<sub>2</sub> exhibited the highest resolving powers and peak capacities across a broad range of masses with a correspondingly high CCS reproducibility (less than 0.40% RSD). CO<sub>2</sub> exhibited relatively low peak capacities, but demonstrated high chemical class and charge state specific selectivity, and in some cases, resulted in inverted analyte elution times as compared to the other drift gases. In general, the reference CCS values and recommended experimental parameters outlined in this manuscript are expected to aid IM researchers in selecting the drift gas appropriate for their analytical problem, as well as providing guidance for incorporating alternative drift gases in future work. Finally, the highly-consistent CCS values measured in this study can be used as a basis for improving fundamental IM theory and developing algorithms which consider the identity of the drift gas in predicting IM elution behavior.

## **Supplementary Material**

Refer to Web version on PubMed Central for supplementary material.

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#### Figure 1.

A conceptual schematic of the commercial drift tube ion mobility-mass spectrometer (6560, Agilent) used in this work. The selected drift gas is supplied to the high-pressure funnel (HPF), trap funnel (TF), and drift tube (DT) regions. The flow controller adjusts the gas flow into TF region based on the DT pressure readback, maintaining a constant pressure in the DT. The DT is isolated from the vacuum system, and gas evacuation from DT is thus facilitated via the entrance and exit apertures of the chamber



## Figure 2.

(a) Mass coverage of each class of molecules investigated. (b) Representative structures for each of the four classes of molecules investigated in this study, with their respective masses indicated by the white dots in panel a



#### Figure 3.

A comparison of the CCS values obtained for Cs<sup>+</sup> (132.91 Da) in multiple drift gases (filled bars) and at various HPF and TF operational pressures (vertical axes labels, in Torr) along with the CCS and corresponding error (outlined bars) reported by Ellis et al. (Ellis, HW; McDaniel, EW; Albritton, DL; Viehland, LA; Lin, SL; Mason, EA; Atomic Data and Nuclear Data Tables 22, 179-217, 1978). Error bars for data obtained in this study correspond to 7 repeat measurements. Red and blue shading in the vertical axes labels indicates high or low relative pressure conditions for HPF and TF



#### Figure 4.

Mass versus CCS plots for each drift gas investigated. The location of all measurements are designated with grey symbols for comparison. The inset in each panel shows an expanded region at low m/z which contains various compound classes. Note that the CCS of carbohydrates (orange squares) relative to the other compounds exhibits a strong dependency on the identity of the drift gas utilized

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## Figure 5.

Relative abundance (R.A.) vs drift time for two ions from each class are shown for the four drift gases at the same field strength. Solid lines show four ions of similar mass, with one ion from each of the four classes (structures are depicted in Figure 2b); dashed lines show four ions of higher mass, with one ion from each class. (b) Two peak resolution ( $R_{pp}$ ) is shown as histogram plots for each closest-mass ion pair of the lower mass group, those depicted with solid lines in the IM spectra. Note the drift time inversion for the high mass carbohydrate (dashed orange) and polyalanine (dashed blue) in CO<sub>2</sub> compared to the other three drift gases.

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#### Figure 6.

(a) The CCS for select ions plotted against the polarizability of the drift gas as given in the CRC Handbook of Chemistry and Physics [45]. All ions are singly charged except polyalanine 14-mer (PAla 14) which is doubly charged. Increasing the polarizability of the drift gas is correlated with an increase in CCS. However, the relative change in CCS with the drift gas polarizability is both class and charge-state specific, as indicated by the different slopes. The slope, intercept, and regression coefficient values are given in Table S5. Error bars are within the marker size. (b) Resolving power ( $R_p$ ) and peak capacity bar graphs for each drift gas. Maximum  $R_p$  values are shown, requiring a different drift field for He (8.3 V/cm) versus the other drift gases (16.0 V/cm). Peak capacity was calculated between HFAP and TAA cations of similar mass. Error bars were based on five measurements conducted on different days.